

# The phase diagrams of the high temperature superconductors

Talk online: [sachdev.physics.harvard.edu](http://sachdev.physics.harvard.edu)

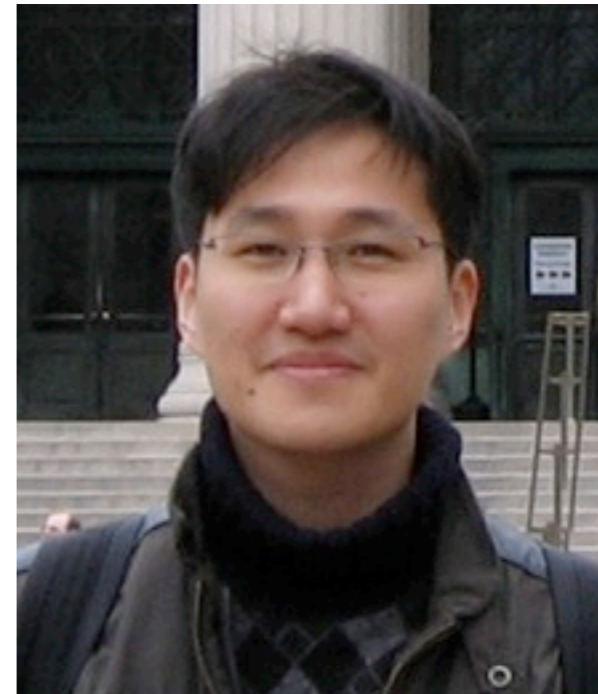
PHYSICS



HARVARD



Max Metlitski, Harvard



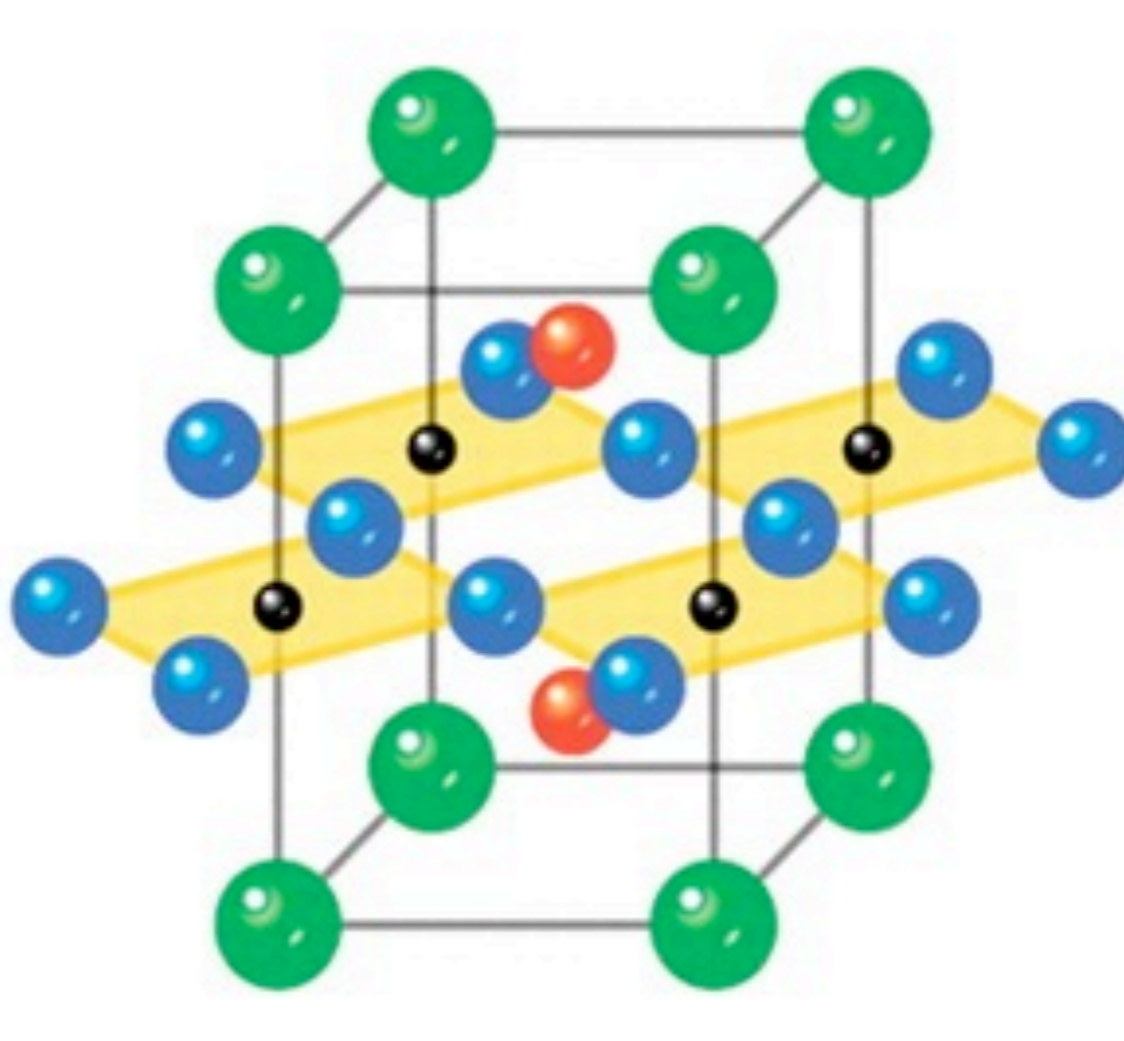
Eun Gook Moon, Harvard



# *The cuprate superconductors*

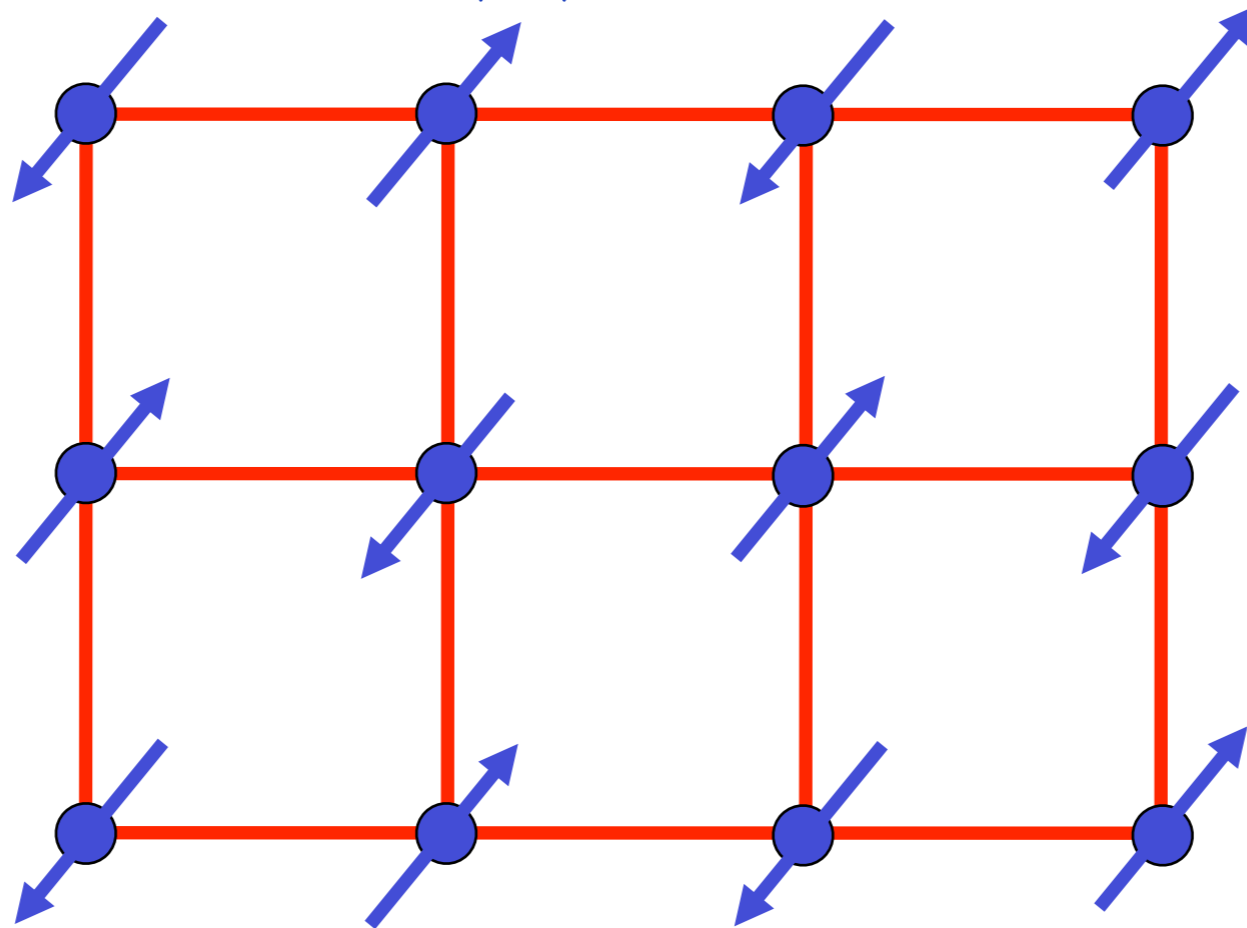
Na-CCOC

- Cu
- Ca/Na
- O
- Cl



# Square lattice antiferromagnet

$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$



Ground state has long-range Néel order

Order parameter is a single vector field  $\vec{\varphi} = \eta_i \vec{S}_i$

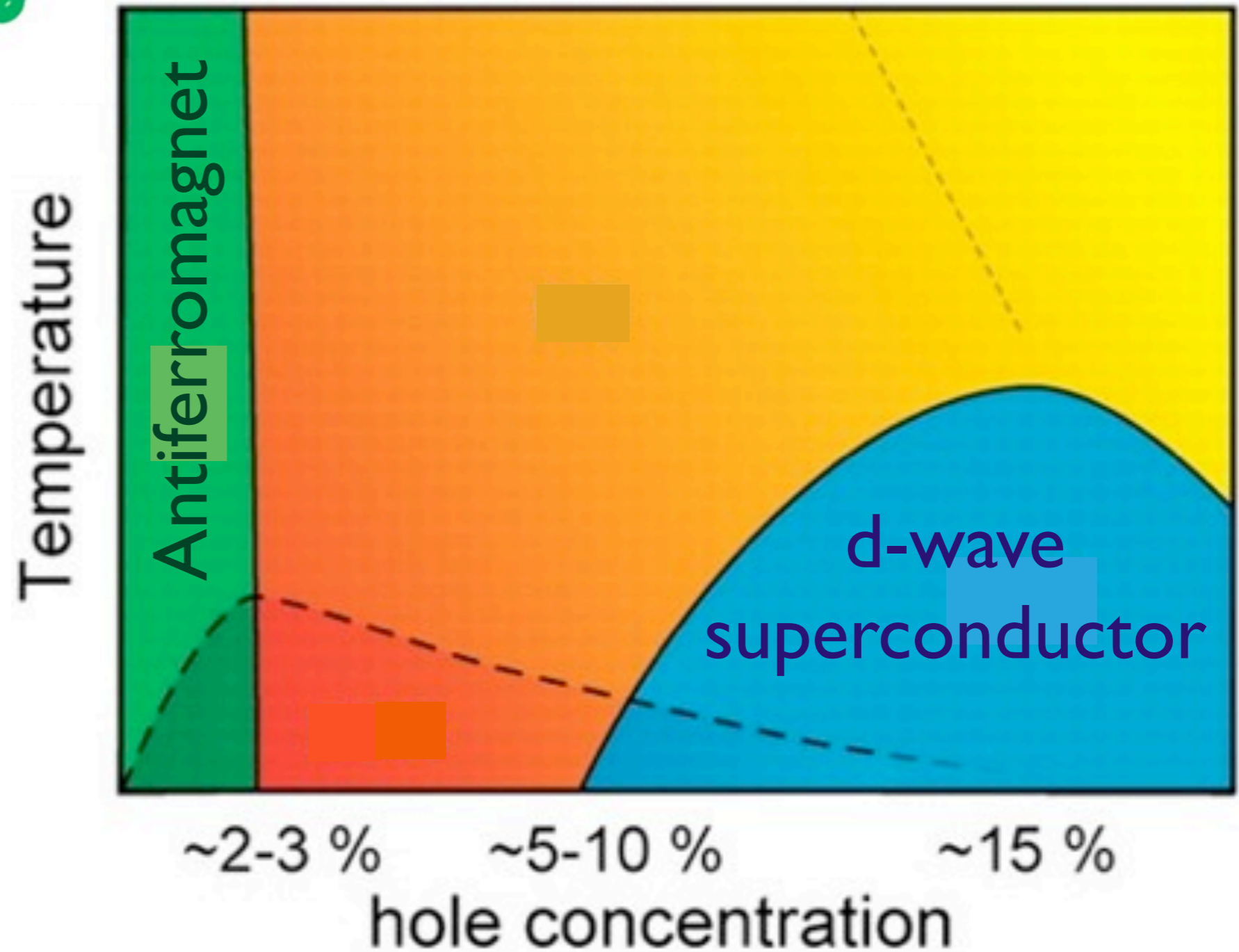
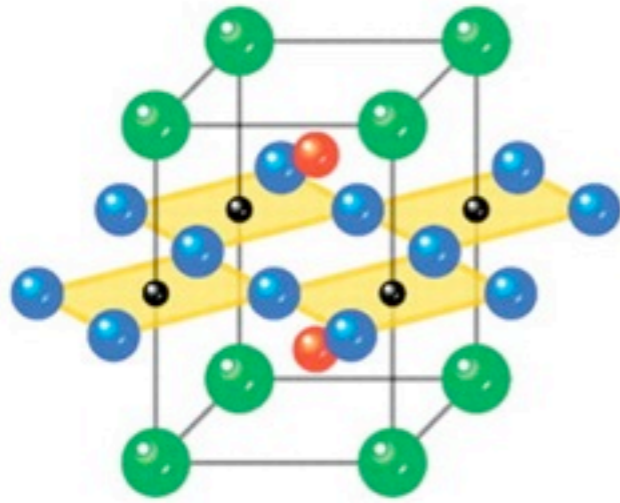
$\eta_i = \pm 1$  on two sublattices

$\langle \vec{\varphi} \rangle \neq 0$  in Néel state.

# The cuprate superconductors

Na-CCOC

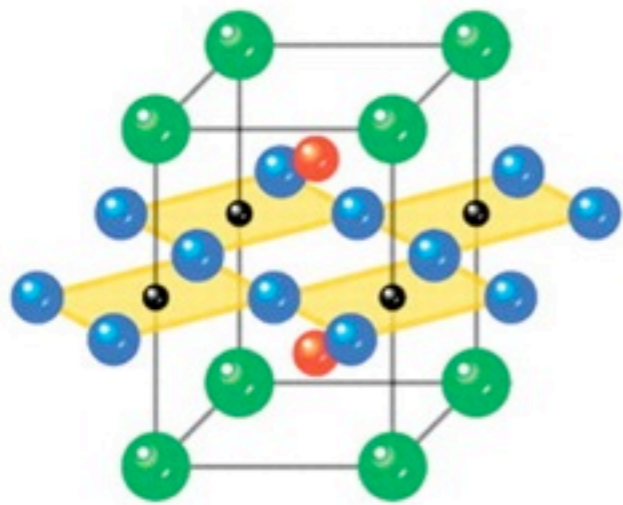
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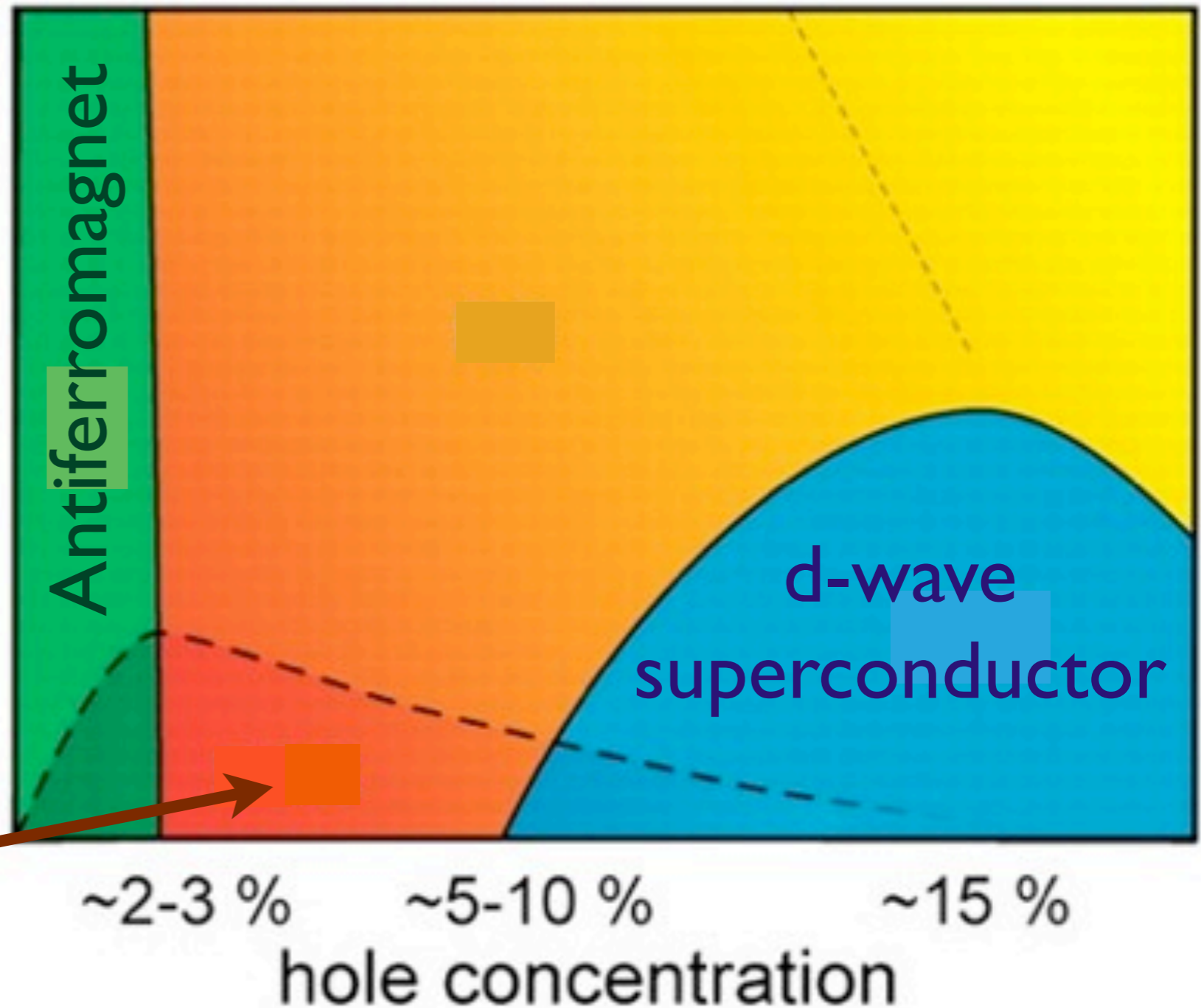
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Temperature

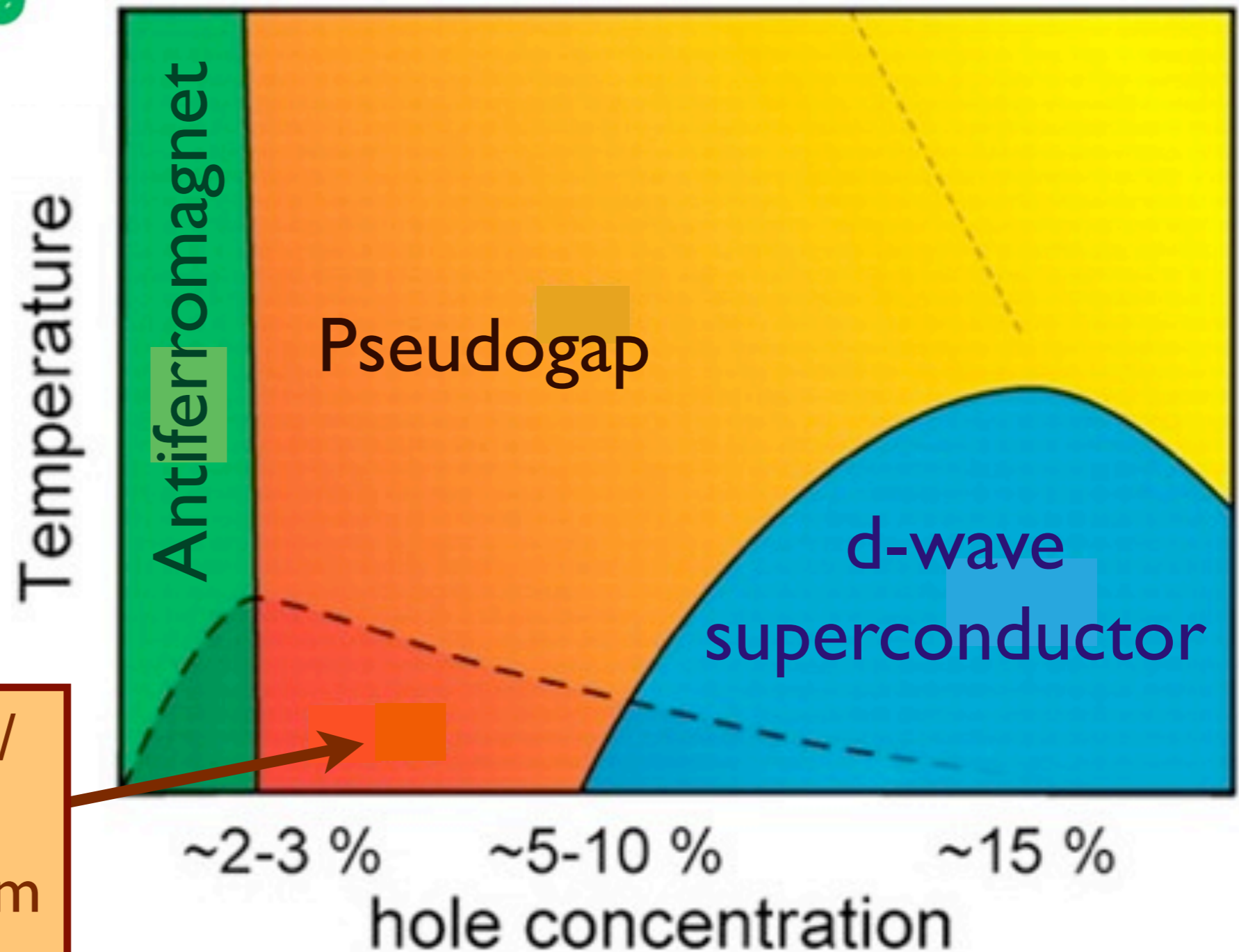
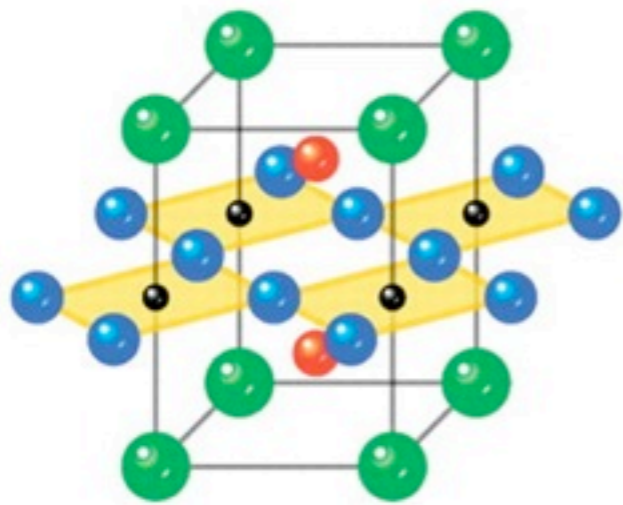


Incommensurate/  
disordered  
antiferromagnetism  
and charge order

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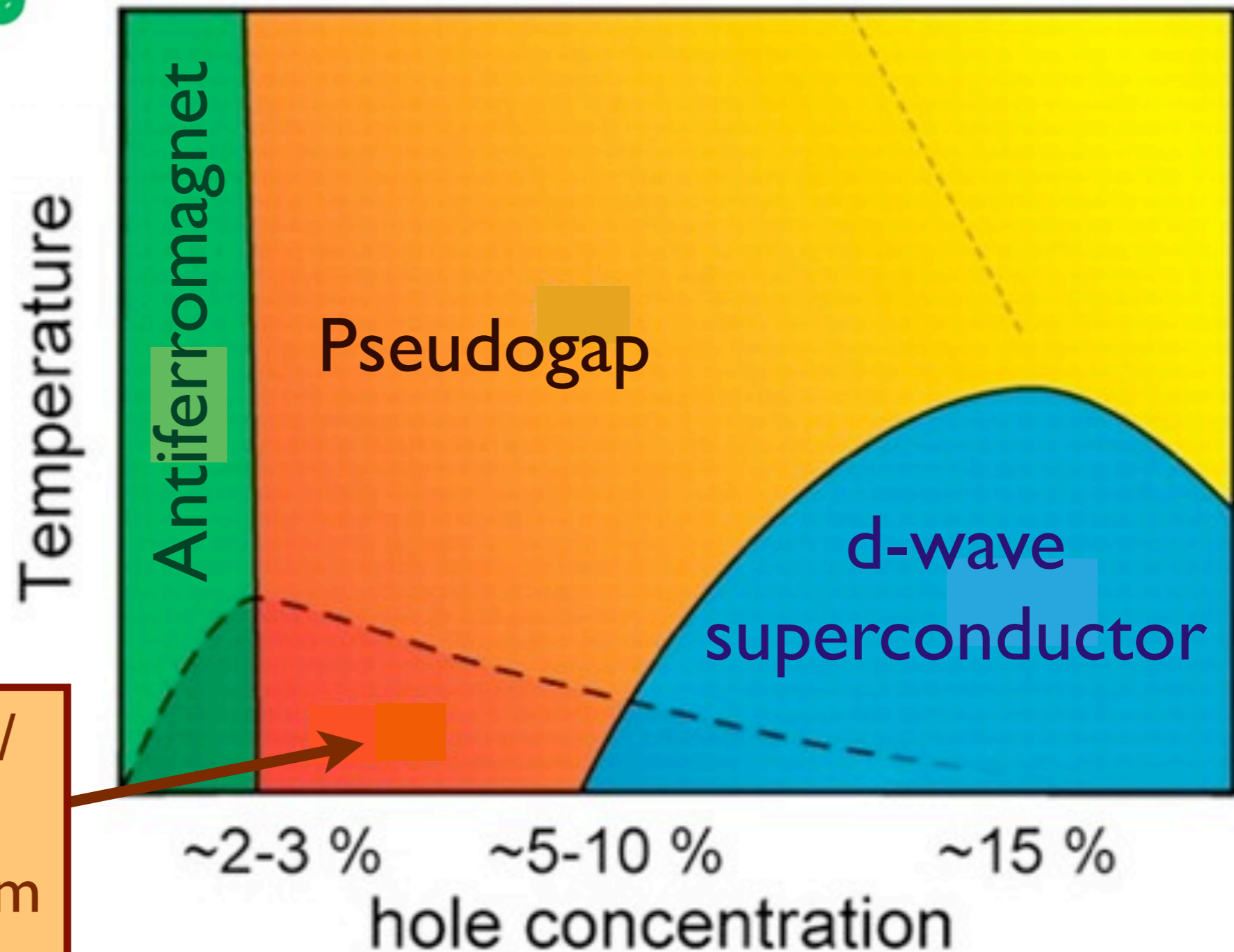
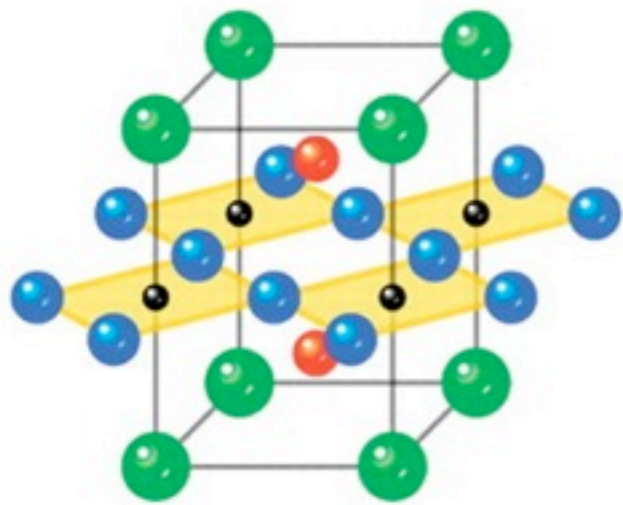


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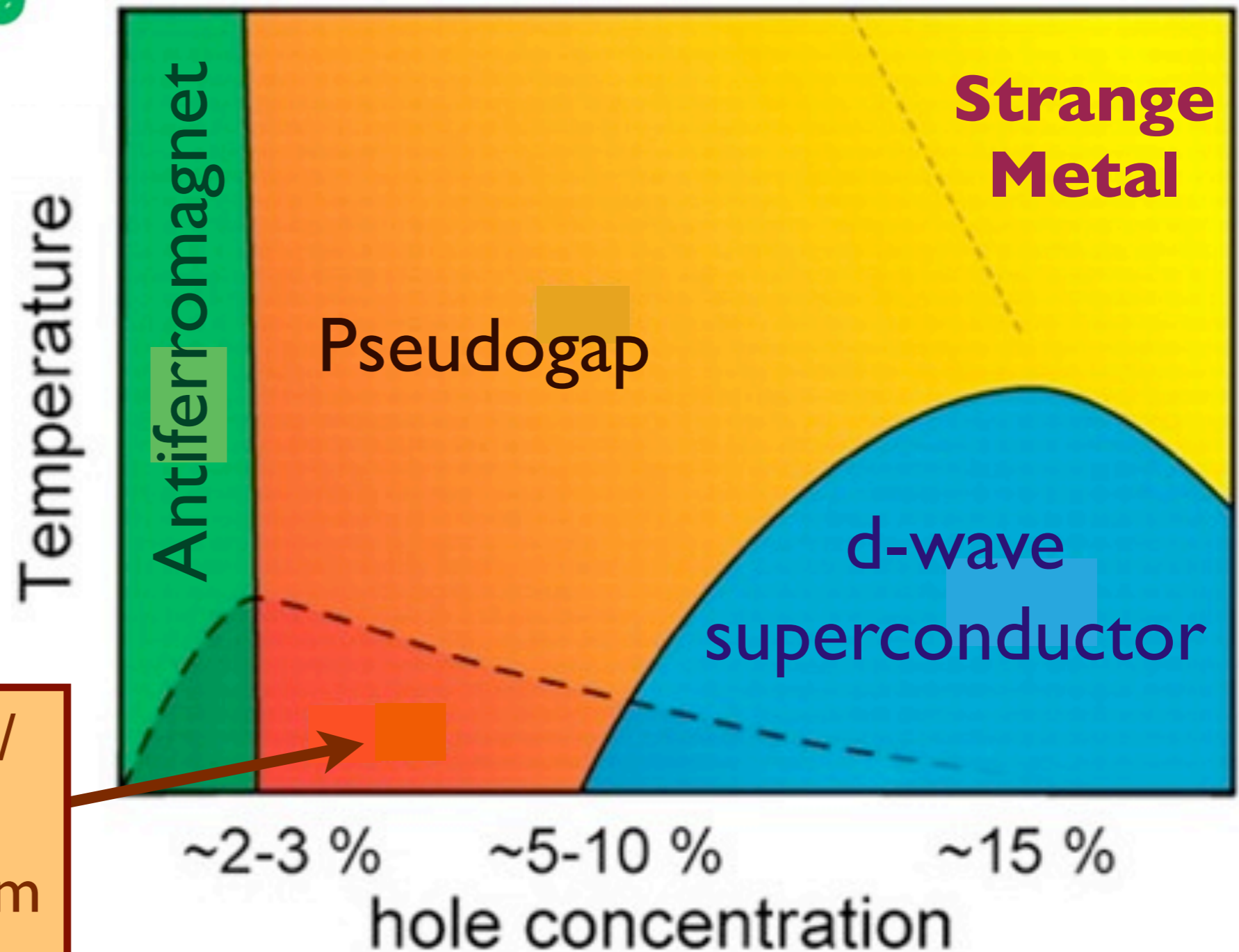
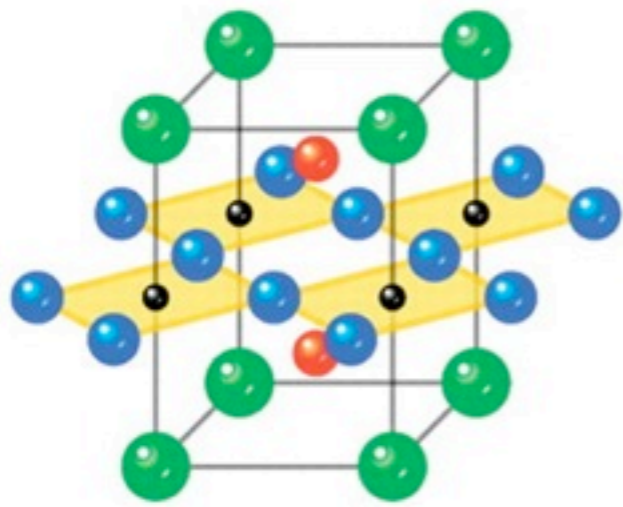
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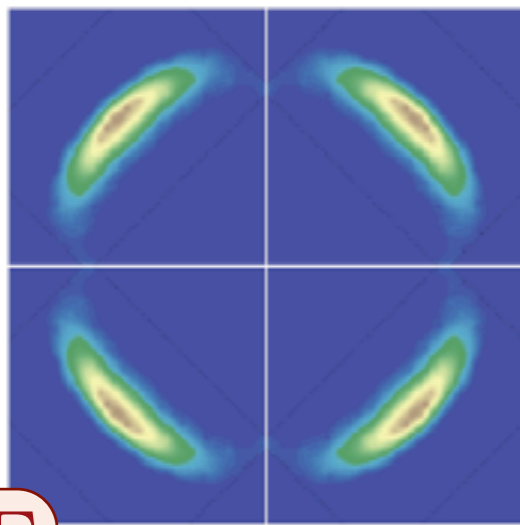
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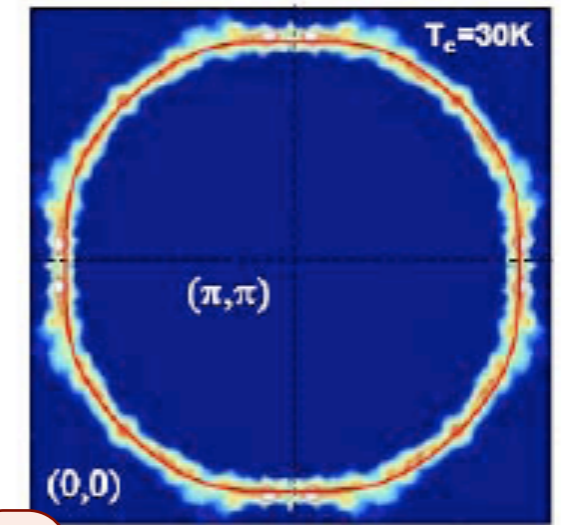
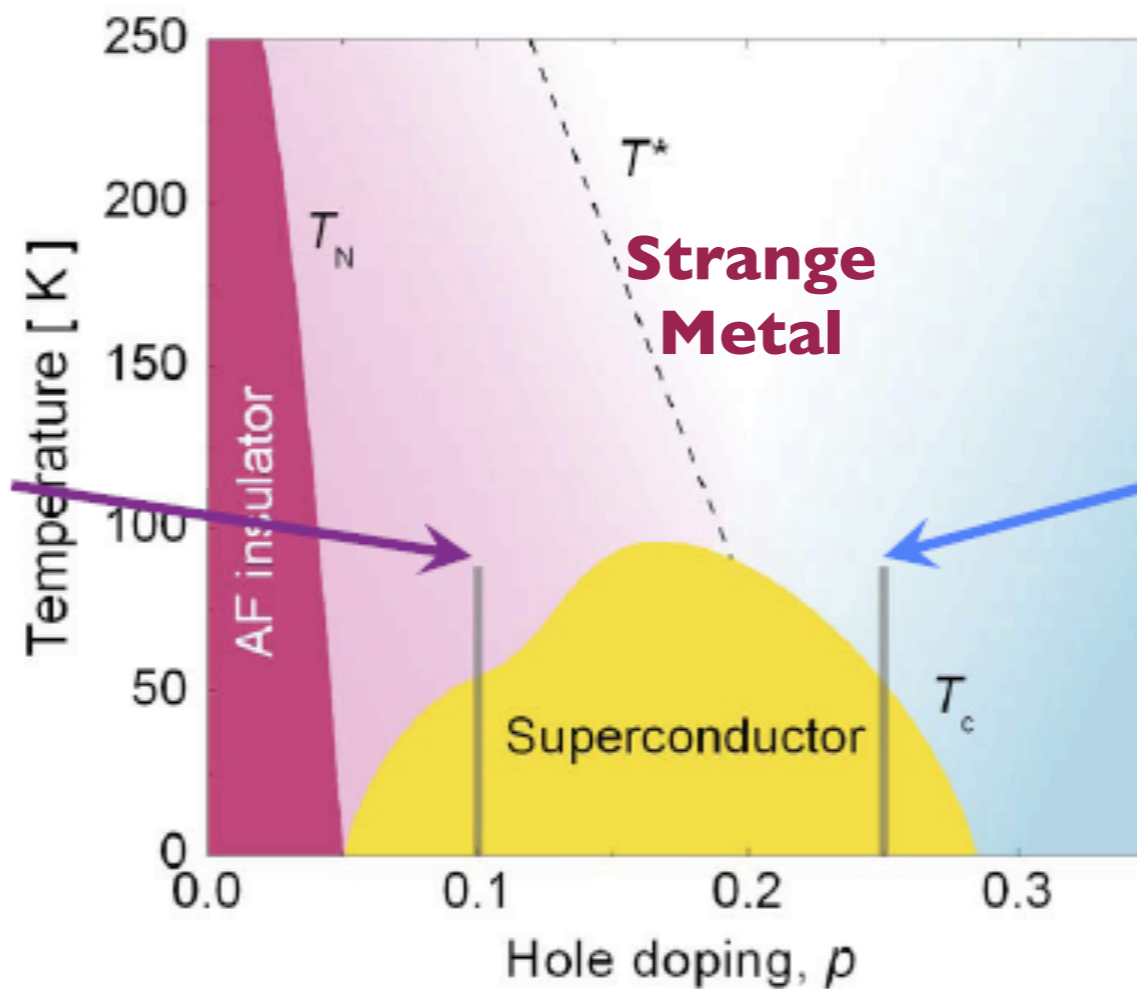


Incommensurate/  
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$\Gamma$

*K.M. Shen et al., Science 2005*



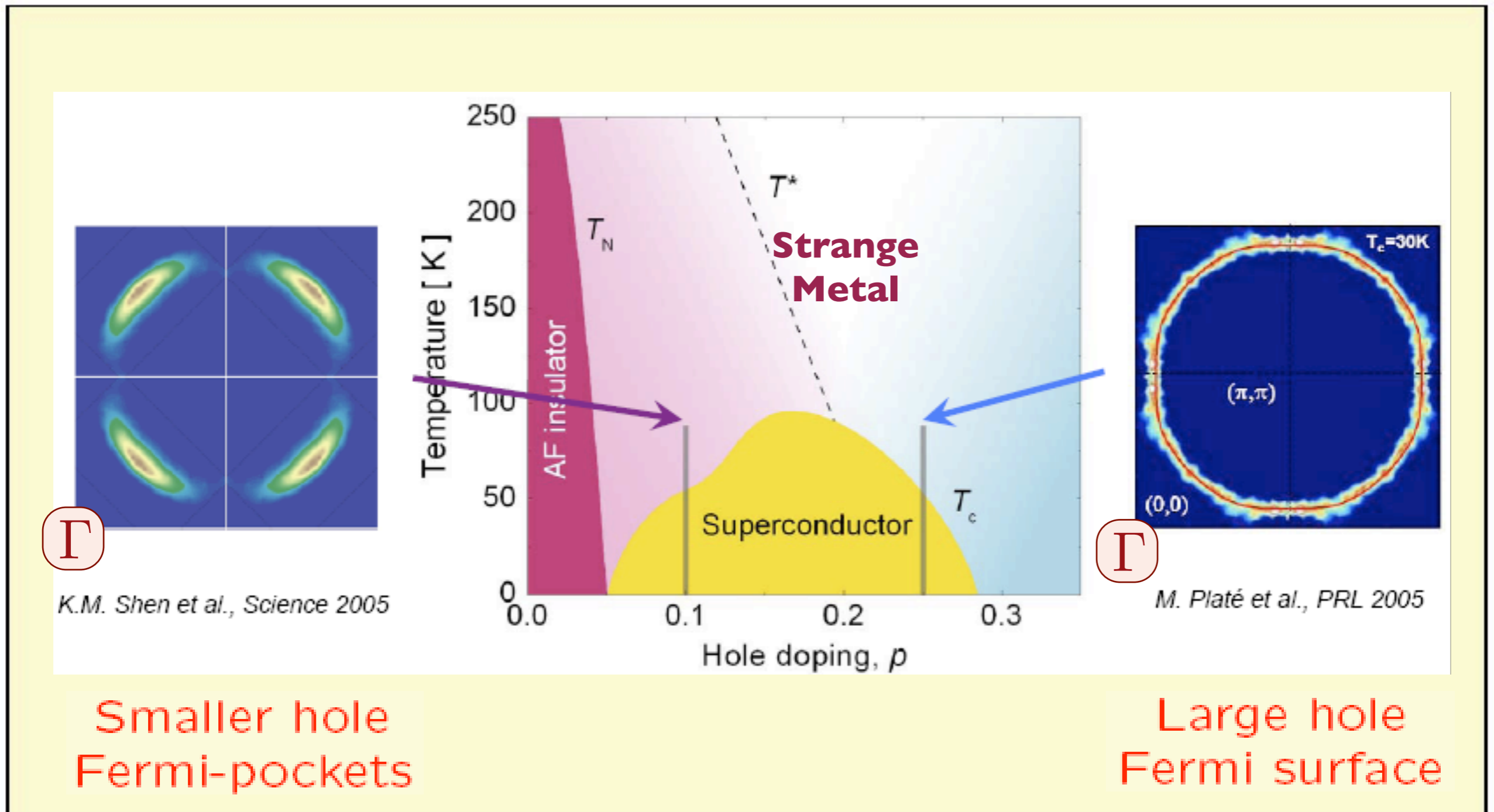
$\Gamma$

*M. Platé et al., PRL 2005*

Smaller hole  
Fermi-pockets

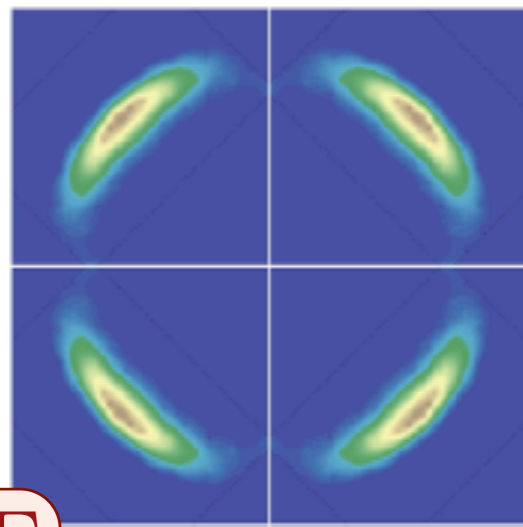
Large hole  
Fermi surface

# Key Ingredients:



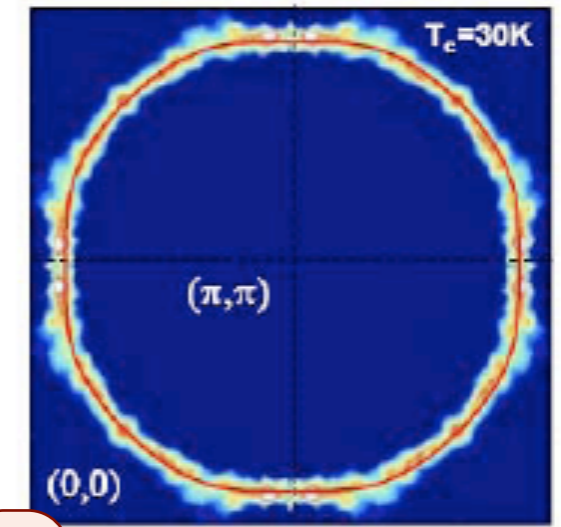
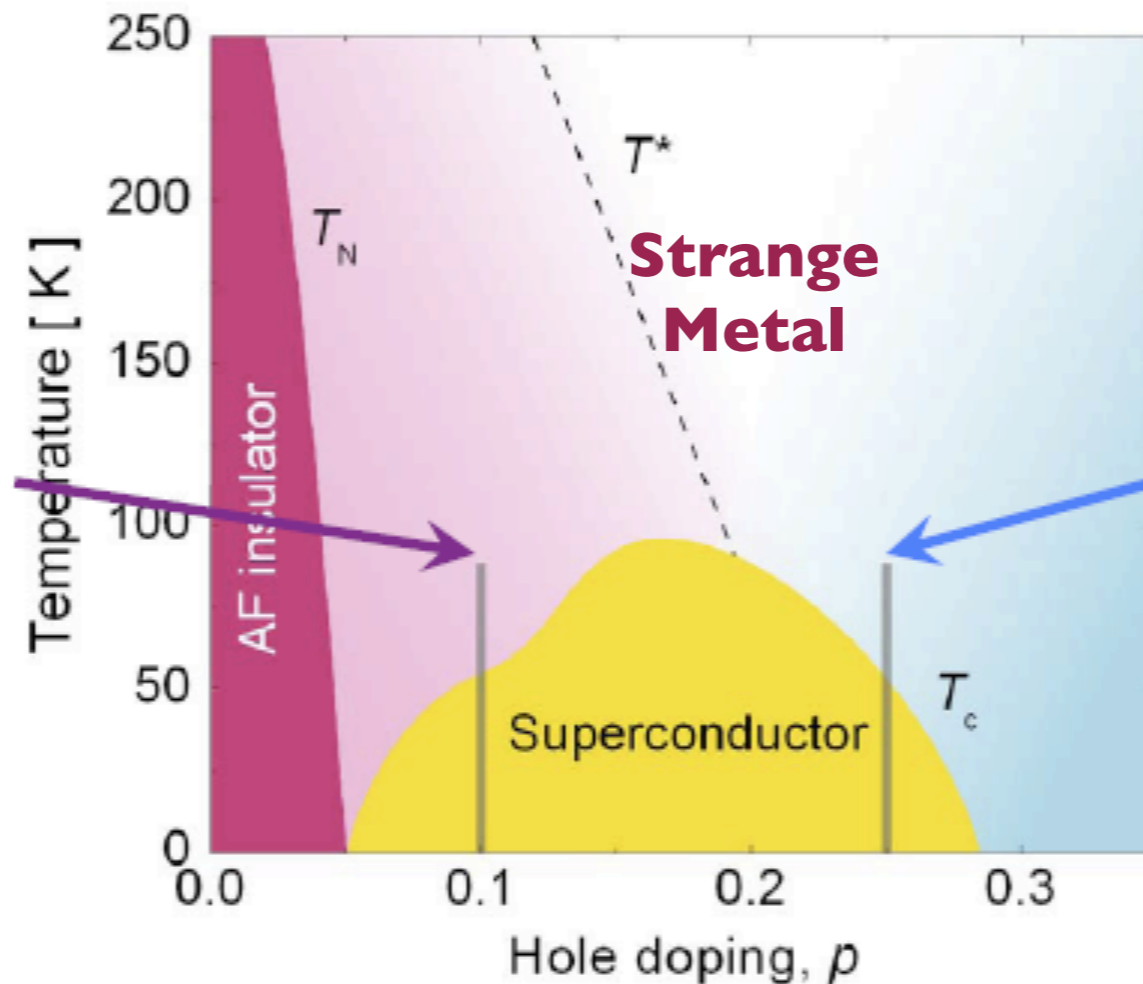
# Key Ingredients:

Antiferromagnetism (AF)  
Spin density wave (SDW)



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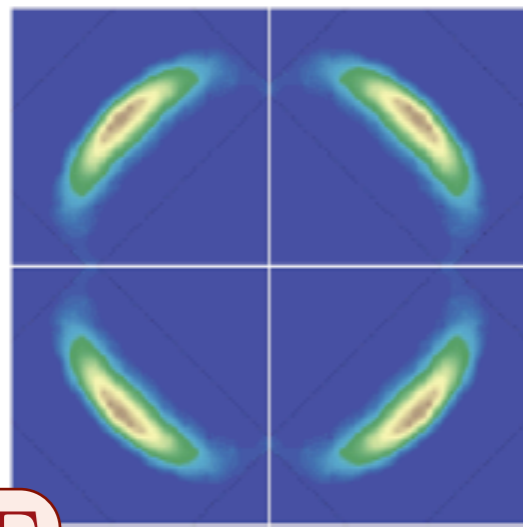
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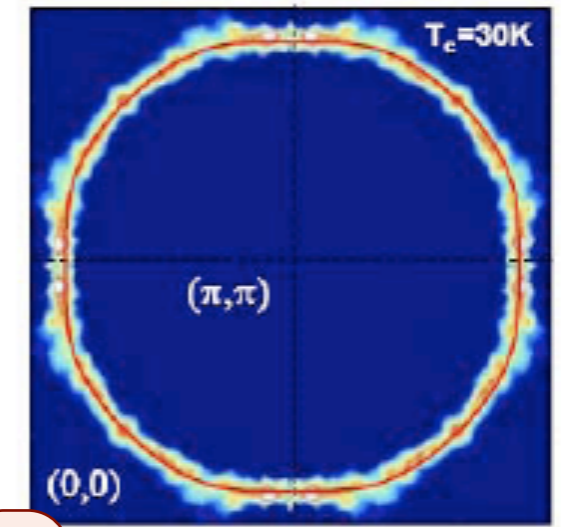
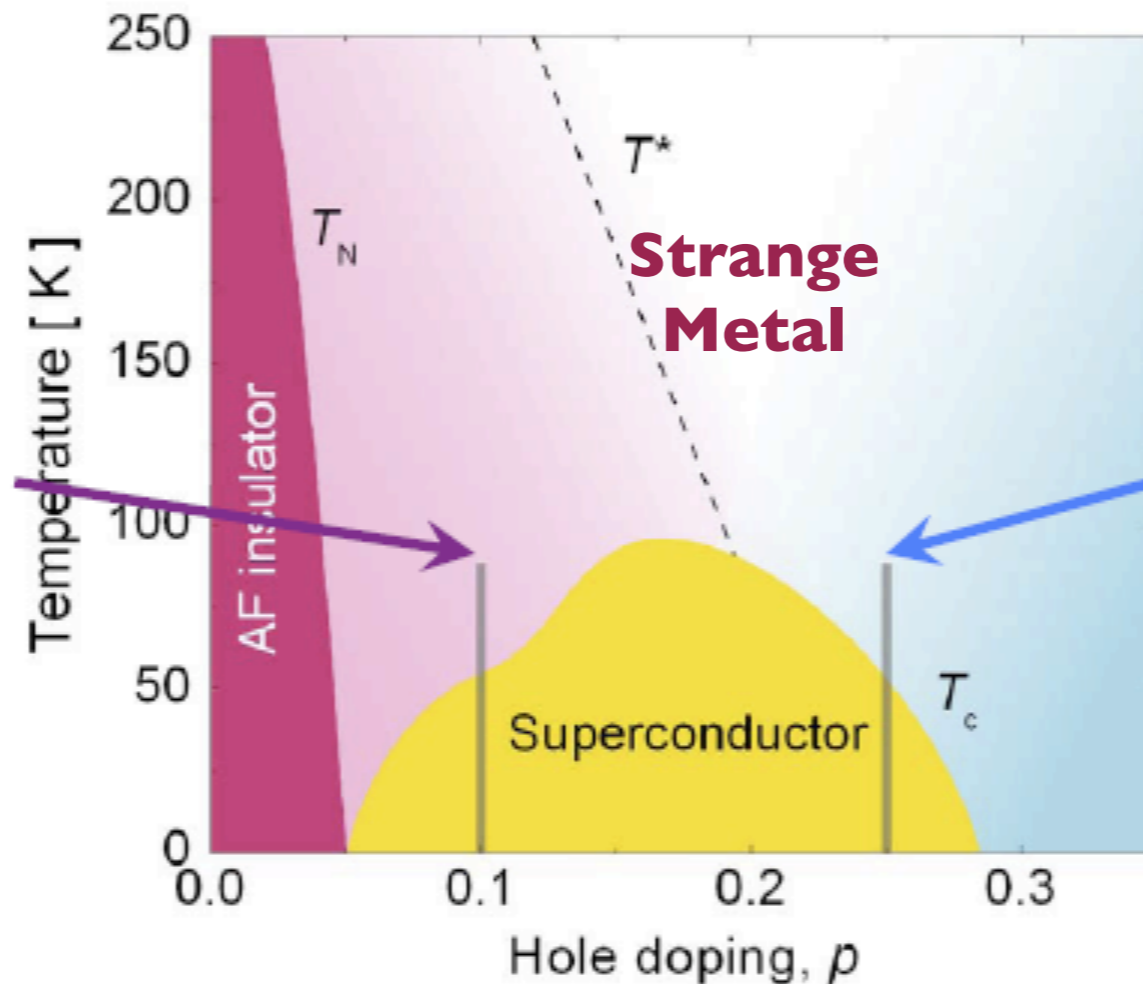
## d-wave superconductivity

Antiferromagnetism (AF)  
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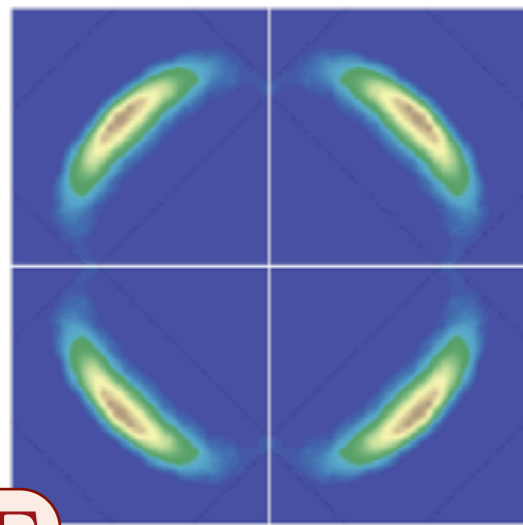
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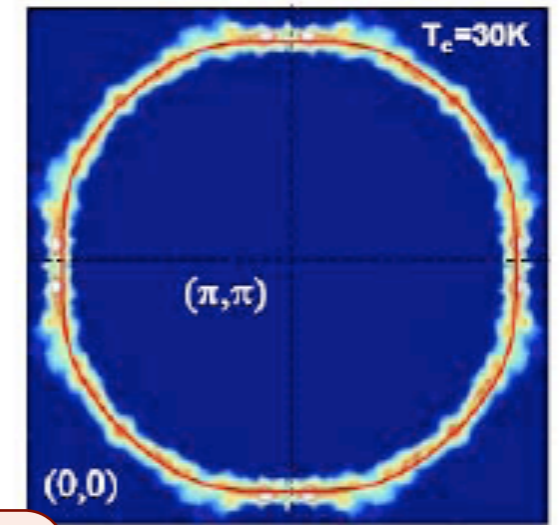
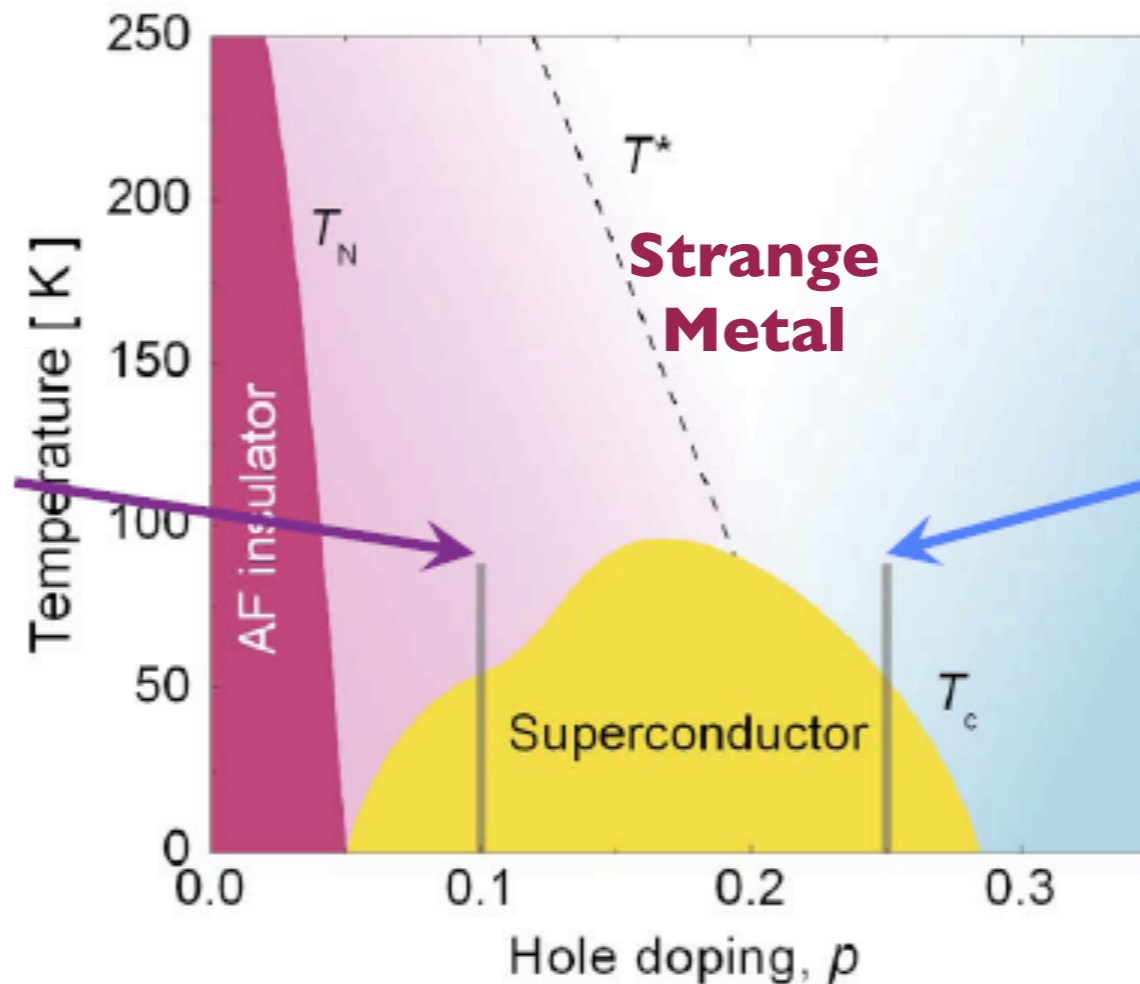
Fermi surface reconstruction



$\Gamma$

K.M. Shen et al., Science 2005

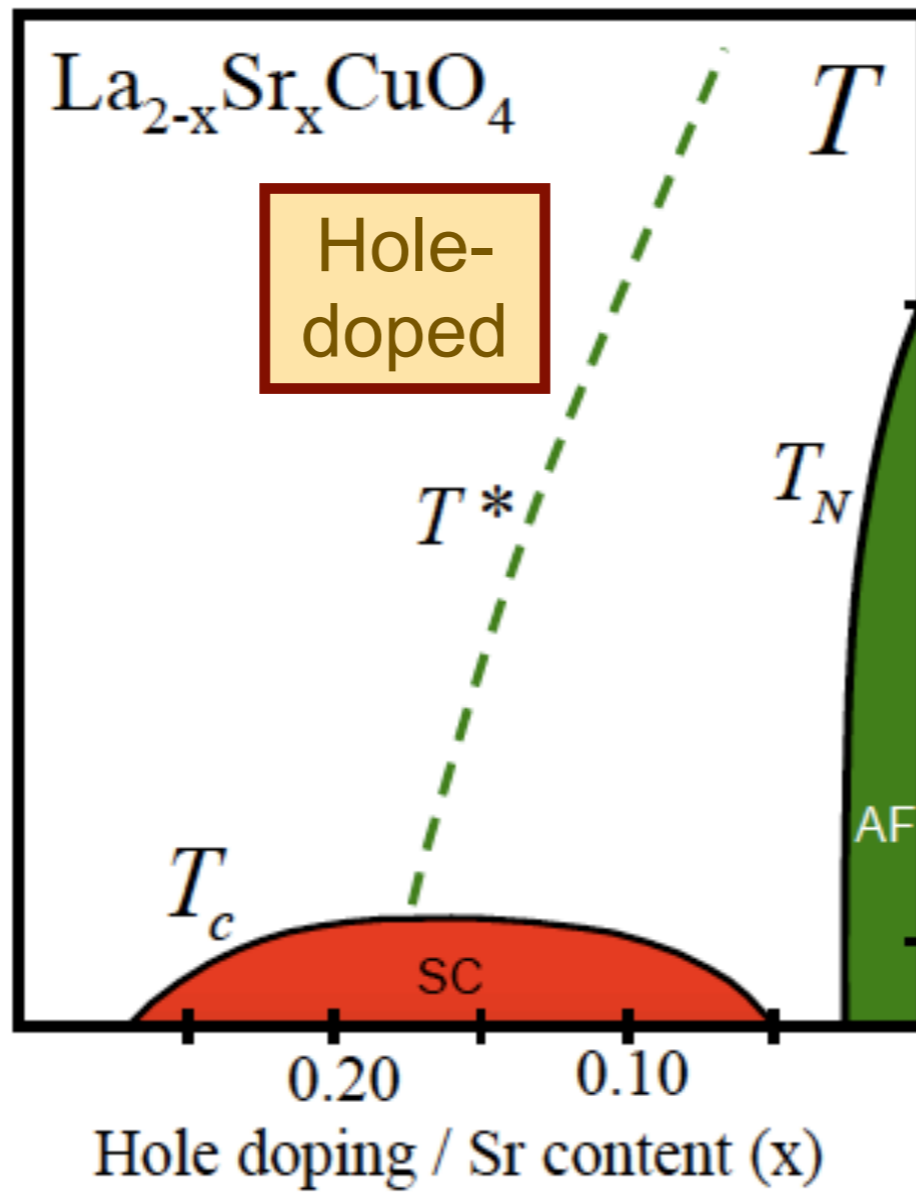
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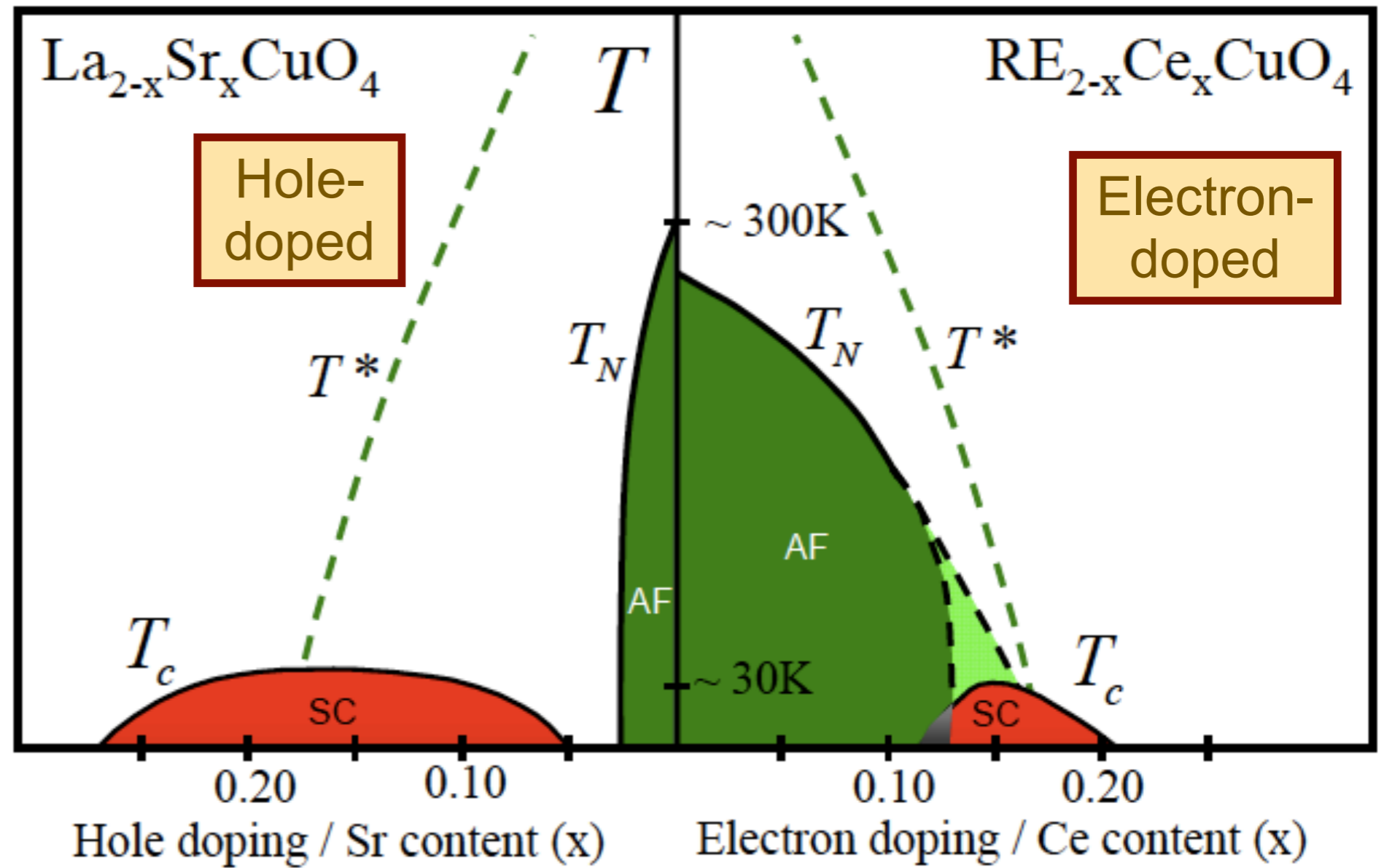
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# Electron-doped cuprate superconductors





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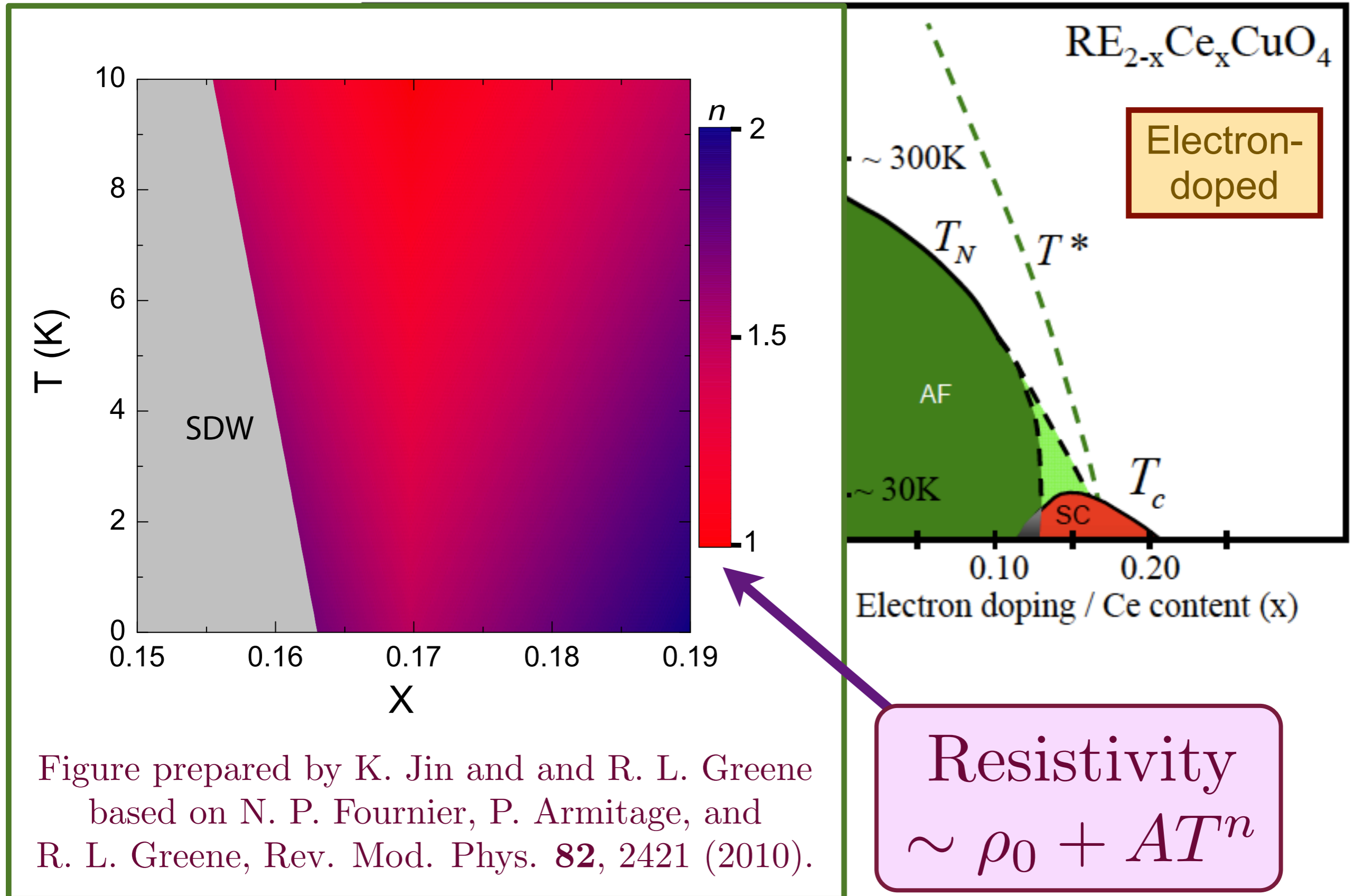


Figure prepared by K. Jin and R. L. Greene based on N. P. Fournier, P. Armitage, and R. L. Greene, Rev. Mod. Phys. **82**, 2421 (2010).

Resistivity  
 $\sim \rho_0 + AT^n$

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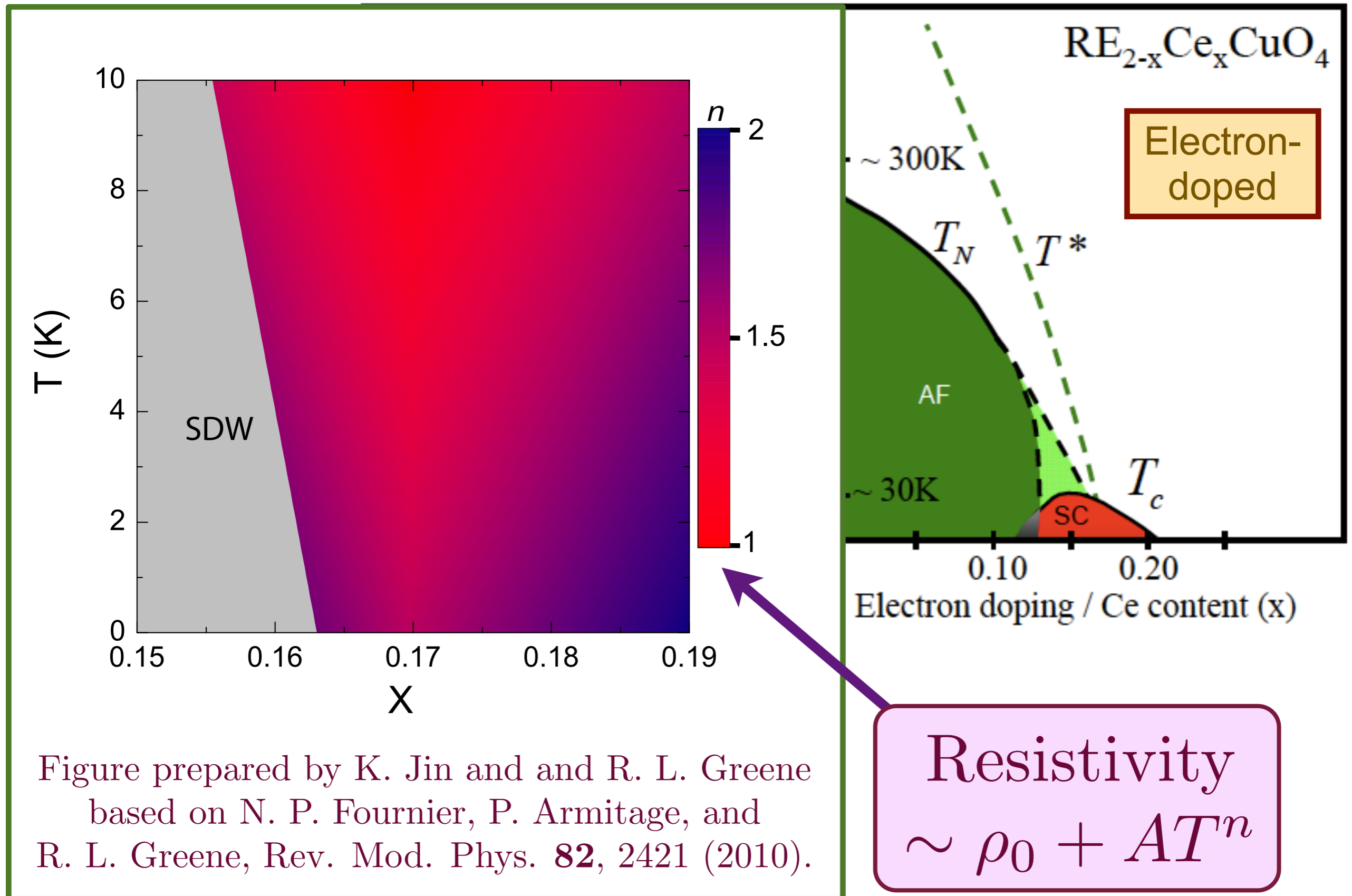


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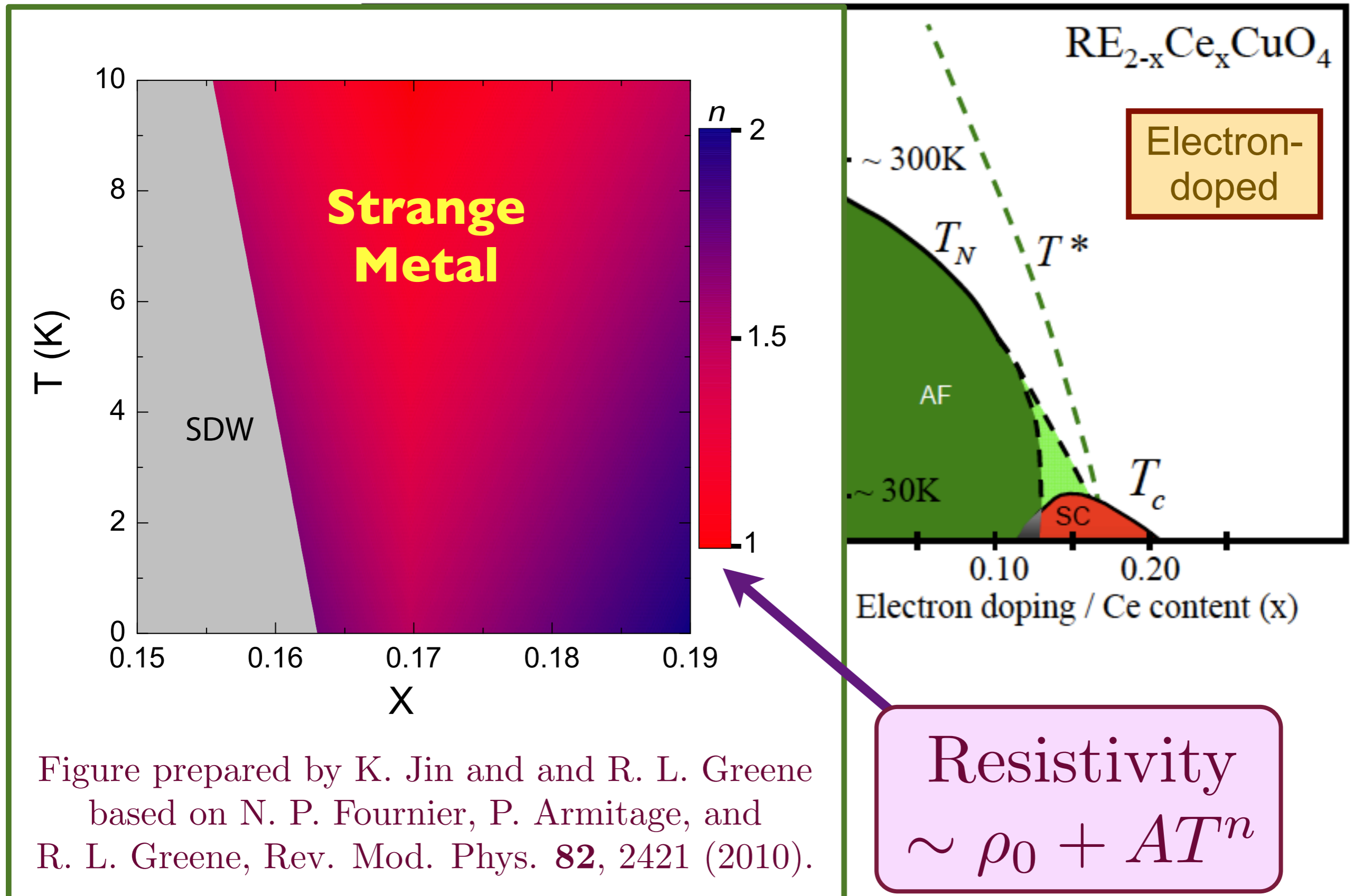
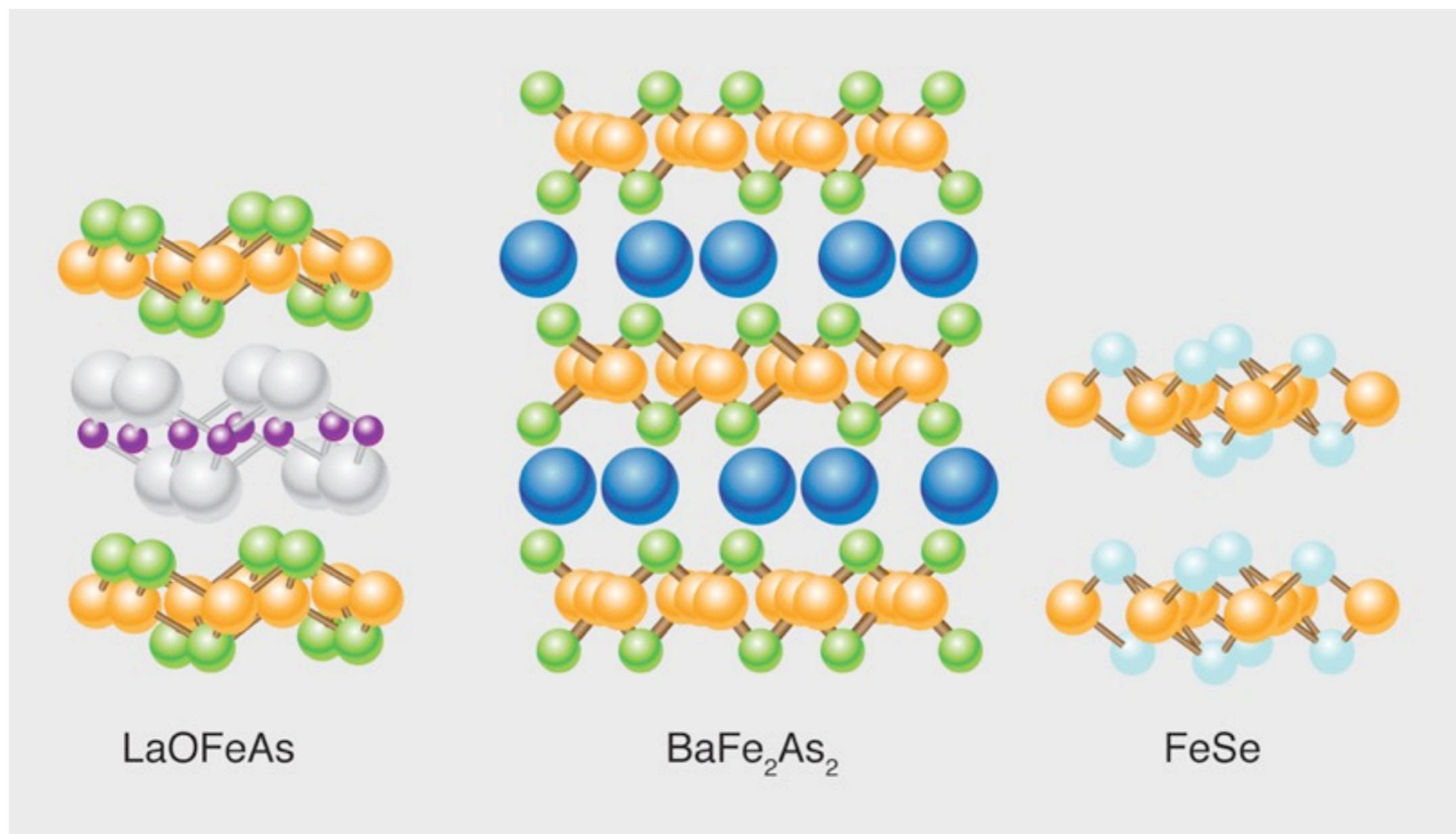


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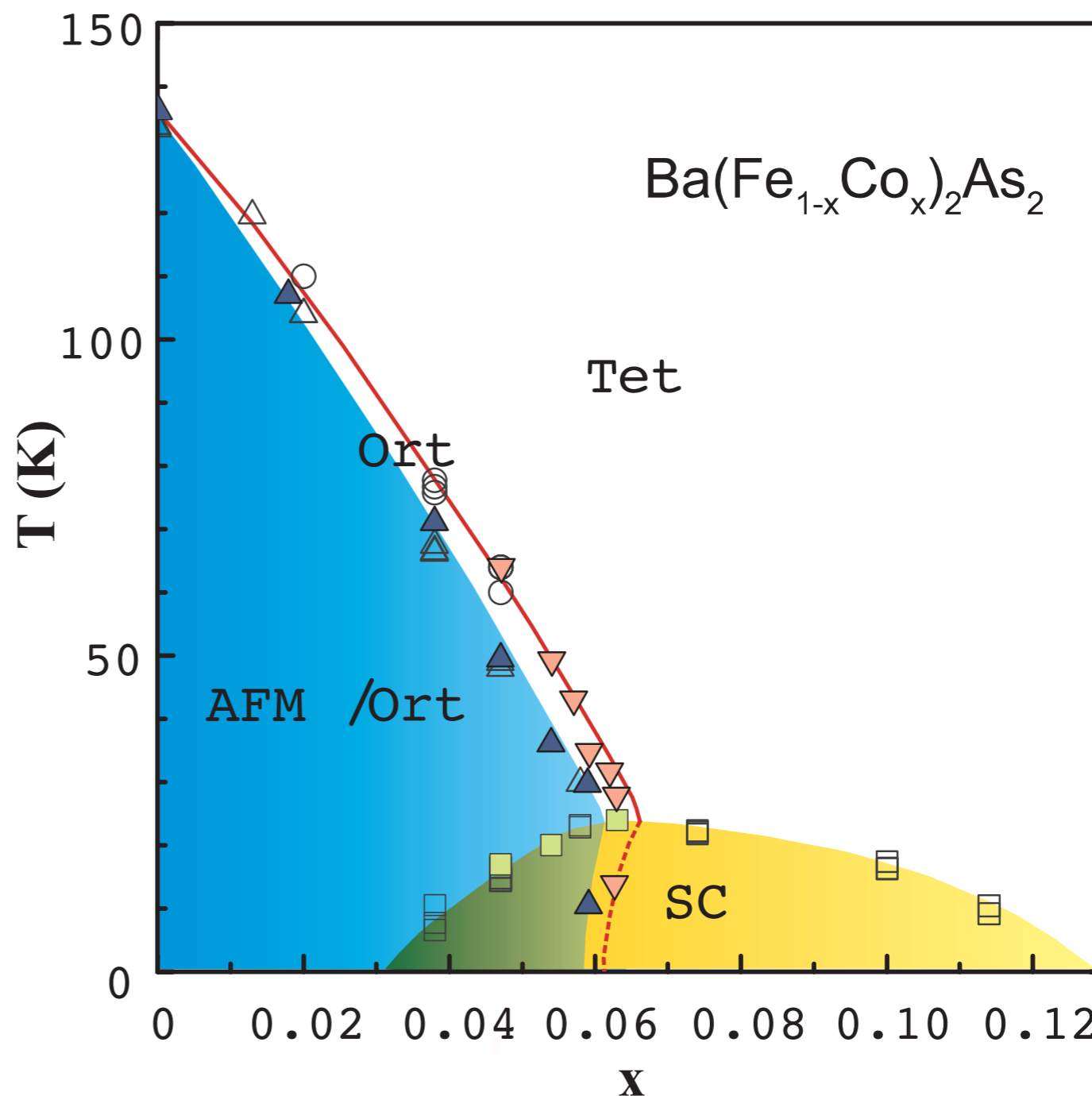
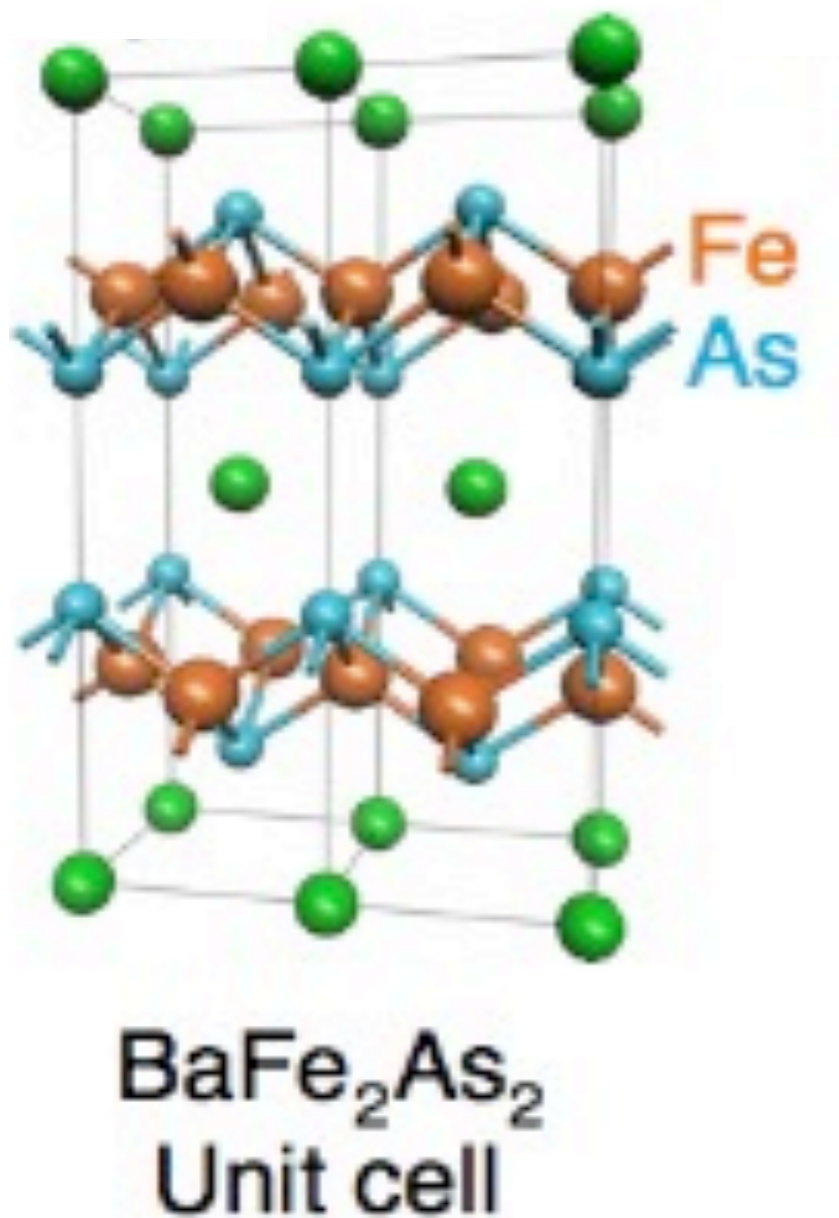
# Iron pnictides:

a new class of high temperature superconductors



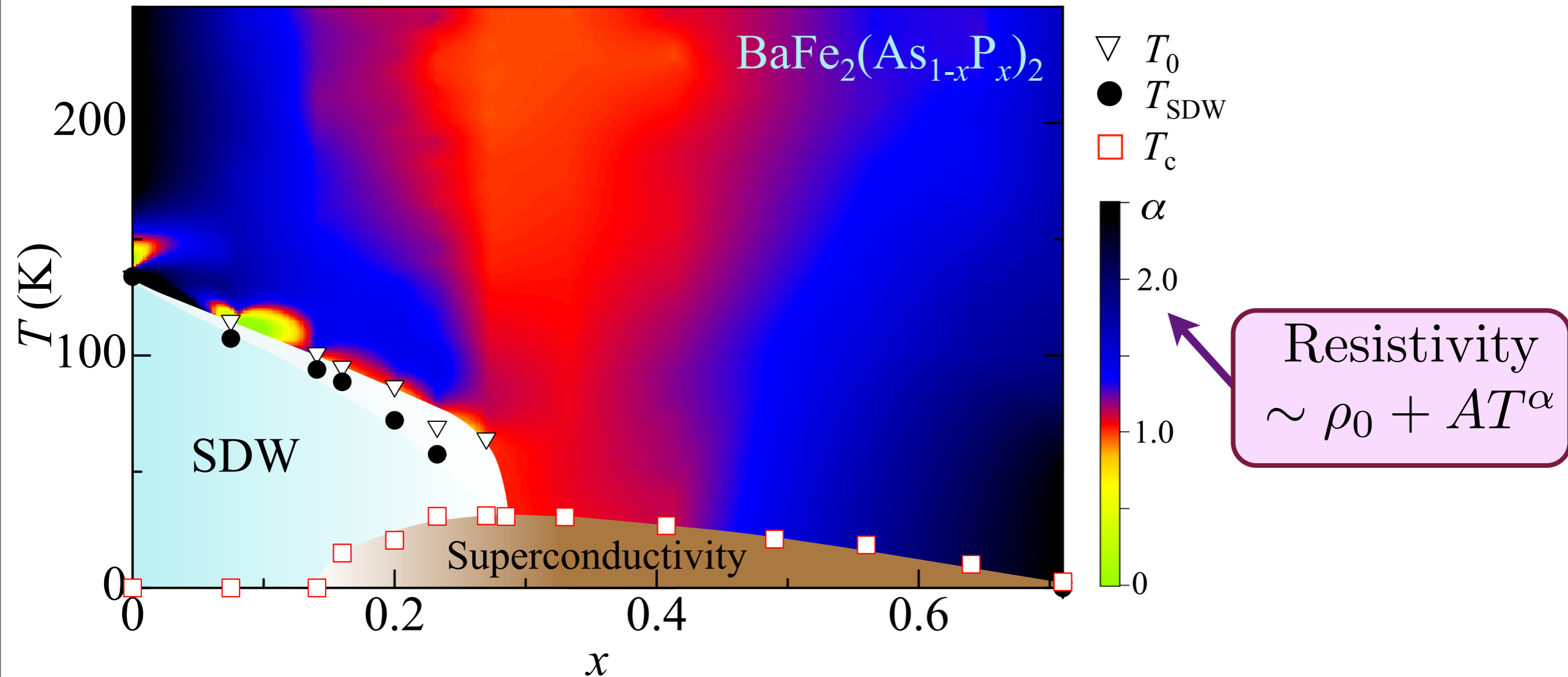
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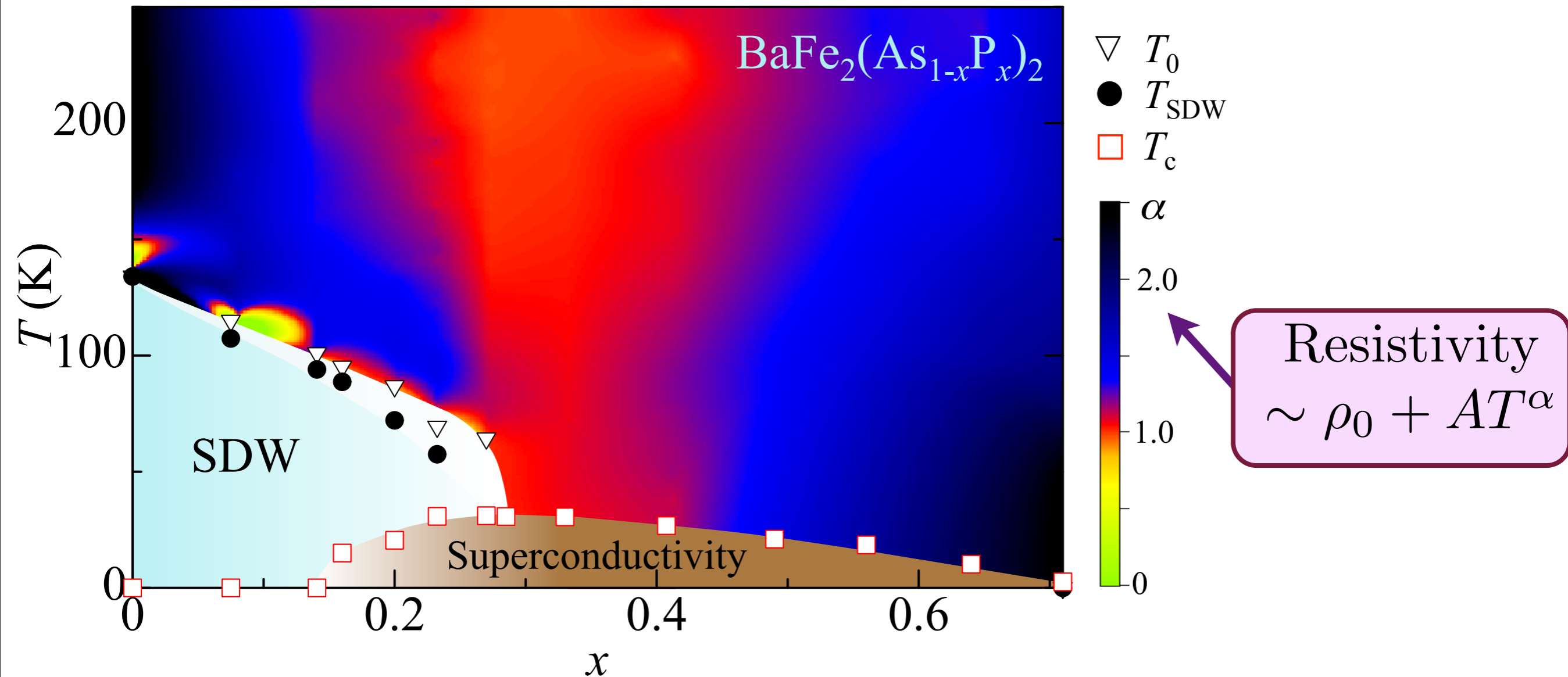
S. Nandi, M. G. Kim, A. Kreyssig, R. M. Fernandes, D. K. Pratt, A. Thaler, N. Ni,  
S. L. Bud'ko, P. C. Canfield, J. Schmalian, R. J. McQueeney, A. I. Goldman,  
*Physical Review Letters* **104**, 057006 (2010).

# Temperature-doping phase diagram of the iron pnictides:



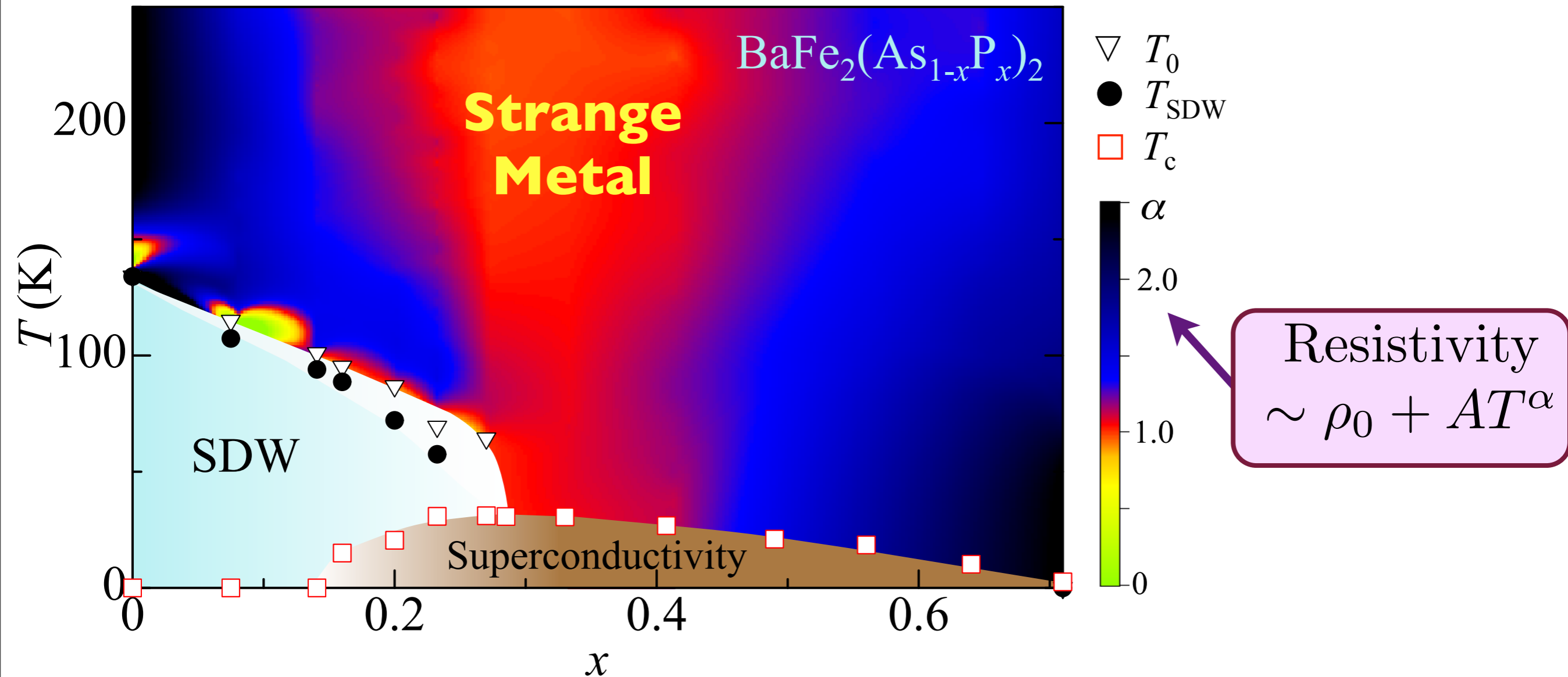
S. Kasahara, T. Shibauchi, K. Hashimoto, K. Ikada, S. Tonegawa, R. Okazaki, H. Shishido, H. Ikeda, H. Takeya, K. Hirata, T. Terashima, and Y. Matsuda, *Physical Review B* **81**, 184519 (2010)

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## Questions

- *Can quantum fluctuations near the loss of antiferromagnetism induce higher temperature superconductivity ?*

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- *If so, why is there no antiferromagnetism in the hole-doped cuprates near the point where the superconductivity is strongest ?*
- *What is the physics of the strange metal ?*

# Outline

## 1. Phenomenology of the onset of antiferromagnetism in a metal

*Quantum criticality of Fermi surface reconstruction, and the phase diagram in a magnetic field*

## 2. Strongly-coupled quantum criticality in metals

*“Mechanism” of higher temperature superconductivity*

## 3. Theory of the competition between superconductivity and antiferromagnetism

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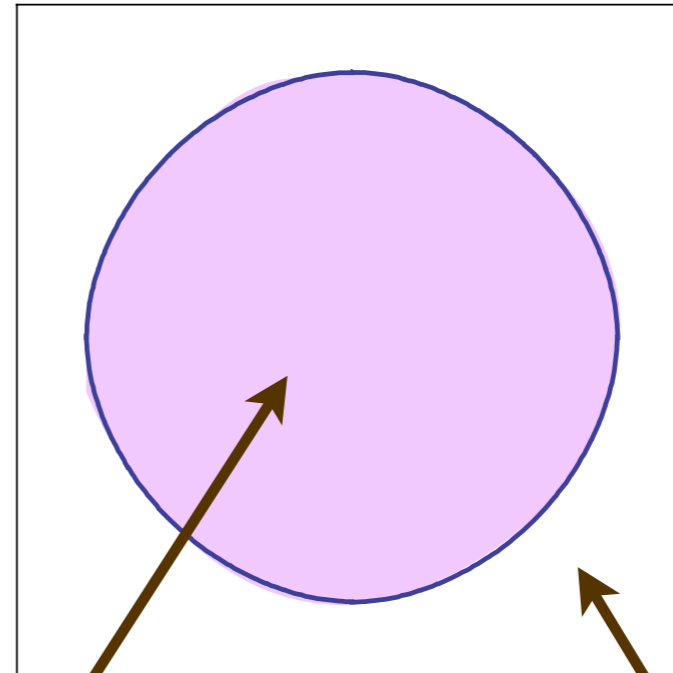
## 2. Strongly-coupled quantum criticality in metals

*“Mechanism” of higher temperature superconductivity*

## 3. Theory of the competition between superconductivity and antiferromagnetism

# Fermi surface

Metal with “large”  
Fermi surface

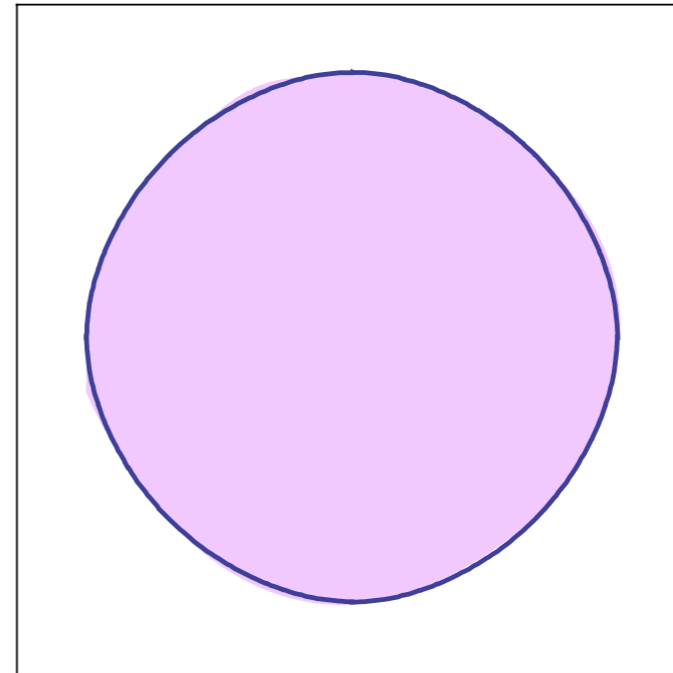


Momenta with  
electronic  
states empty

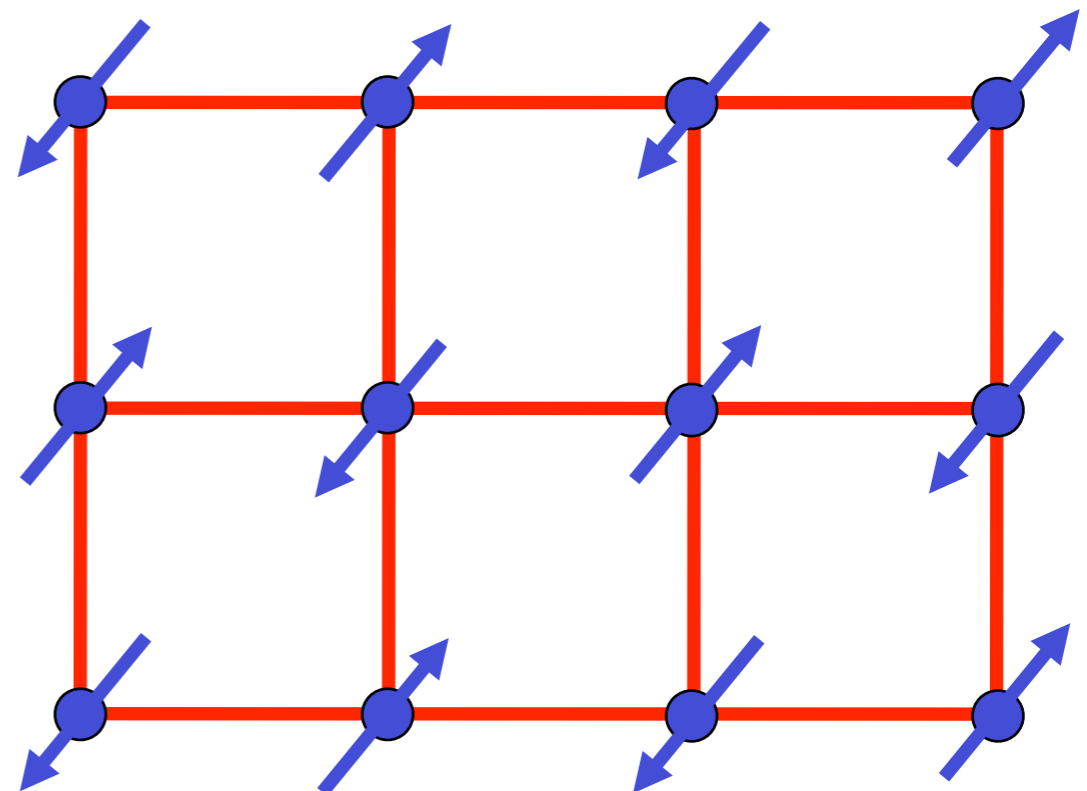
Momenta with  
electronic  
states  
occupied

# Fermi surface+antiferromagnetism

Metal with “large”  
Fermi surface



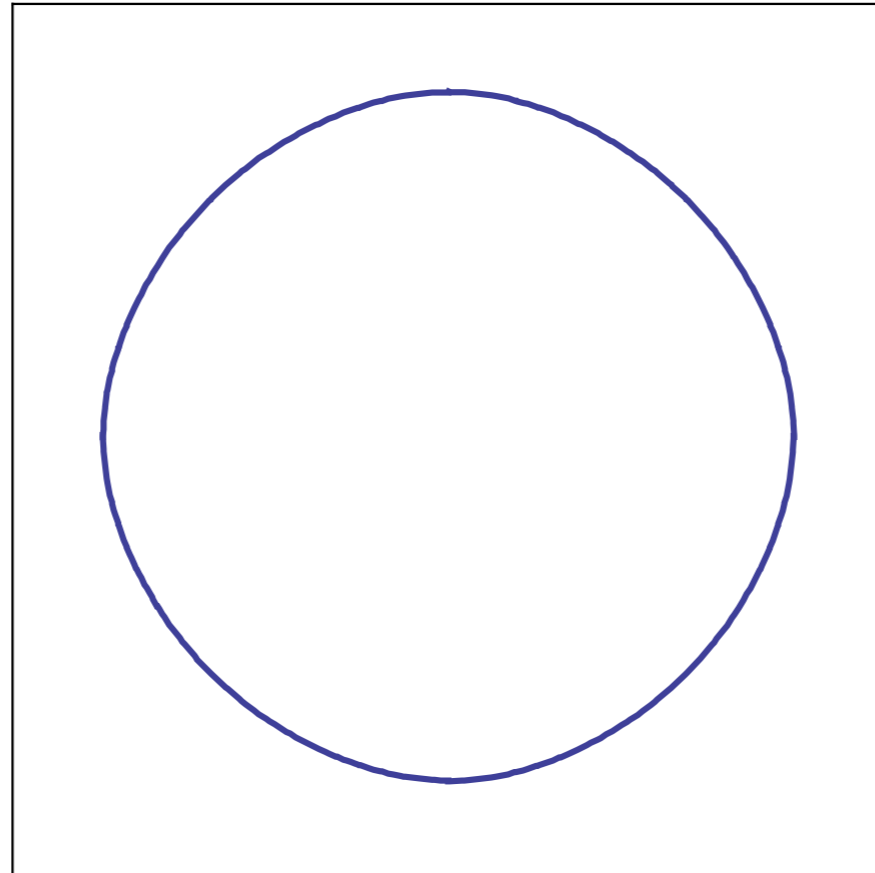
+



The electron spin polarization obeys

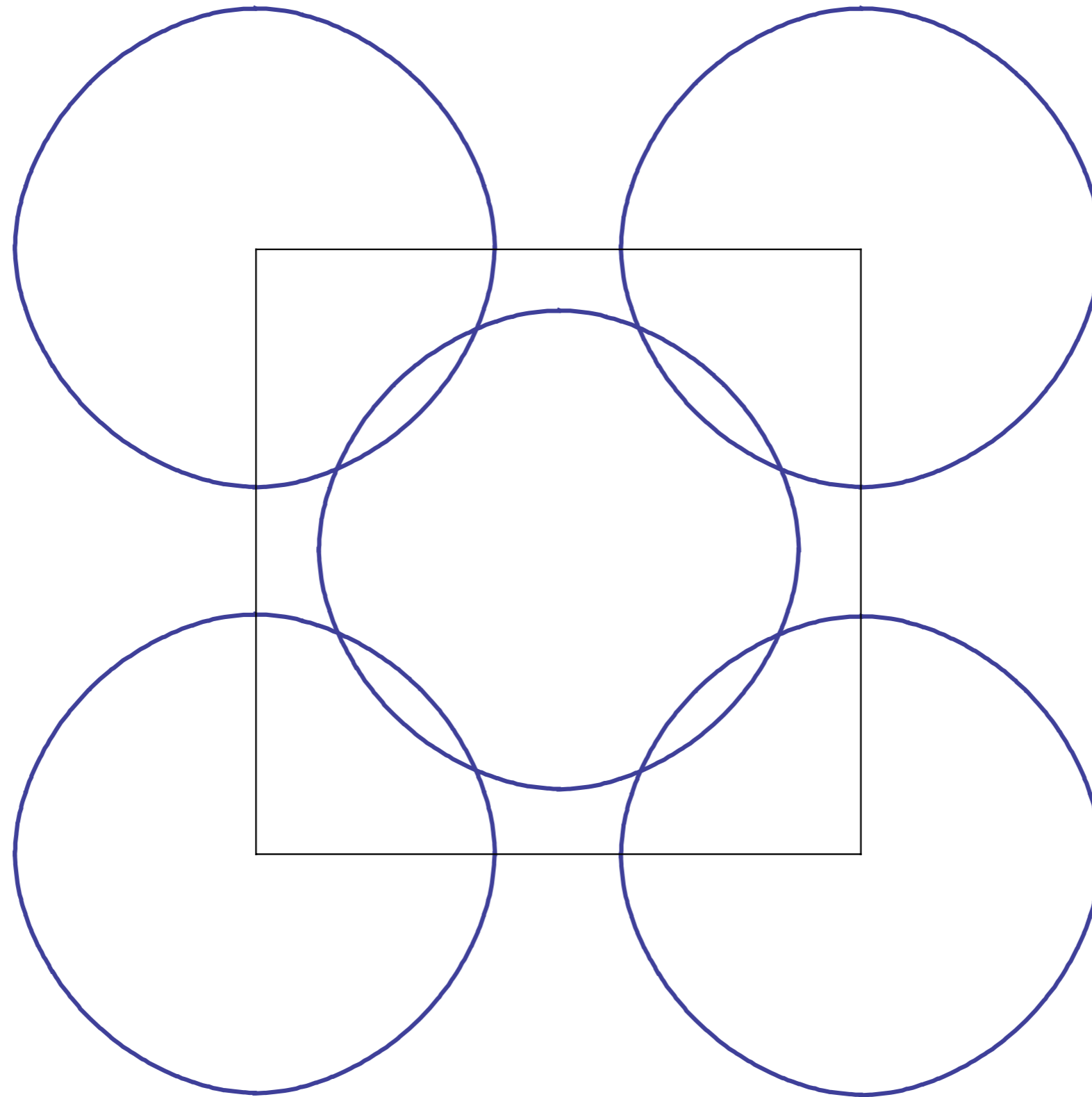
$$\langle \vec{S}(\mathbf{r}, \tau) \rangle = \vec{\varphi}(\mathbf{r}, \tau) e^{i\mathbf{K} \cdot \mathbf{r}}$$

where  $\mathbf{K}$  is the ordering wavevector.

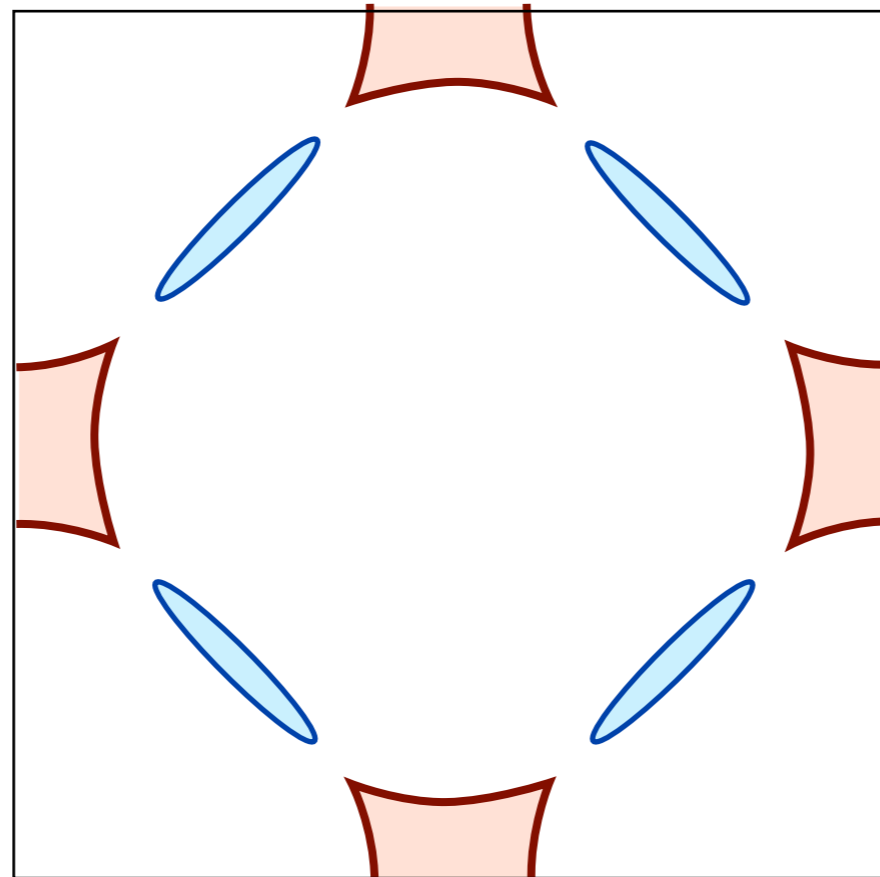


**Metal with “large” Fermi surface**



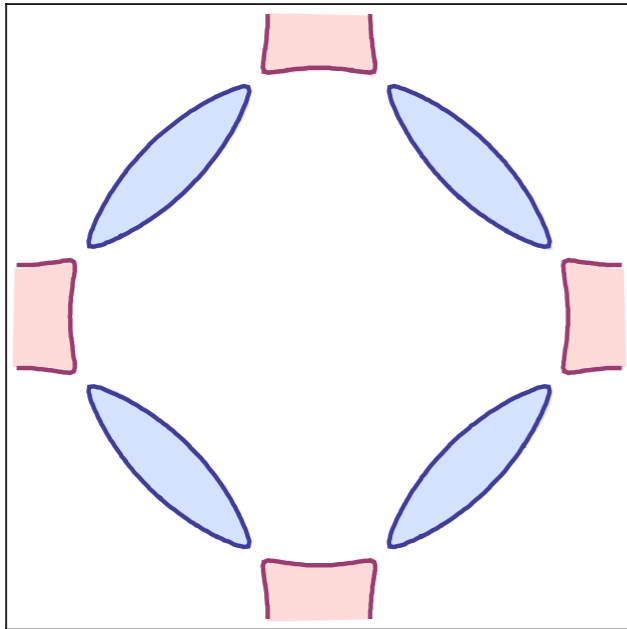


Fermi surfaces translated by  $\mathbf{K} = (\pi, \pi)$ .



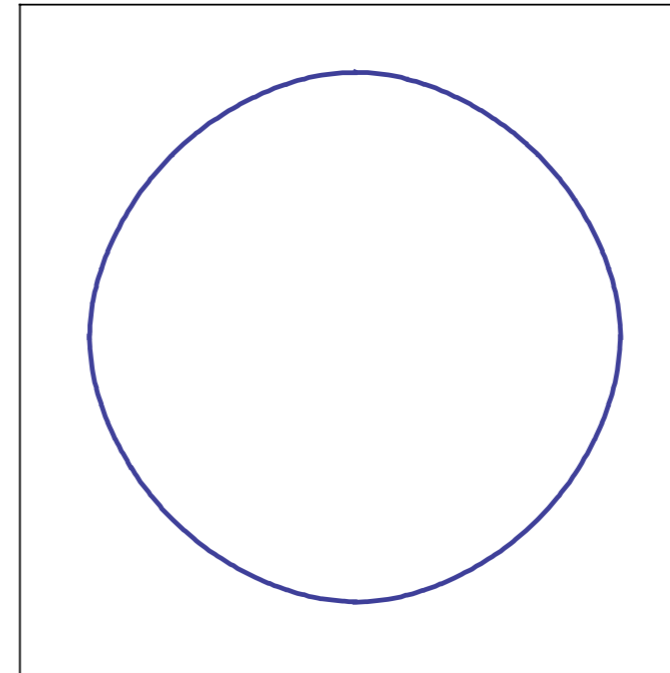
Electron and hole pockets in  
antiferromagnetic phase with  $\langle \vec{\varphi} \rangle \neq 0$

# Fermi surface+antiferromagnetism



$$\langle \vec{\varphi} \rangle \neq 0$$

Metal with electron  
and hole pockets



$$\langle \vec{\varphi} \rangle = 0$$

Metal with “large”  
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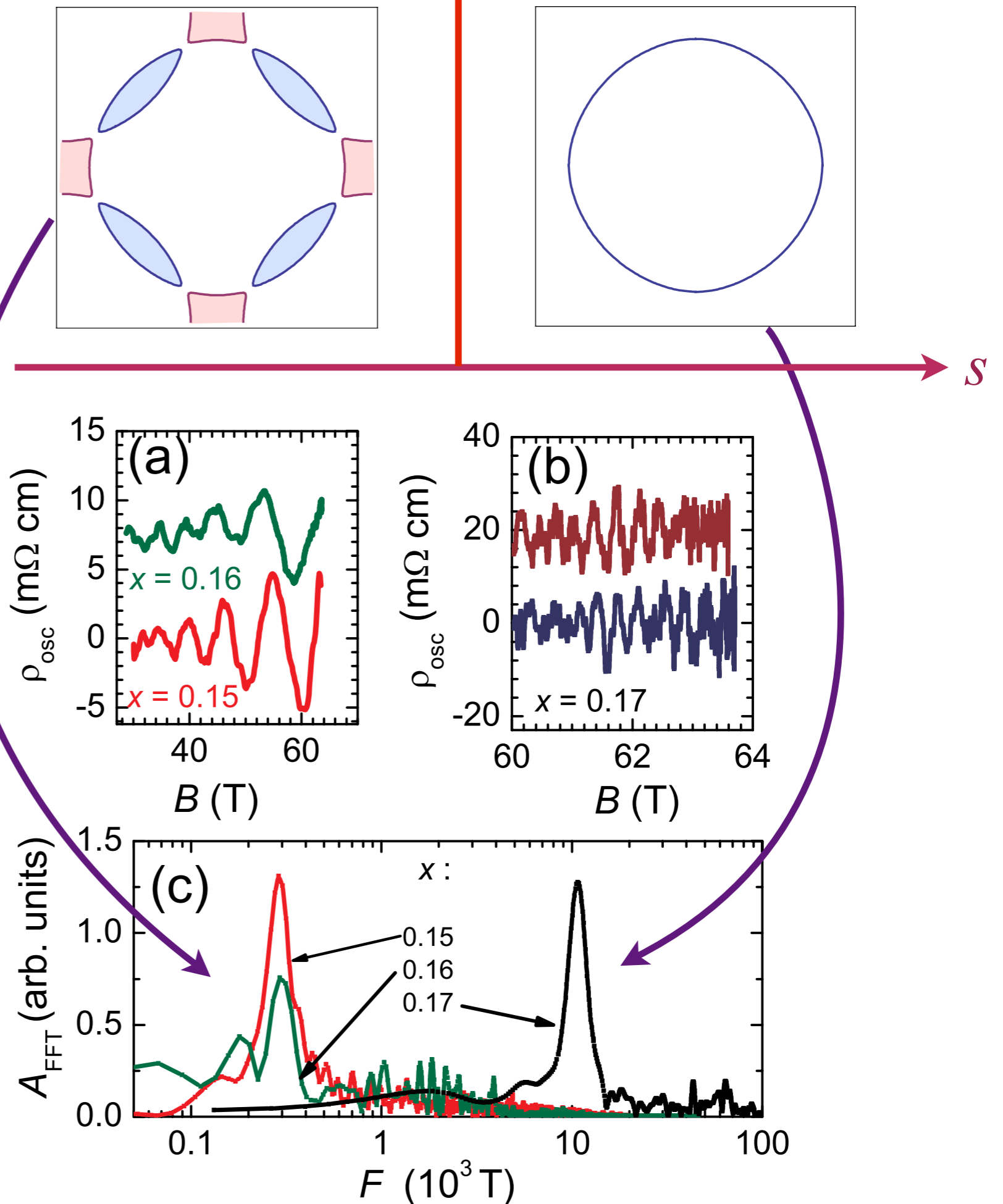
← Increasing interaction

S. Sachdev, A. V. Chubukov, and A. Sokol, *Phys. Rev. B* **51**, 14874 (1995).  
A. V. Chubukov and D. K. Morr, *Physics Reports* **288**, 355 (1997).

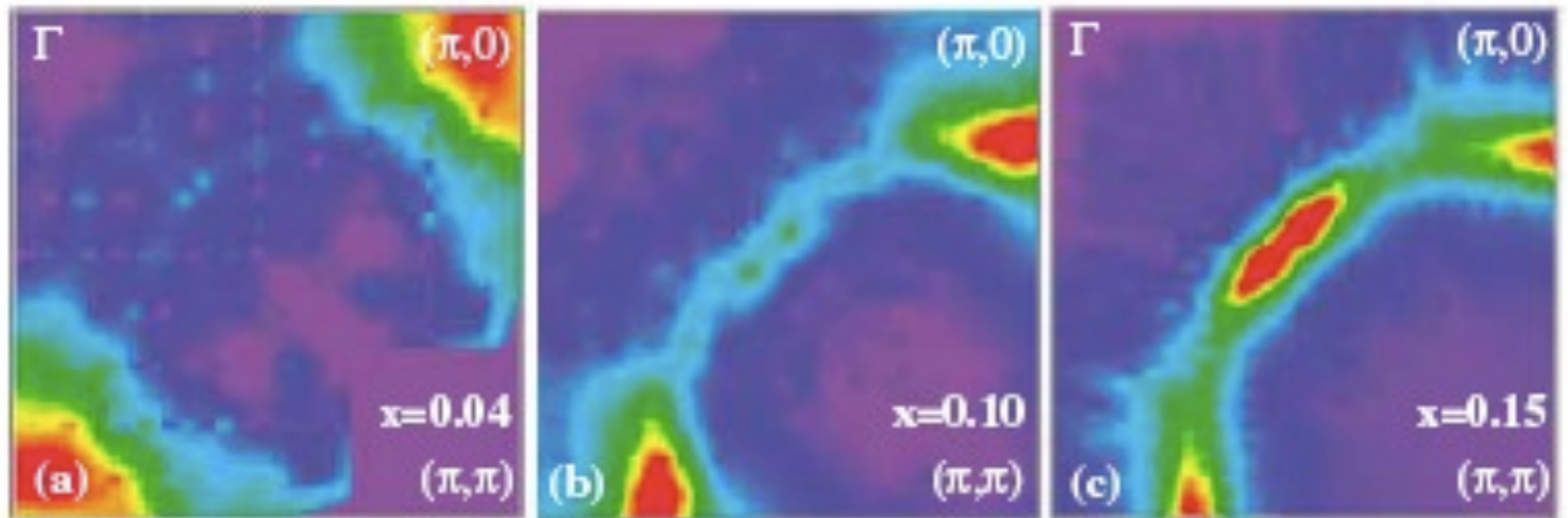
# Quantum oscillations



T. Helm, M.V. Kartsovnik,  
M. Bartkowiak, N. Bittner,  
M. Lambacher, A. Erb, J. Wosnitza,  
and R. Gross,  
*Phys. Rev. Lett.* **103**, 157002 (2009).

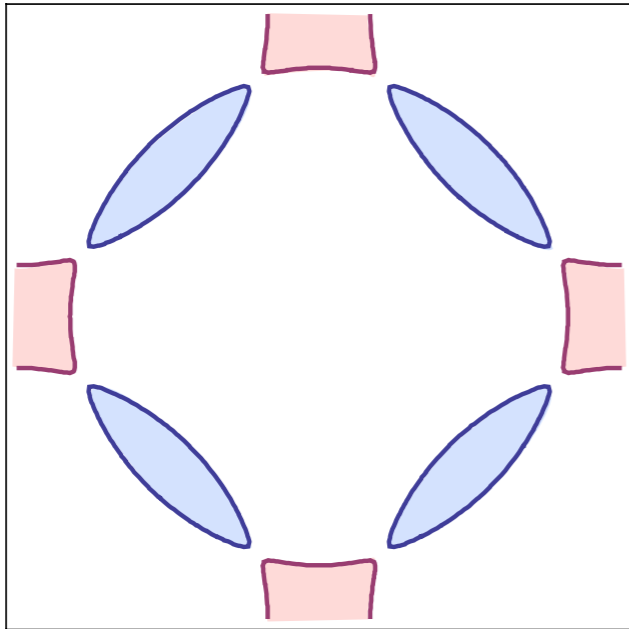


# Photoemission in $\text{Nd}_{2-x}\text{Ce}_x\text{CuO}_4$



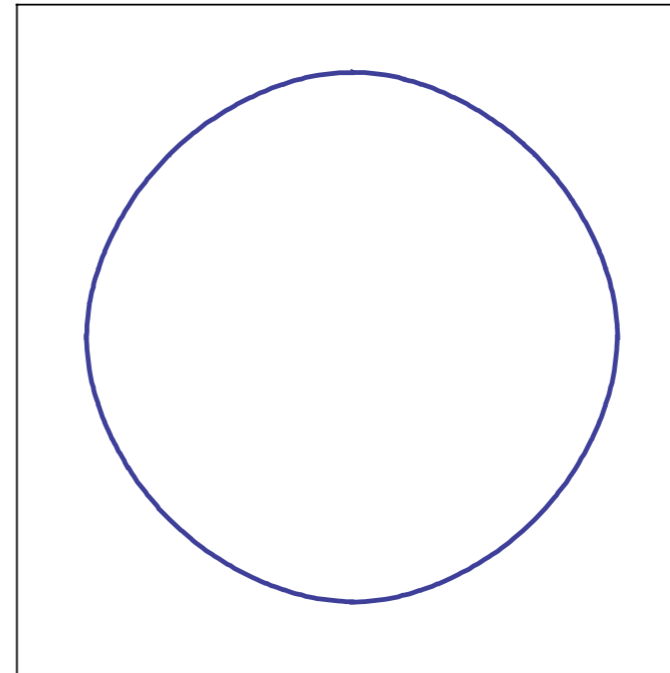
N. P. Armitage *et al.*, Phys. Rev. Lett. **88**, 257001 (2002).

# Fermi surface+antiferromagnetism



$$\langle \vec{\varphi} \rangle \neq 0$$

Metal with electron  
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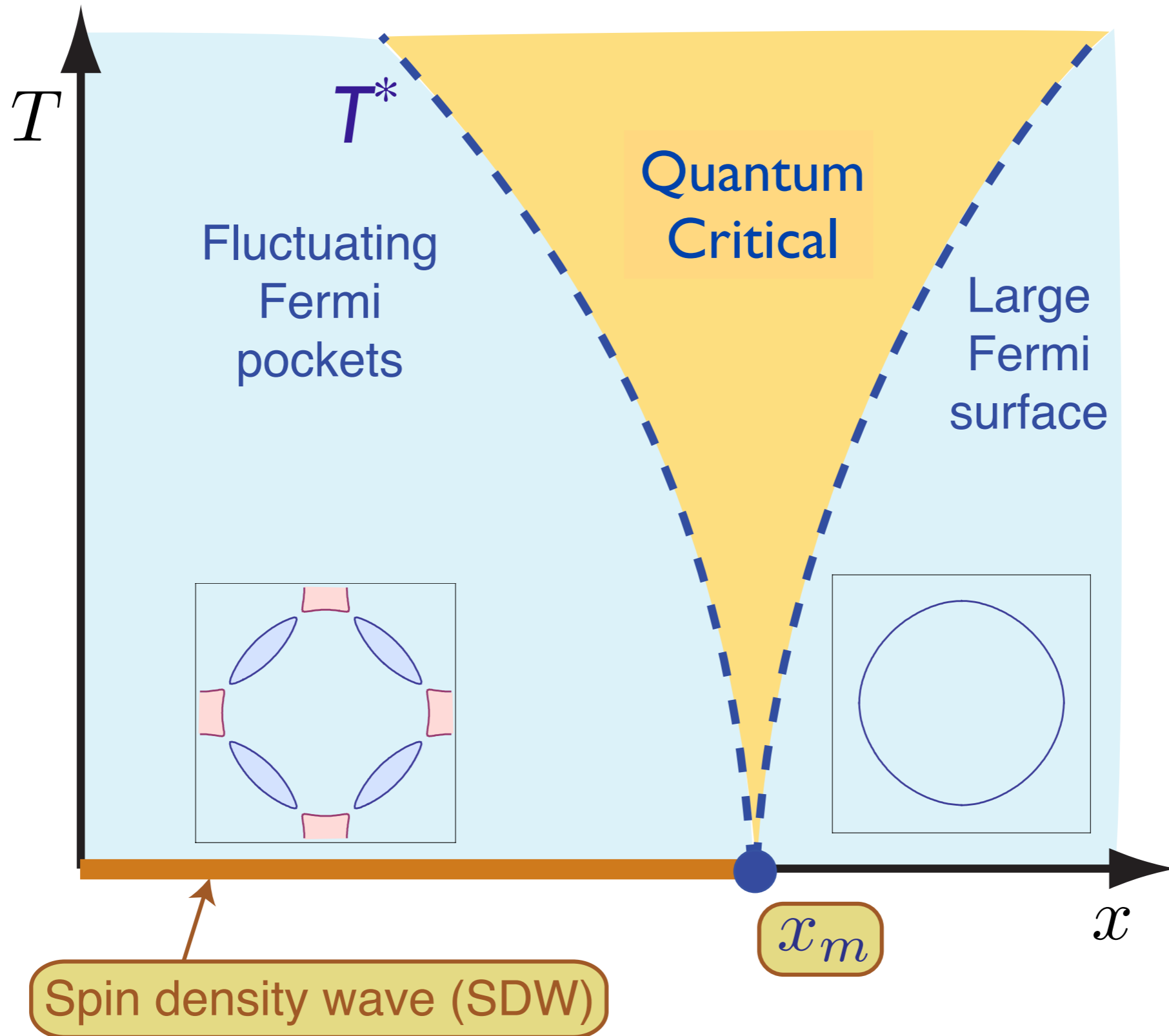
$$\langle \vec{\varphi} \rangle = 0$$

Metal with “large”  
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$S$

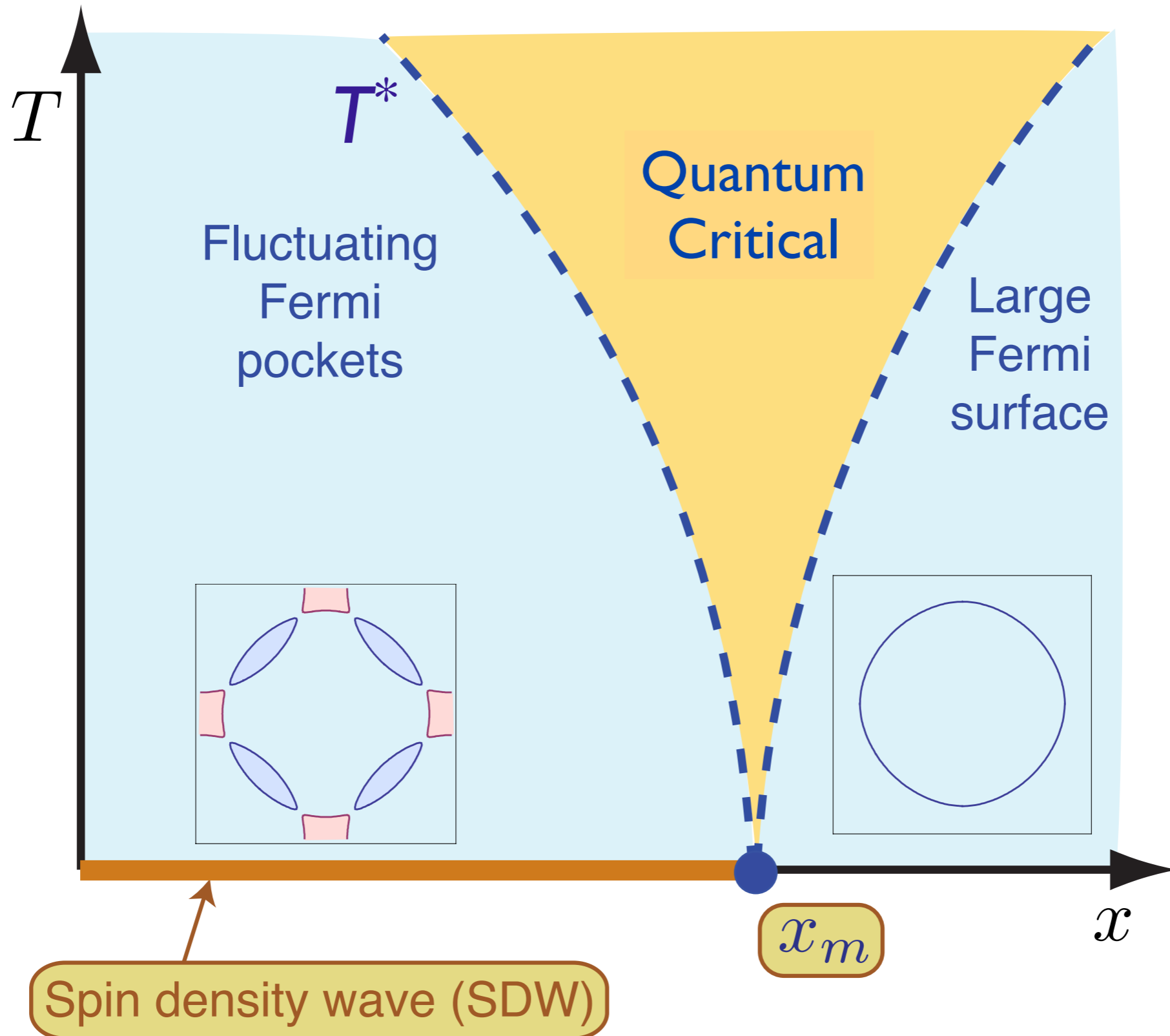
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# Theory of quantum criticality in the cuprates



Underlying SDW ordering quantum critical point  
in metal at  $x = x_m$

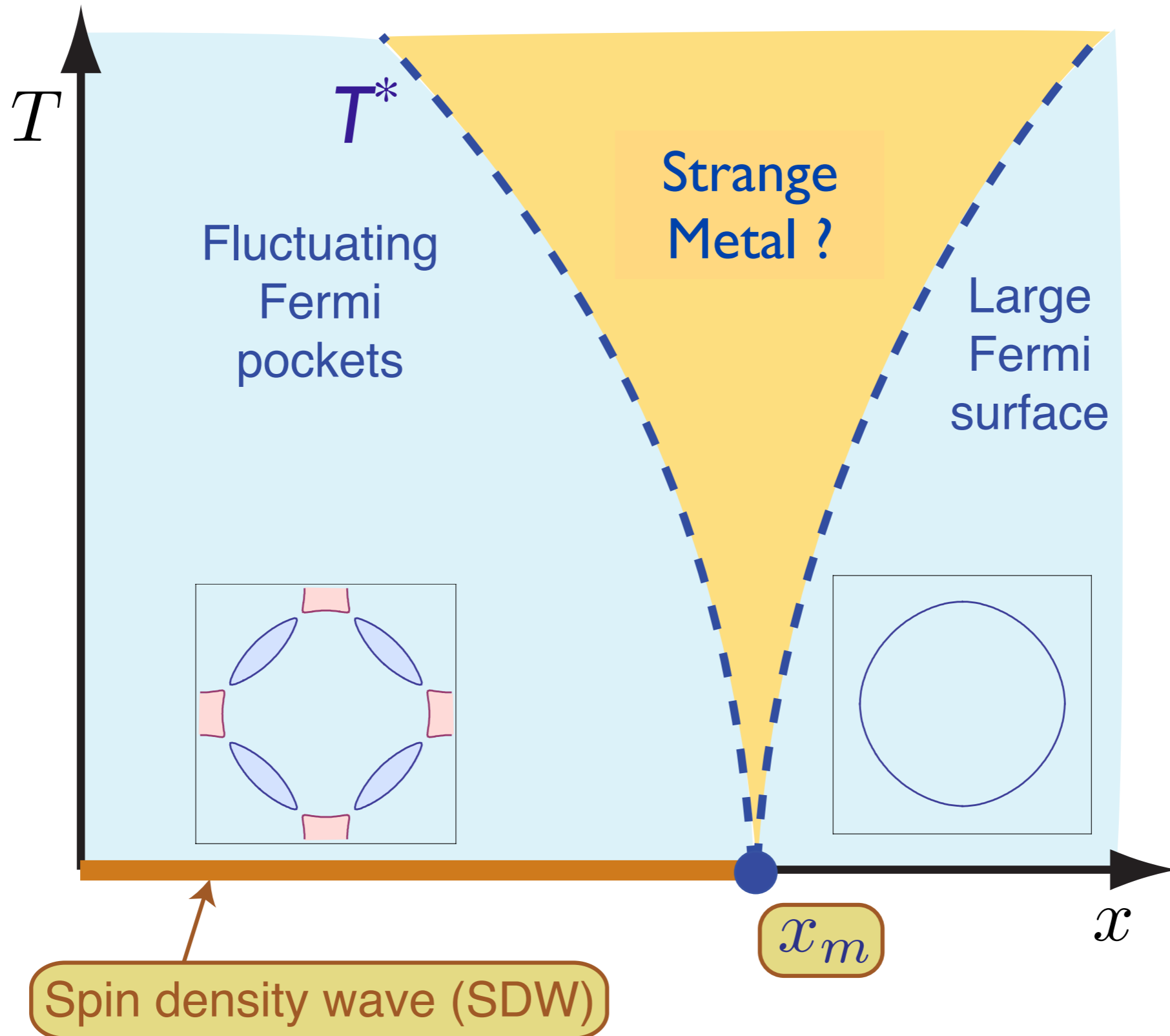
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Relaxation and equilibration times  $\sim \hbar/k_B T$  are robust properties of strongly-coupled quantum criticality

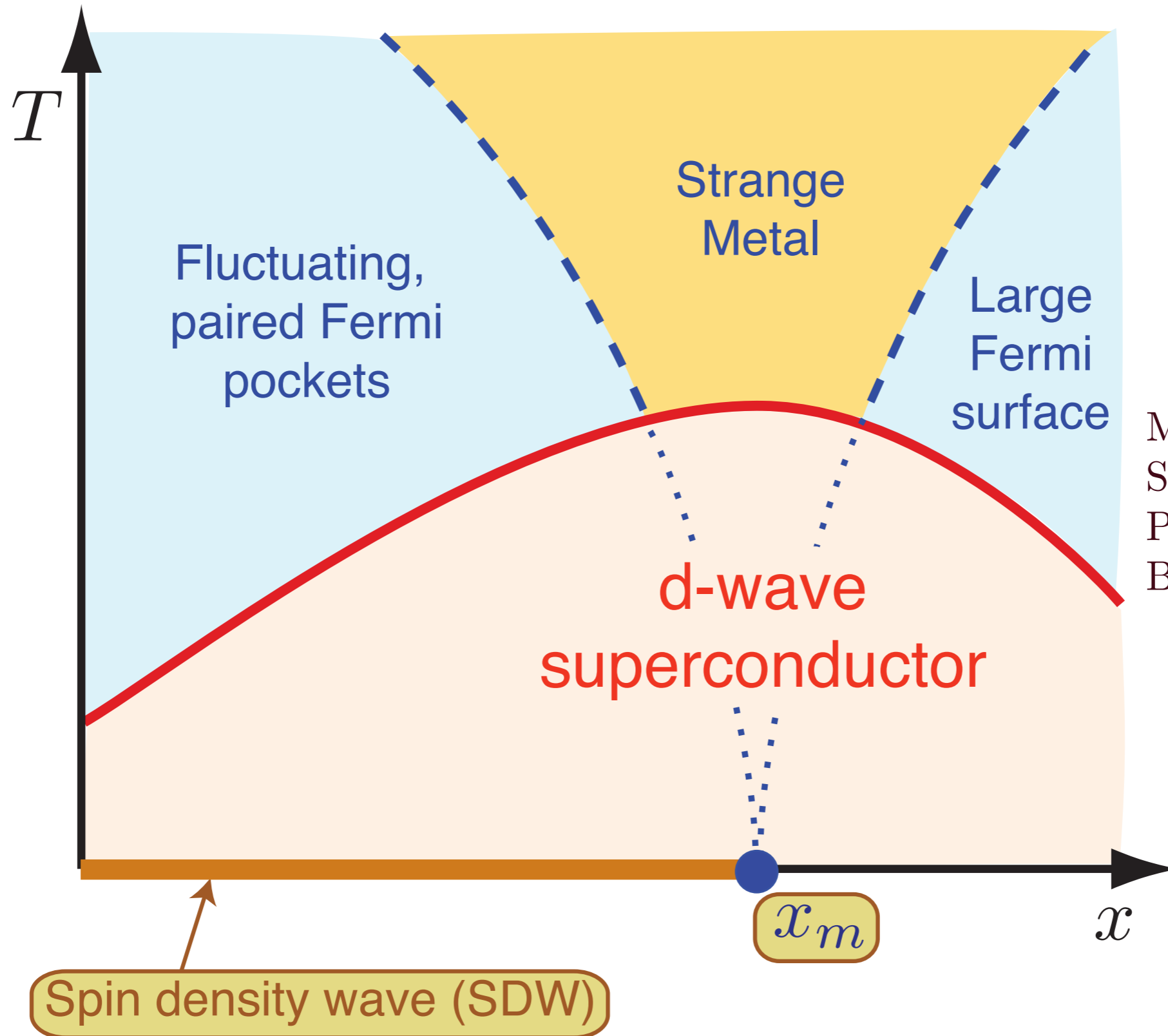


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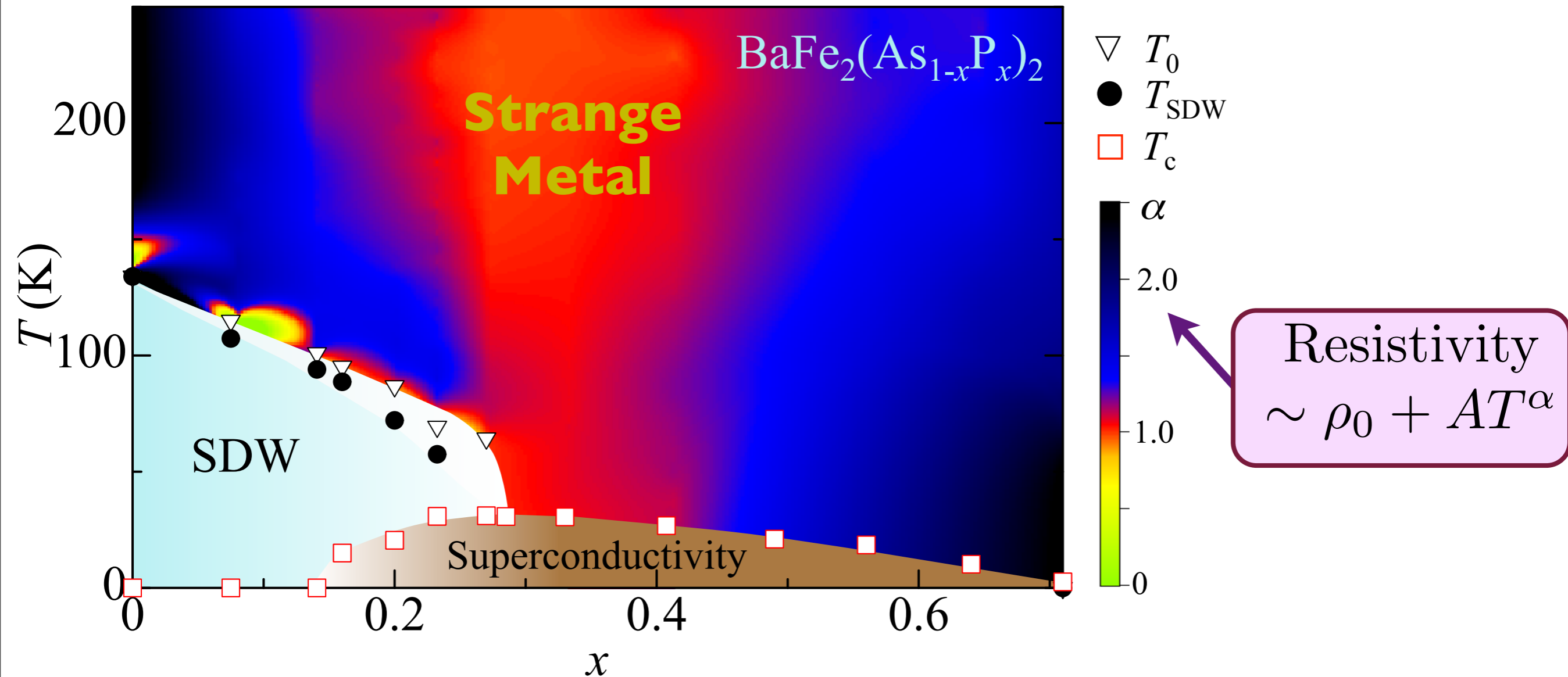
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M. A. Metlitski and  
S. Sachdev,  
Physical Review  
B **82**, 075128 (2010)

SDW quantum critical point is unstable to  $d$ -wave superconductivity  
This instability is stronger than that in the BCS theory

# Temperature-doping phase diagram of the iron pnictides:

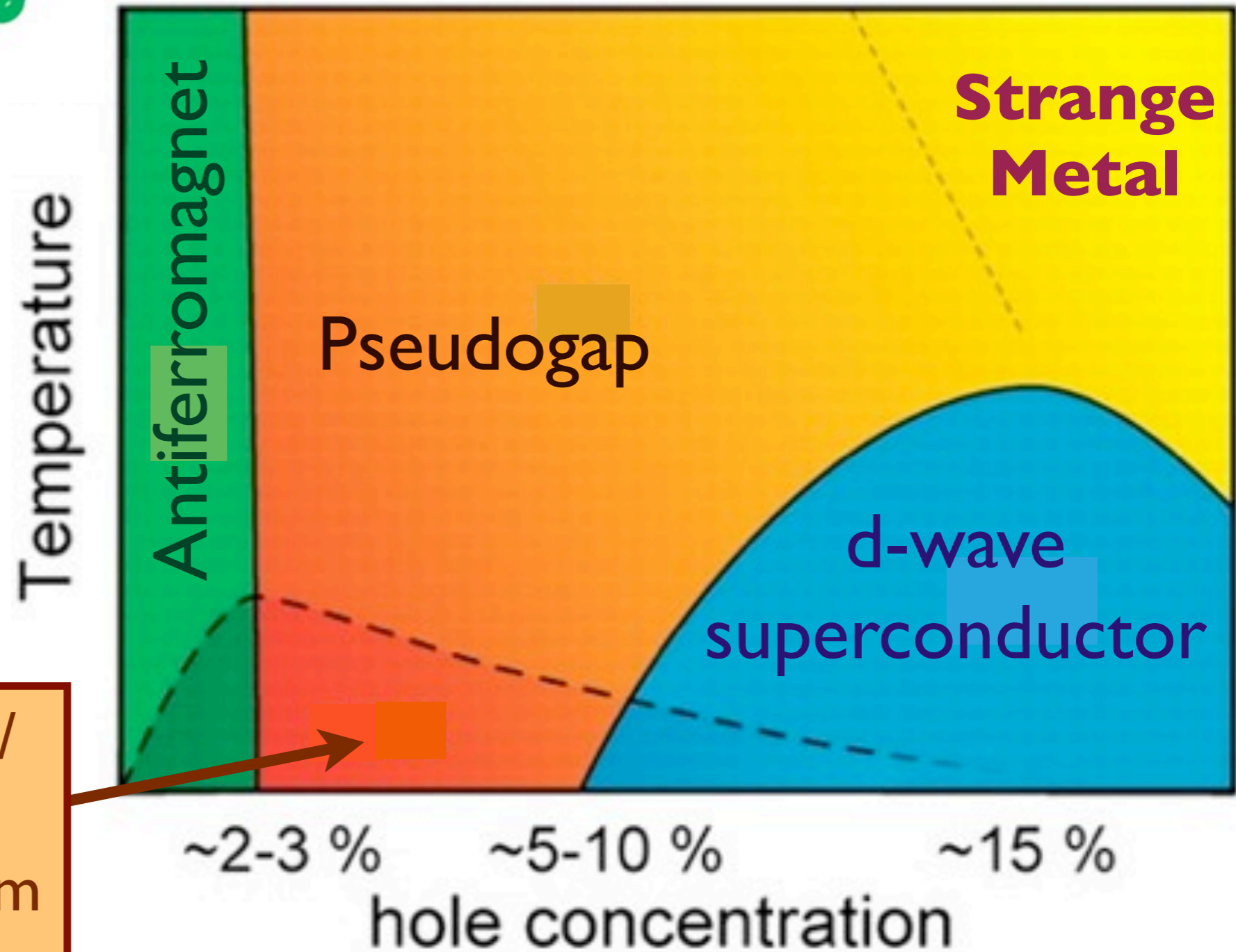
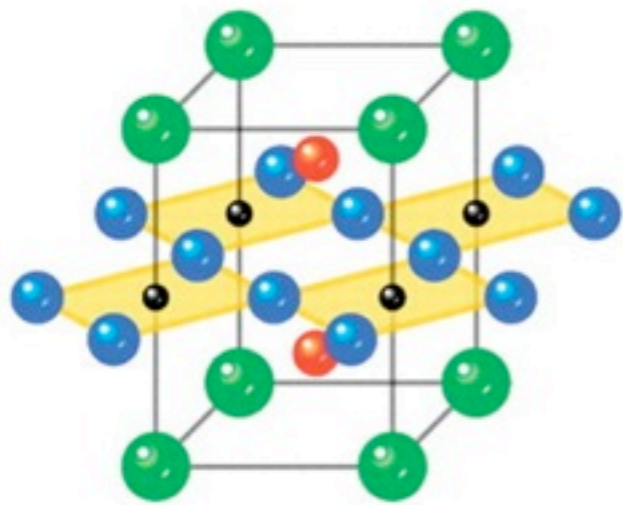


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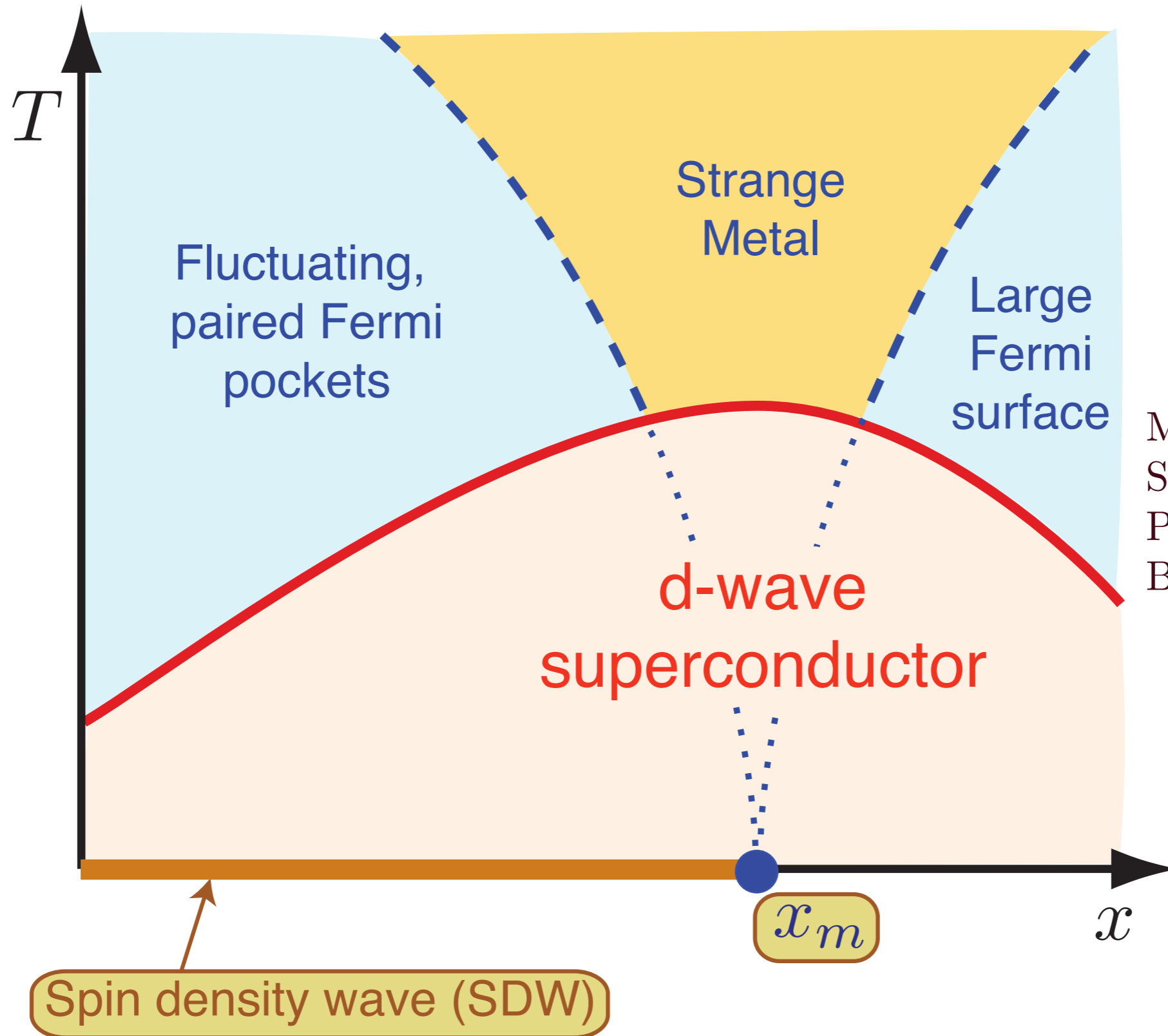
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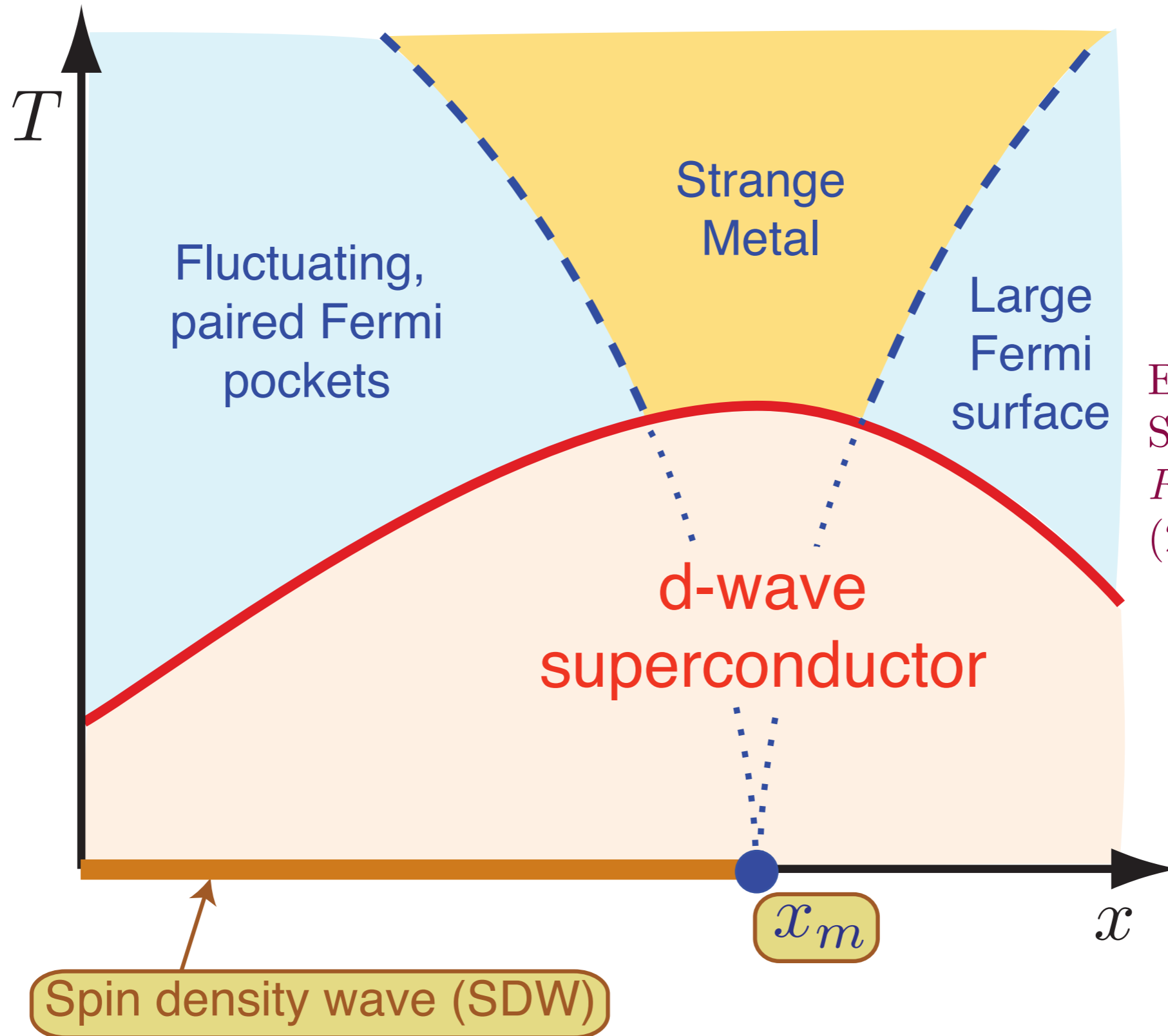
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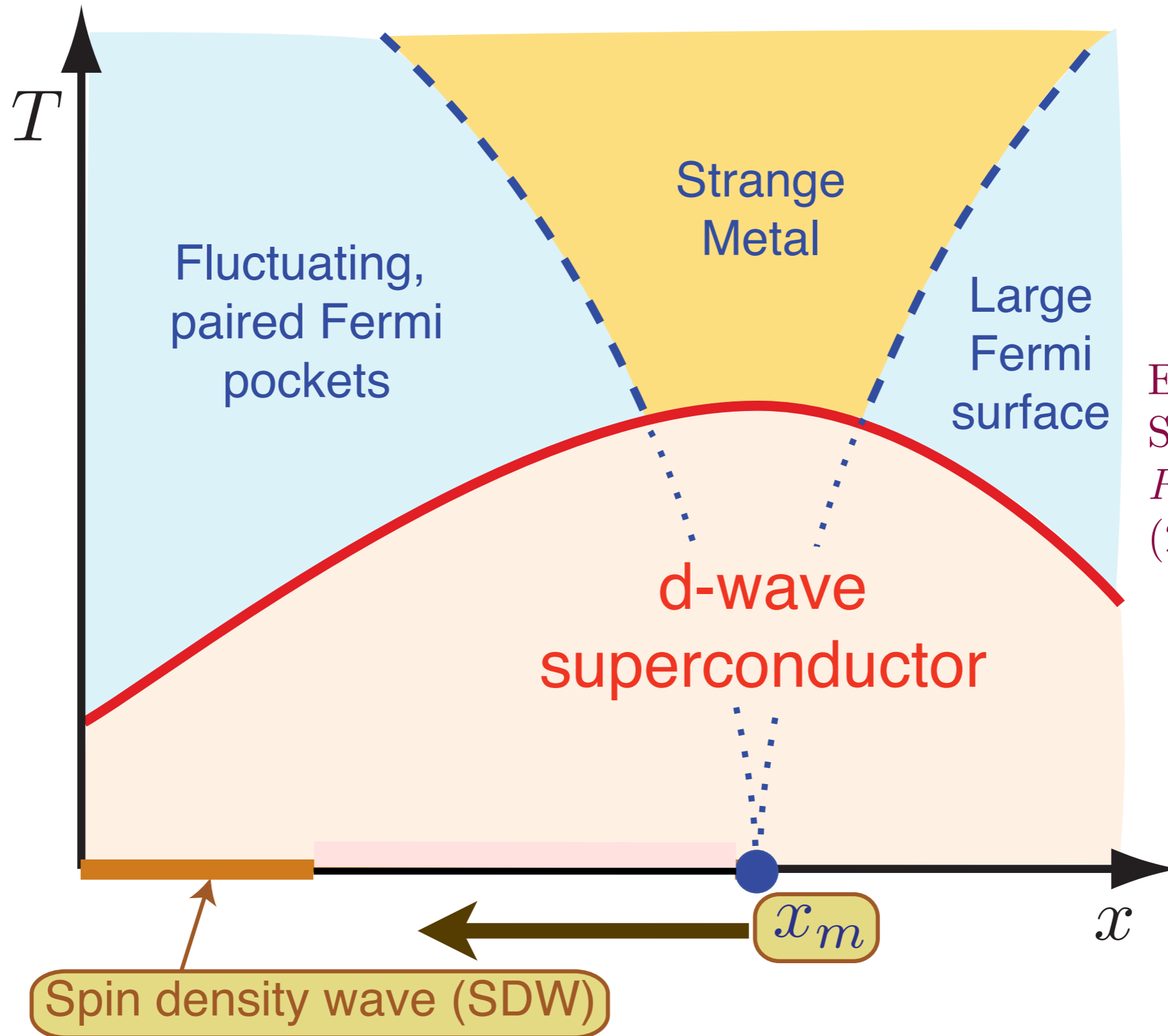
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E. G. Moon and S. Sachdev, *Phy. Rev. B* **80**, 035117 (2009)

Competition between SDW order and superconductivity moves the actual quantum critical point to  $x = x_s < x_m$ .

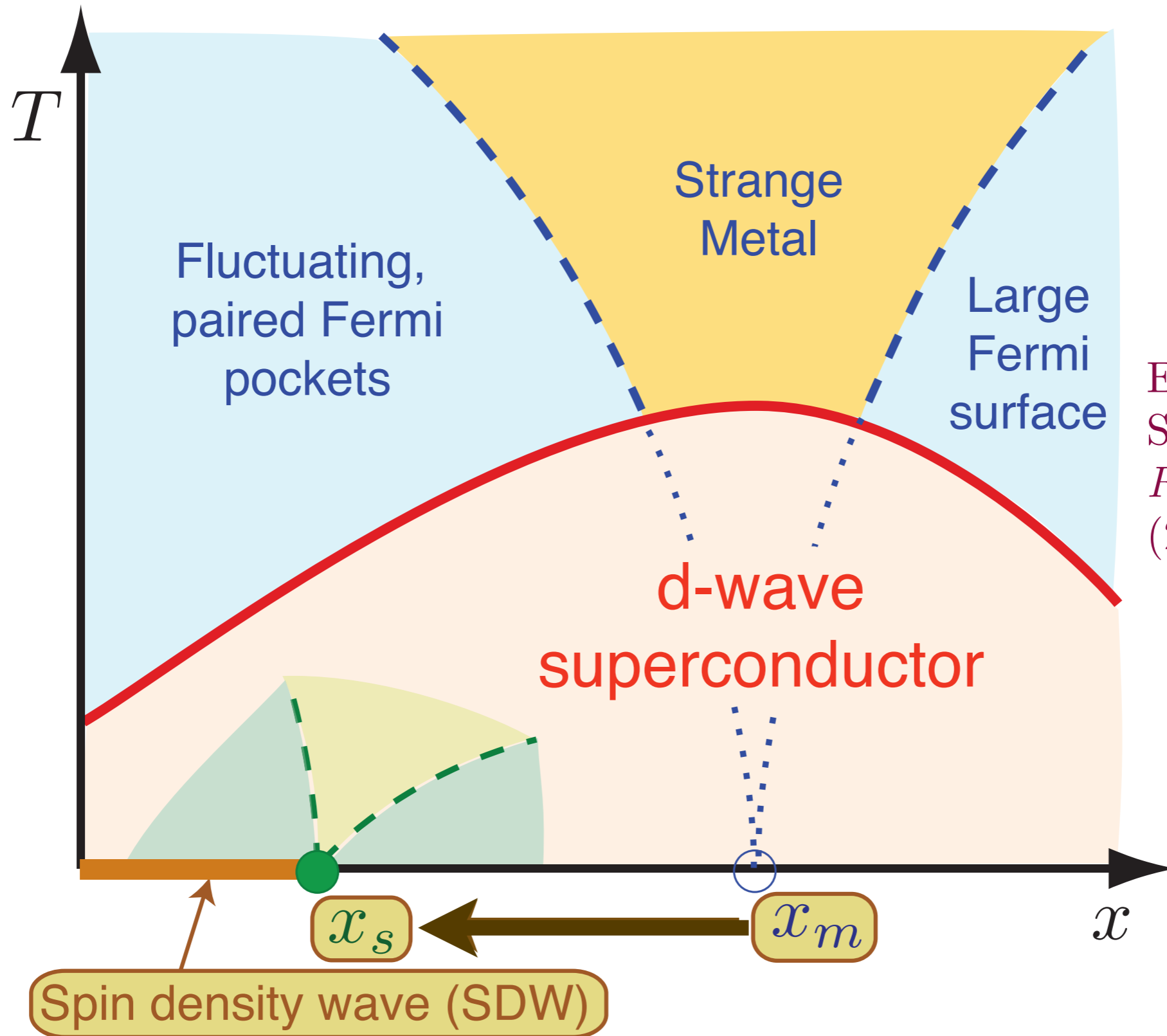
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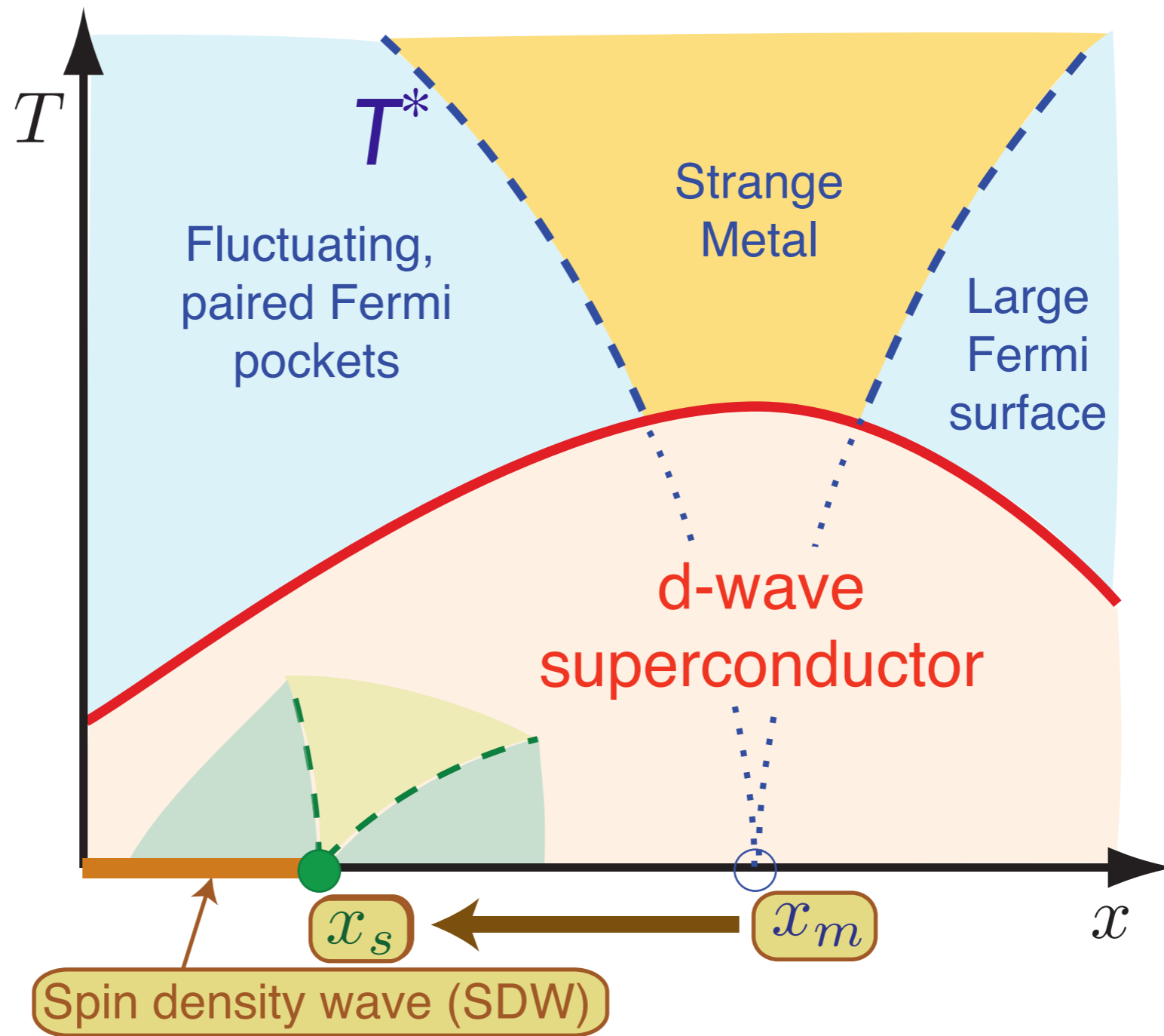
# Theory of quantum criticality in the cuprates



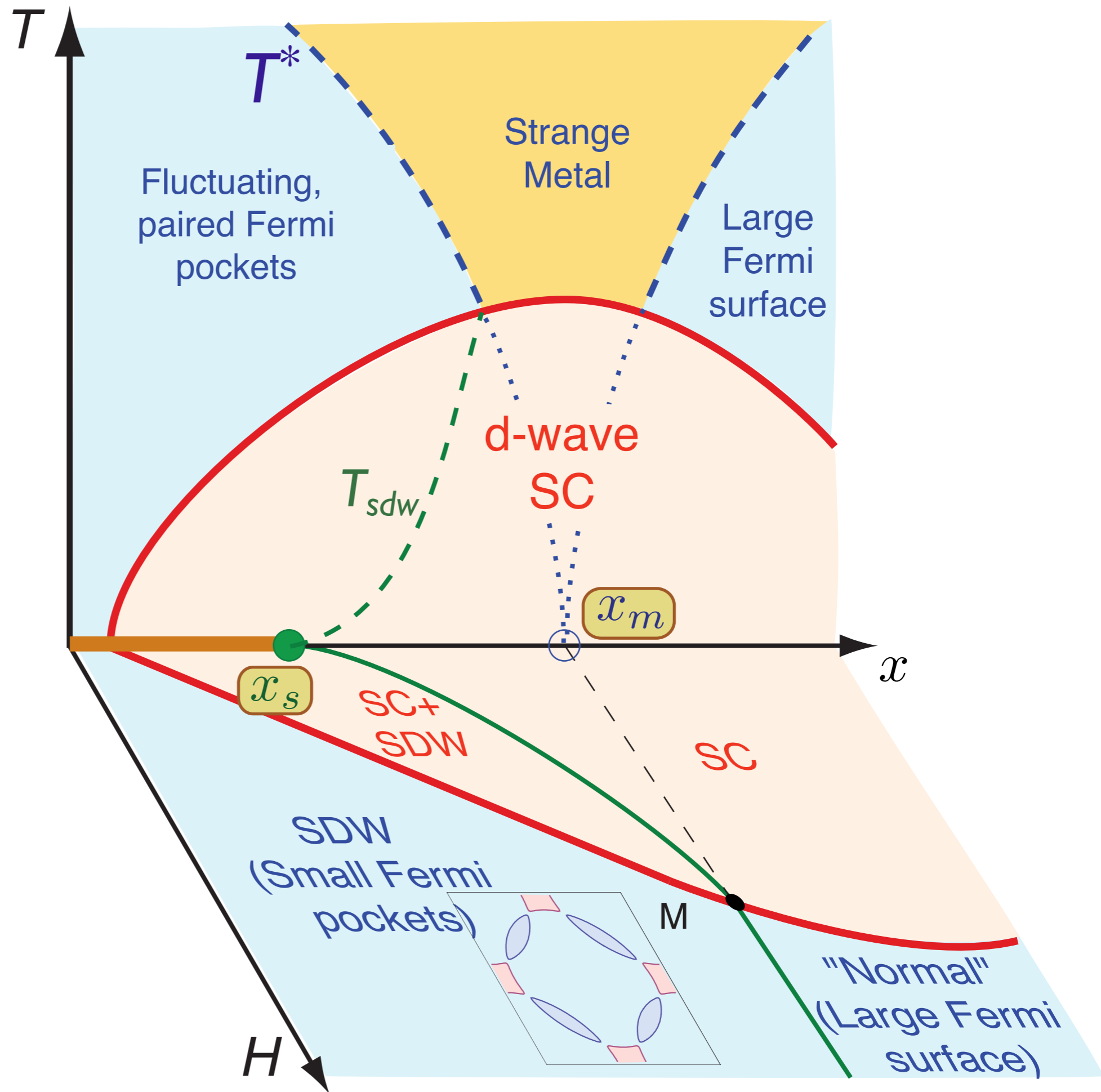
E. G. Moon and S. Sachdev, *Phy. Rev. B* **80**, 035117 (2009)

Competition between SDW order and superconductivity moves the actual quantum critical point to  $x = x_s < x_m$ .



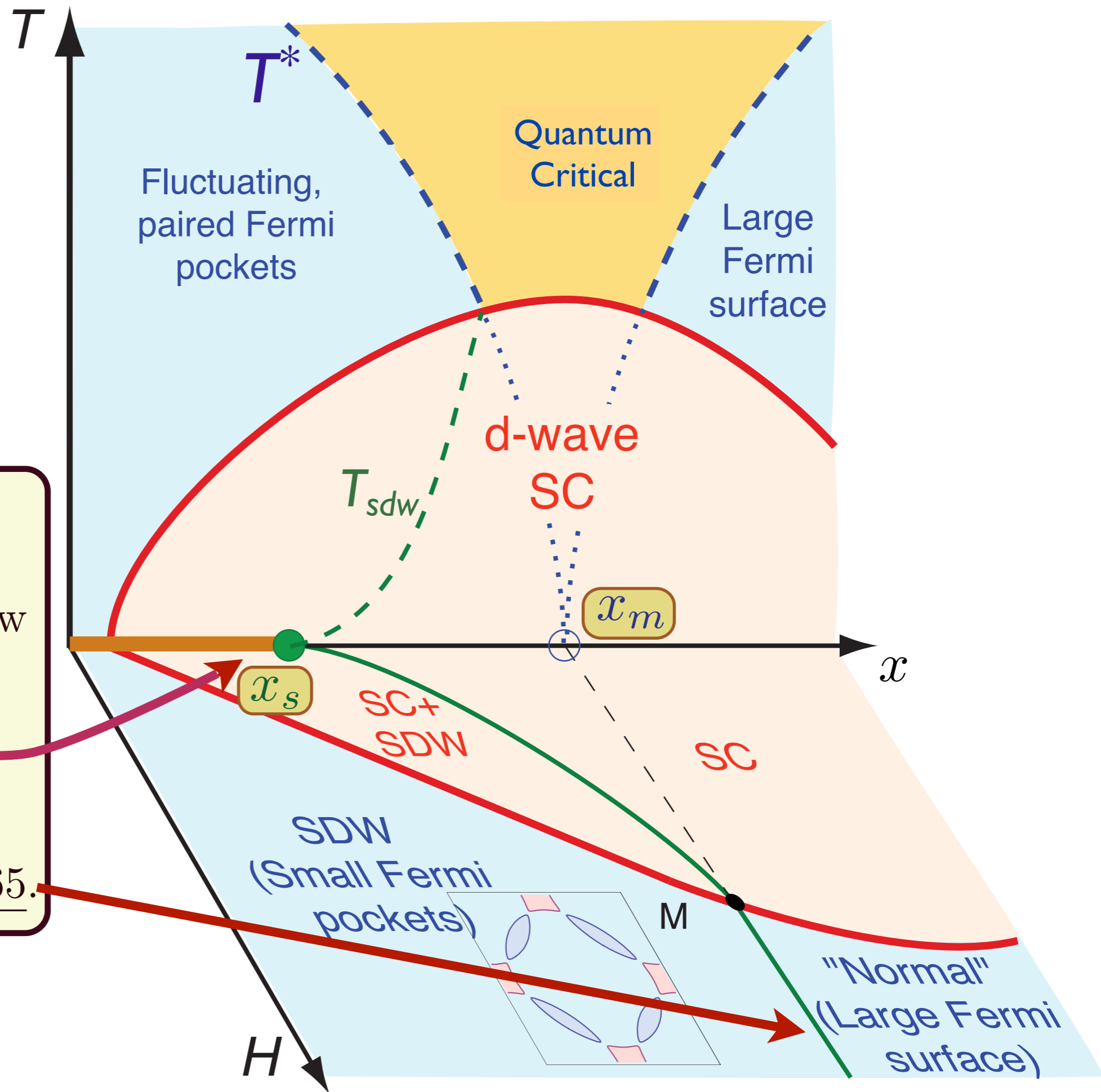


E. Demler, S. Sachdev  
and Y. Zhang, *Phys.  
Rev. Lett.* **87**,  
067202 (2001).



E. Demler, S. Sachdev  
and Y. Zhang, *Phys. Rev. Lett.* **87**,  
067202 (2001).

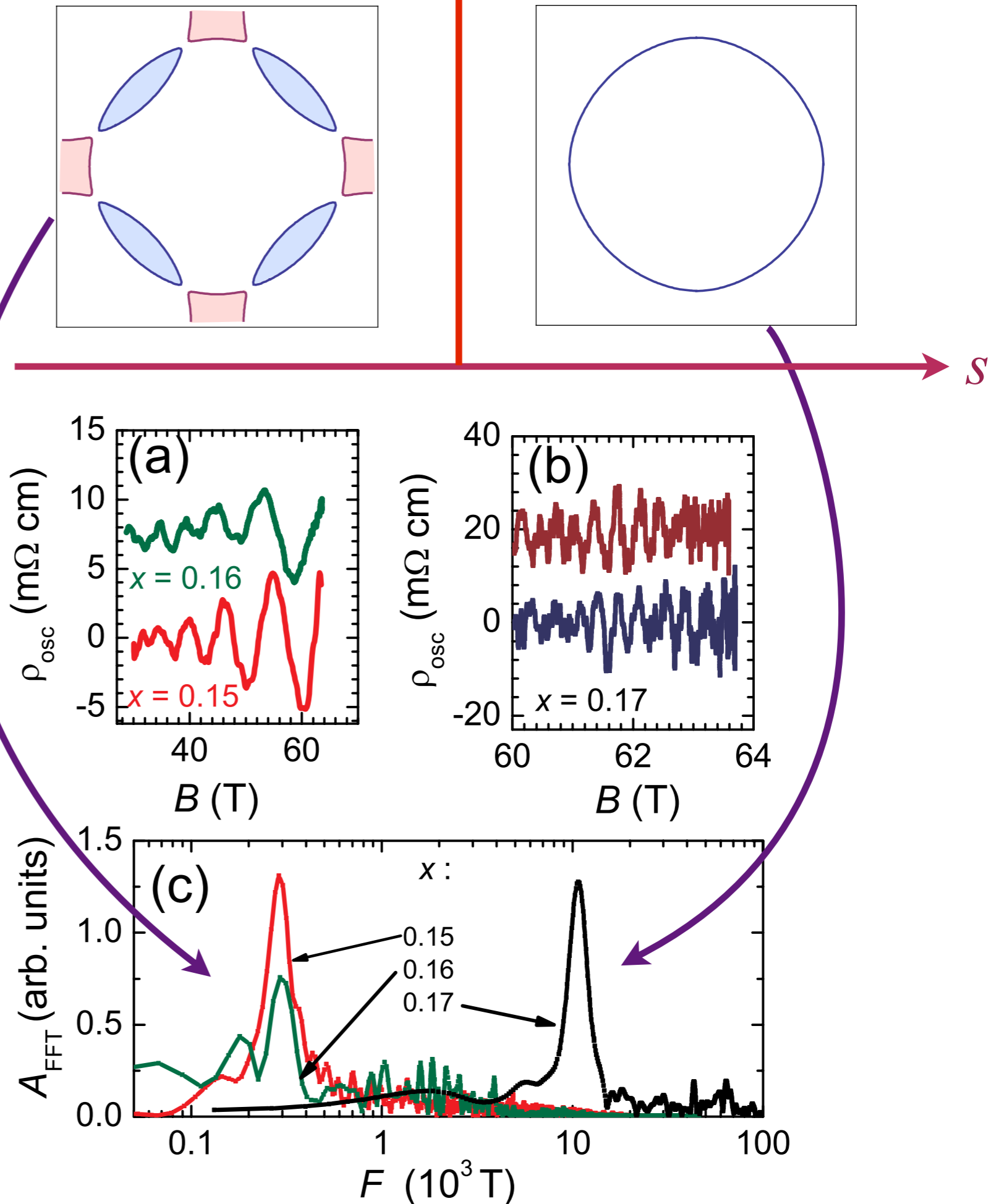
Neutron scattering experiments on  $\text{Nd}_{2-x}\text{Ce}_x\text{CuO}_4$  show that at low fields  $x_s = 0.14$ , while quantum oscillations at high fields show that  $x_m = 0.165$ .

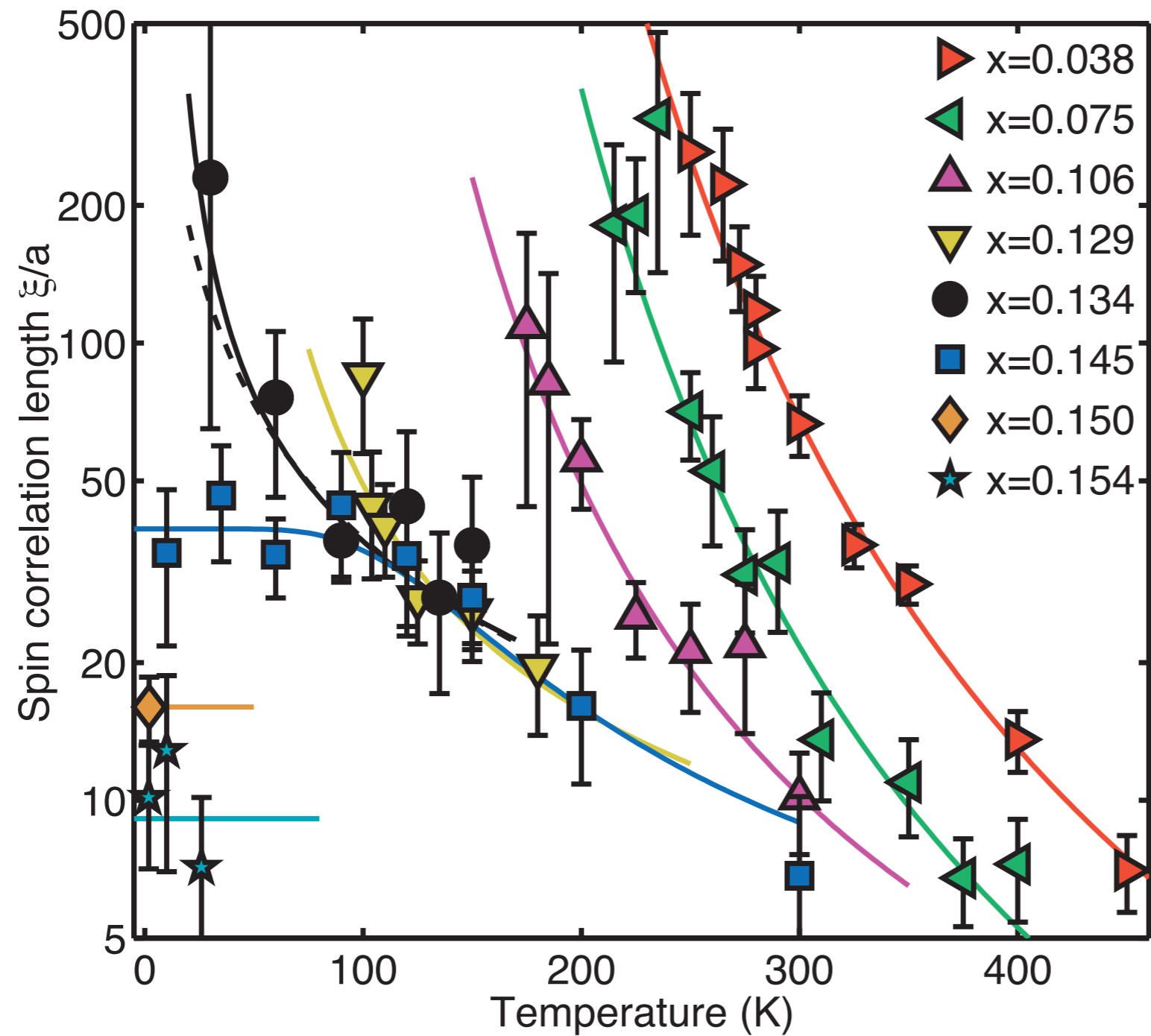


# Quantum oscillations



T. Helm, M.V. Kartsovnik,  
M. Bartkowiak, N. Bittner,  
M. Lambacher, A. Erb, J. Wosnitza,  
and R. Gross,  
*Phys. Rev. Lett.* **103**, 157002 (2009).

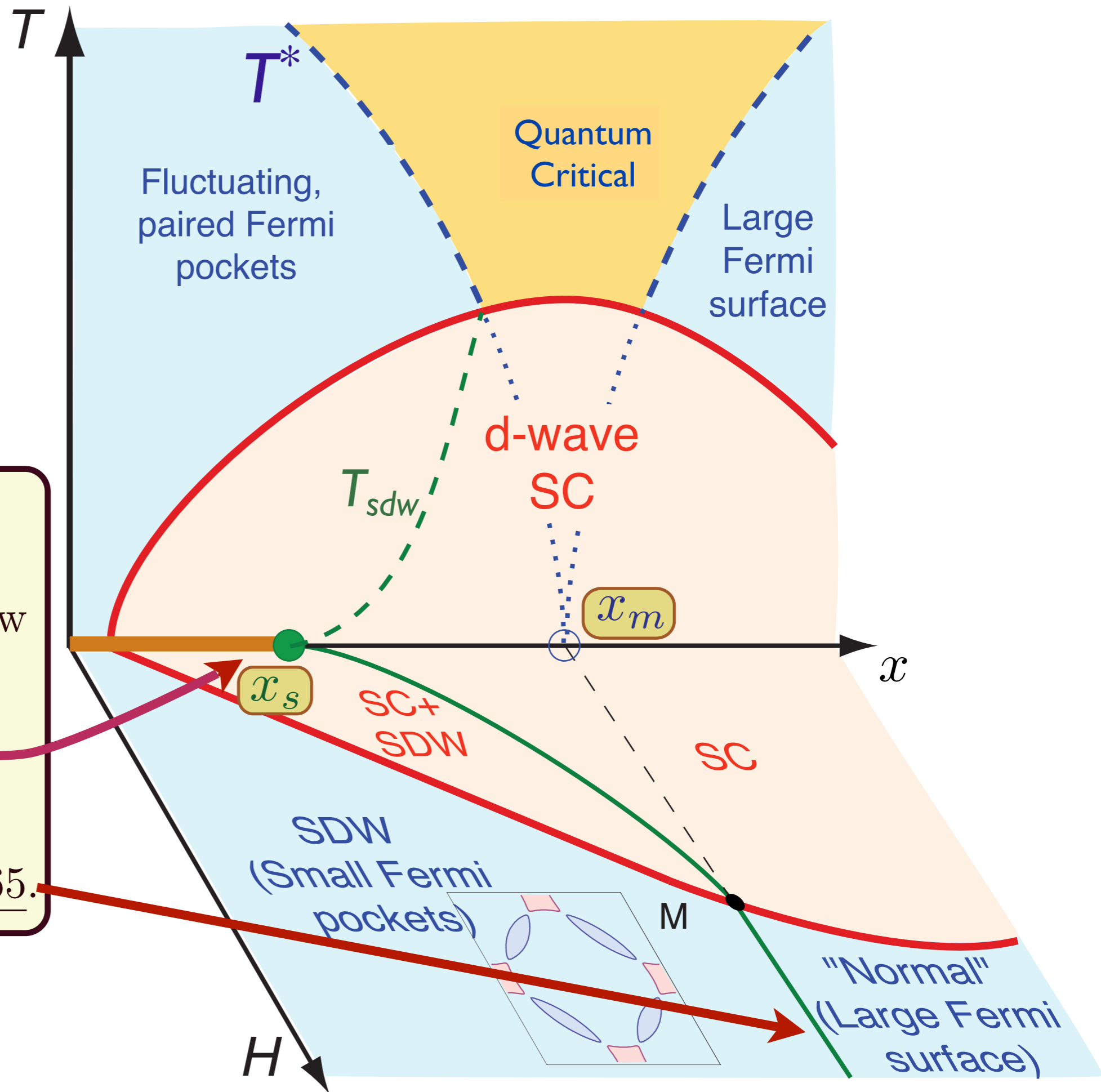




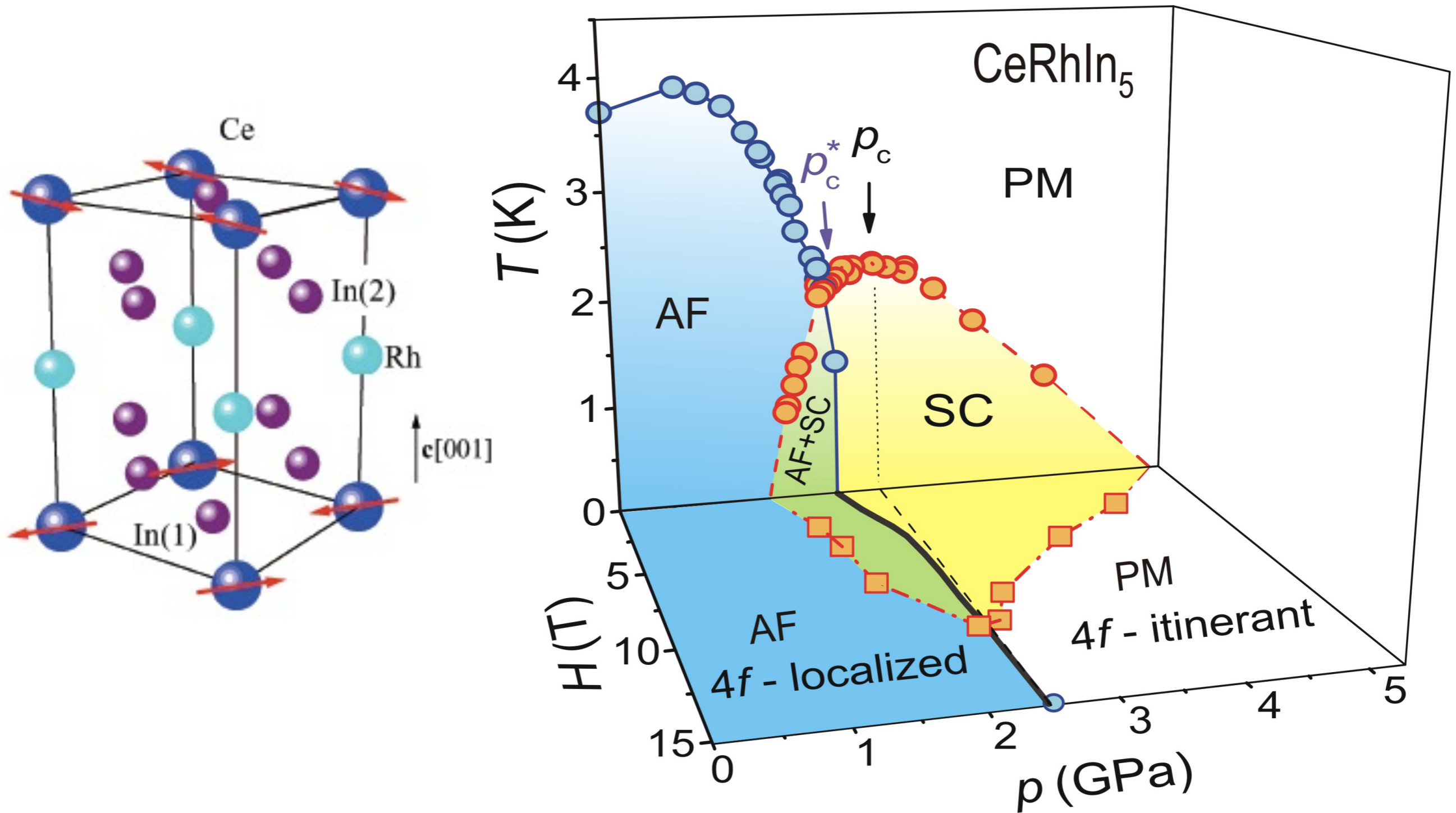
E. M. Motoyama, G. Yu, I. M. Vishik, O. P. Vajk, P. K. Mang, and M. Greven,  
*Nature* **445**, 186 (2007).

E. Demler, S. Sachdev  
and Y. Zhang, *Phys. Rev. Lett.* **87**,  
067202 (2001).

Neutron scattering experiments on  $\text{Nd}_{2-x}\text{Ce}_x\text{CuO}_4$  show that at low fields  $x_s = 0.14$ , while quantum oscillations at high fields show that  $x_m = 0.165$ .



# Similar phase diagram for CeRhIn<sub>5</sub>

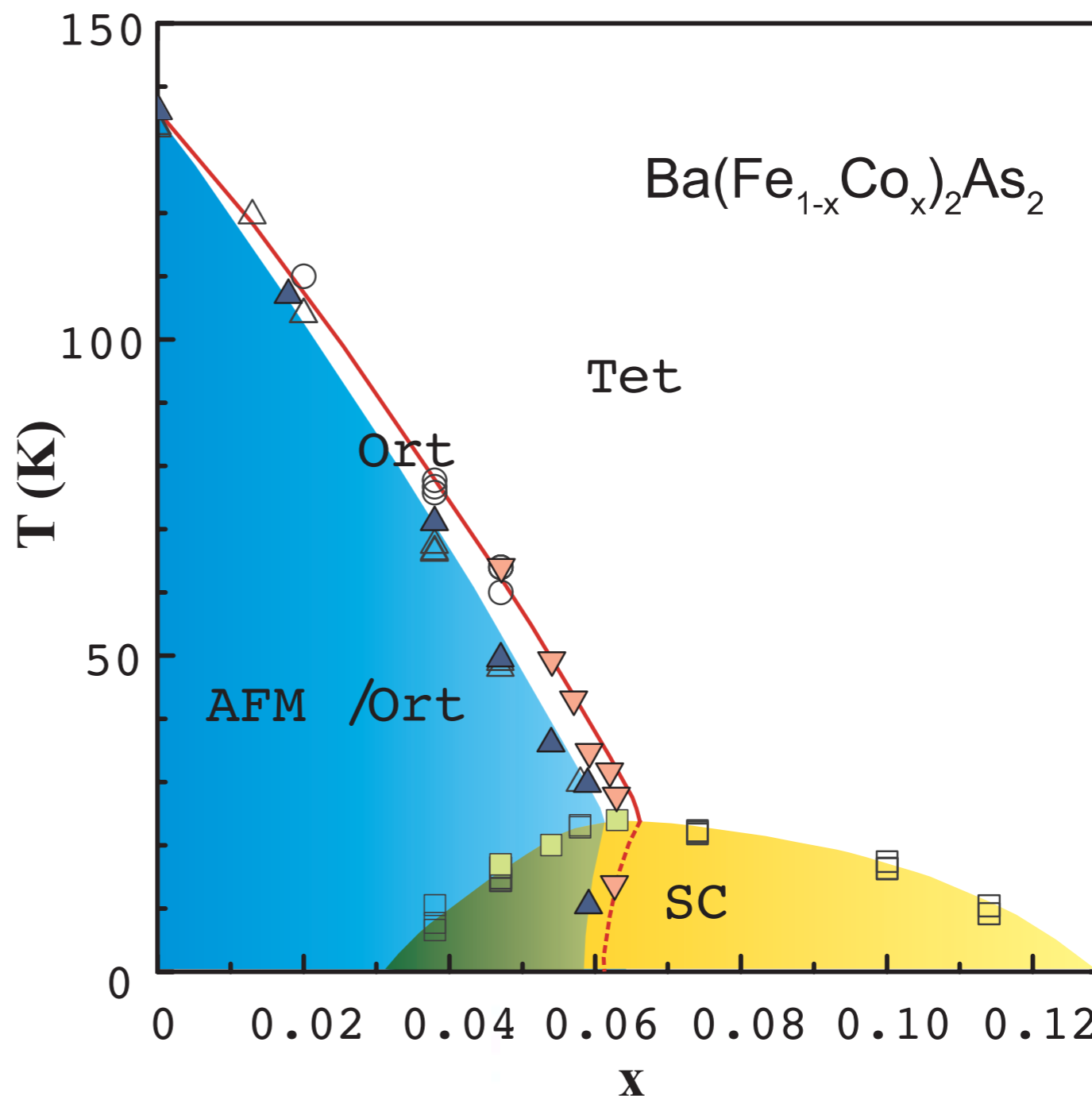
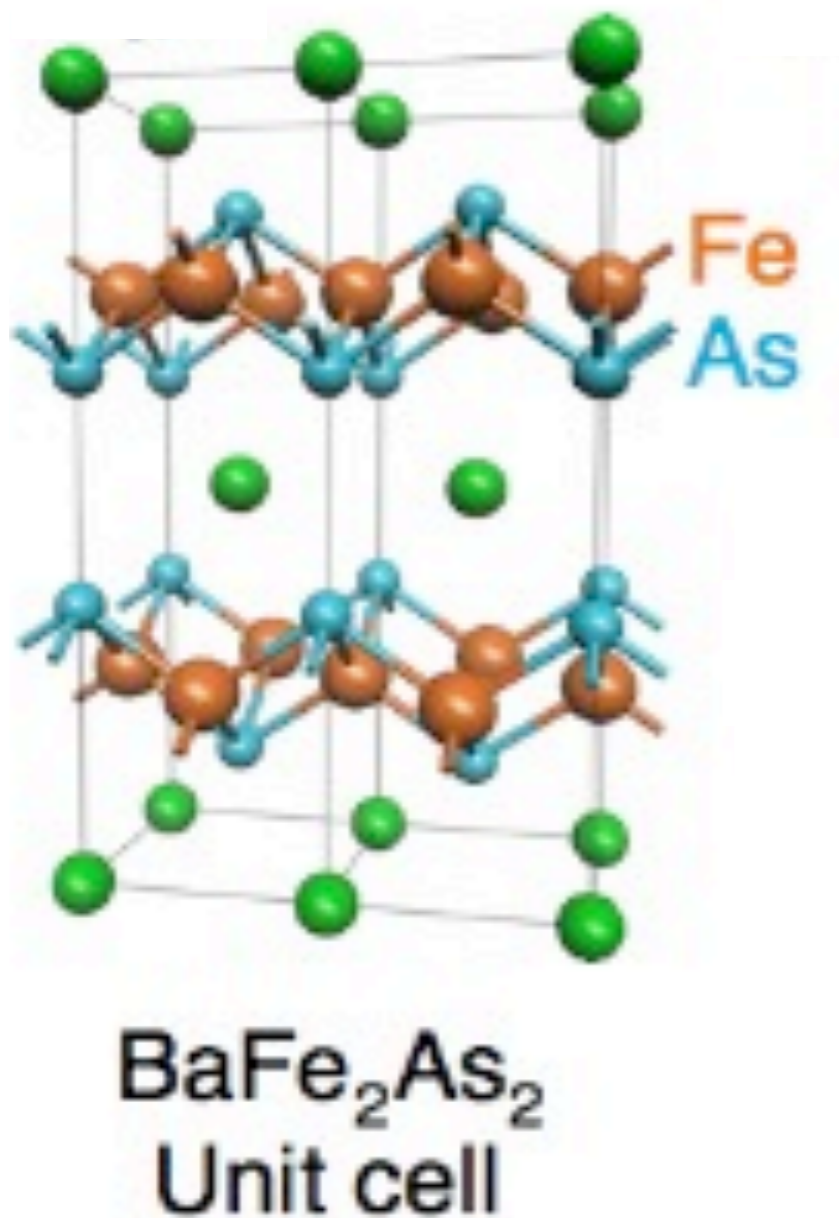


G. Knebel, D. Aoki, and J. Flouquet, arXiv:0911.5223.

Tuson Park, F. Ronning, H. Q. Yuan, M. B. Salamon, R. Movshovich, J. L. Sarrao, and J. D. Thompson, *Nature* **440**, 65 (2006)

# Iron pnictides:

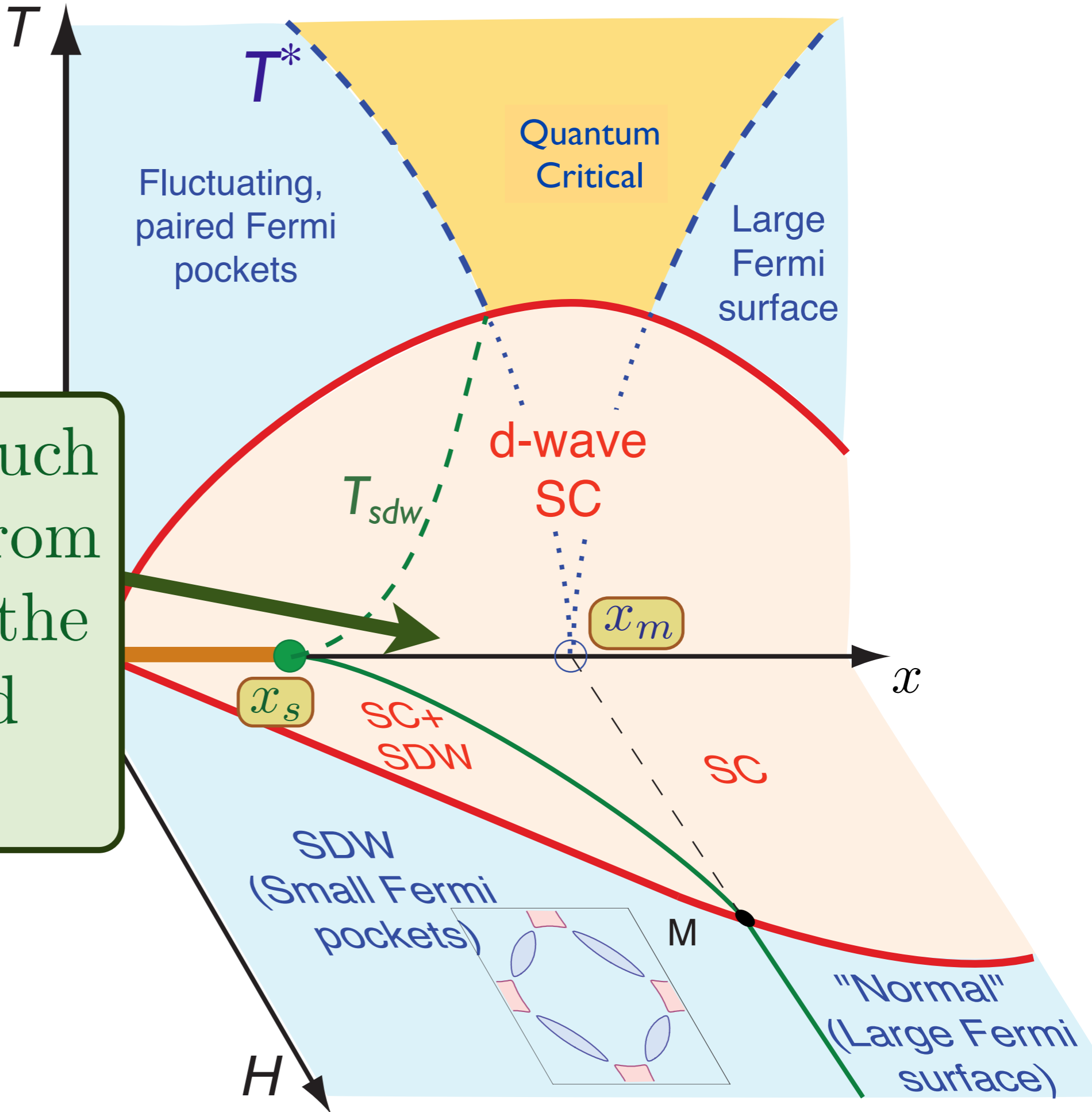
a new class of high temperature superconductors



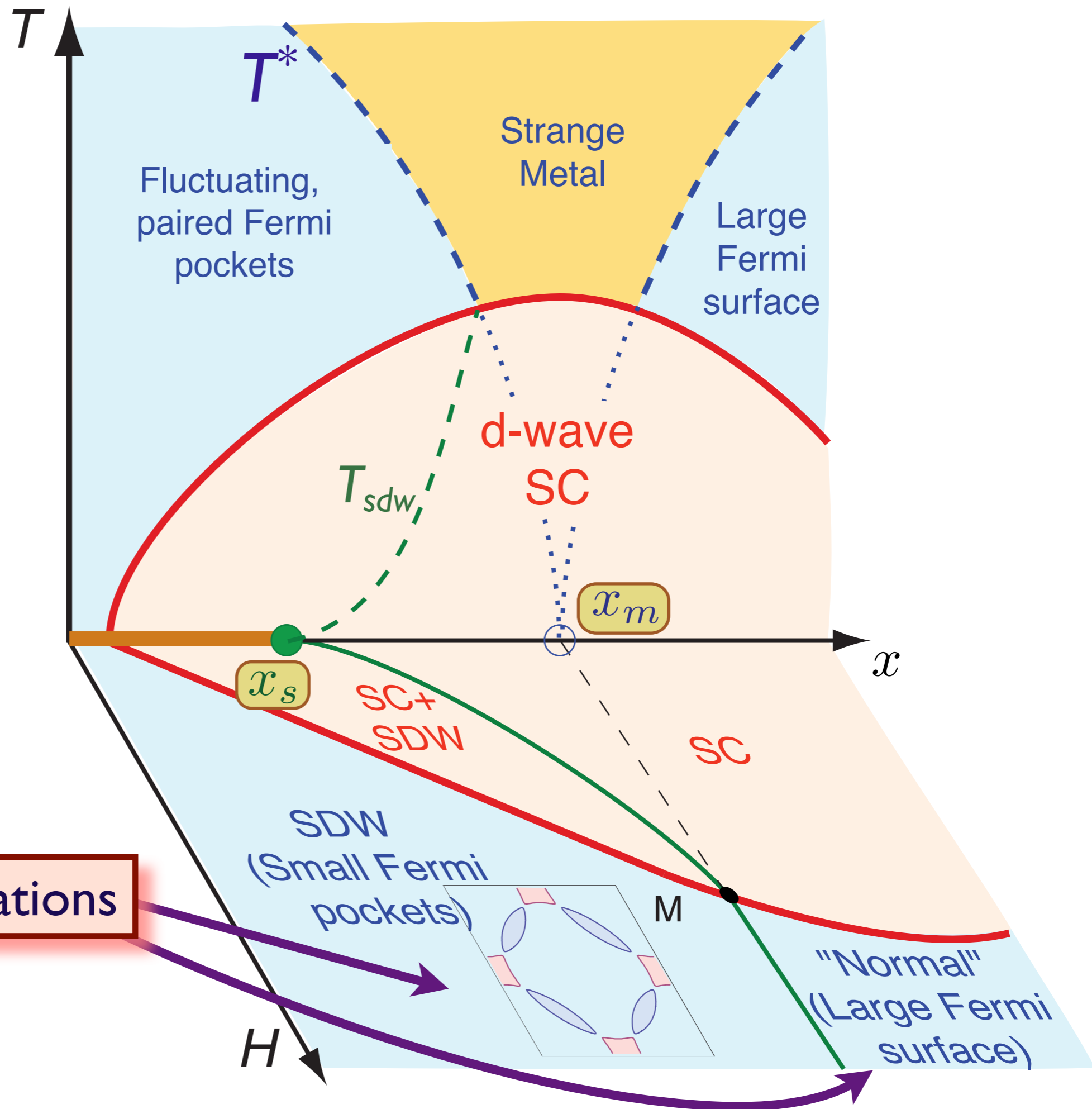
S. Nandi, M. G. Kim, A. Kreyssig, R. M. Fernandes, D. K. Pratt, A. Thaler, N. Ni,  
S. L. Bud'ko, P. C. Canfield, J. Schmalian, R. J. McQueeney, A. I. Goldman,  
*Physical Review Letters* **104**, 057006 (2010).



There is a much larger shift from  $x_m$  to  $x_s$  in the hole-doped cuprates.

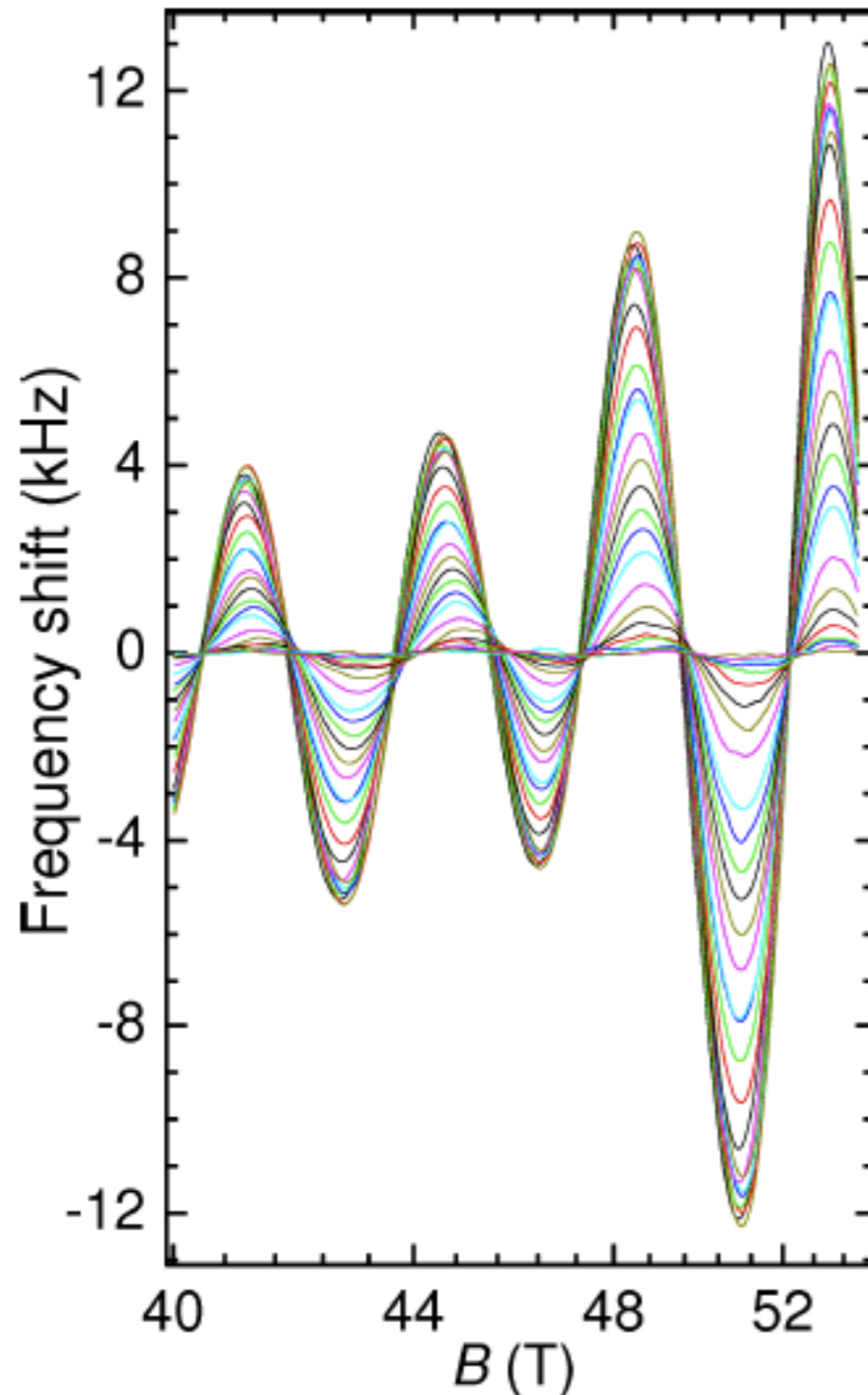


E. Demler, S. Sachdev  
and Y. Zhang, *Phys.  
Rev. Lett.* **87**,  
067202 (2001).



Quantum oscillations

# Evidence for small Fermi pockets in hole-doped cuprates



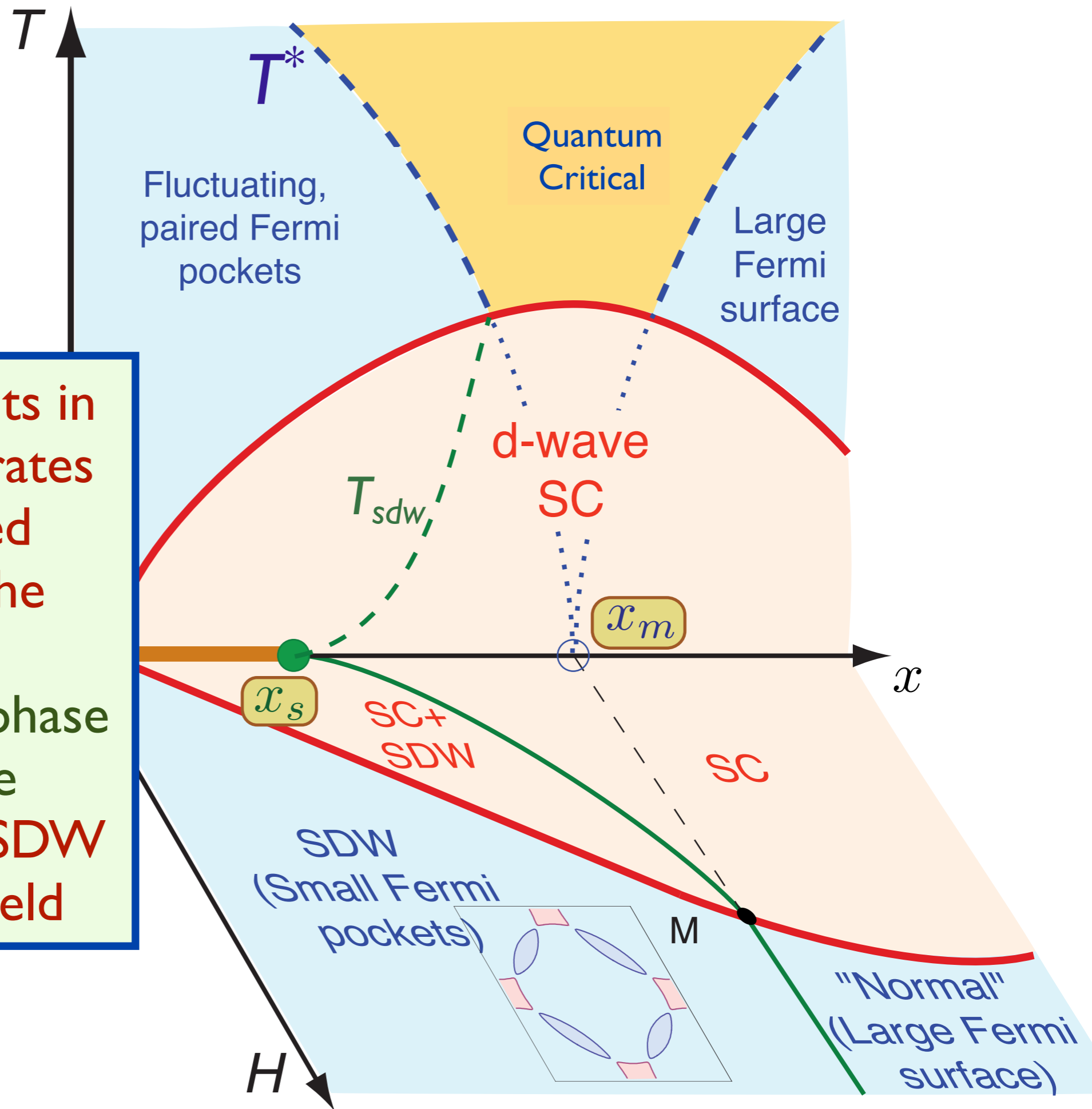
Suchitra E. Sebastian, N. Harrison,  
M. M. Altarawneh, Ruixing Liang, D.A. Bonn,  
W. N. Hardy, and G. G. Lonzarich  
*Physical Review B* **81**, 140505(R) (2010)

Original observation:  
N. Doiron-Leyraud, C. Proust,  
D. LeBoeuf, J. Levallois,  
J.-B. Bonnemaïson, R. Liang,  
D.A. Bonn, W. N. Hardy,  
and L. Taillefer,  
*Nature* **447**, 565 (2007)

FIG. 2: Magnetic quantum oscillations measured in  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  with  $x \approx 0.56$  (after background polynomial subtraction). This restricted interval in  $B = |\mathbf{B}|$  furnishes a dynamic range of  $\sim 50$  dB between  $T = 1$  and 18 K. The actual  $T$  values are provided in Fig. 3.

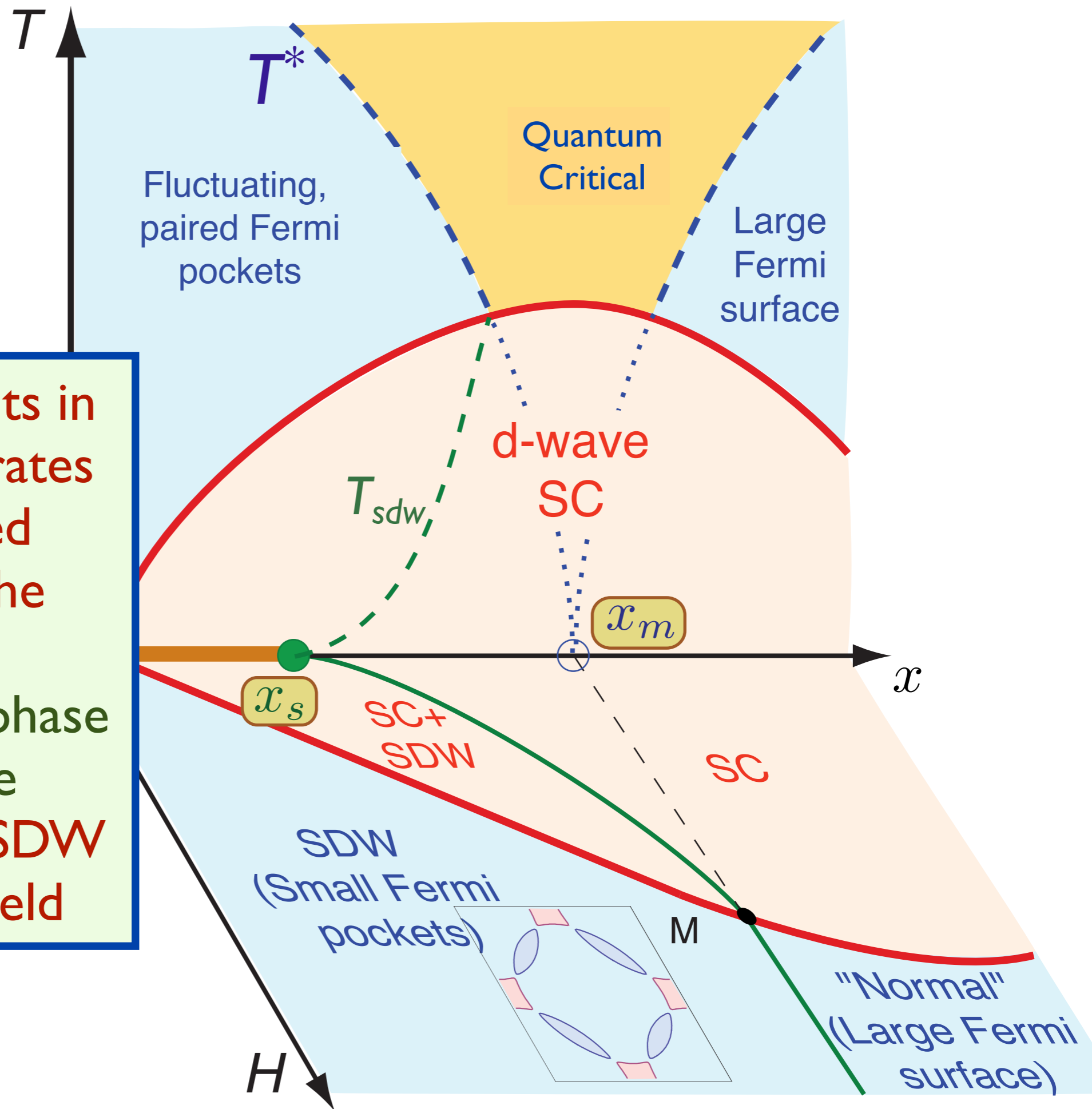
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Many experiments in  
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in a magnetic field



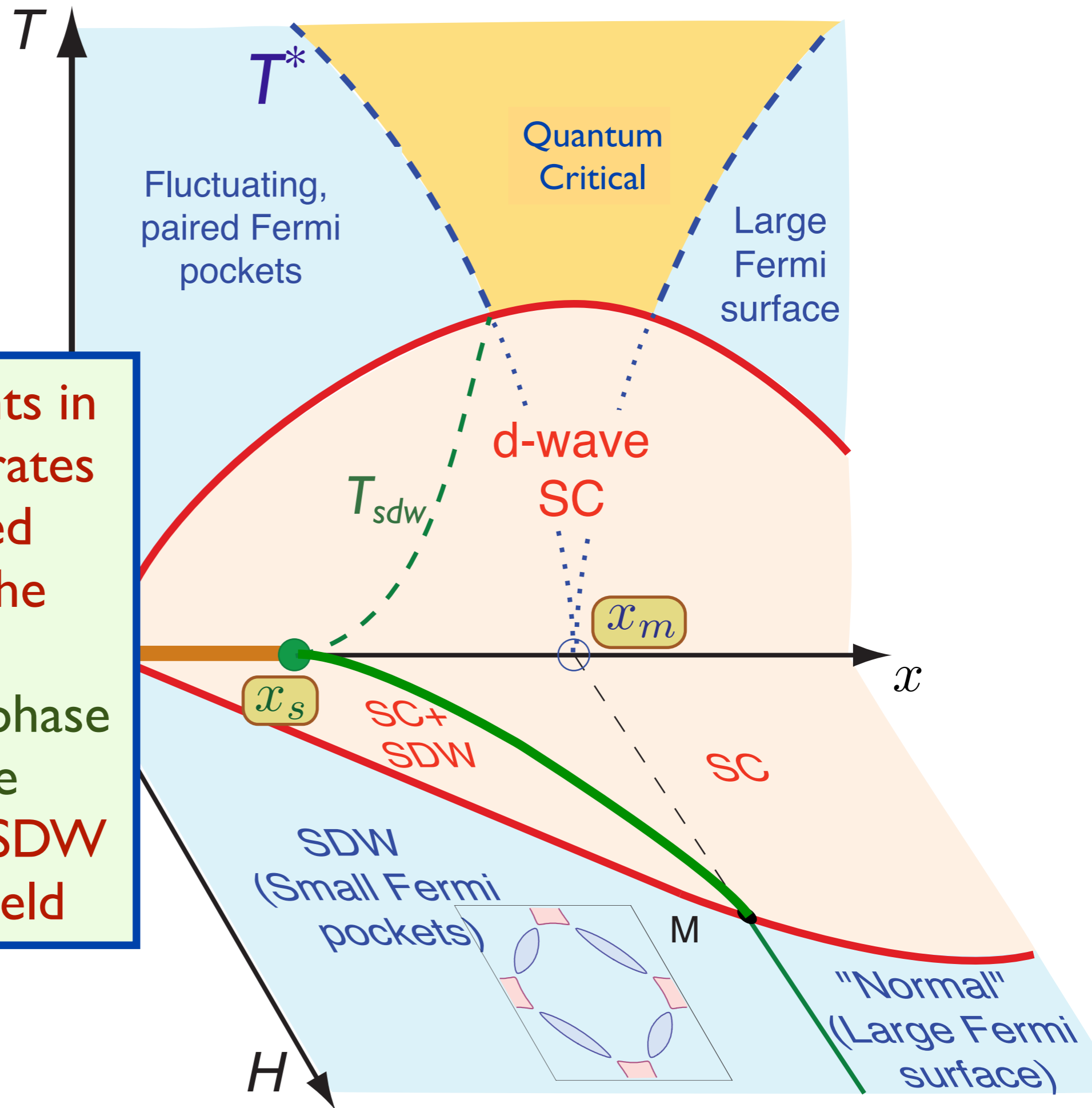
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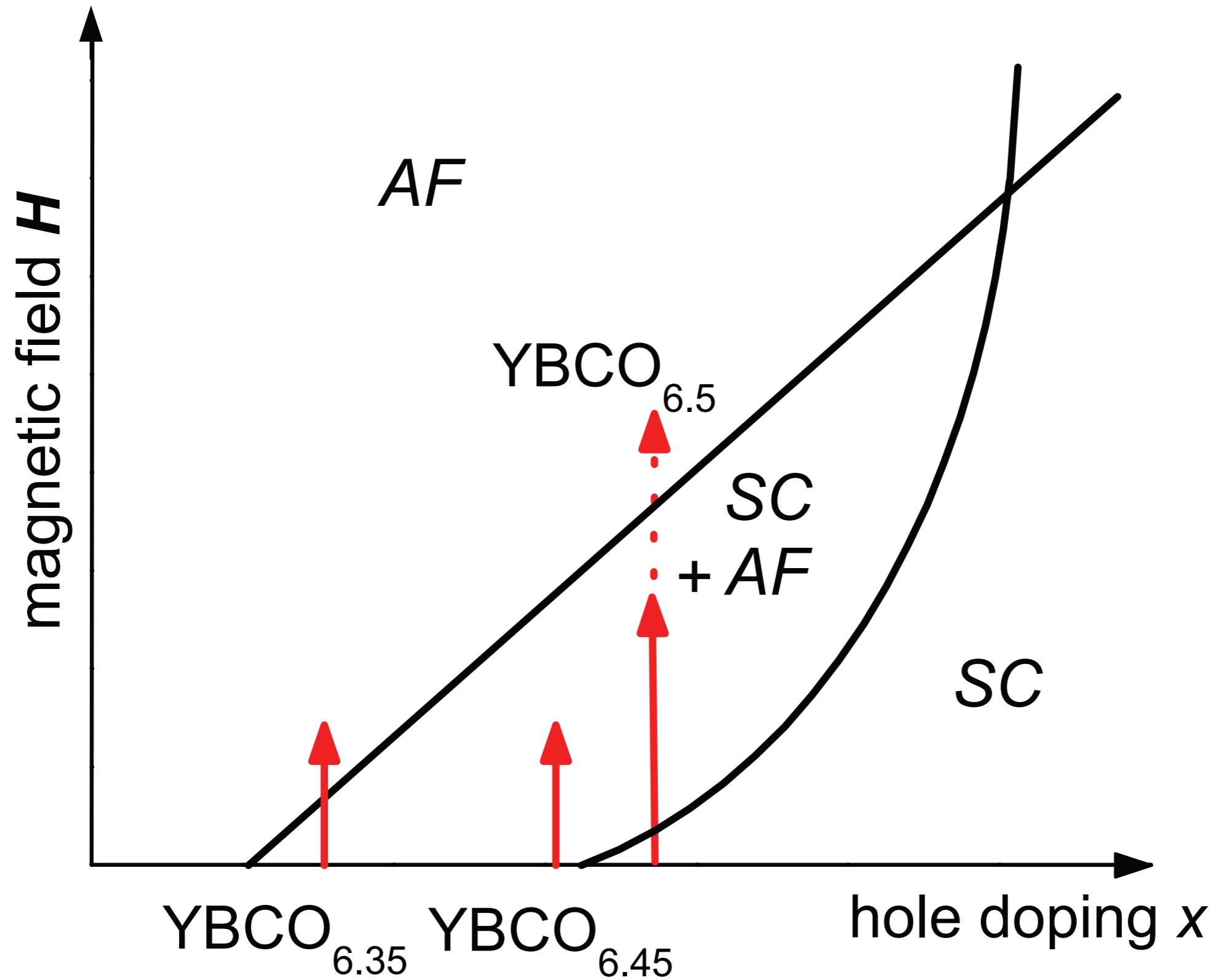
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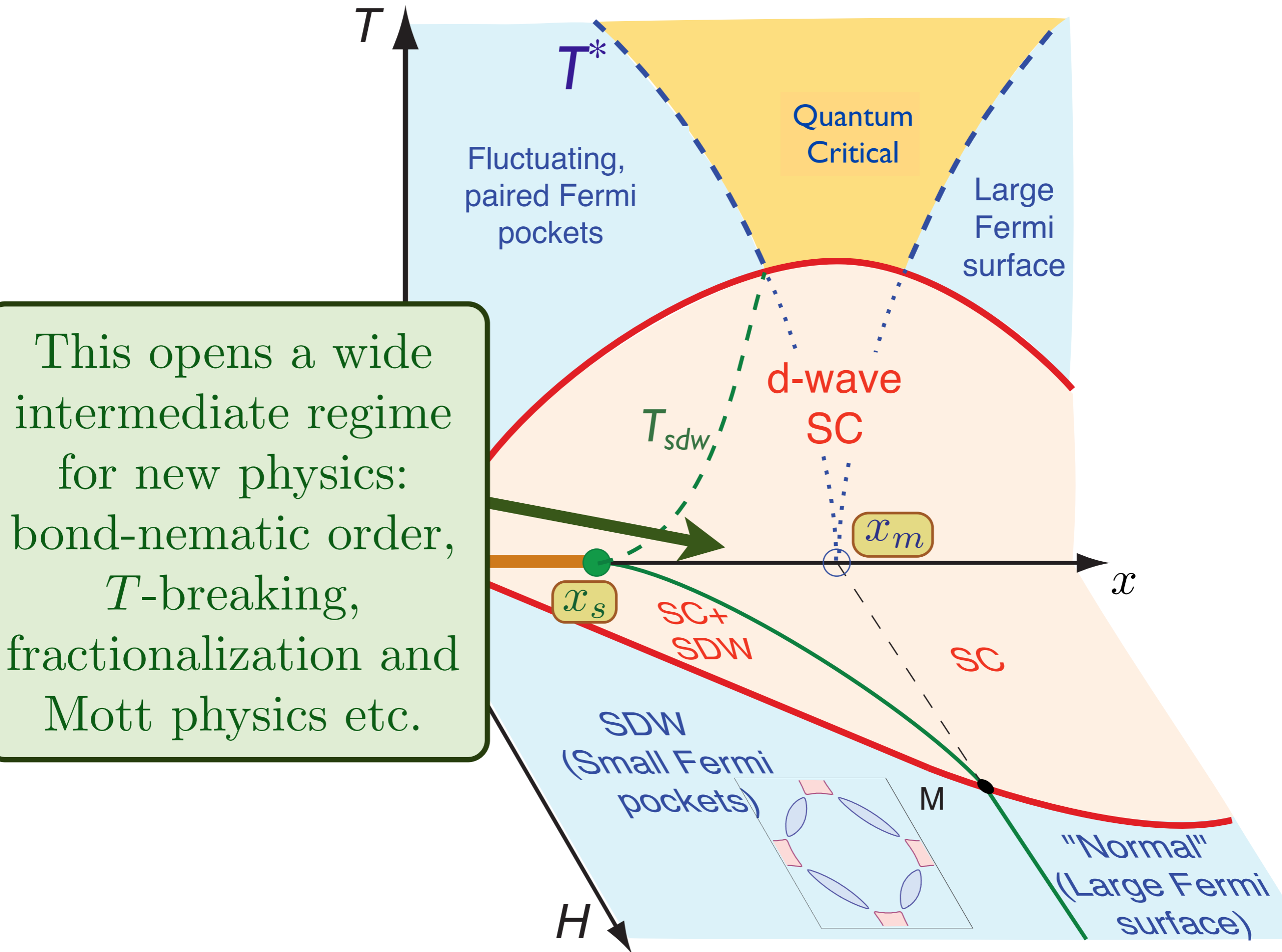
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D. Haug, V. Hinkov, Y. Sidis, P. Bourges, N. B. Christensen, A. Ivanov, T. Keller, C. T. Lin, and B. Keimer, *New J. Phys.* **12**, 105006 (2010)



This opens a wide intermediate regime for new physics: bond-nematic order,  $T$ -breaking, fractionalization and Mott physics etc.



# Outline

## 1. Phenomenology of the onset of antiferromagnetism in a metal

*Quantum criticality of Fermi surface reconstruction, and the phase diagram in a magnetic field*

## 2. Strongly-coupled quantum criticality in metals

*“Mechanism” of higher temperature superconductivity*

## 3. Theory of the competition between superconductivity and antiferromagnetism

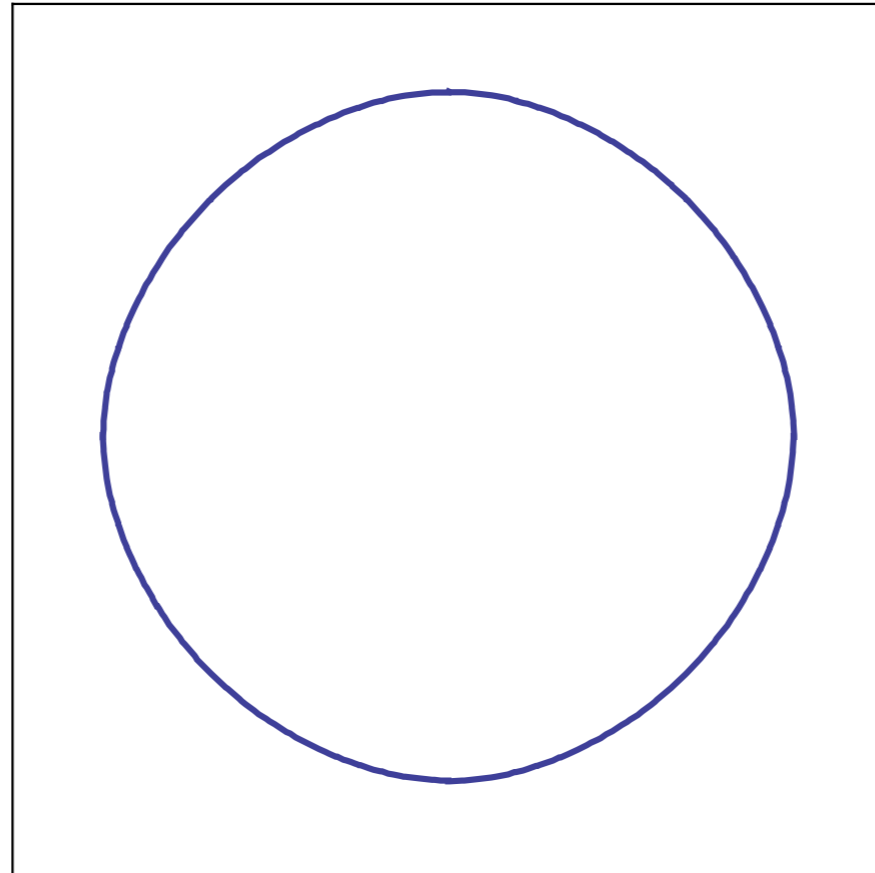
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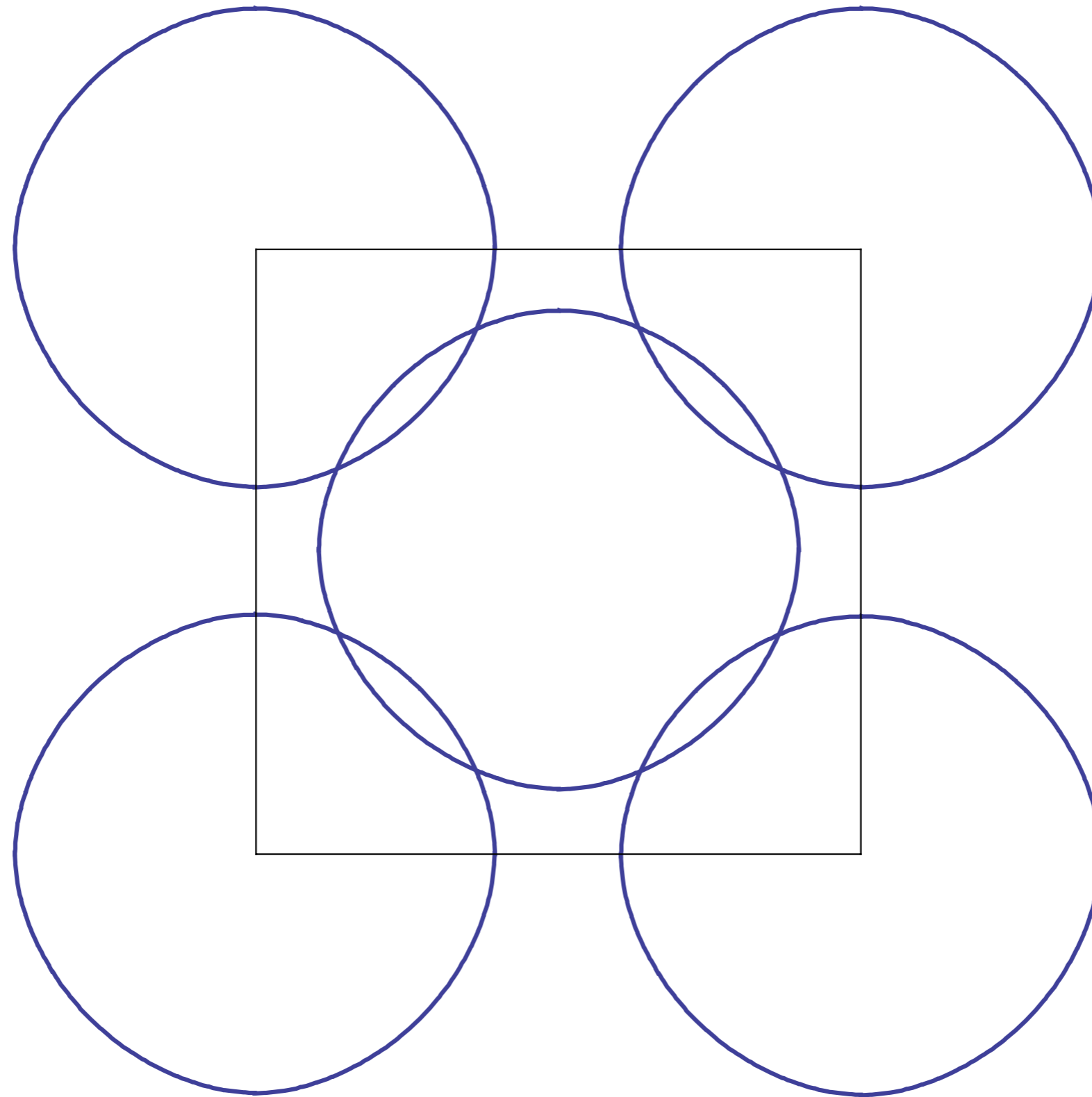
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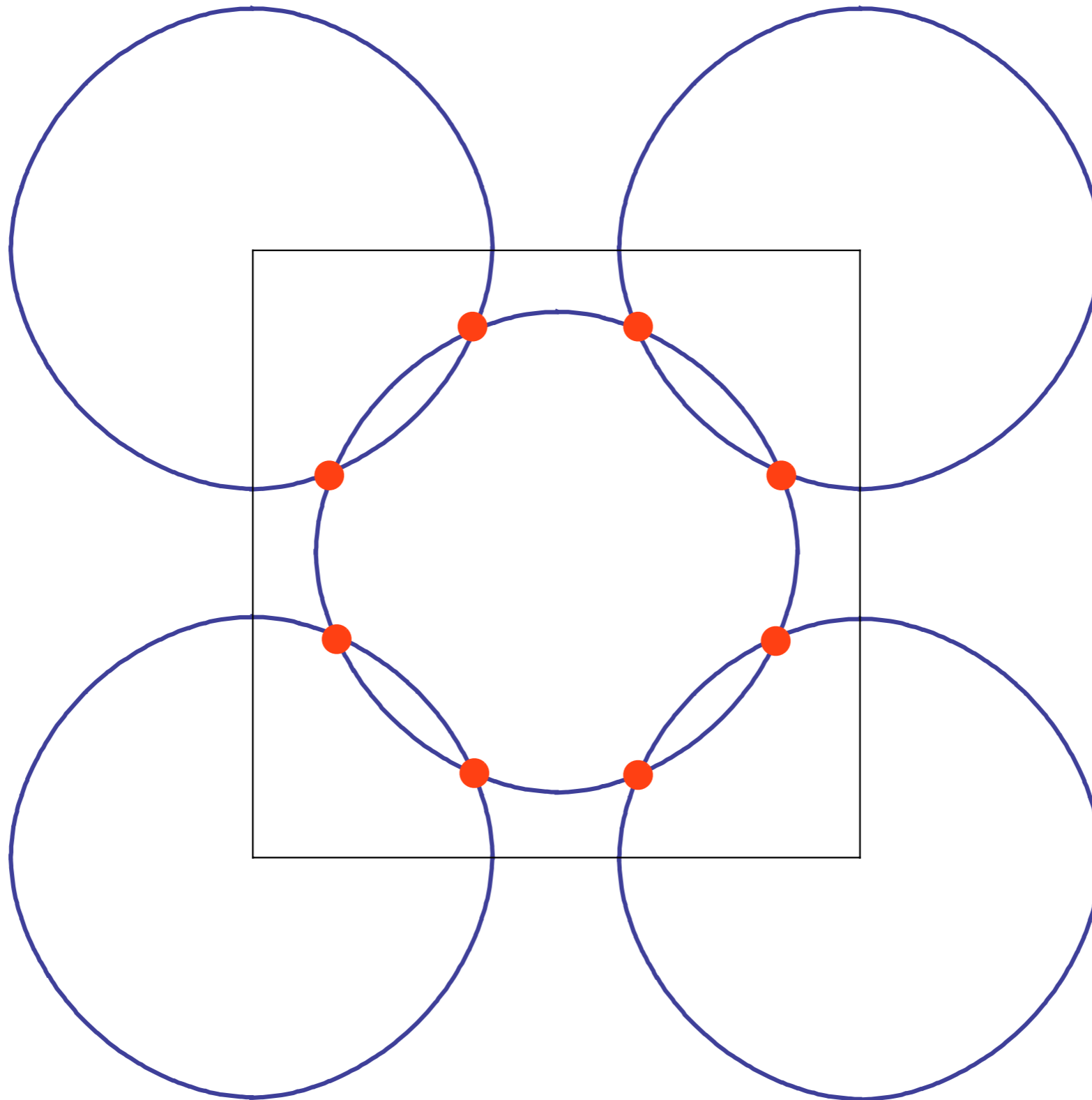
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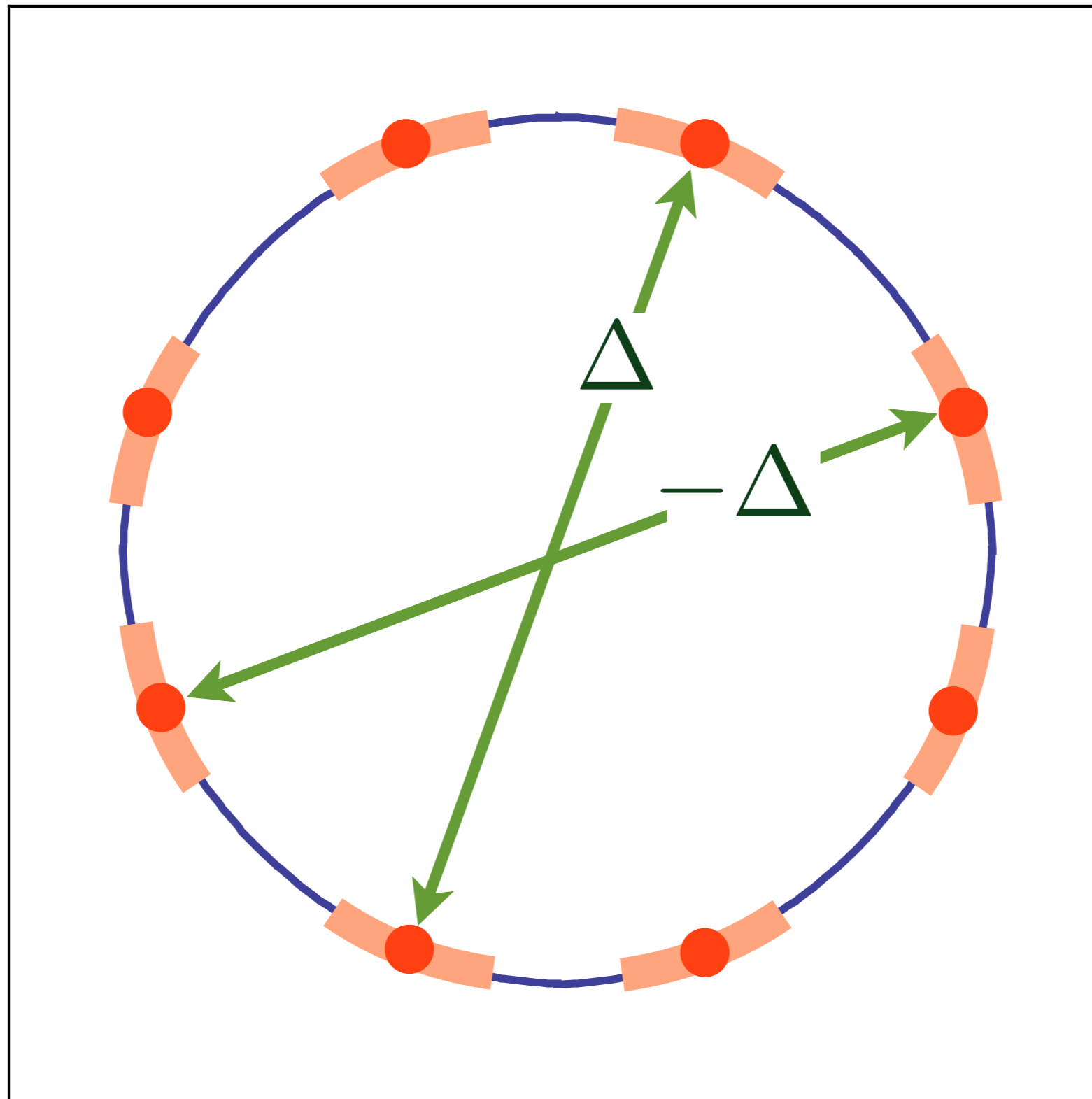
**Metal with “large” Fermi surface**



Fermi surfaces translated by  $\mathbf{K} = (\pi, \pi)$ .

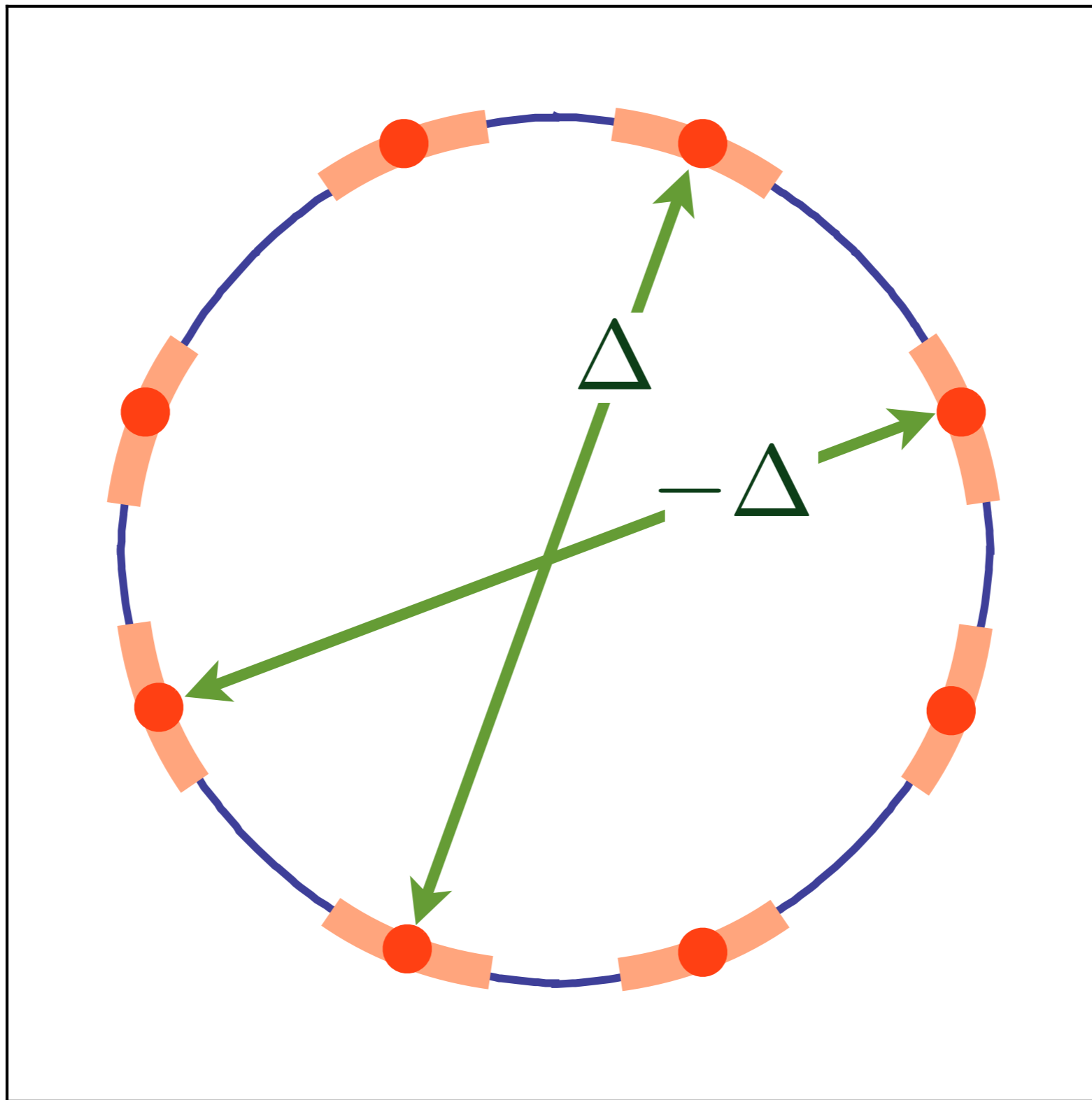


**“Hot” spots**



Unconventional pairing at and near hot spots

$$\langle c_{\mathbf{k}\alpha}^\dagger c_{-\mathbf{k}\beta}^\dagger \rangle = \varepsilon_{\alpha\beta} \Delta (\cos k_x - \cos k_y)$$



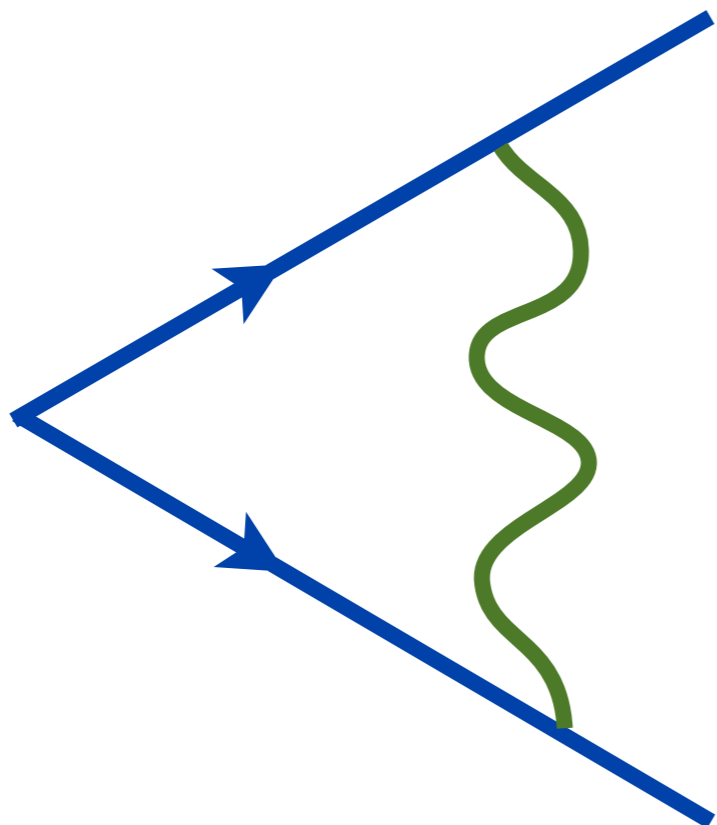
Unconventional pairing at and near hot spots

# BCS theory

$$1 + \lambda_{\text{e-ph}} \log \left( \frac{\omega_D}{\omega} \right)$$



Cooper  
logarithm





# BCS theory

$$1 + \lambda_{\text{e-ph}} \log \left( \frac{\omega_D}{\omega} \right)$$

Electron-phonon  
coupling



# BCS theory

$$1 + \lambda_{\text{e-ph}} \log \left( \frac{\omega_D}{\omega} \right)$$

Electron-phonon  
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Debye  
frequency

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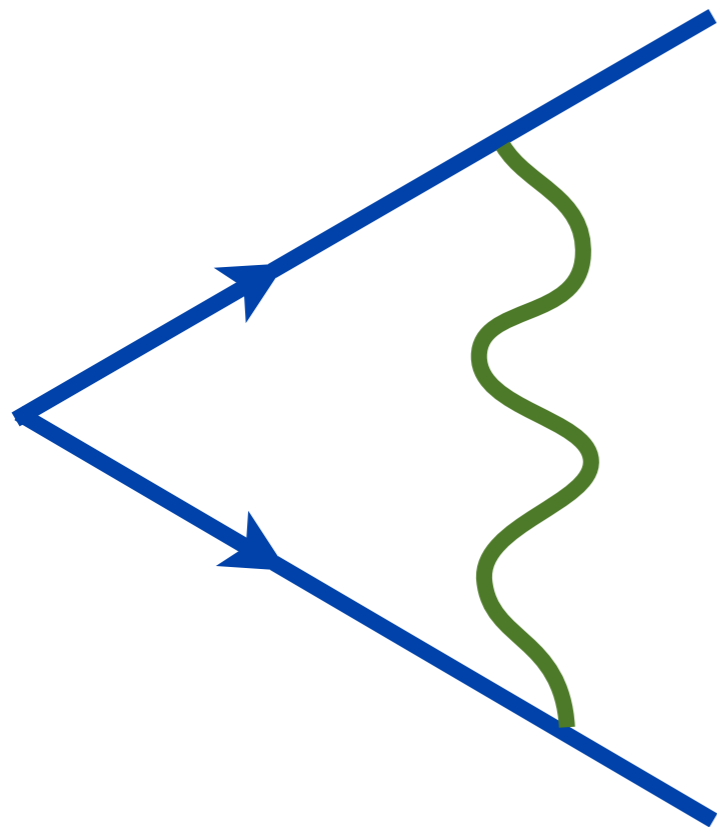
Implies

$$T_c \sim \omega_D \exp(-1/\lambda)$$

# Enhancement of pairing susceptibility by interactions

## Antiferromagnetic fluctuations: weak-coupling

$$1 + \left(\frac{U}{t}\right)^2 \log\left(\frac{E_F}{\omega}\right)$$



Cooper  
logarithm

V. J. Emery, *J. Phys. (Paris) Colloq.* **44**, C3-977 (1983)

D. J. Scalapino, E. Loh, and J. E. Hirsch, *Phys. Rev. B* **34**, 8190 (1986)

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S. Raghu, S. A. Kivelson, and D. J. Scalapino, *Phys. Rev. B* **81**, 224505 (2010)

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V.J. Emery, *J. Phys. (Paris) Colloq.* **44**, C3-977 (1983)

D.J. Scalapino, E. Loh, and J.E. Hirsch, *Phys. Rev. B* **34**, 8190 (1986)

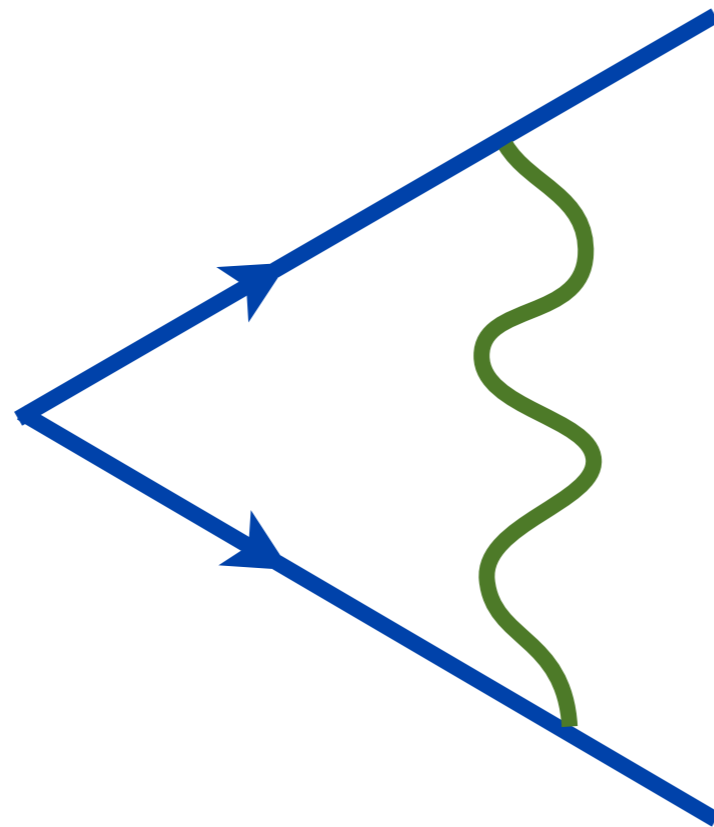
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# Enhancement of pairing susceptibility by interactions

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M.A. Metlitski and S. Sachdev, *Phys. Rev. B* **85**, 075127 (2010)



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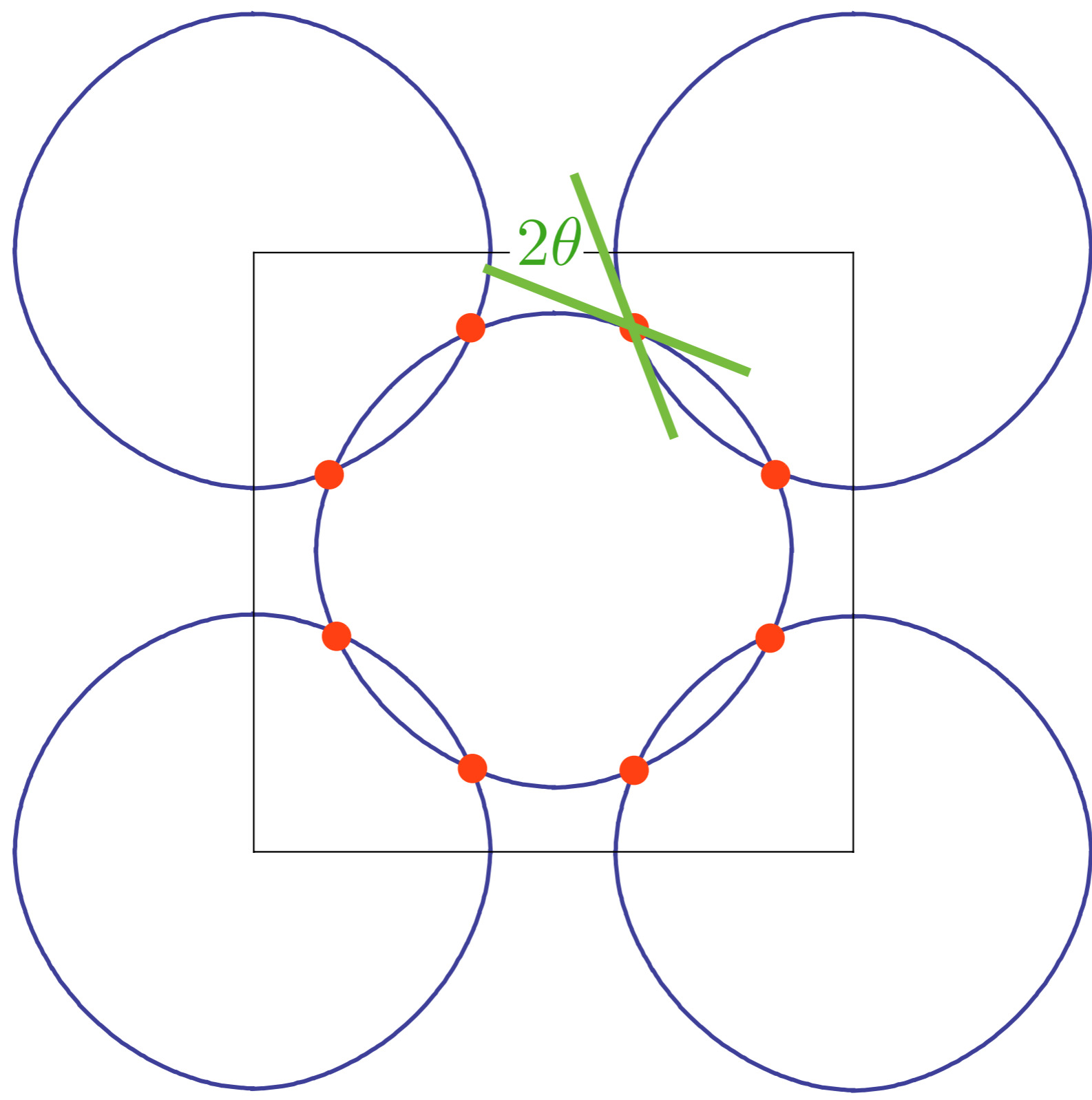
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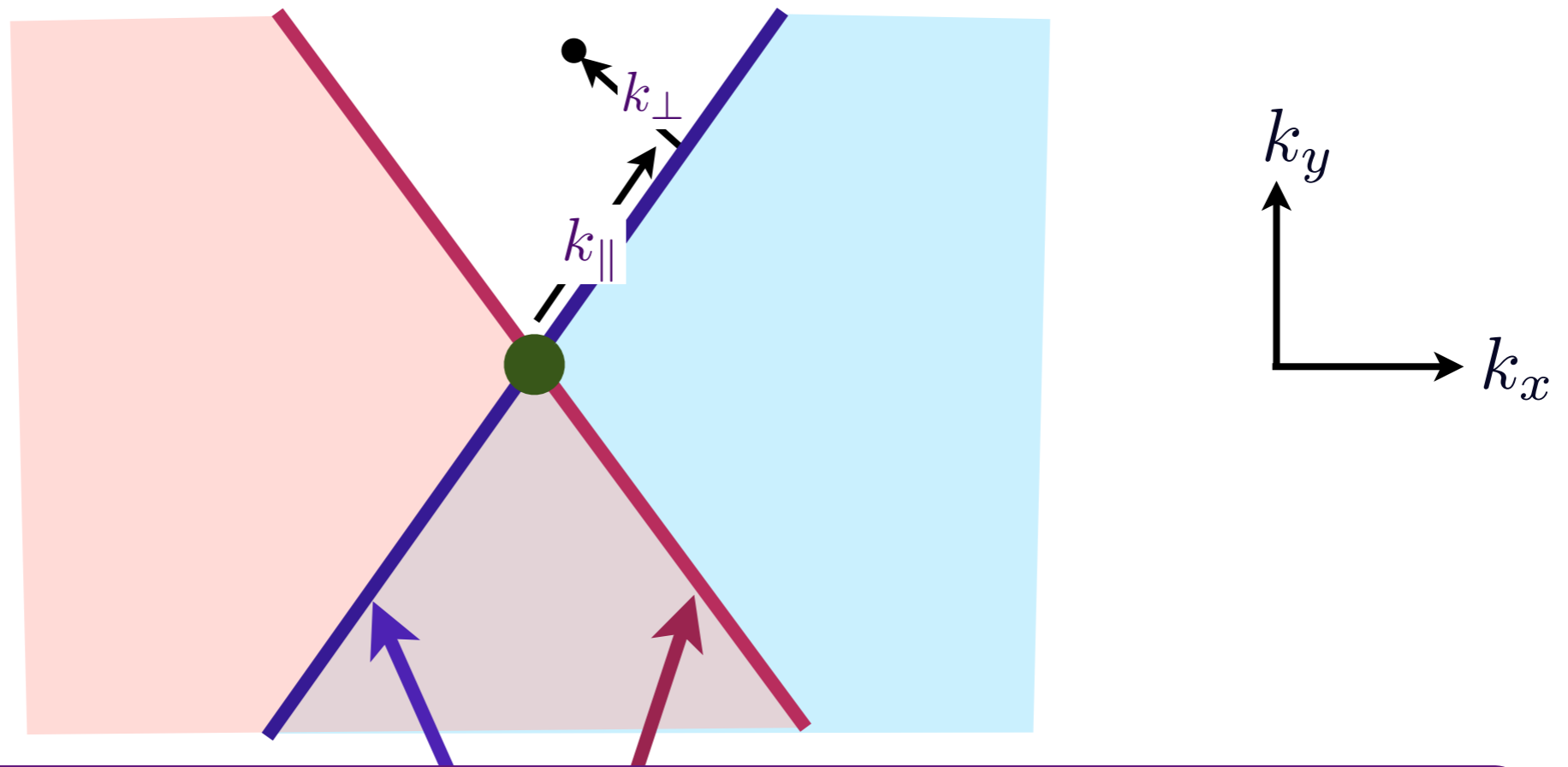
Fermi  
energy

$\alpha = \tan \theta$ , where  $2\theta$  is  
the angle between Fermi lines.  
Independent of interaction strength  
 $U$  in 2 dimensions.

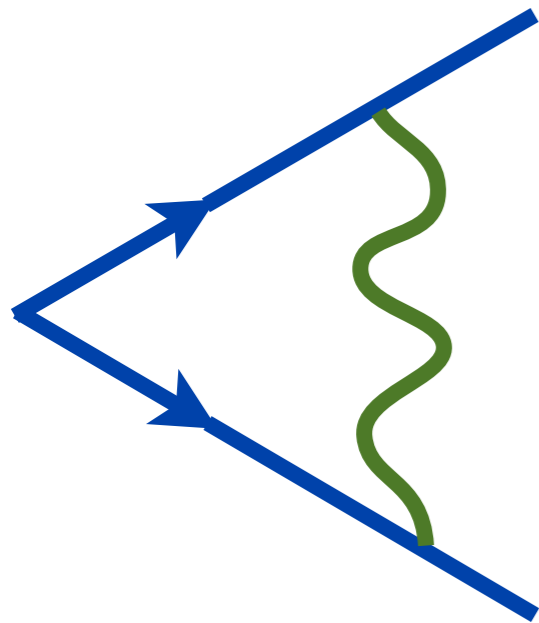
(see also Ar. Abanov, A. V. Chubukov, and A. M. Finkel'stein, *Europhys. Lett.* **54**, 488 (2001))  
M. A. Metlitski and S. Sachdev, *Phys. Rev. B* **85**, 075127 (2010)



M.A. Metlitski  
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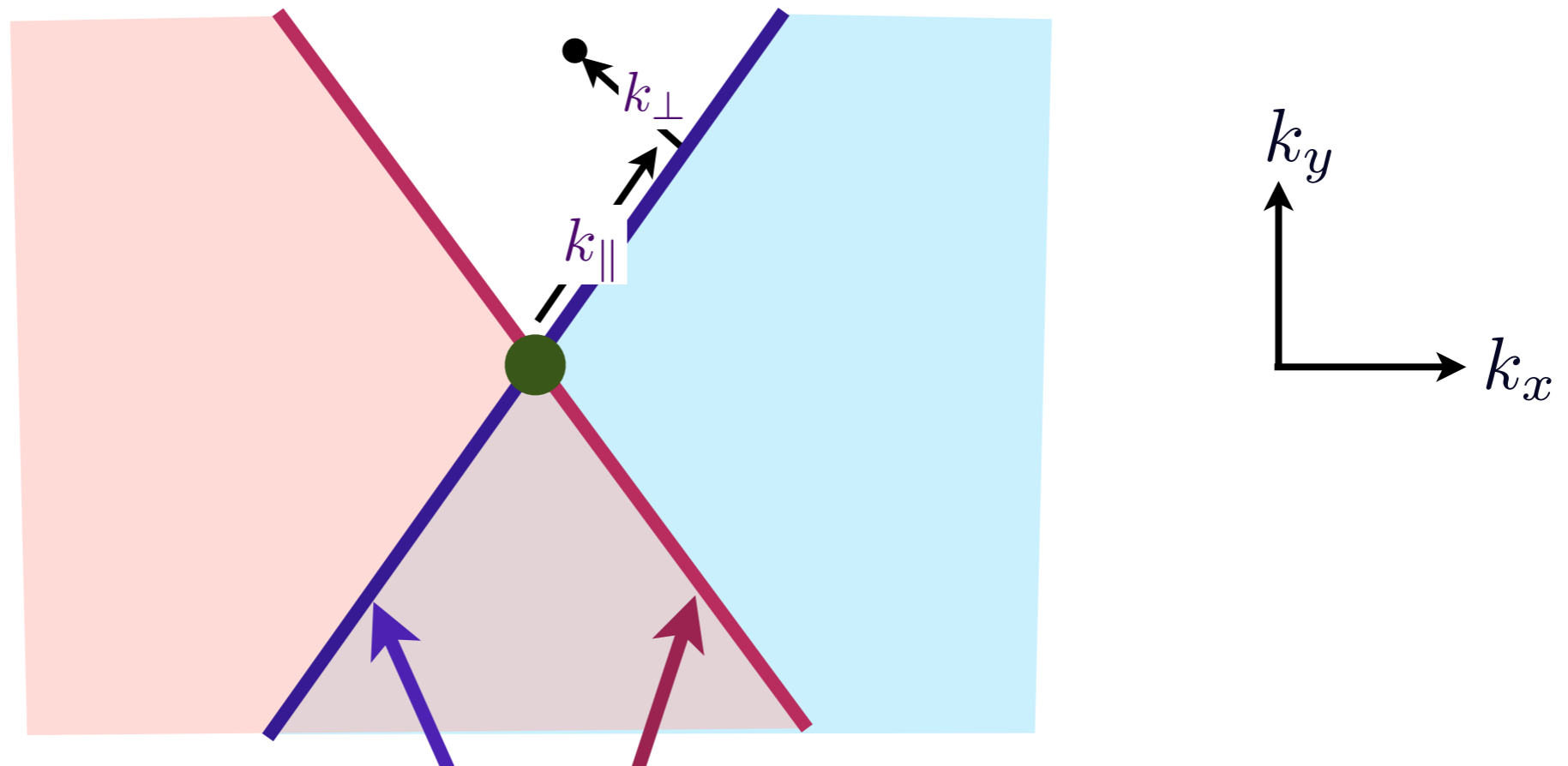


$$G_{\text{fermion}} = \frac{Z(k_{\parallel})}{i\omega - v_F(k_{\parallel})k_{\perp}}, \quad Z(k_{\parallel}) \sim v_F(k_{\parallel}) \sim k_{\parallel}$$

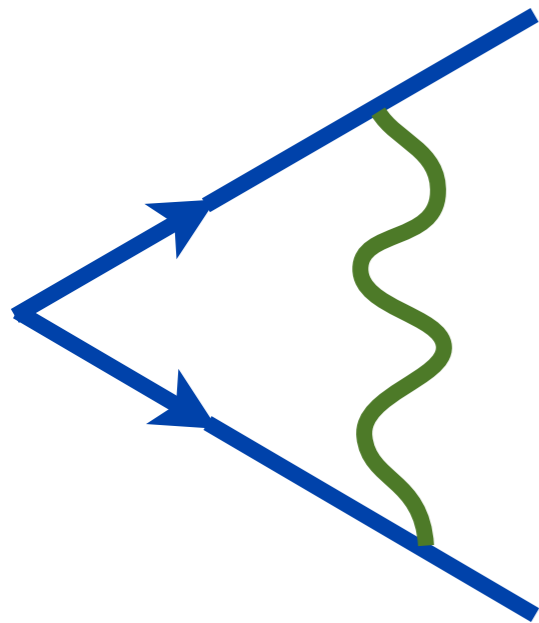


$$\int dk_{\parallel} \frac{1}{k_{\parallel}^2} \left( \frac{Z^2(k_{\parallel})}{v_F(k_{\parallel})} \right) \log \frac{k_{\parallel}^2}{\omega}$$

M.A. Metlitski  
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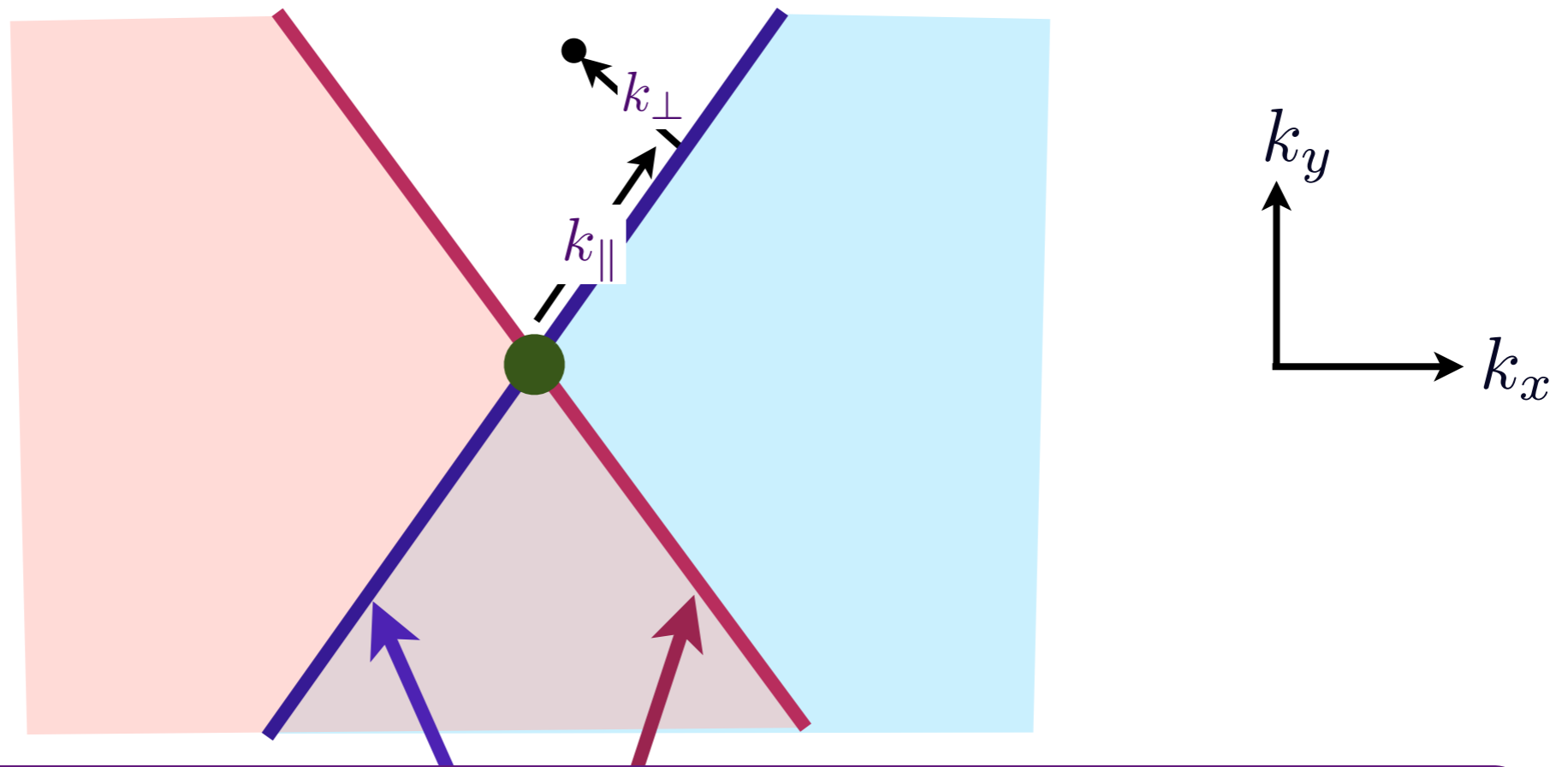
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Cooper  
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M.A. Metlitski  
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Spin fluctuation propagator

Cooper logarithm



# Enhancement of pairing susceptibility by interactions

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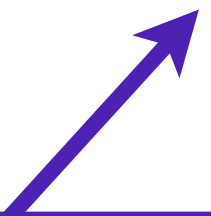
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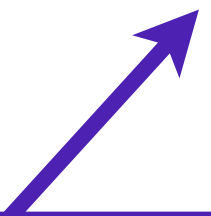
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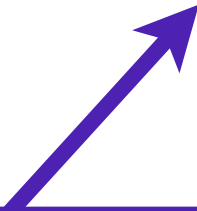
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Ar. Abanov, A. V. Chubukov, and A. M. Finkel'stein, *Europhys. Lett.* **54**, 488 (2001)

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# Phenomenological quantum theory of competing orders

Competition between superconductivity (SC) and spin-density wave (SDW) order

Write down a Landau-Ginzburg action for the quantum fluctuations of the SDW order ( $\vec{\varphi}$ ) and superconductivity ( $\Delta$ ):

$$\mathcal{S} = \int d^2r d\tau \left[ \frac{1}{2} (\partial_\tau \vec{\varphi})^2 + \frac{c^2}{2} (\nabla_x \vec{\varphi})^2 + \frac{s}{2} \vec{\varphi}^2 + \frac{u}{4} (\vec{\varphi}^2)^2 + \kappa \vec{\varphi}^2 |\Delta|^2 \right] + \int d^2r \left[ |(\nabla_x - i(2e/\hbar c)\mathcal{A})\Delta|^2 - |\Delta|^2 + \frac{|\Delta|^4}{2} \right]$$

where  $\kappa > 0$  is the repulsion between the two order parameters, and  $\nabla \times \mathcal{A} = H$  is the applied magnetic field.

E. Demler, S. Sachdev and Y. Zhang, *Phys. Rev. Lett.* **87**, 067202 (2001).



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Write down a Landau-Ginzburg action for the quantum fluctuations of the SDW order ( $\vec{\varphi}$ ) and superconductivity ( $\Delta$ ):

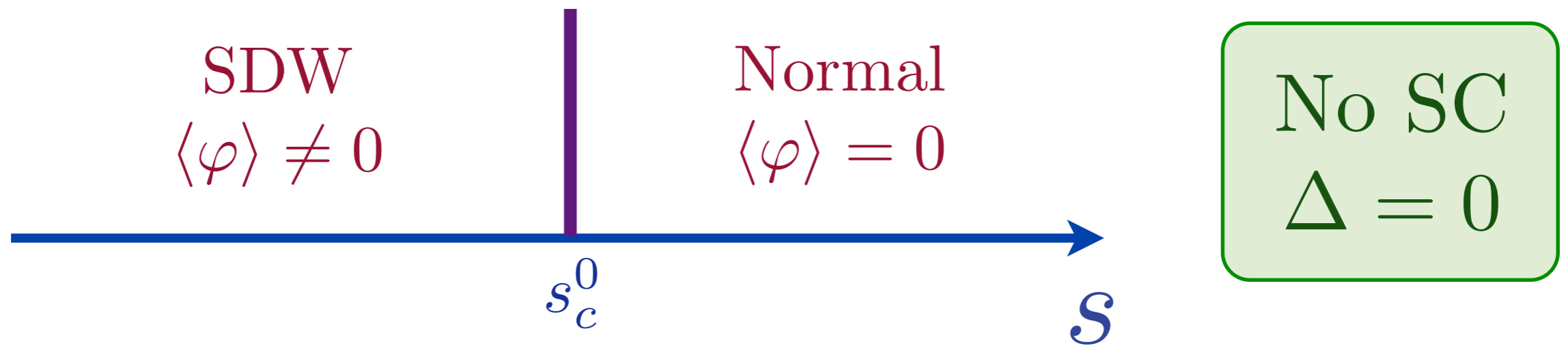
$$\mathcal{S} = \int d^2r d\tau \left[ \frac{1}{2} (\partial_\tau \vec{\varphi})^2 + \frac{c^2}{2} (\nabla_x \vec{\varphi})^2 + \frac{s}{2} \vec{\varphi}^2 + \frac{u}{4} (\vec{\varphi}^2)^2 + \kappa \vec{\varphi}^2 |\Delta|^2 \right] + \int d^2r \left[ |(\nabla_x - i(2e/\hbar c)\mathcal{A})\Delta|^2 - |\Delta|^2 + \frac{|\Delta|^4}{2} \right]$$

where  $\kappa > 0$  is the repulsion between the two order parameters, and  $\nabla \times \mathcal{A} = H$  is the applied magnetic field.

E. Demler, S. Sachdev and Y. Zhang, *Phys. Rev. Lett.* **87**, 067202 (2001).

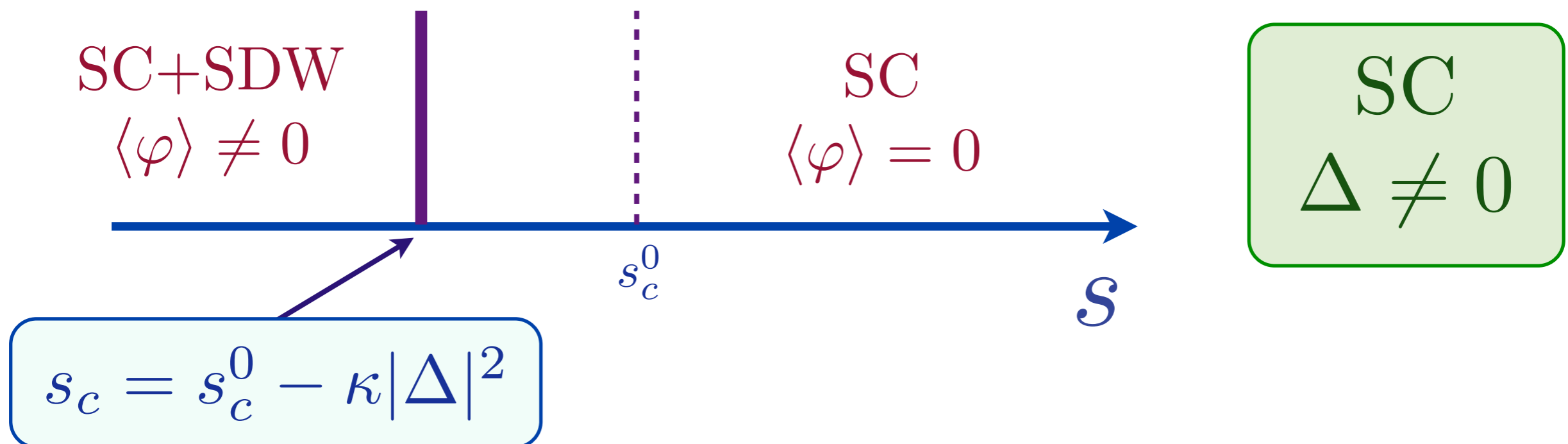
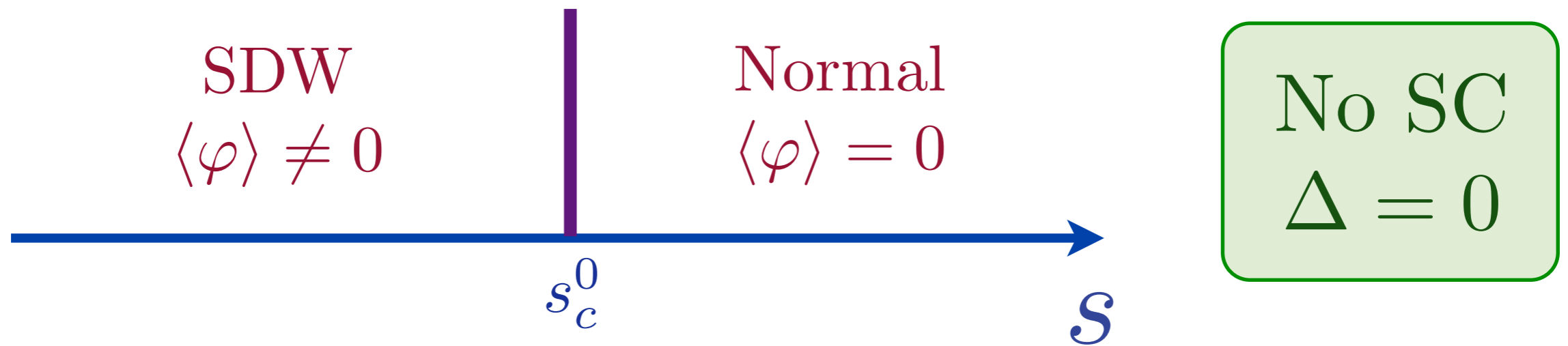
# Phenomenological quantum theory of competing orders

Competition between superconductivity (SC) and spin-density wave (SDW) order



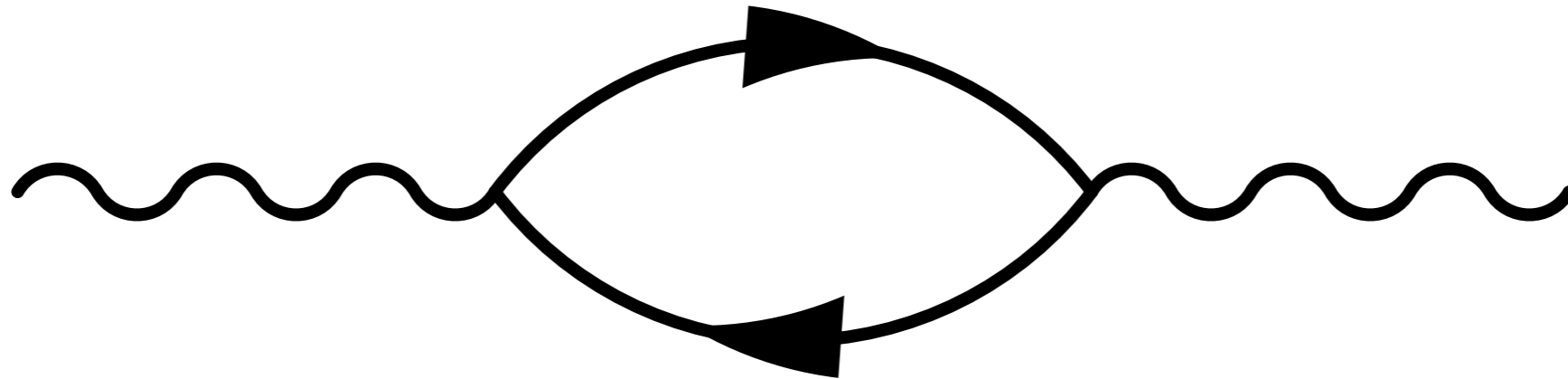
# Phenomenological quantum theory of competing orders

Competition between superconductivity (SC) and spin-density wave (SDW) order



# Fermi surface theory of competing orders

Competition between superconductivity (SC) and spin-density wave (SDW) order



Compute the SDW susceptibility,  $\chi$ , in the superconducting state. As  $\Delta \rightarrow 0$ , we find

$$\chi(\Delta) = \chi(0) - C|\Delta|$$

where  $C$  is a universal constant dominated by the vicinity of the hot spots.

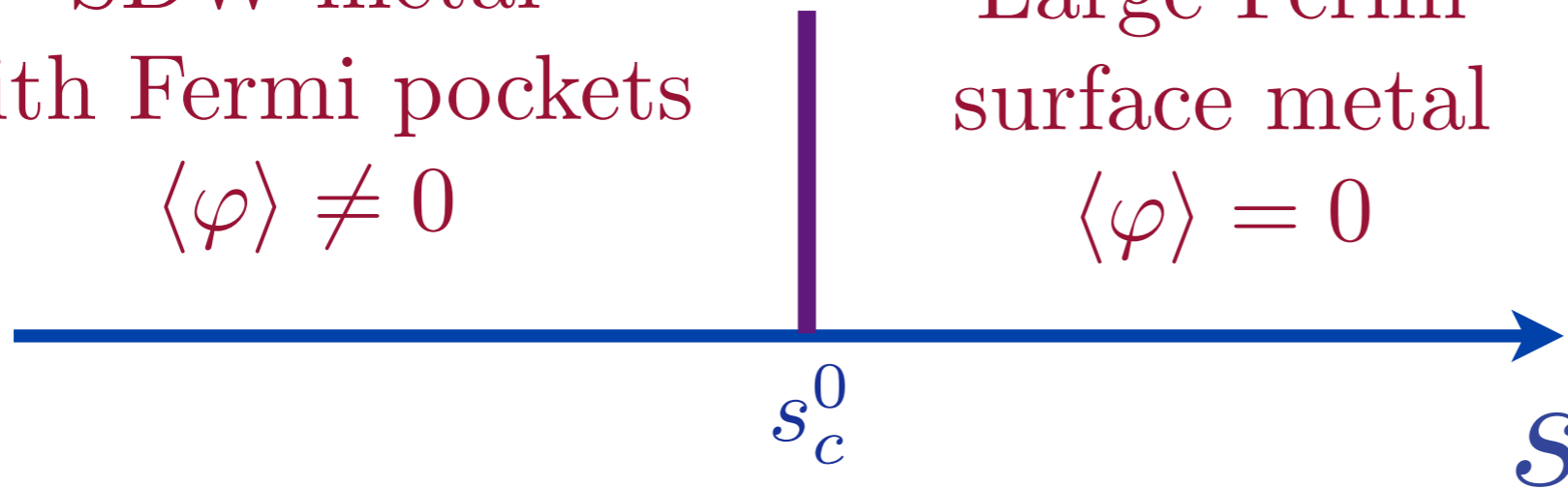
E. G. Moon and S. Sachdev, *Phy. Rev. B* **82**, 104516 (2010)

# Fermi surface theory of competing orders

Competition between superconductivity (SC) and spin-density wave (SDW) order

SDW metal  
with Fermi pockets  
 $\langle \varphi \rangle \neq 0$

Large Fermi  
surface metal  
 $\langle \varphi \rangle = 0$



No SC  
 $\Delta = 0$

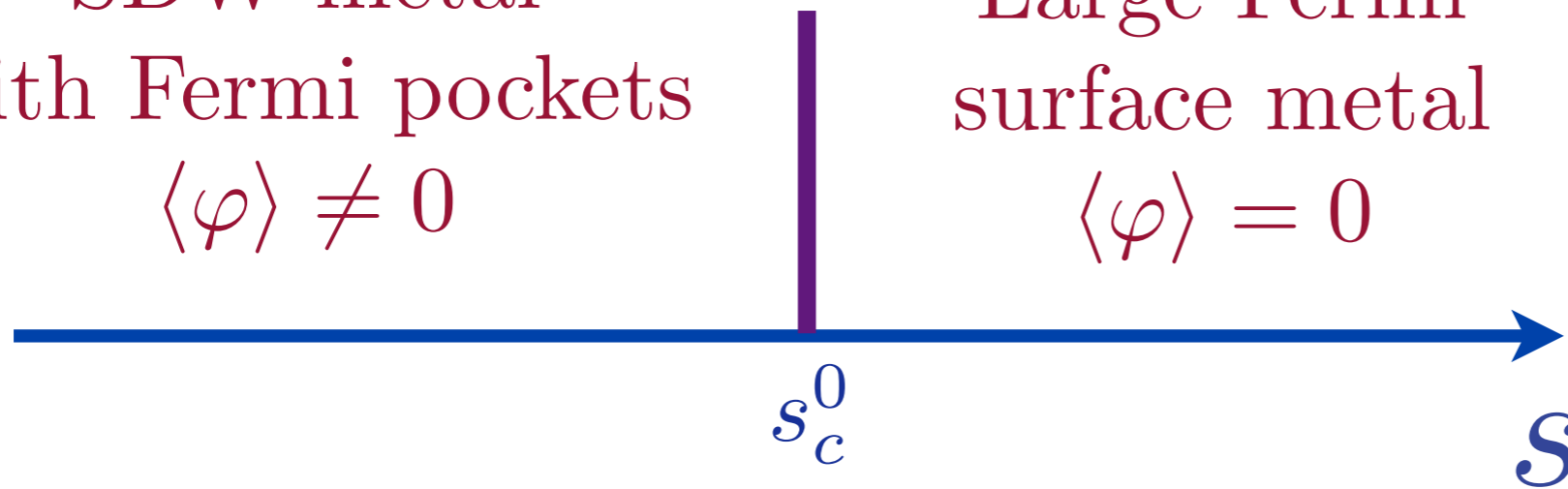
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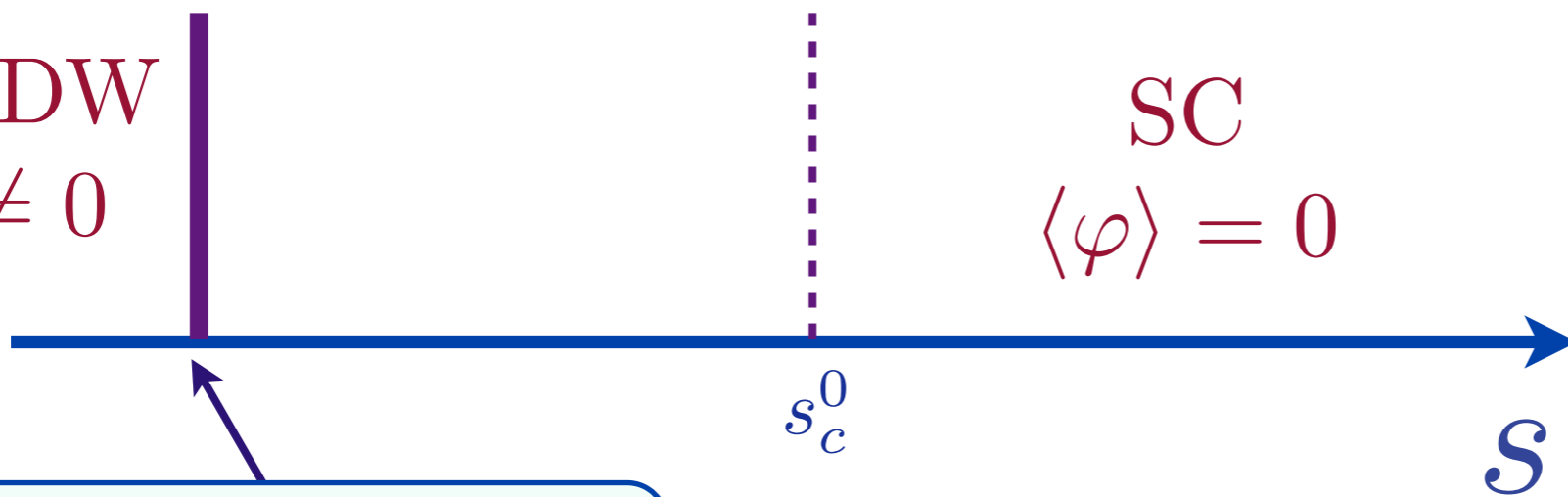
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No SC  
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SC+SDW  
 $\langle \varphi \rangle \neq 0$

SC  
 $\langle \varphi \rangle = 0$

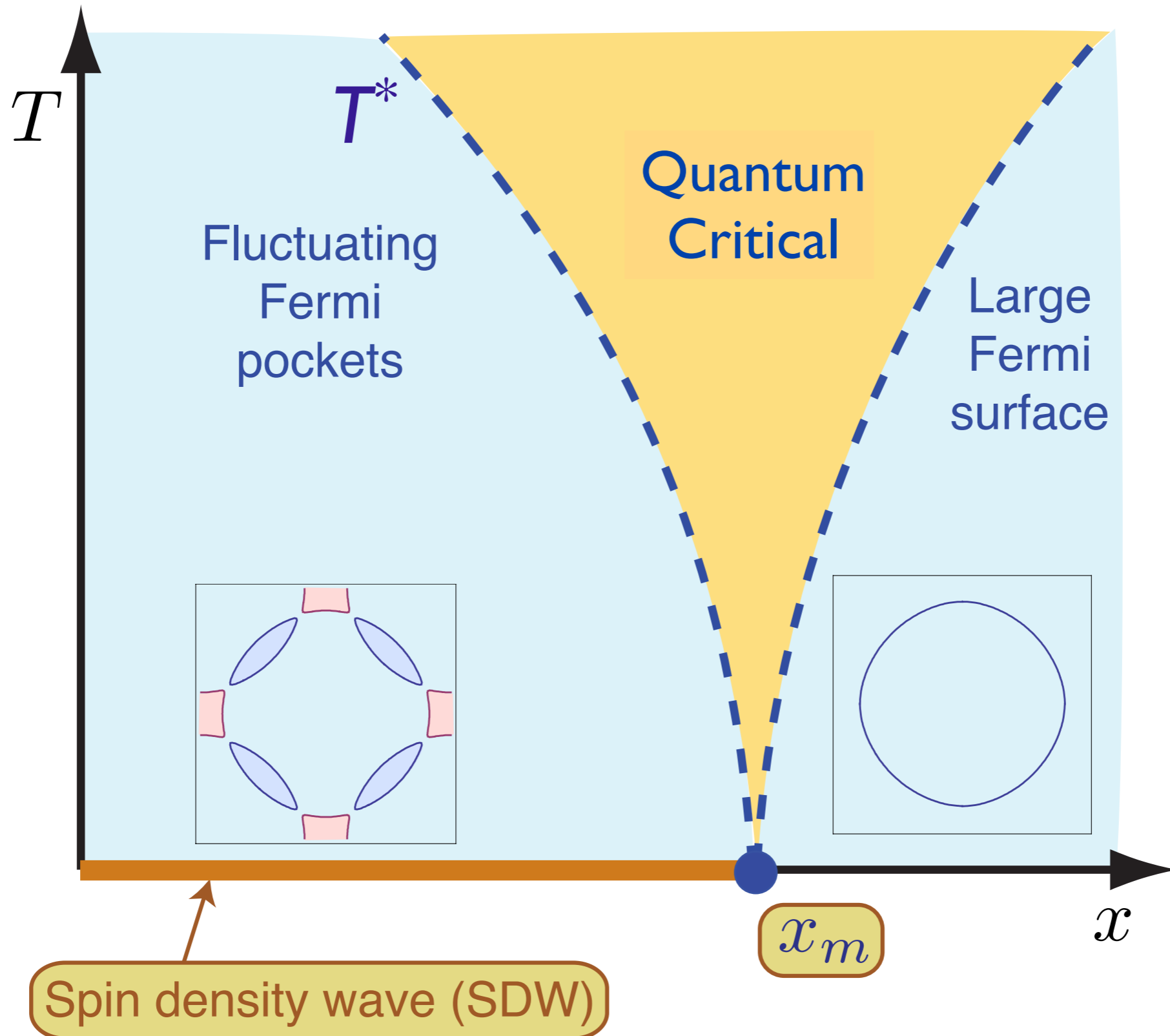


SC  
 $\Delta \neq 0$

$$s_c^0 - s_c \sim |\Delta|$$

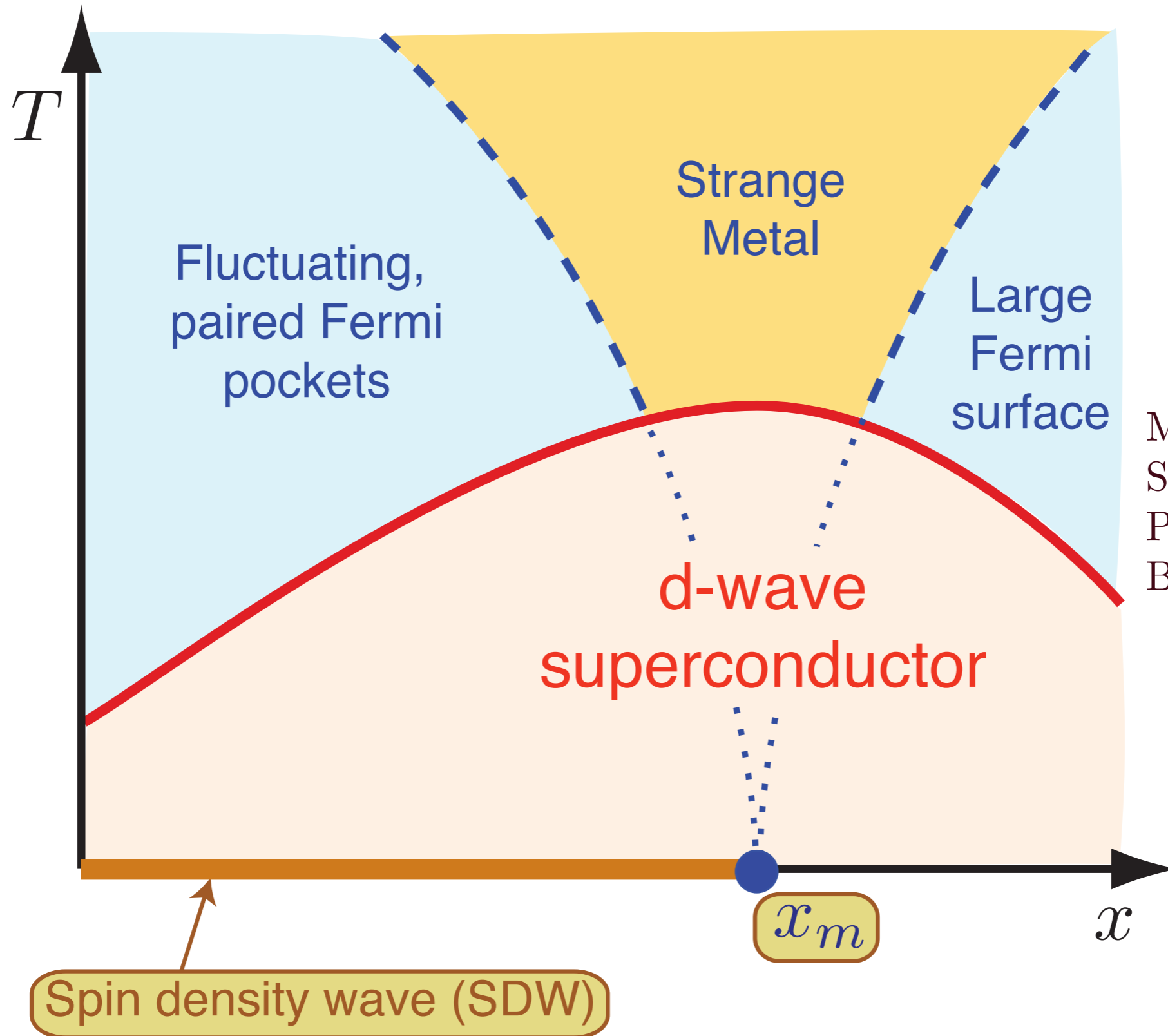
E. G. Moon and S. Sachdev, *Phy. Rev. B* **82**, 104516 (2010)

# Theory of quantum criticality in the cuprates



Relaxation and equilibration times  $\sim \hbar/k_B T$  are robust properties of strongly-coupled quantum criticality

# Theory of quantum criticality in the cuprates

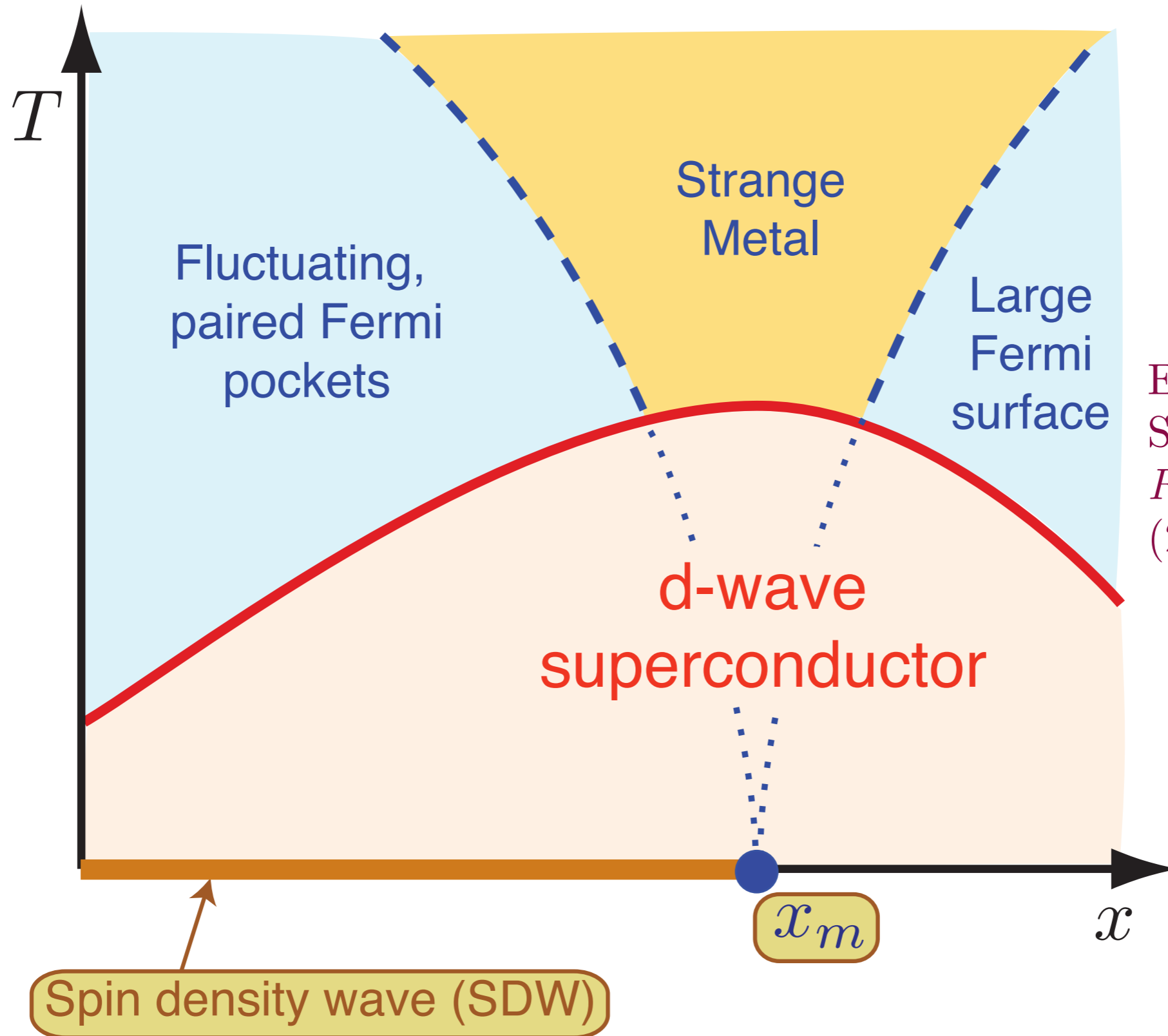


M. A. Metlitski and  
S. Sachdev,  
Physical Review  
B **82**, 075128 (2010)

SDW quantum critical point is unstable to  $d$ -wave superconductivity  
This instability is stronger than that in the BCS theory



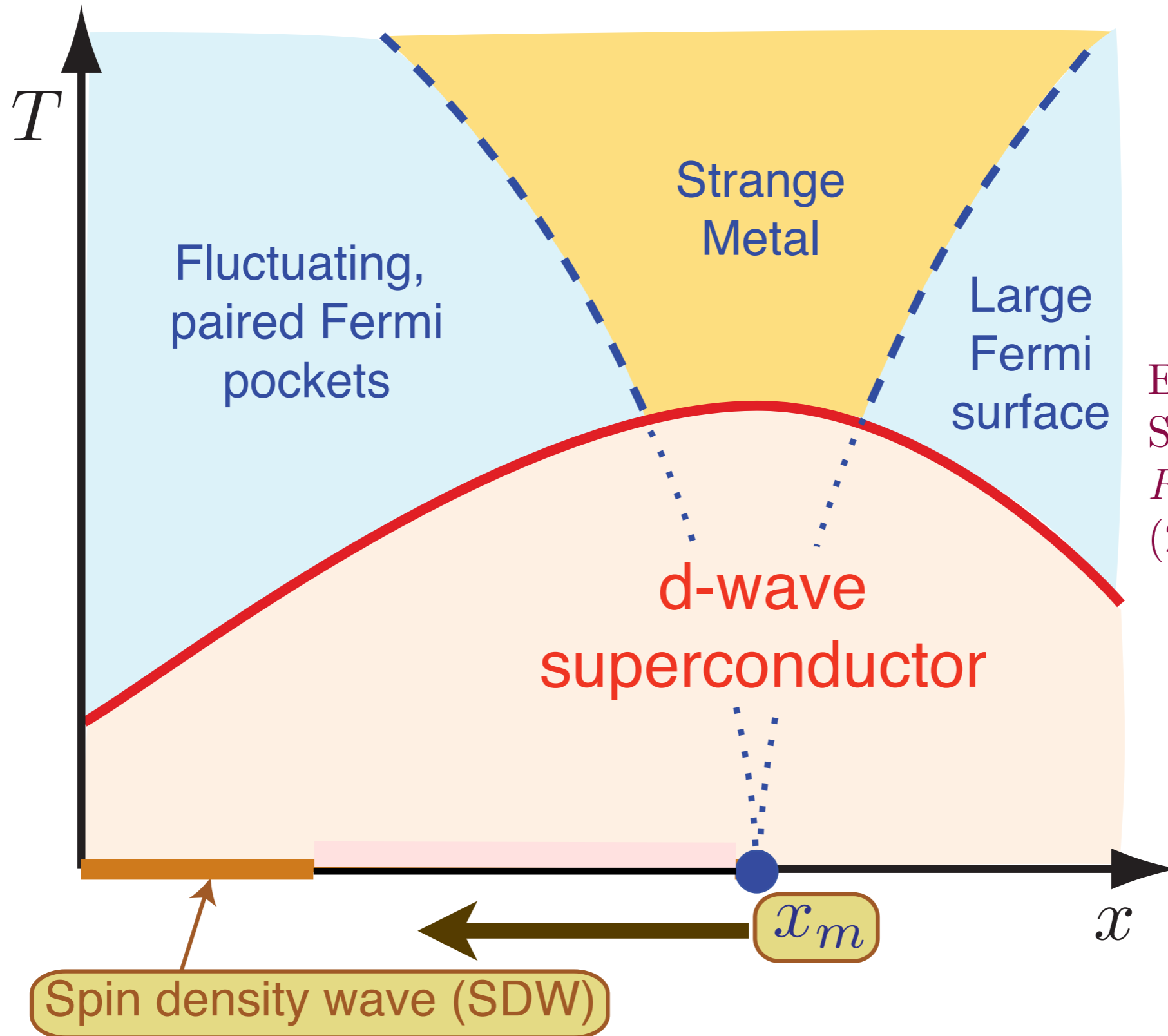
# Theory of quantum criticality in the cuprates



E. G. Moon and S. Sachdev, *Phy. Rev. B* **80**, 035117 (2009)

Competition between SDW order and superconductivity moves the actual quantum critical point to  $x = x_s < x_m$ .

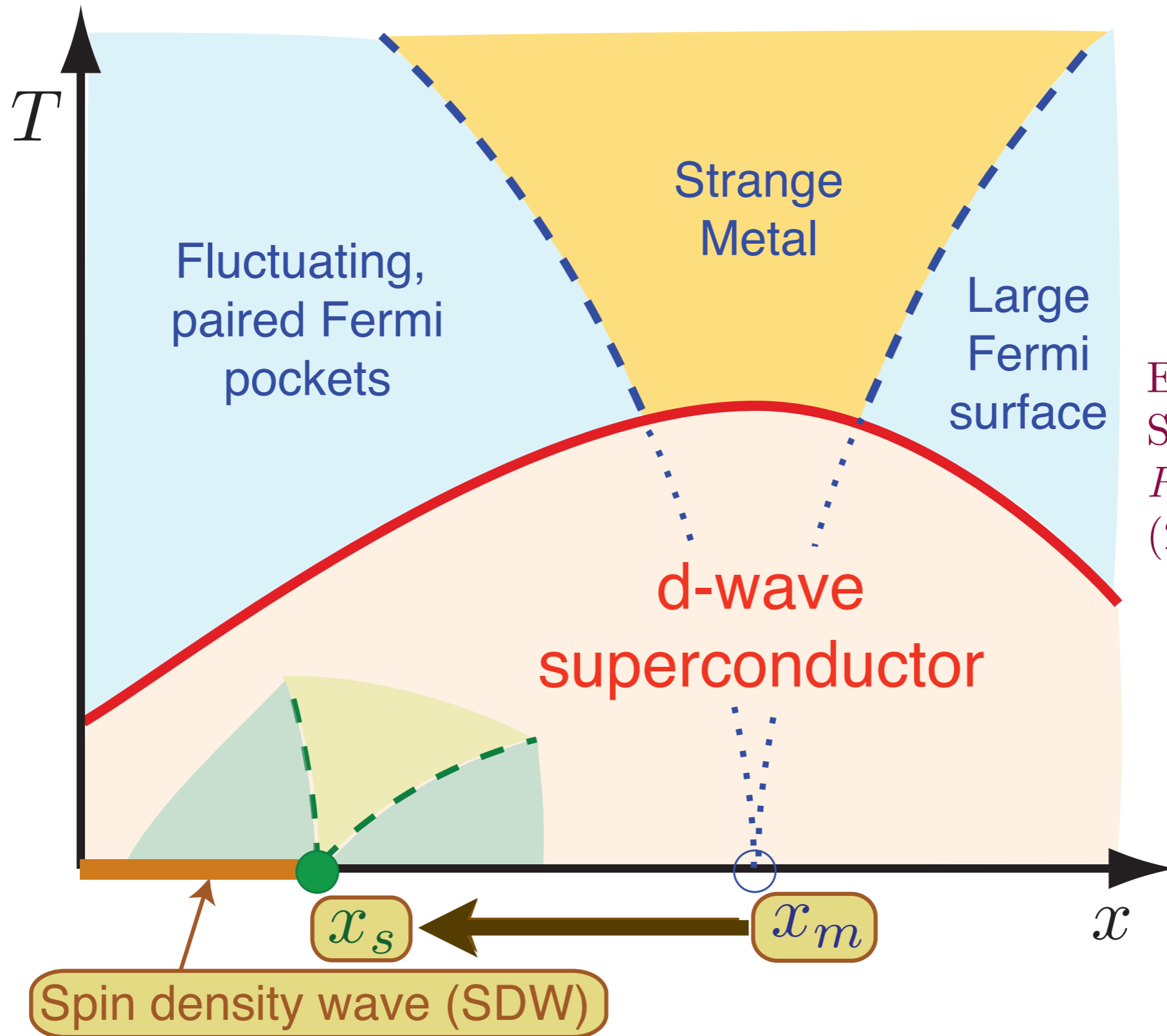
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## Questions

- *Can quantum fluctuations near the loss of antiferromagnetism induce higher temperature superconductivity ?*
- *If so, why is there no antiferromagnetism in the hole-doped cuprates near the point where the superconductivity is strongest ?*
- *What is the physics of the strange metal ?*

## Questions and answers

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**Competition between antiferromagnetism and superconductivity has shifted the antiferromagnetic quantum-critical point (QCP), and shrunk the region of antiferromagnetism. This QCP shift is largest in the cuprates**

● *What is the physics of the strange metal ?*

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Proposal: strongly-coupled quantum criticality of Fermi surface change in a metal