

# The phase diagrams of the high temperature superconductors

Talk online: [sachdev.physics.harvard.edu](http://sachdev.physics.harvard.edu)

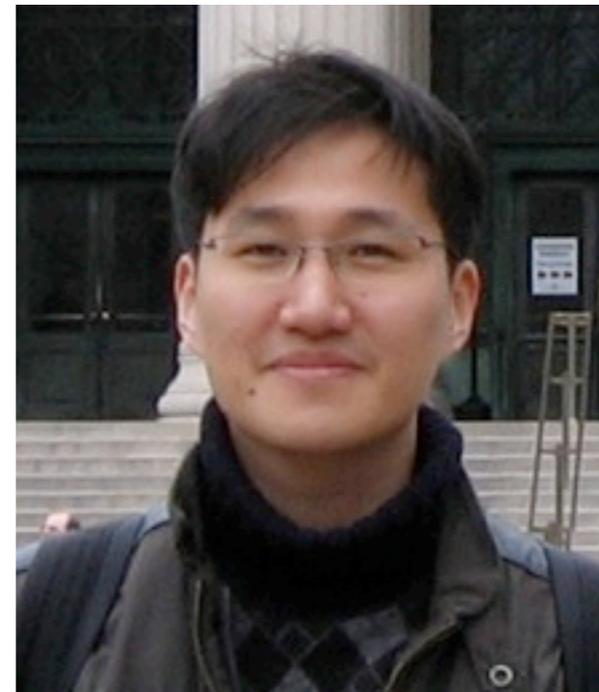
PHYSICS



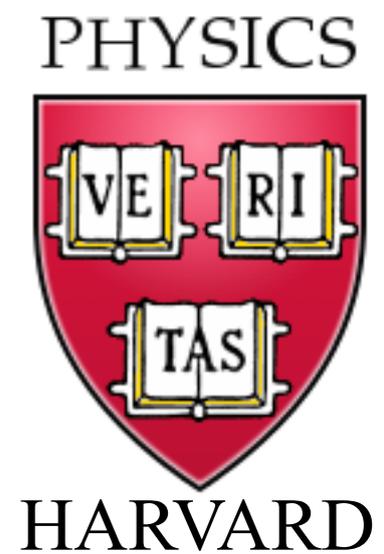
HARVARD



Max Metlitski, Harvard



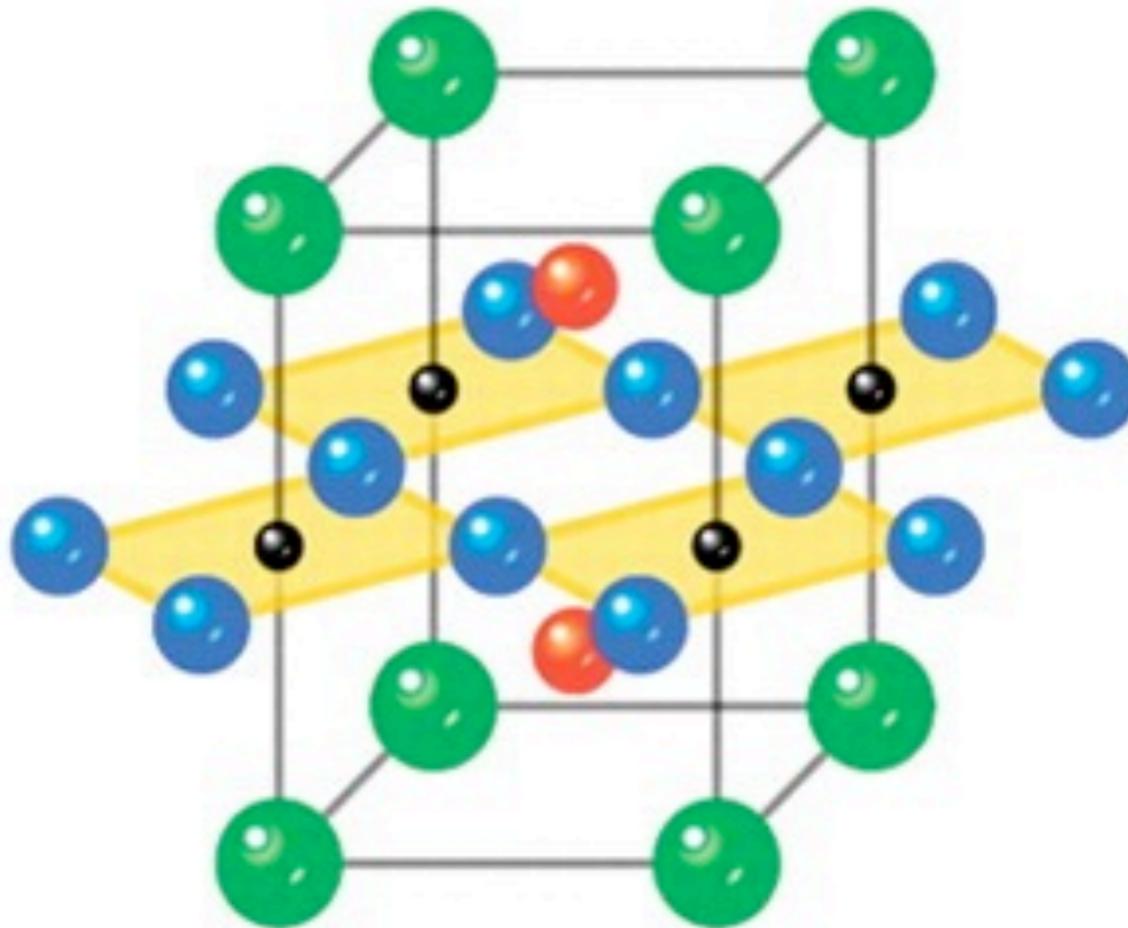
Eun Gook Moon, Harvard



# *The cuprate superconductors*

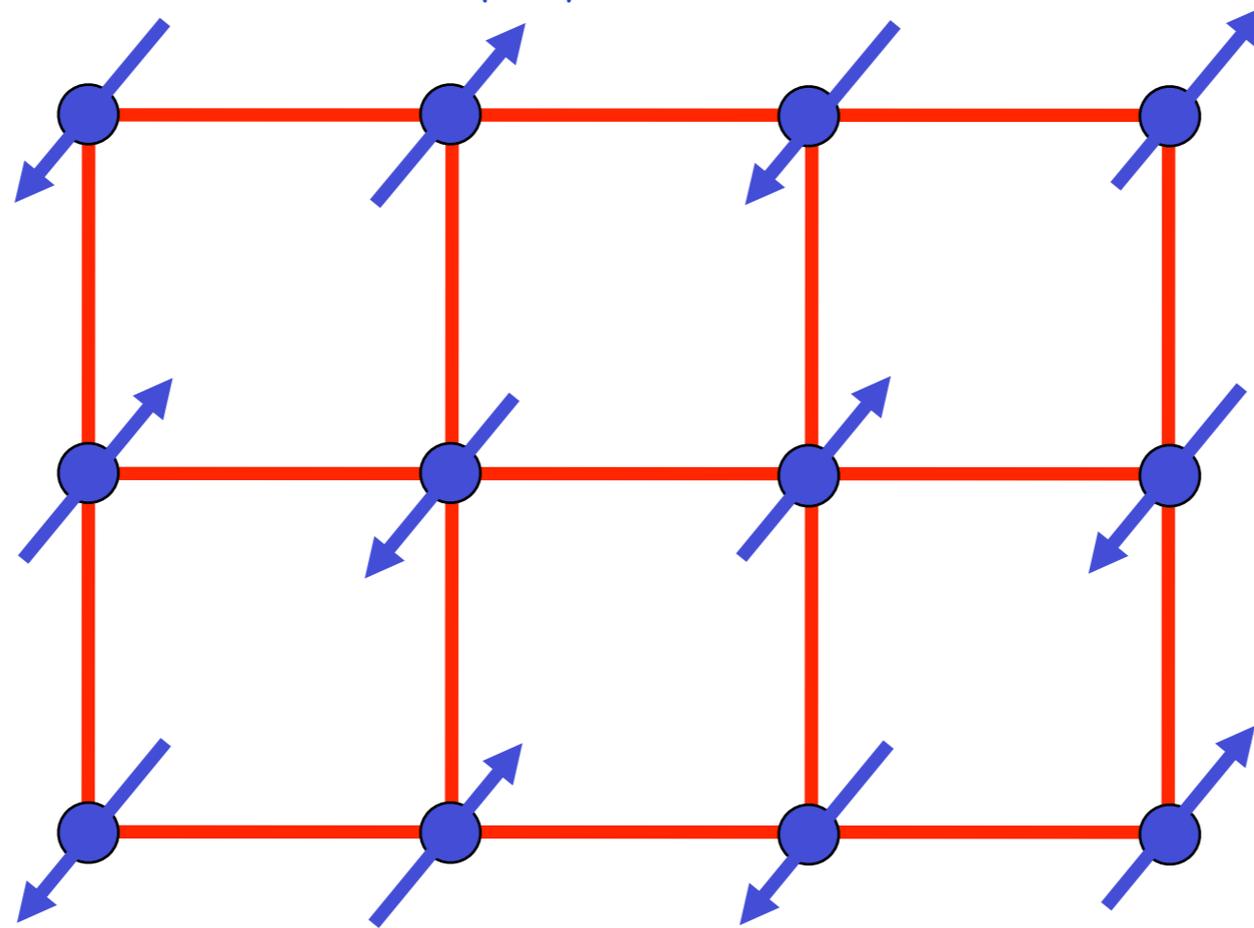
Na-CCOC

- Cu
- Ca/Na
- O
- Cl



# Square lattice antiferromagnet

$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$



Ground state has long-range Néel order

Order parameter is a single vector field  $\vec{\varphi} = \eta_i \vec{S}_i$

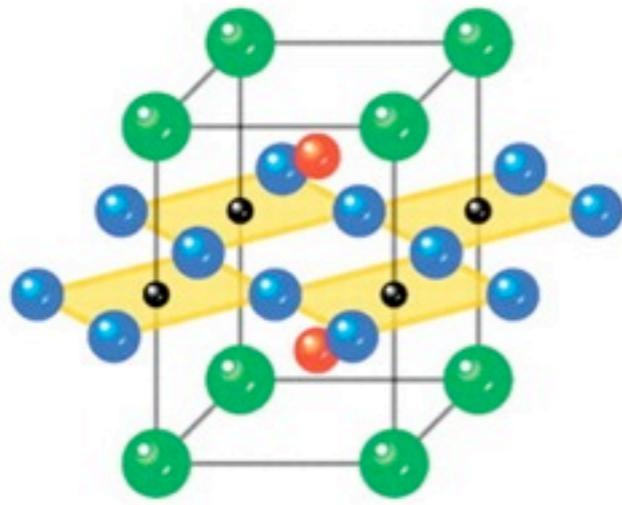
$\eta_i = \pm 1$  on two sublattices

$\langle \vec{\varphi} \rangle \neq 0$  in Néel state.

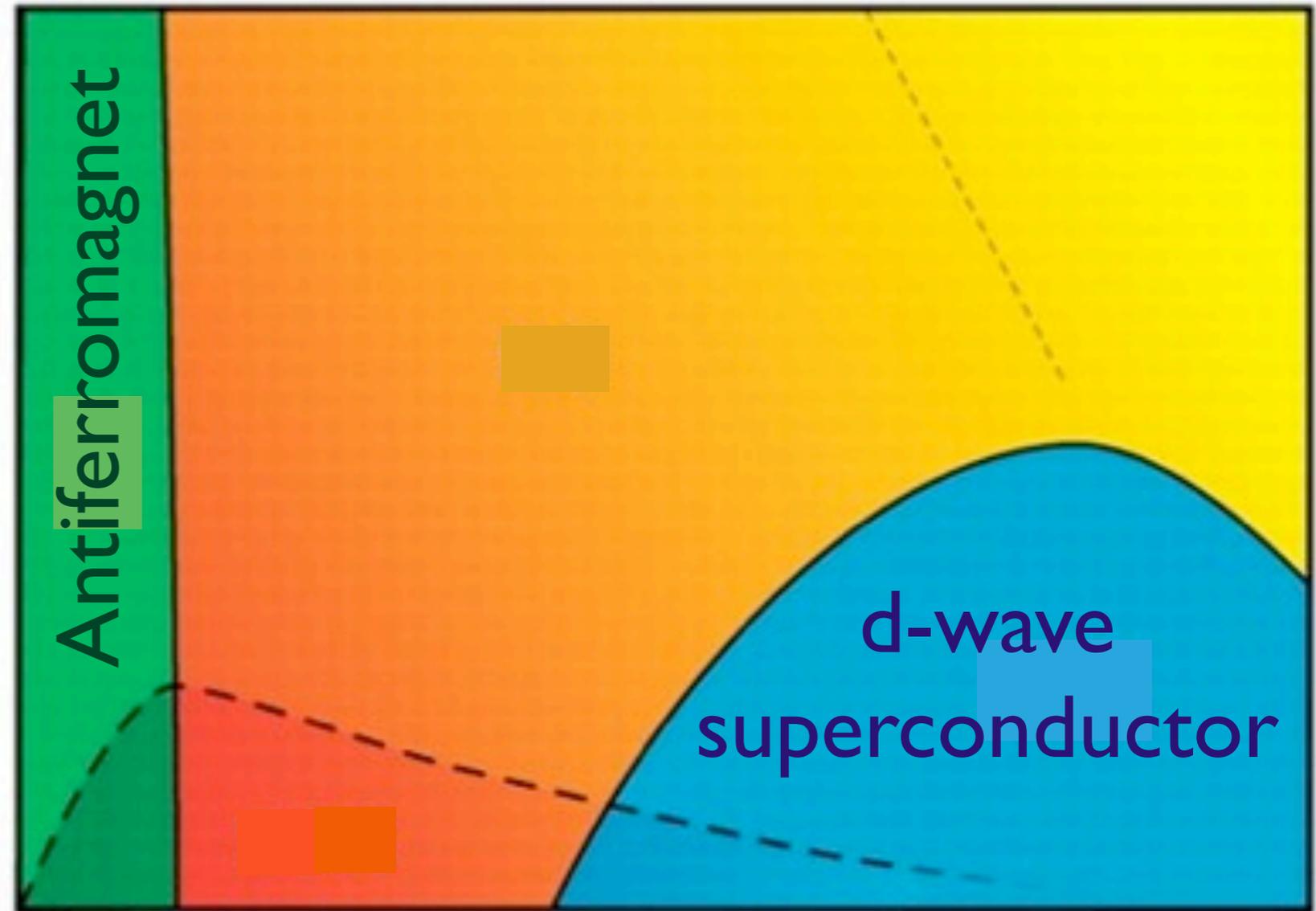
# The cuprate superconductors

Na-CCOC

- Cu
- Ca/Na
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Temperature



~2-3 %

~5-10 %

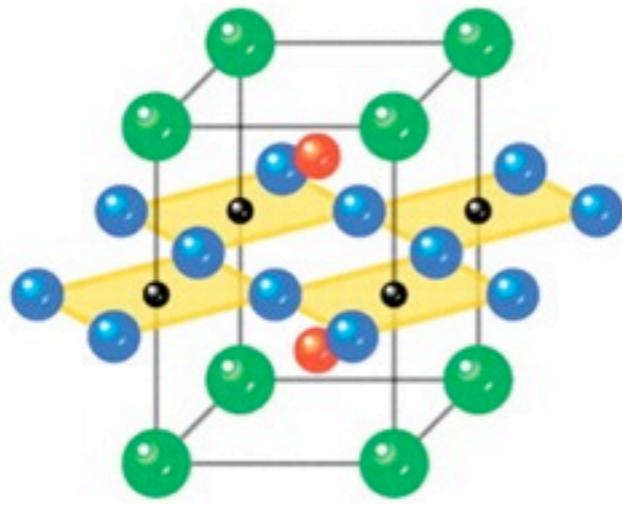
~15 %

hole concentration

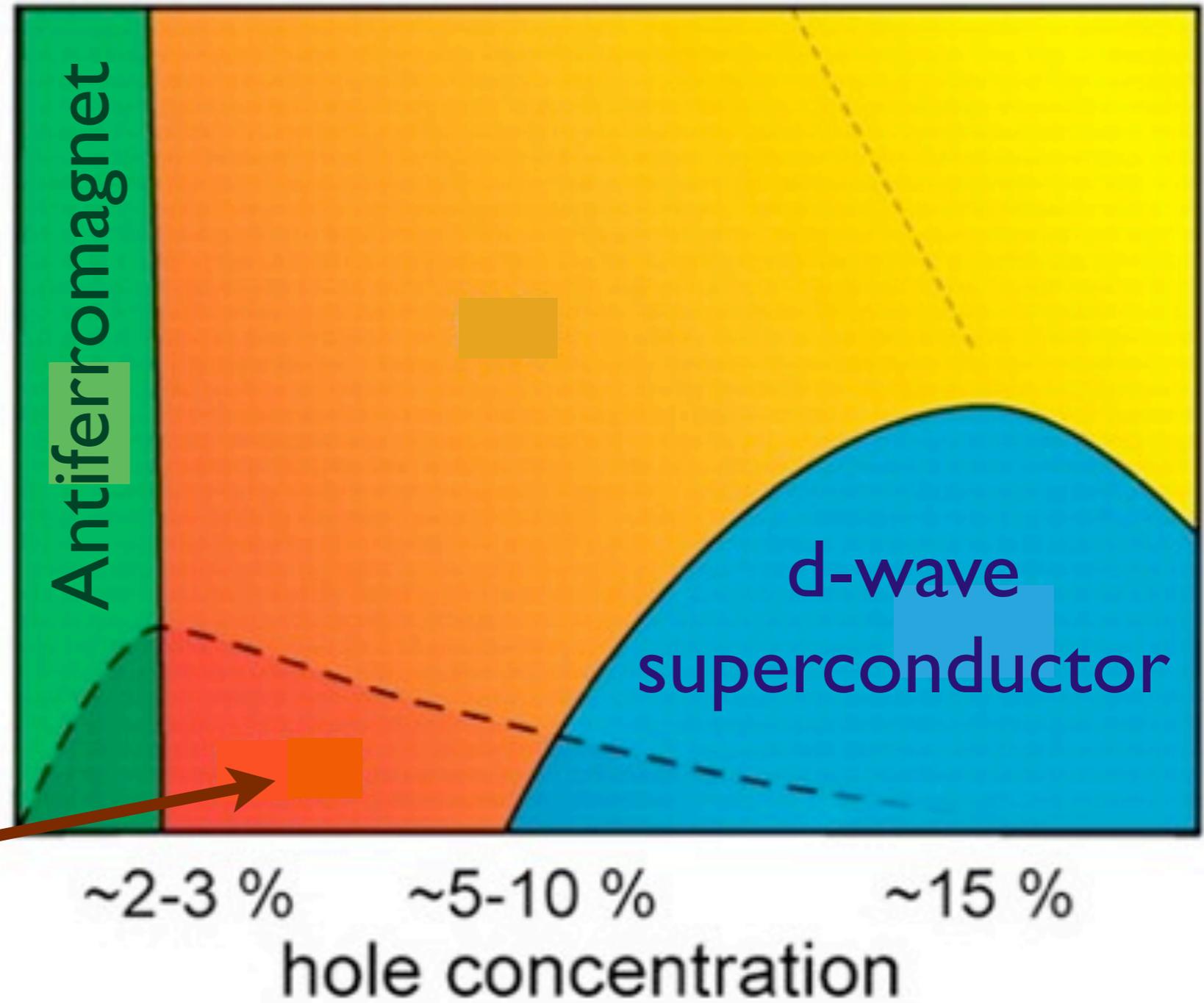
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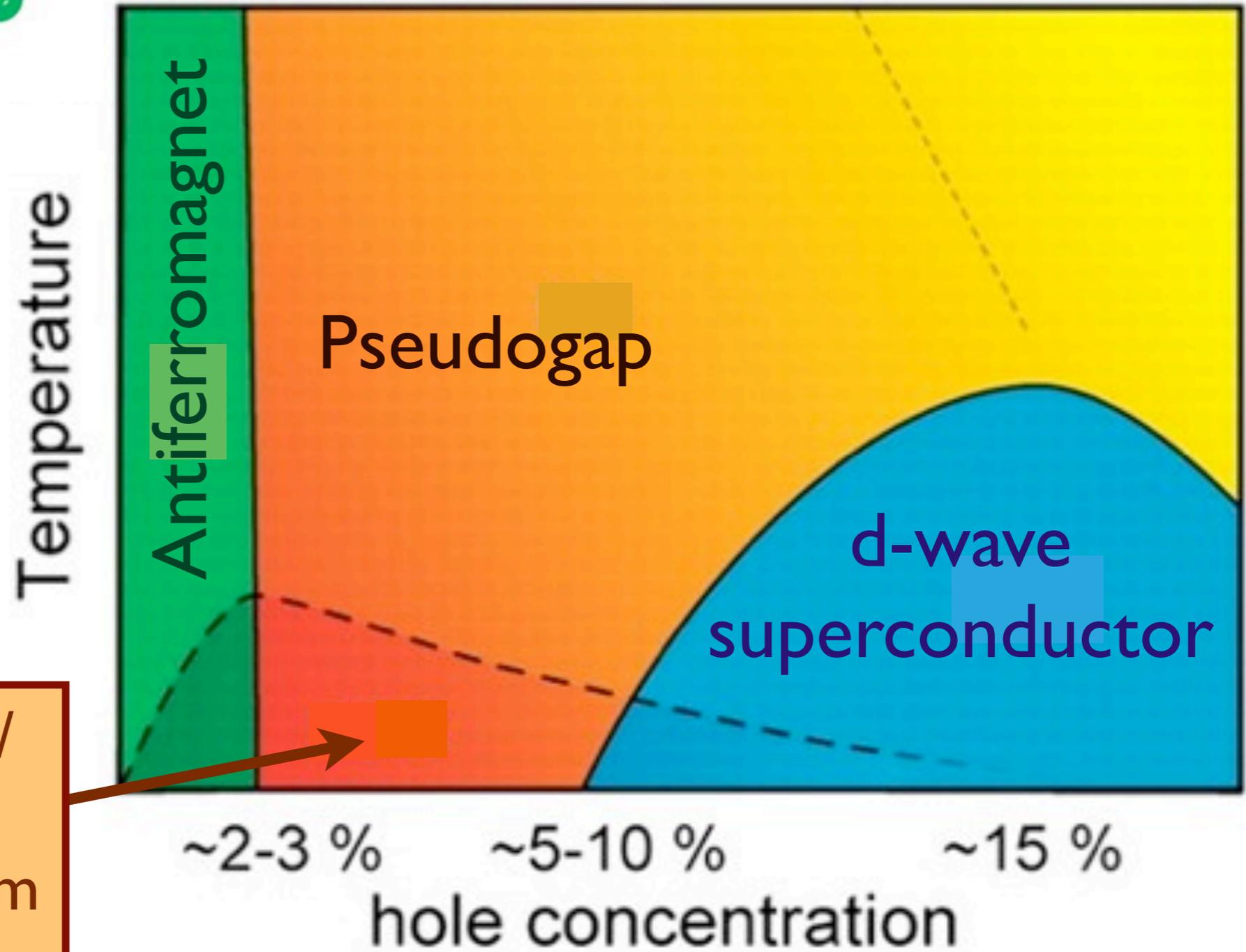
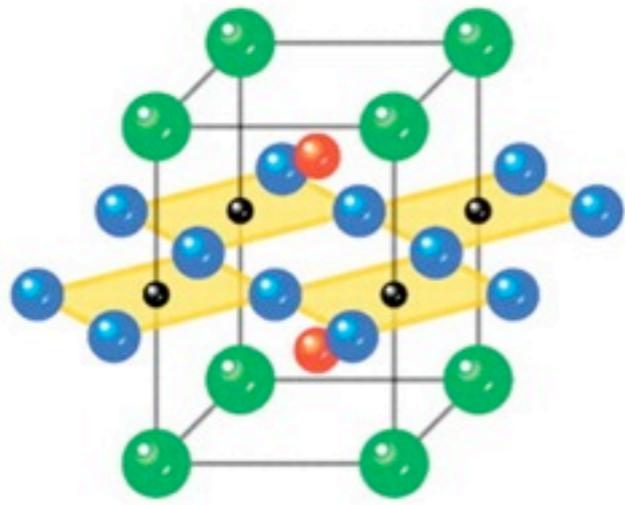


Incommensurate/  
disordered  
antiferromagnetism  
and charge order

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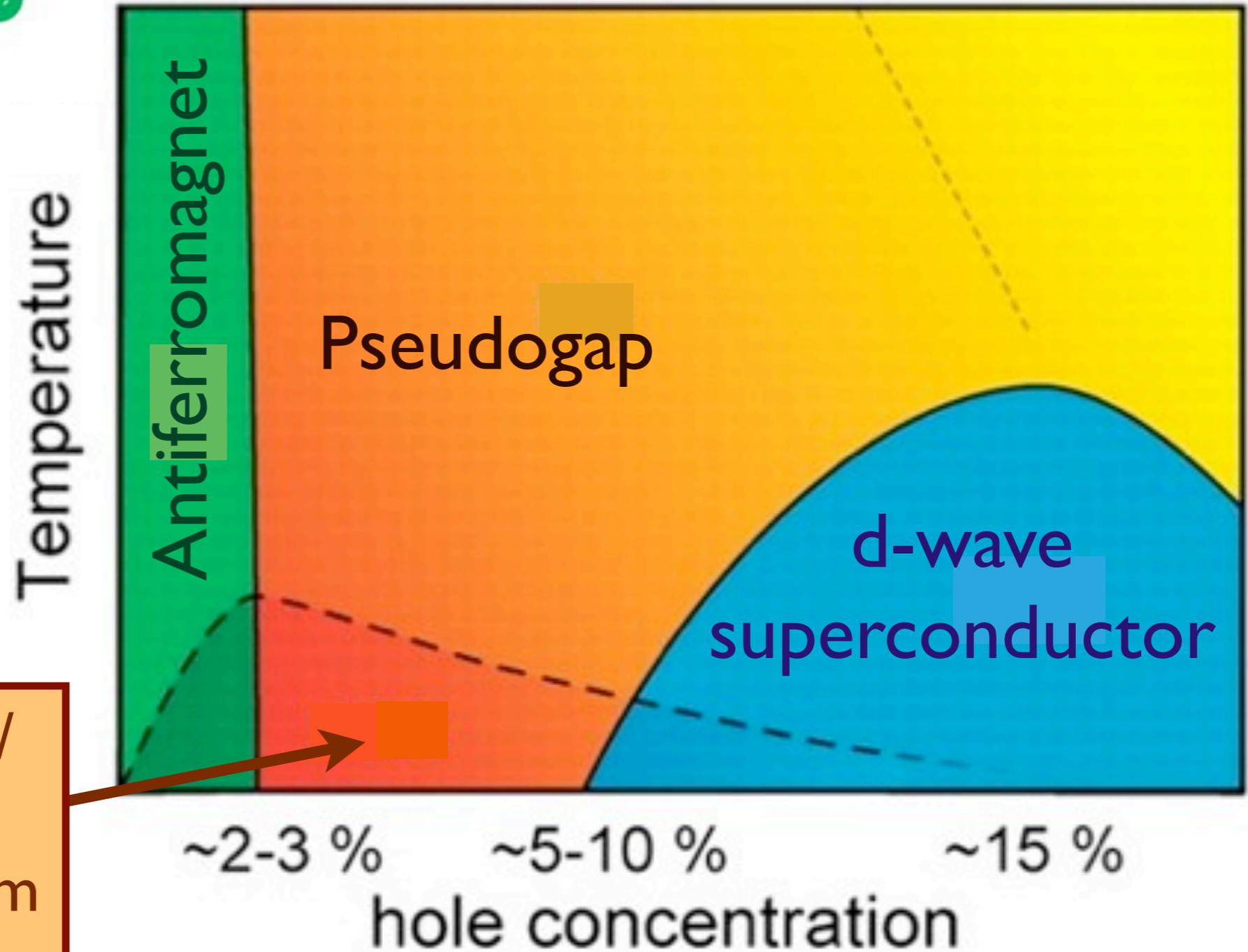
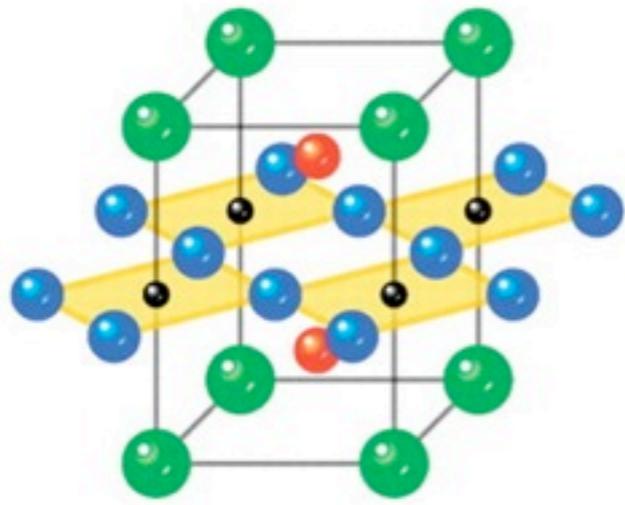


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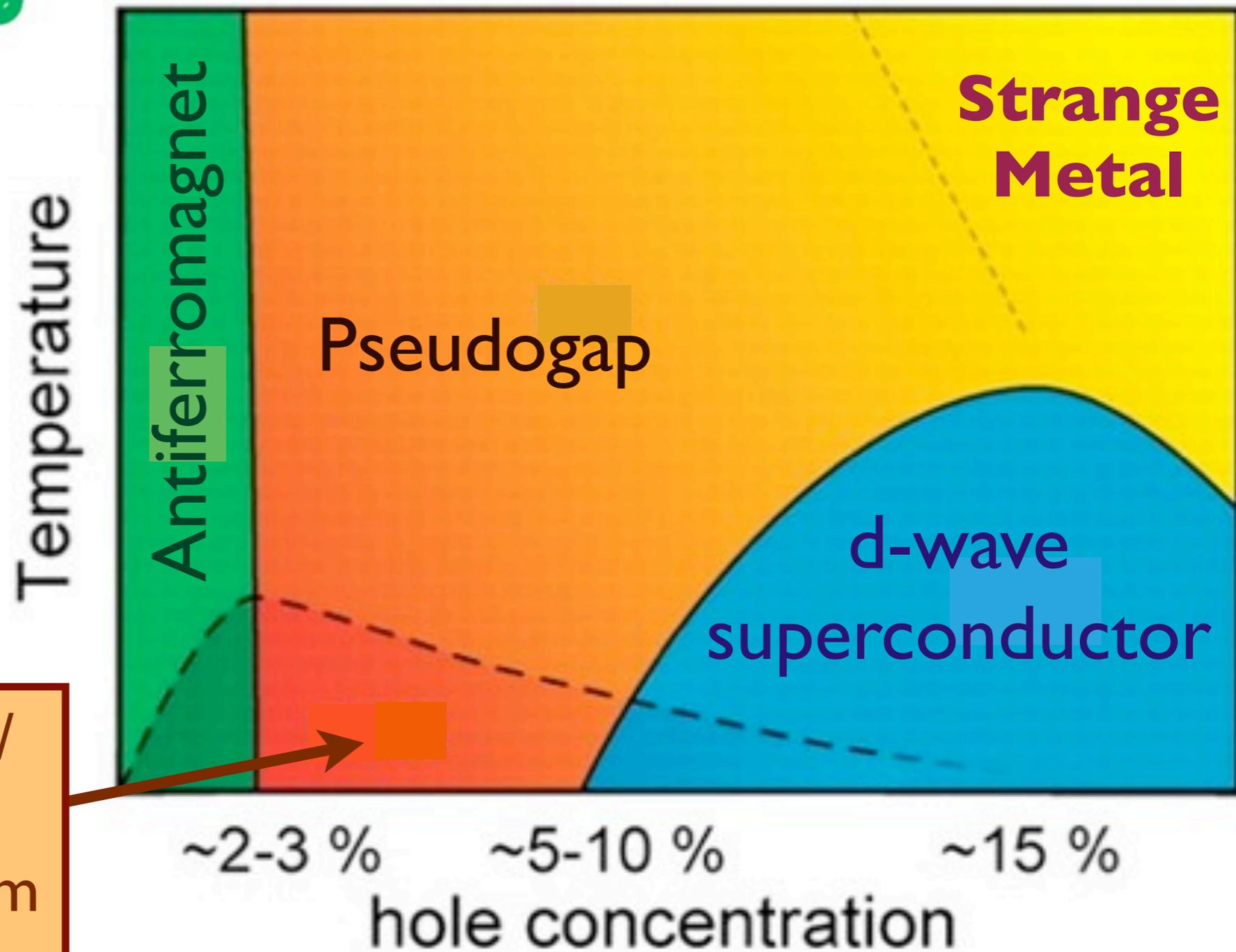
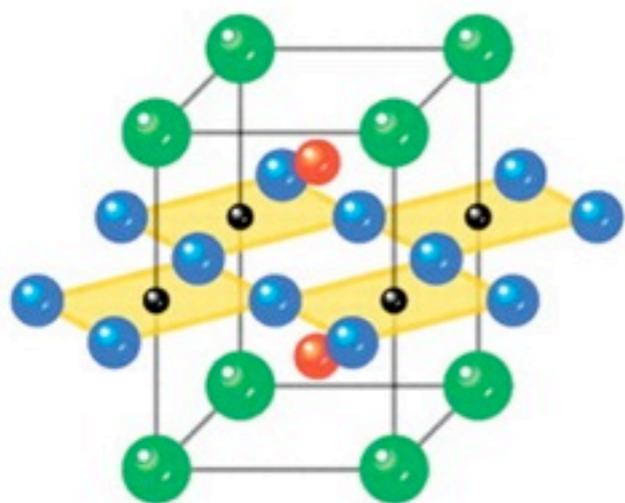
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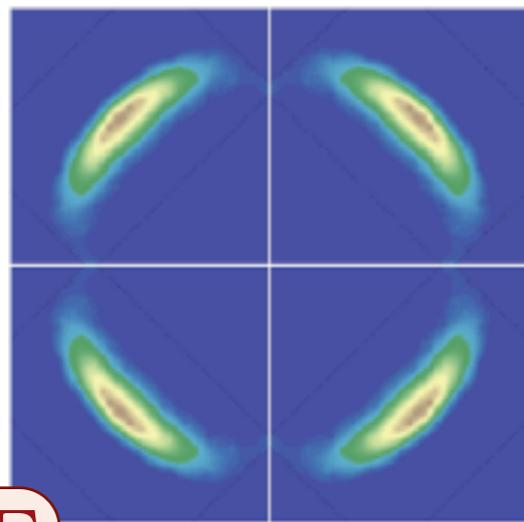
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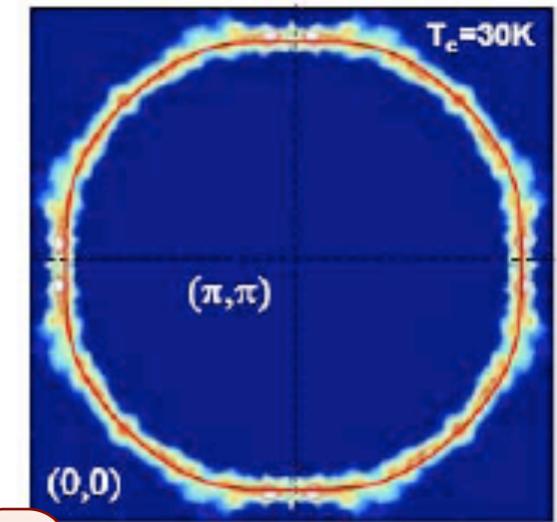
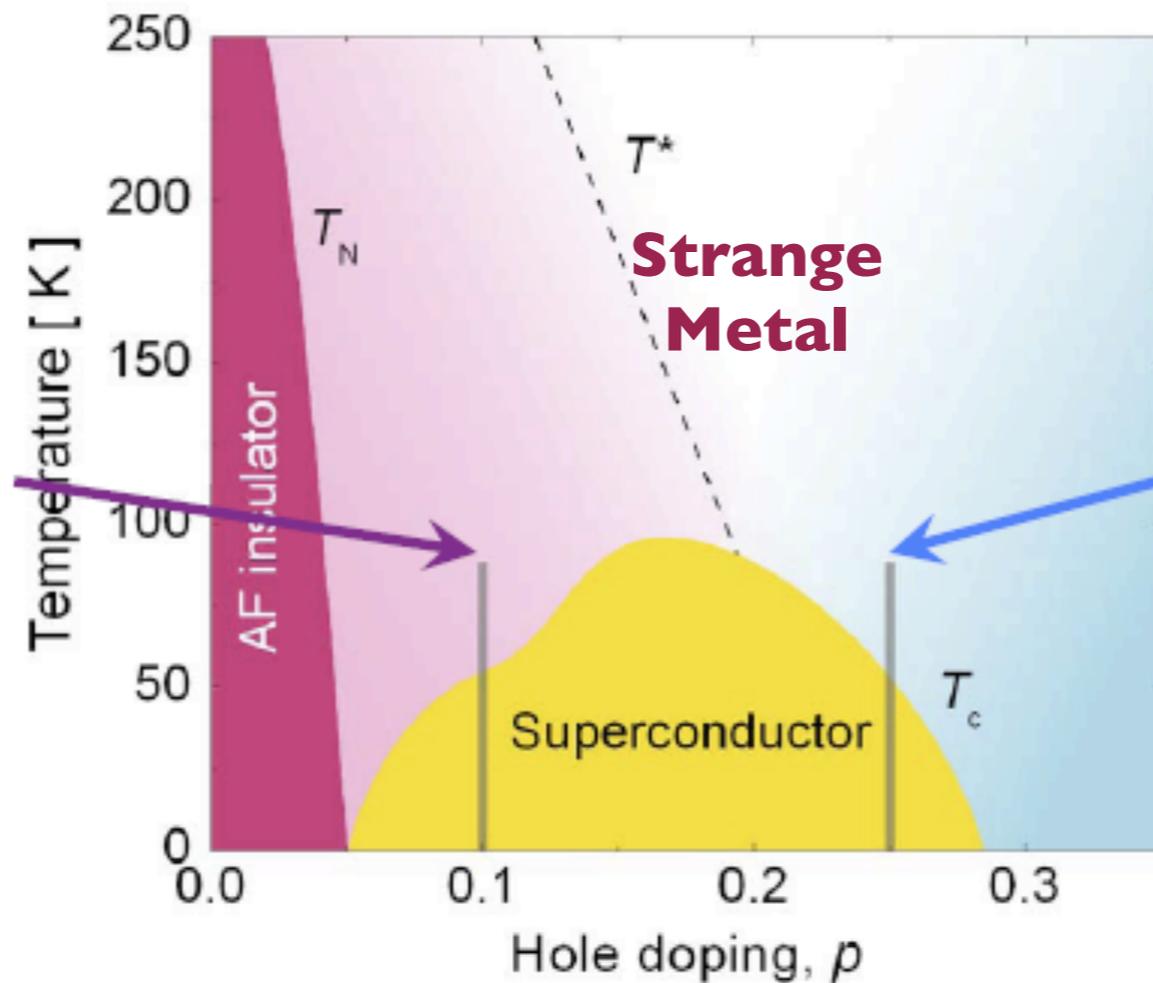
Incommensurate/  
disordered  
antiferromagnetism  
and charge order

# Central ingredients in cuprate phase diagram: antiferromagnetism, superconductivity, and change in Fermi surface



$\Gamma$

*K.M. Shen et al., Science 2005*



$\Gamma$

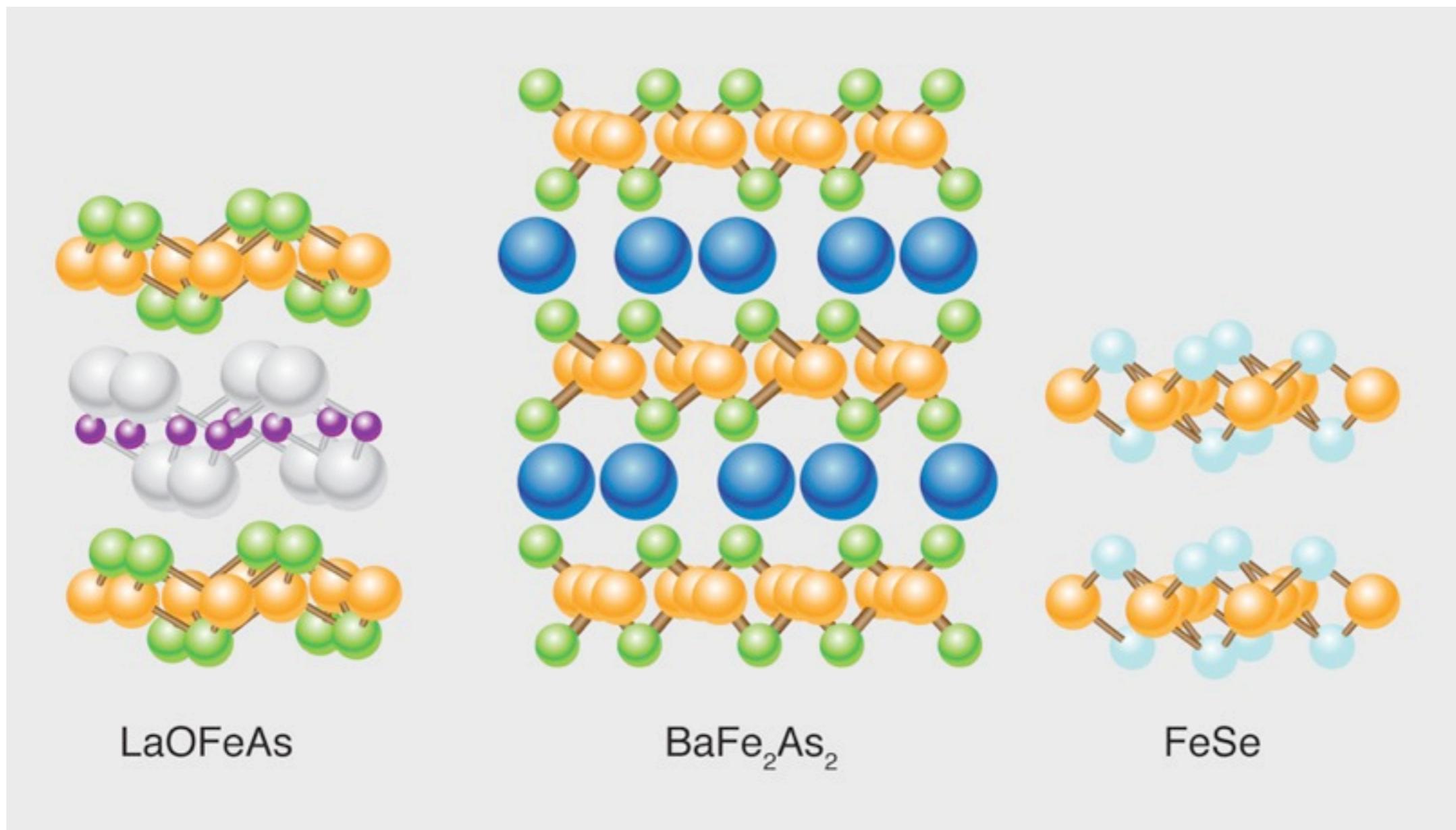
*M. Platé et al., PRL 2005*

Smaller hole  
Fermi-pockets

Large hole  
Fermi surface

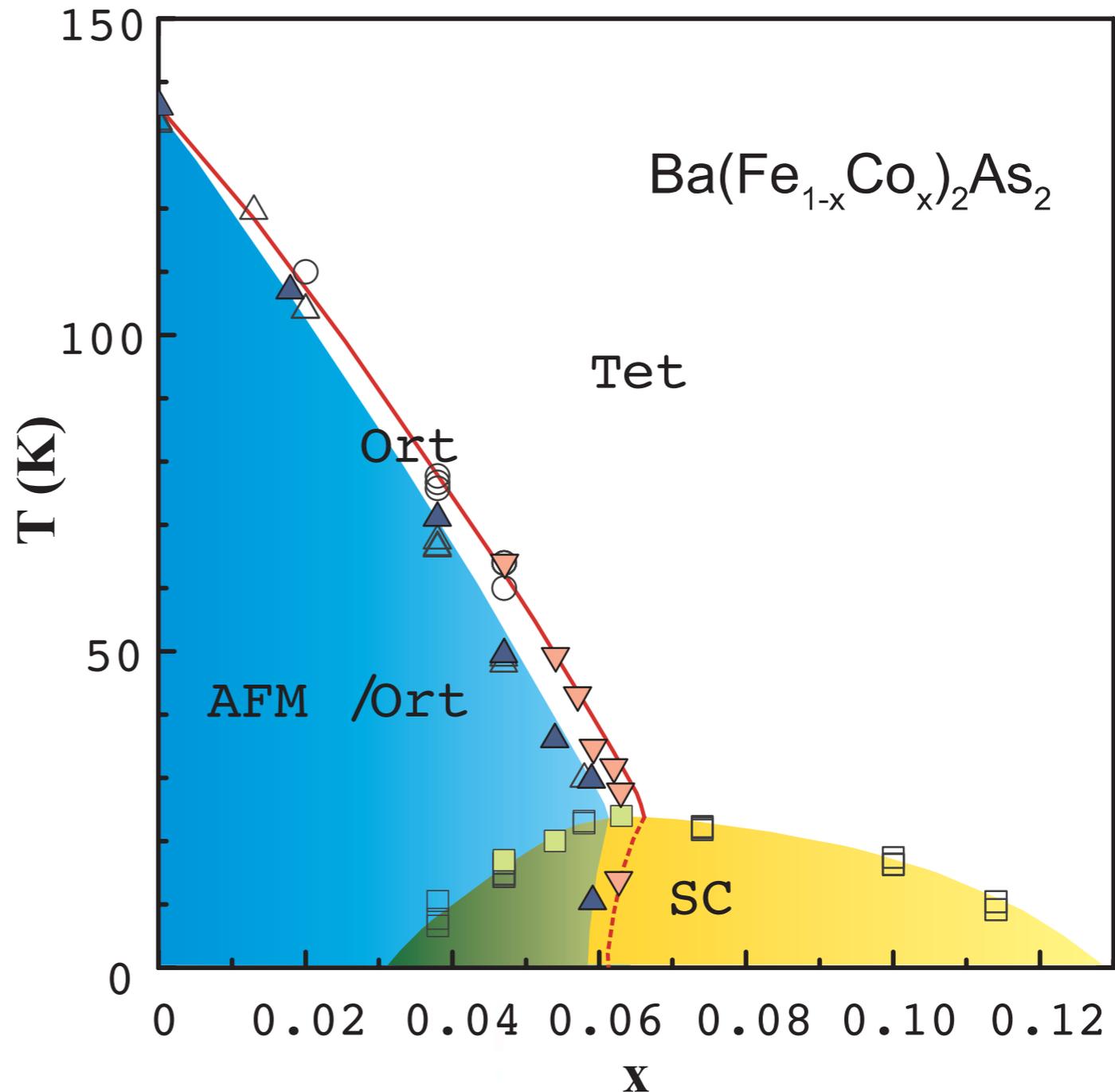
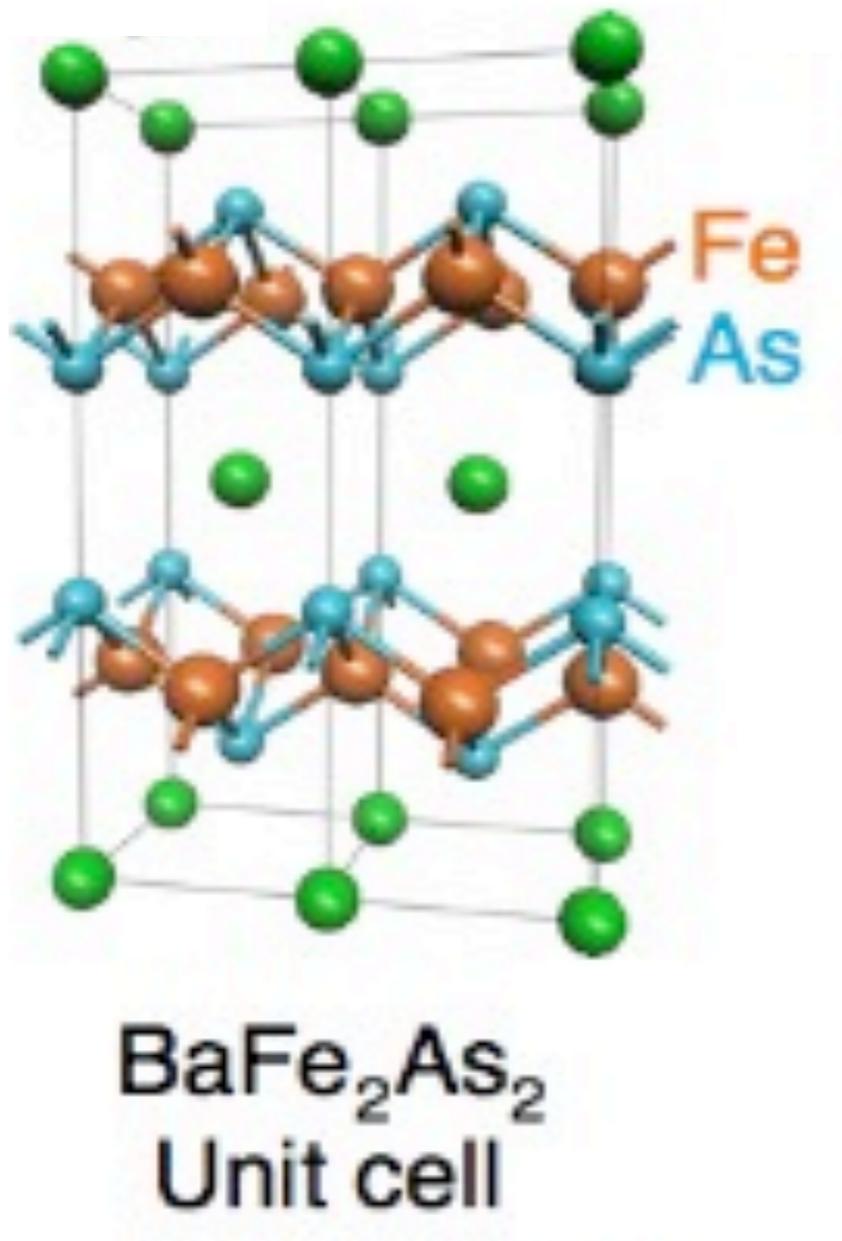
# Iron pnictides:

a new class of high temperature superconductors



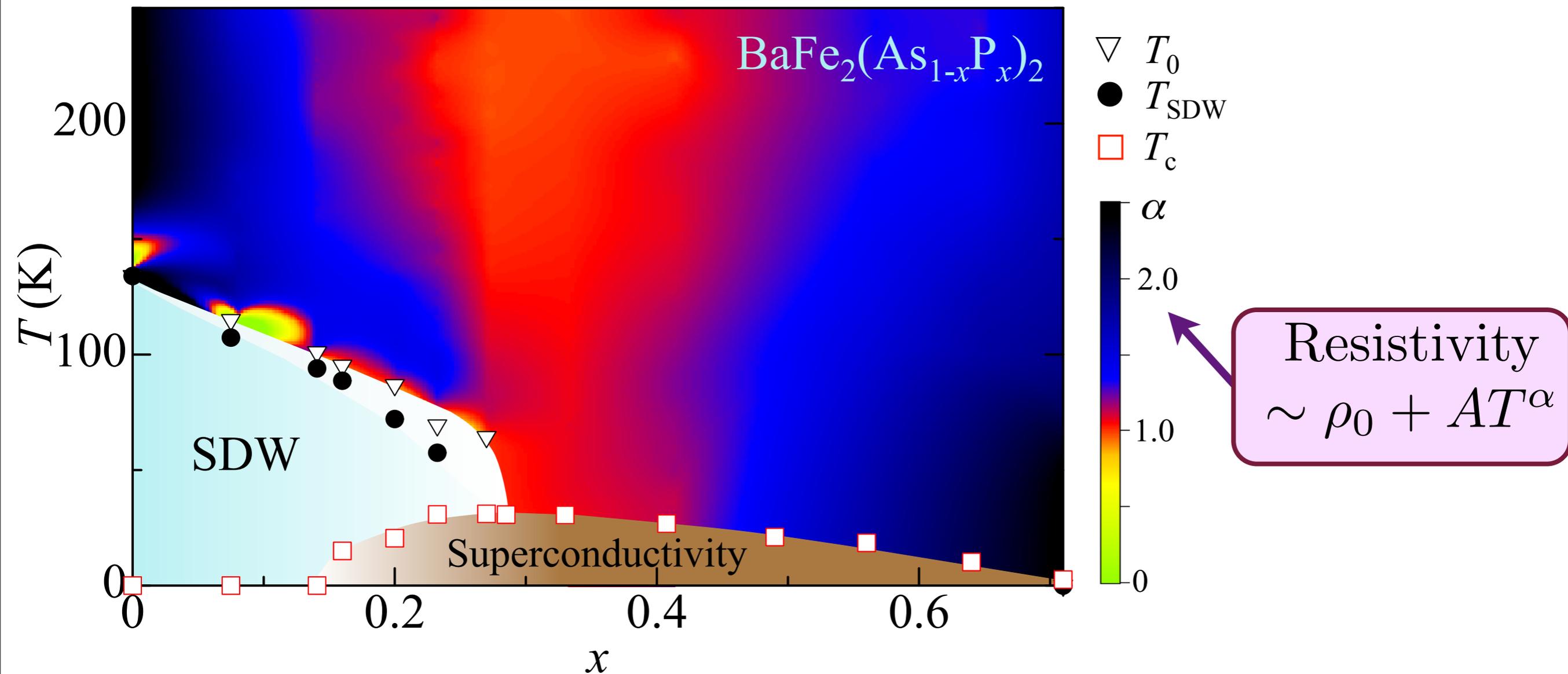
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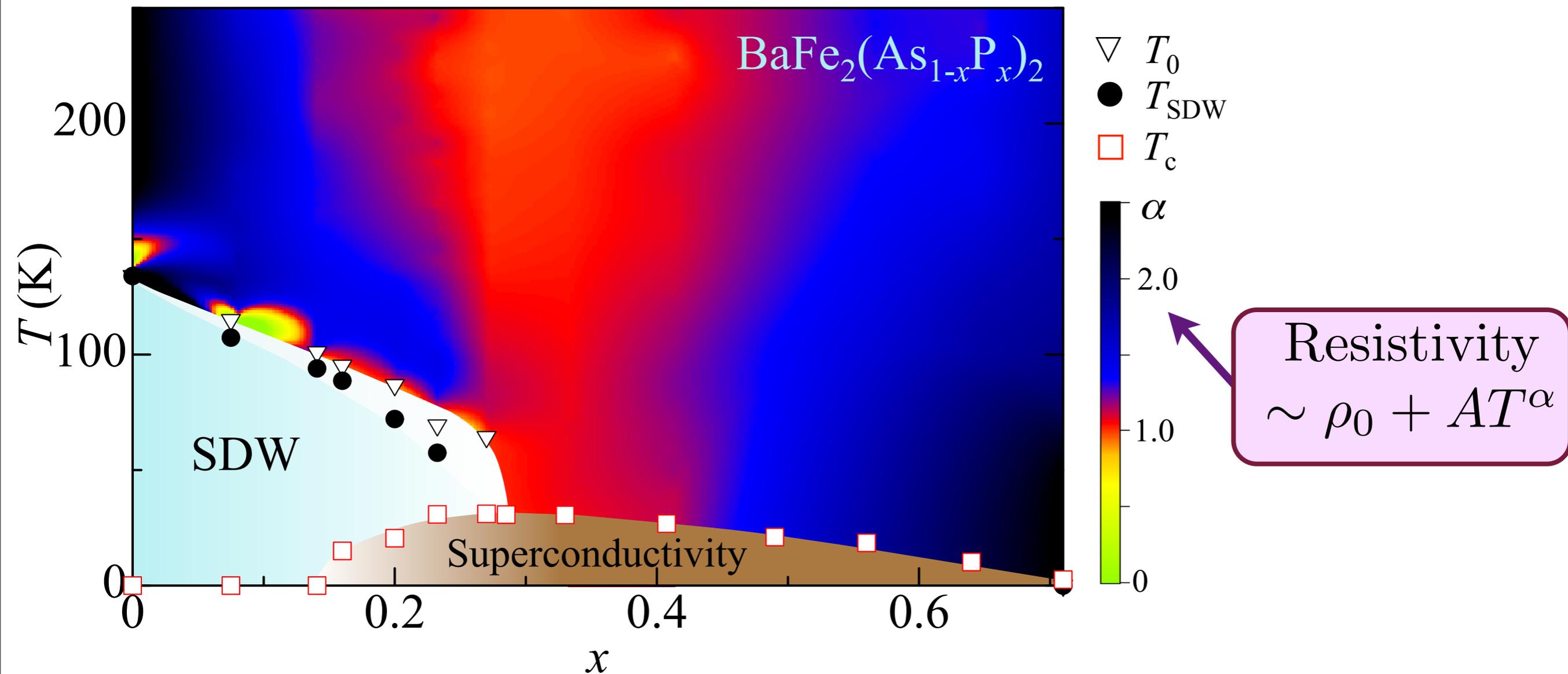
S. Nandi, M. G. Kim, A. Kreyssig, R. M. Fernandes, D. K. Pratt, A. Thaler, N. Ni,  
S. L. Bud'ko, P. C. Canfield, J. Schmalian, R. J. McQueeney, A. I. Goldman,  
*Physical Review Letters* **104**, 057006 (2010).

# Temperature-doping phase diagram of the iron pnictides:



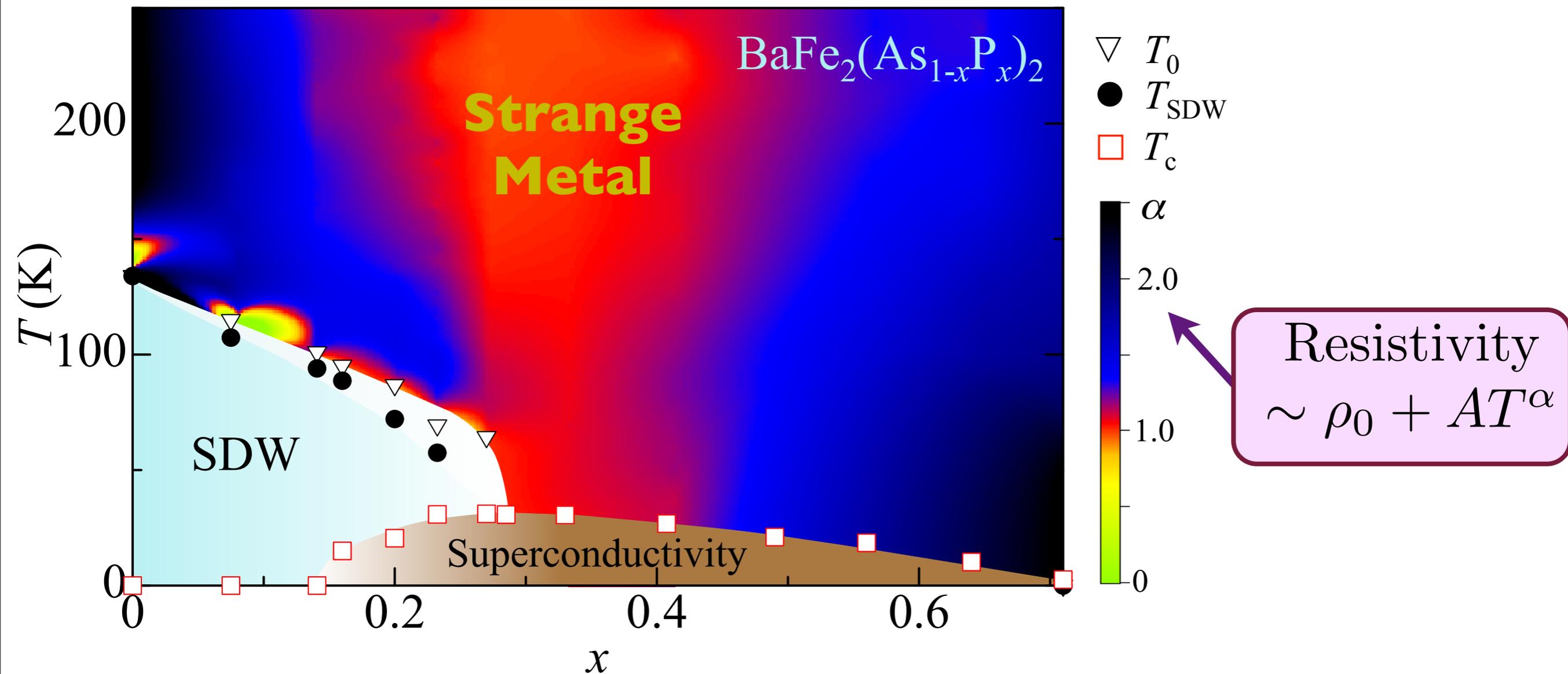
S. Kasahara, T. Shibauchi, K. Hashimoto, K. Ikada, S. Tonegawa, R. Okazaki, H. Shishido, H. Ikeda, H. Takeya, K. Hirata, T. Terashima, and Y. Matsuda, *Physical Review B* **81**, 184519 (2010)

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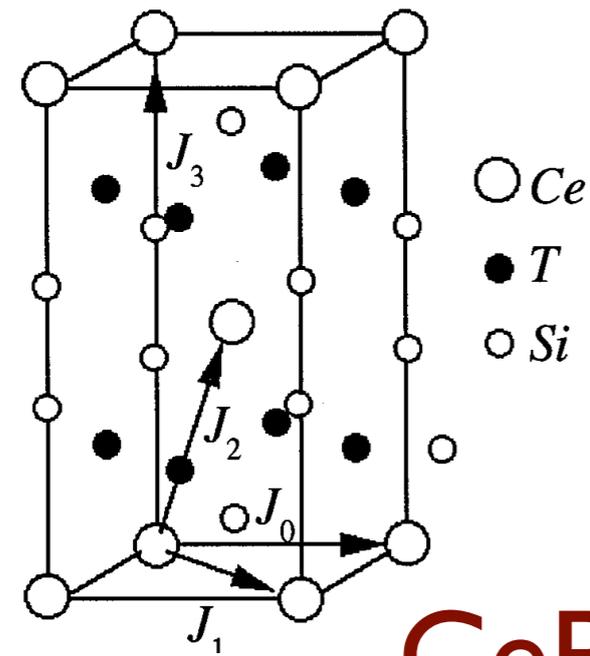
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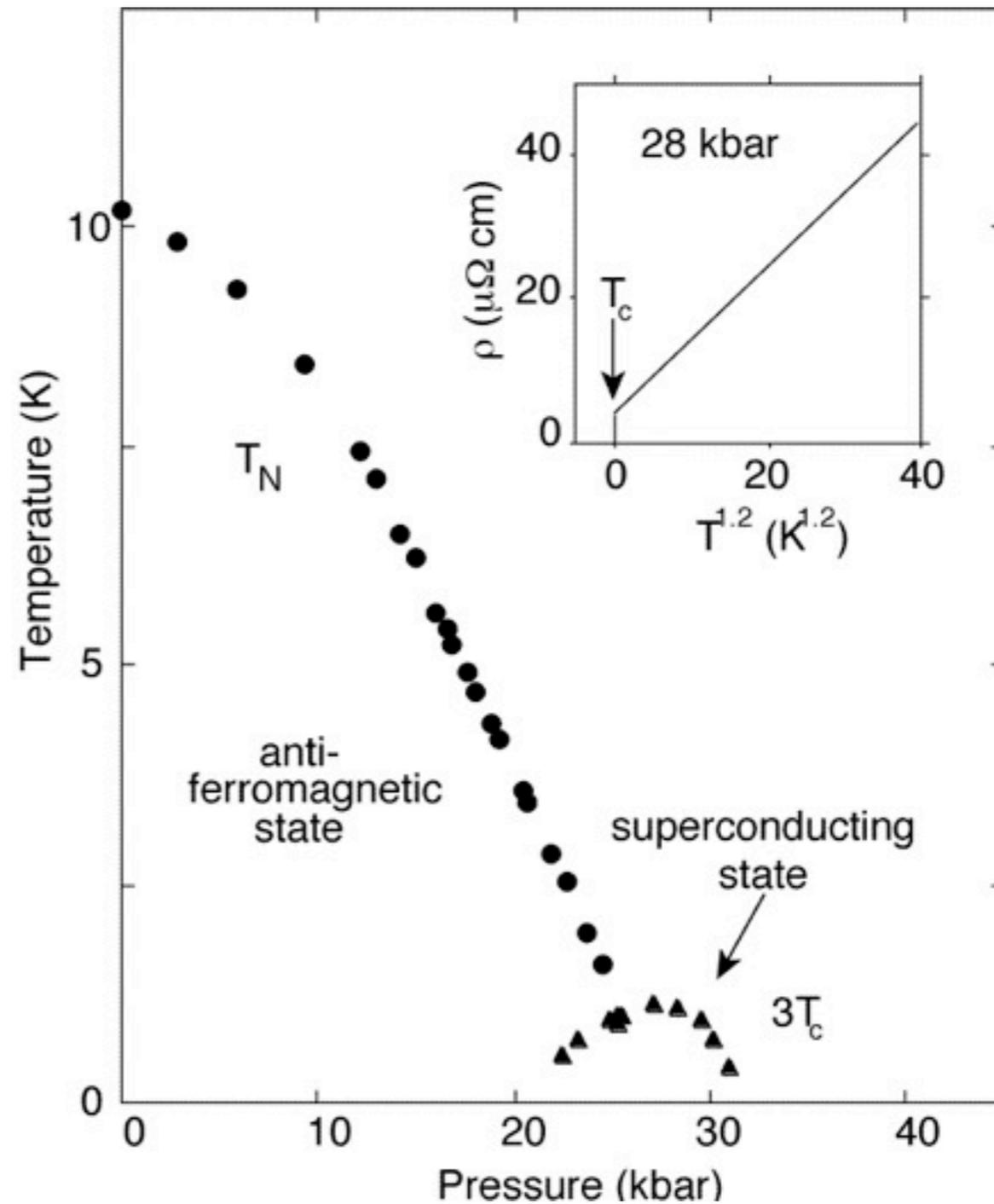
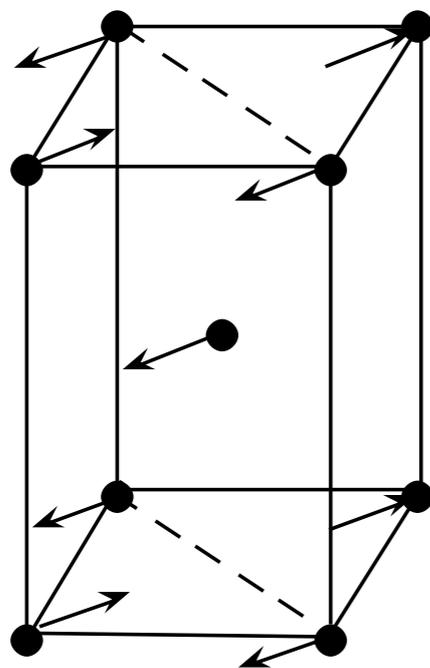


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# Lower $T_c$ superconductivity in the heavy fermion compounds



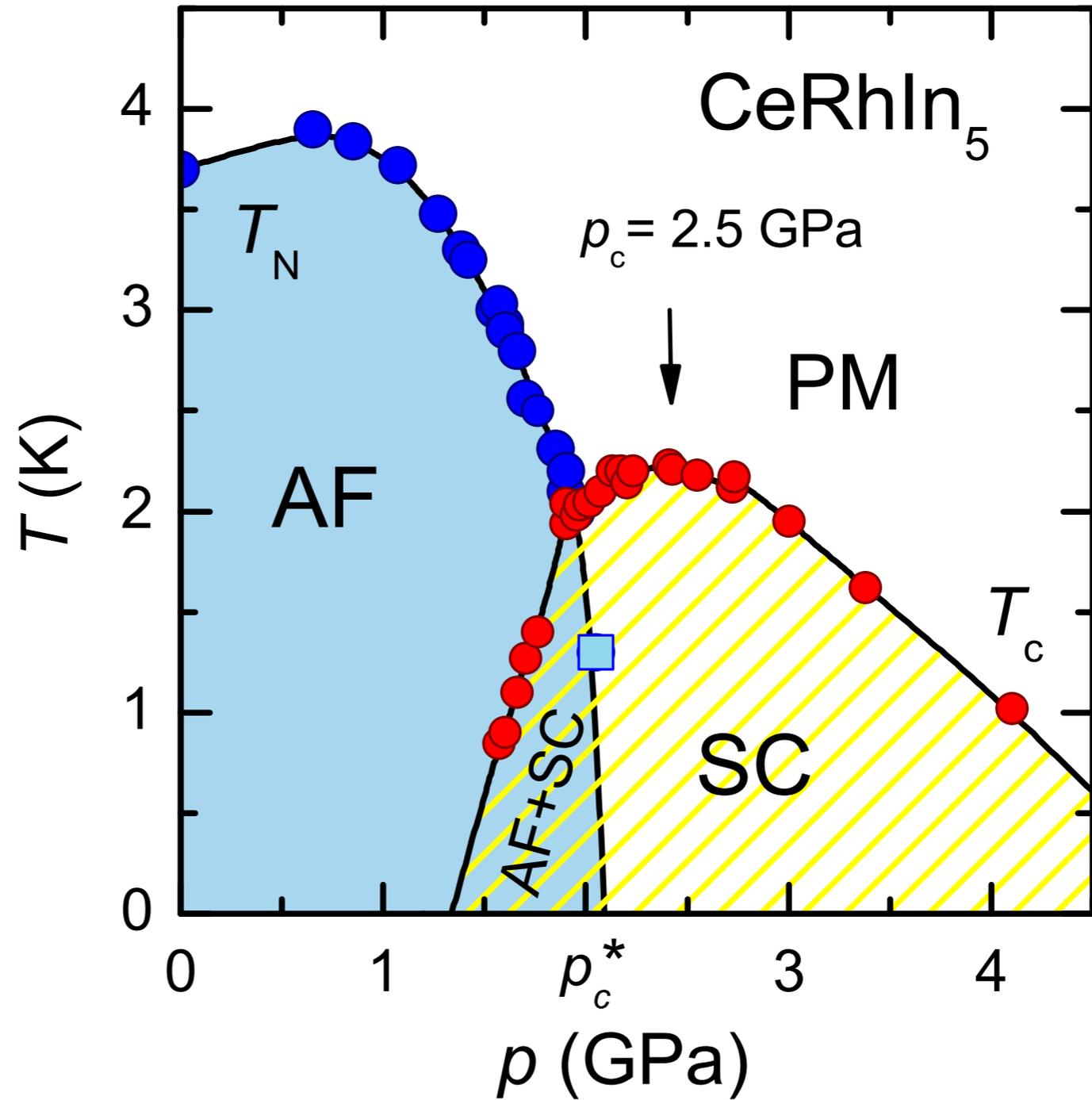
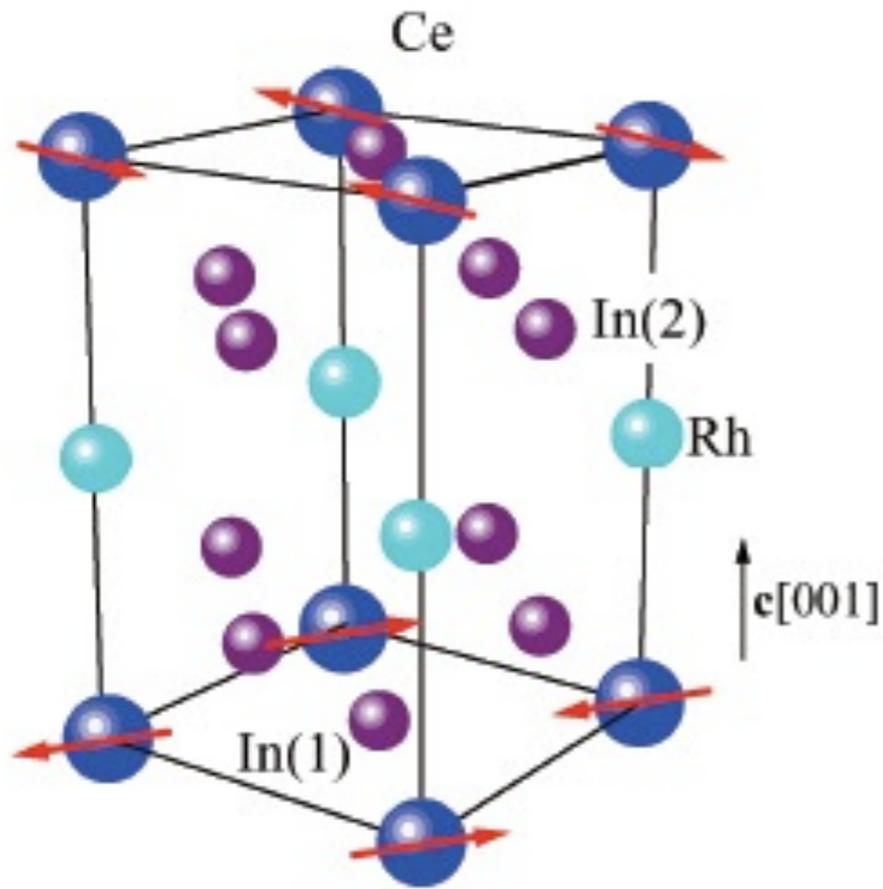
**a**



N. D. Mathur, F. M. Grosche, S. R. Julian, I. R. Walker, D. M. Freye, R. K.W. Haselwimmer, and G. G. Lonzarich, Nature **394**, 39 (1998)



# Lower $T_c$ superconductivity in the heavy fermion compounds



G. Knebel, D. Aoki, and J. Flouquet, arXiv:0911.5223

## Questions

- *Can quantum fluctuations near the loss of antiferromagnetism induce higher temperature superconductivity ?*

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- *If so, why is there no antiferromagnetism in the cuprates near the point where the superconductivity is strongest ?*
- *What is the physics of the strange metal ?*

# Outline

1. Loss of antiferromagnetism in an insulator  
*Coupled-dimer antiferromagnets and quantum criticality*
2. Onset of antiferromagnetism in a metal  
*From large Fermi surfaces to Fermi pockets*
3. Unconventional superconductivity  
*Pairing from antiferromagnetic fluctuations*
4. Competing orders  
*Phase diagram in a magnetic field*
5. Strongly-coupled quantum criticality in metals  
*Fluctuating antiferromagnetism and Fermi surfaces*

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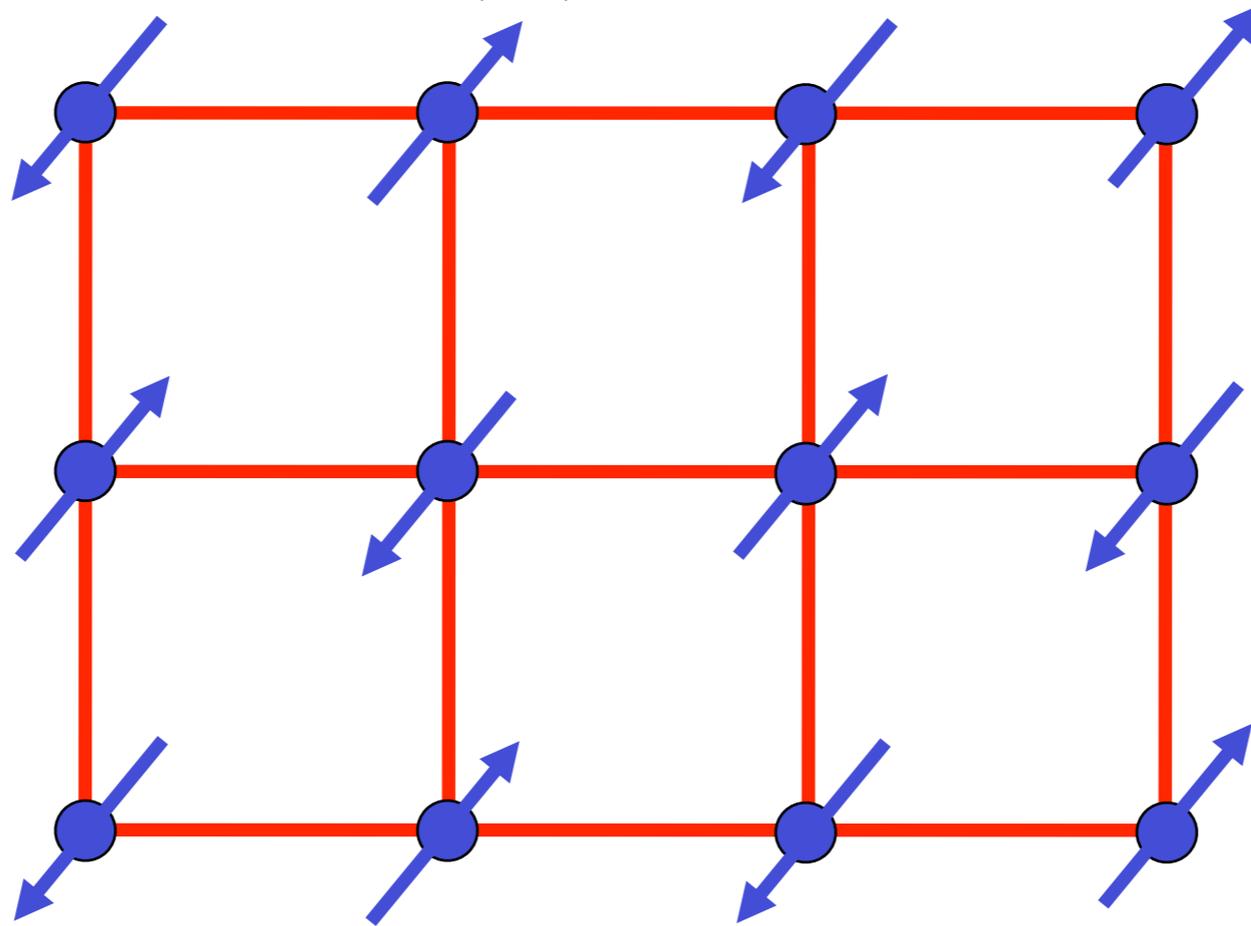
*Phase diagram in a magnetic field*

## 5. Strongly-coupled quantum criticality in metals

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# Square lattice antiferromagnet

$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$



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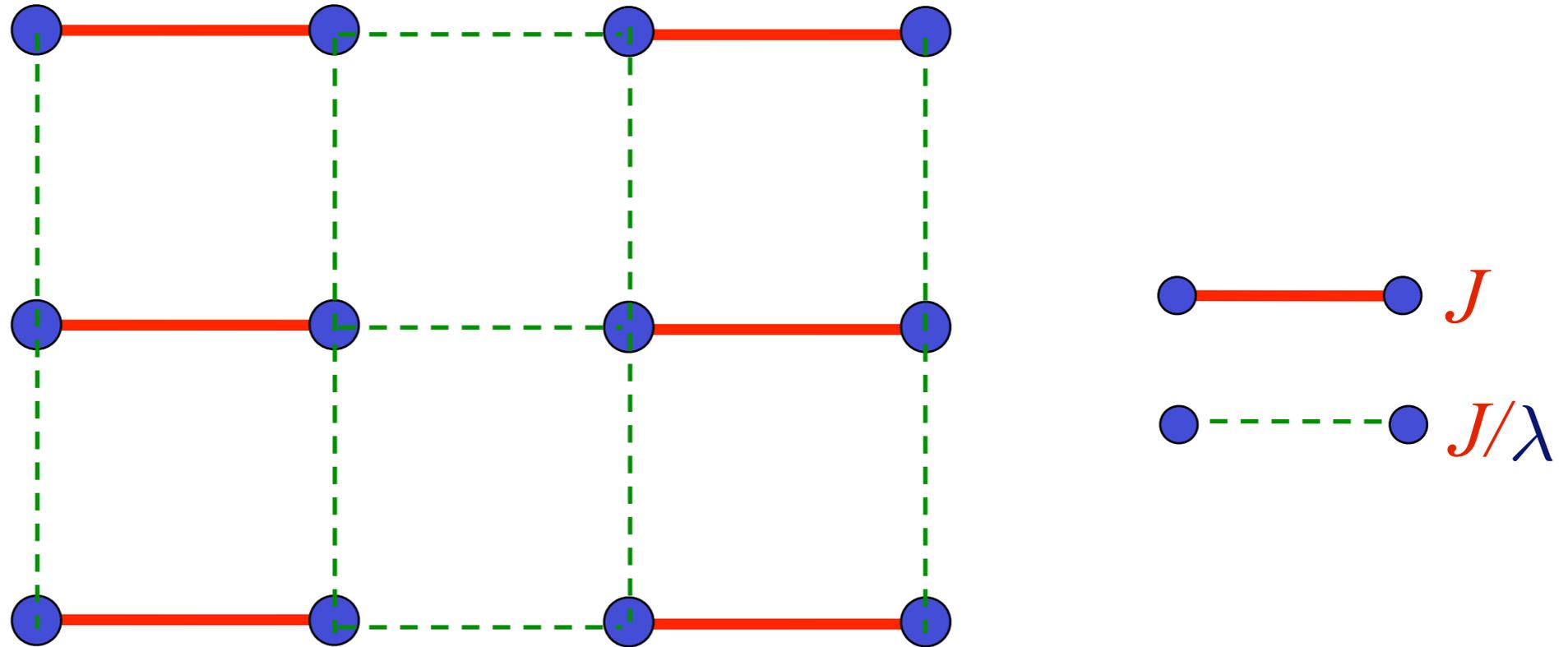
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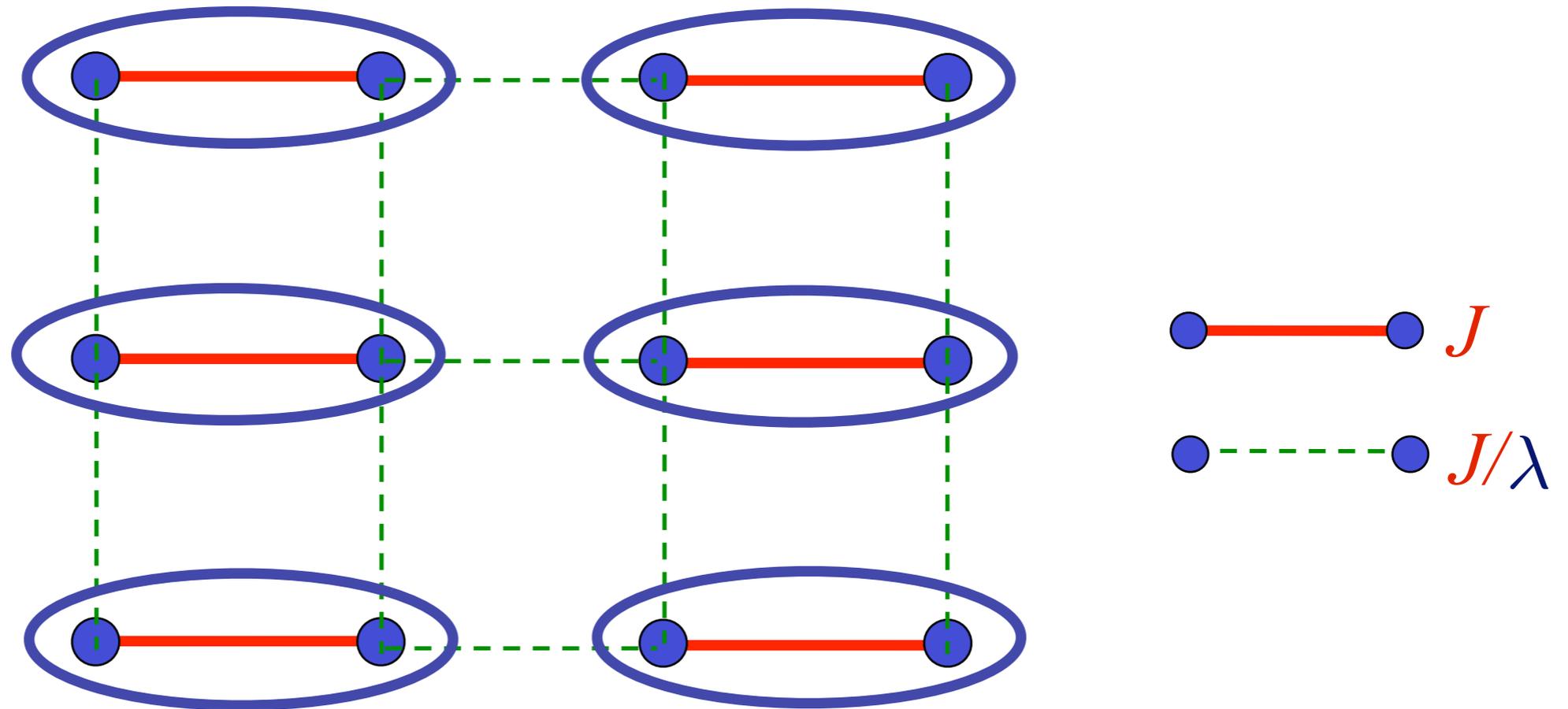


Weaken some bonds to induce spin entanglement in a new quantum phase



# Square lattice antiferromagnet

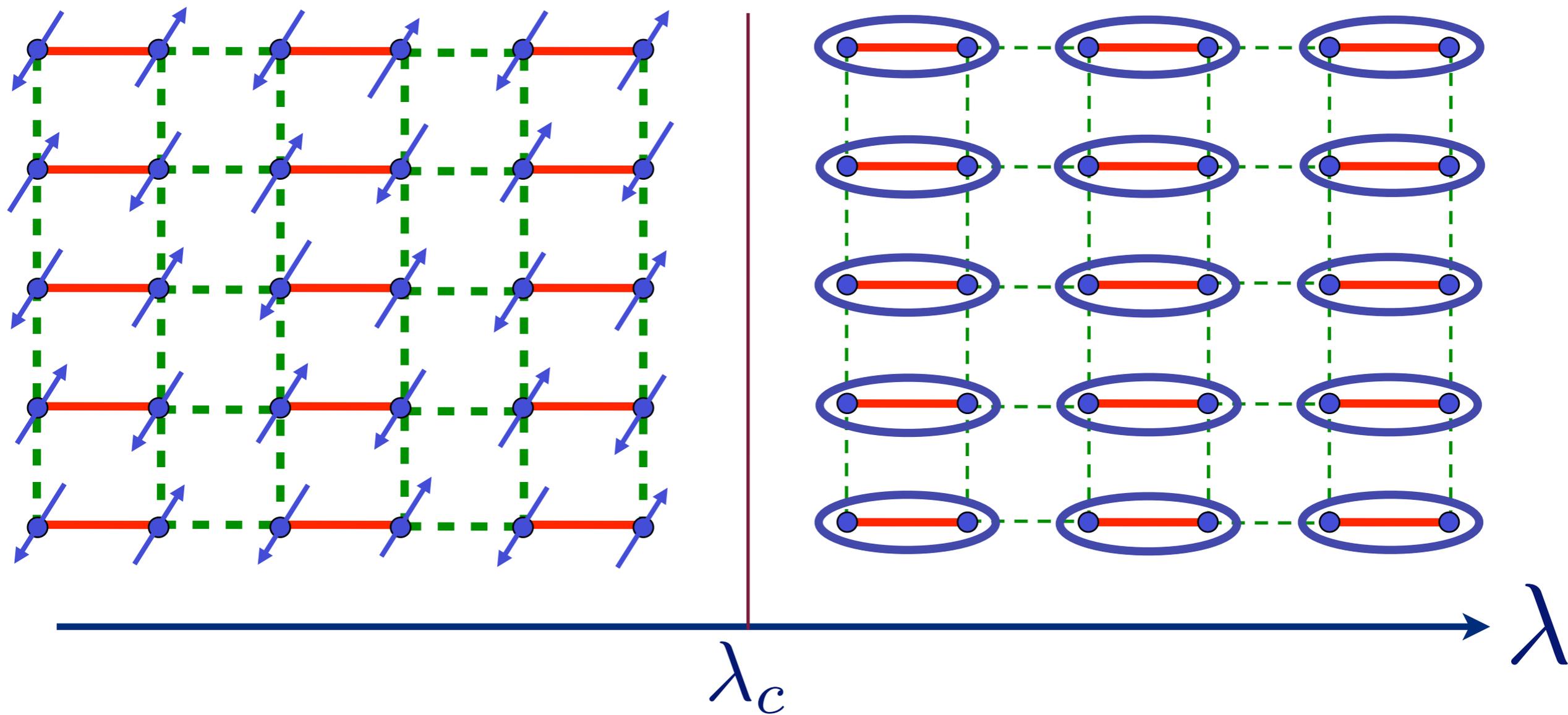
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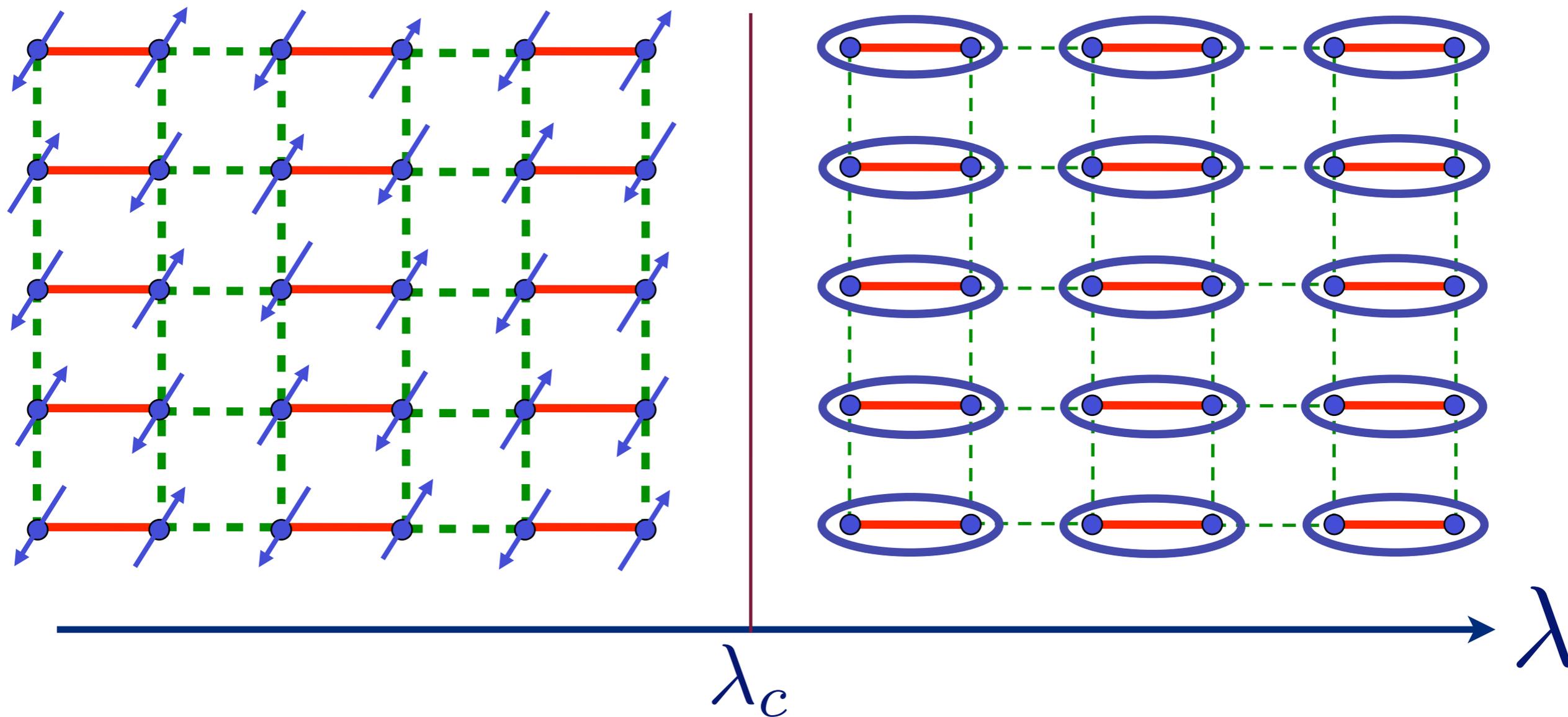
Ground state is a “quantum paramagnet”  
with spins locked in valence bond singlets

$$\text{[Diagram of a valence bond singlet]} = \frac{1}{\sqrt{2}} \left( |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$

$$\text{Diagram of two blue dots connected by a red line, enclosed in a blue oval} = \frac{1}{\sqrt{2}} \left( |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$



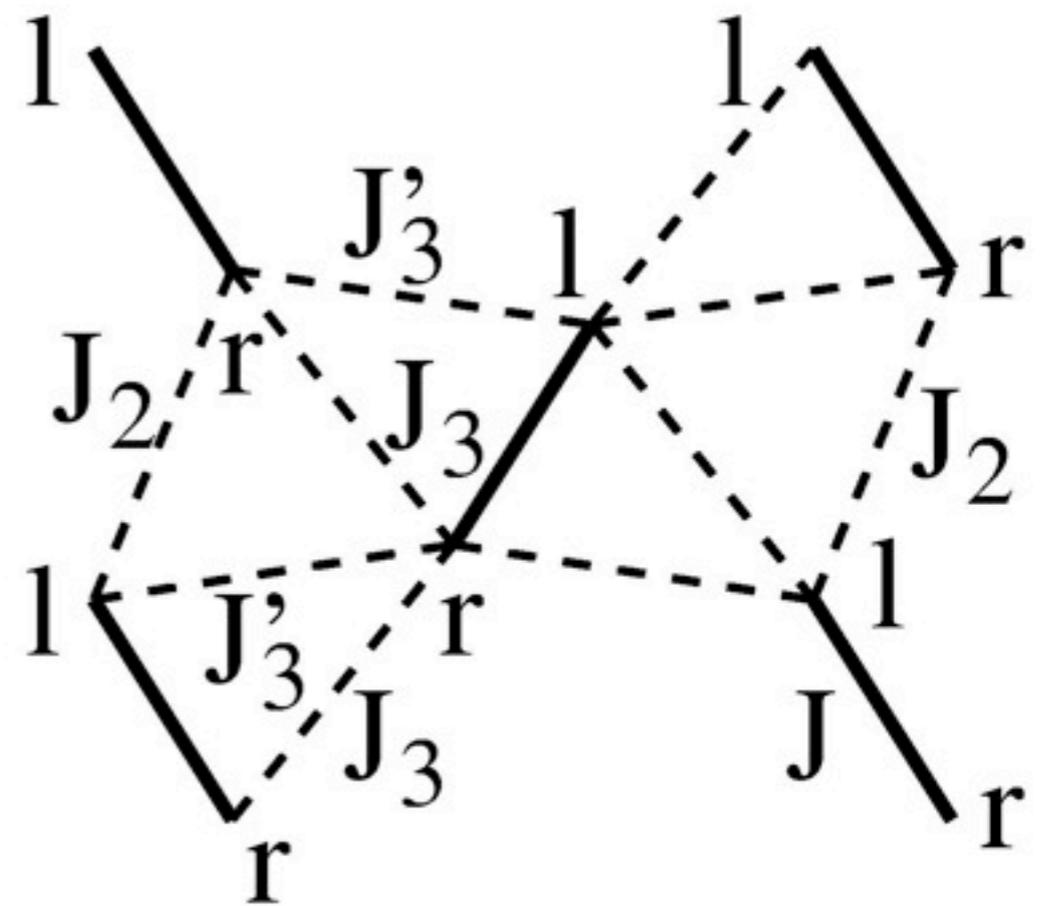
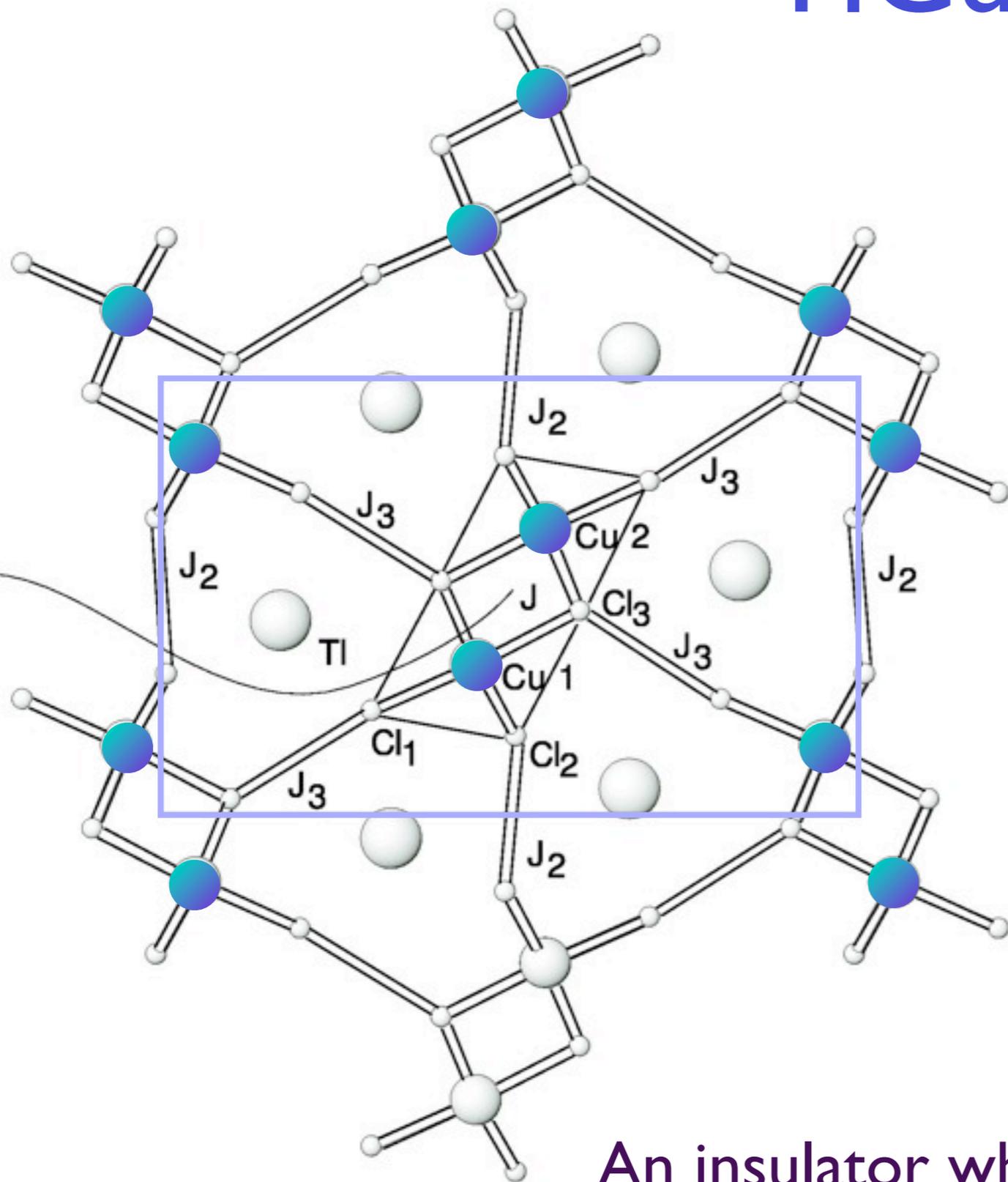
$$\text{Diagram of two blue spheres connected by a red line, enclosed in a blue oval} = \frac{1}{\sqrt{2}} \left( |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$



← Pressure in  $\text{TiCuCl}_3$

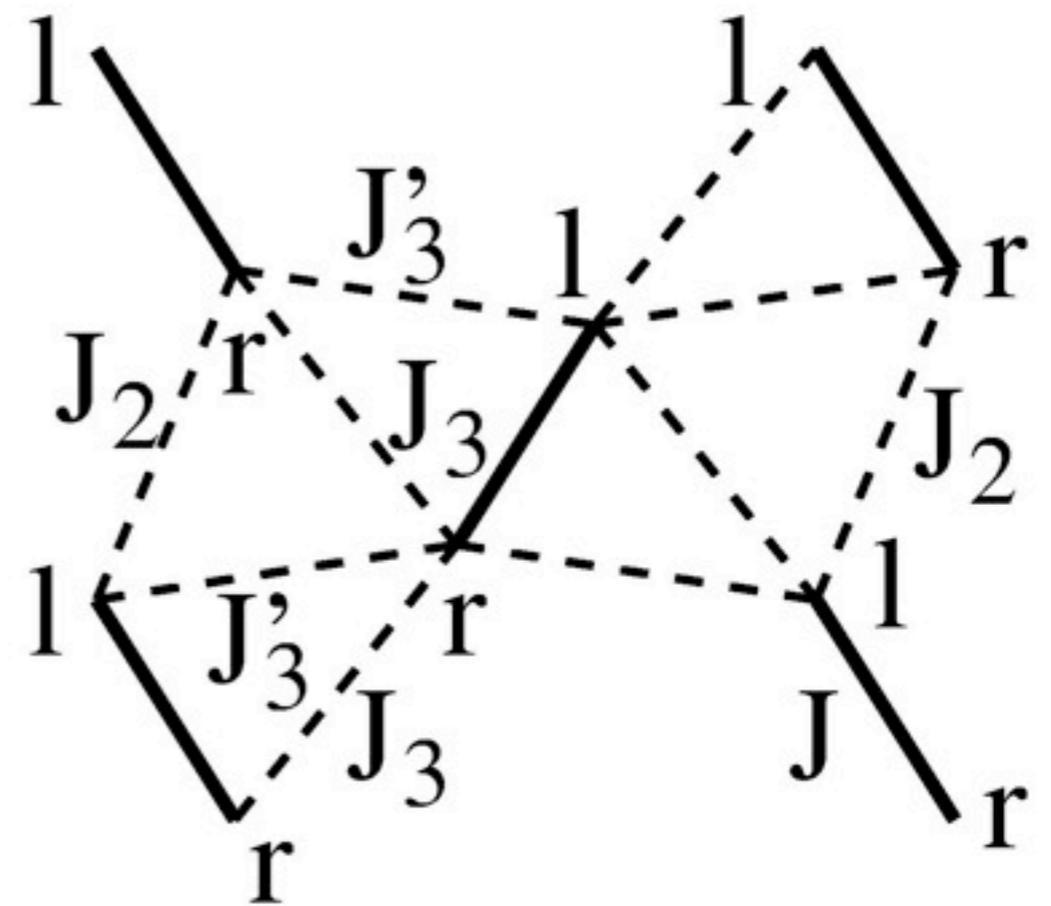
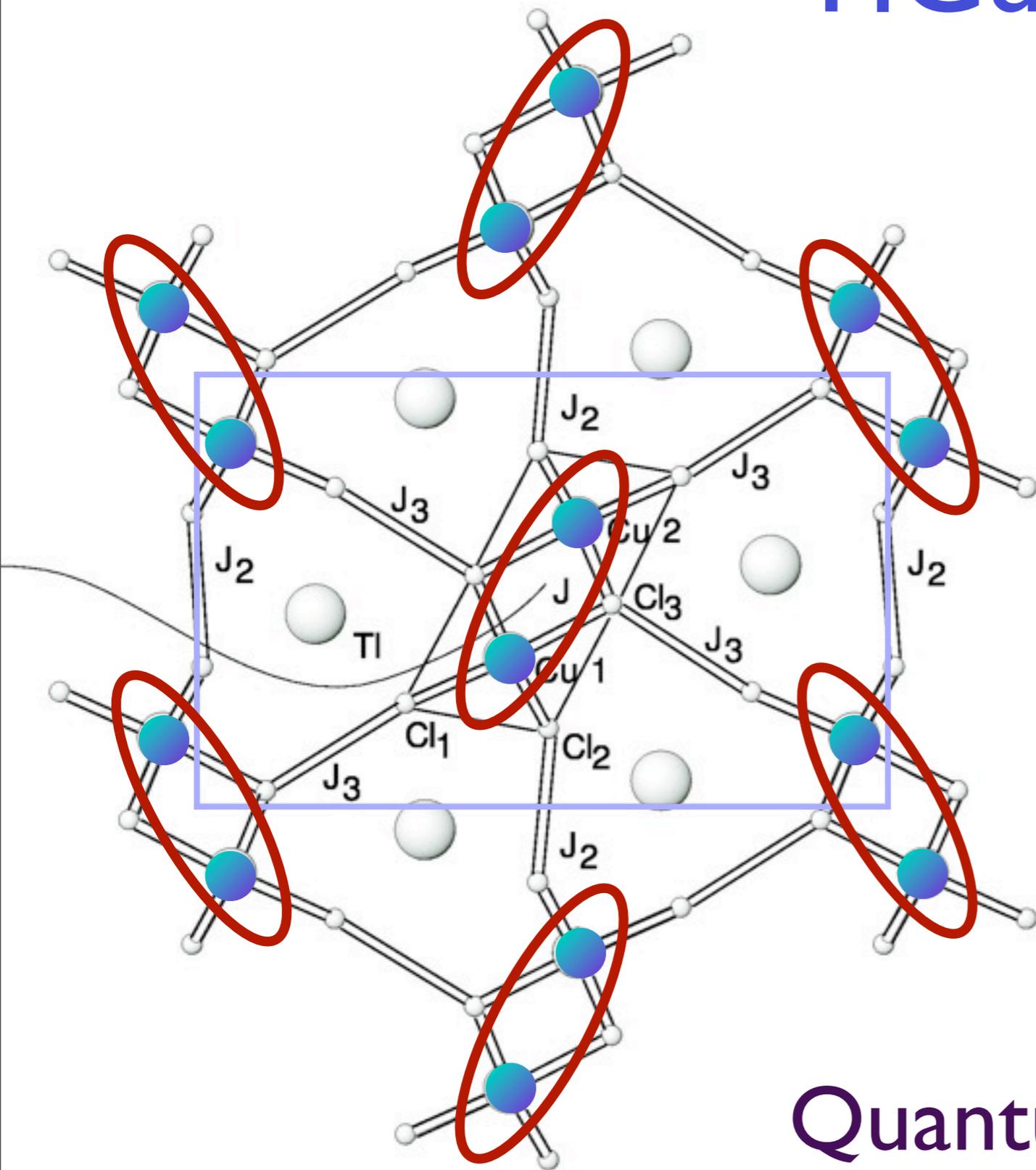
A. Oosawa, K. Kakurai, T. Osakabe, M. Nakamura, M. Takeda, and H. Tanaka, *Journal of the Physical Society of Japan*, **73**, 1446 (2004).

# TlCuCl<sub>3</sub>



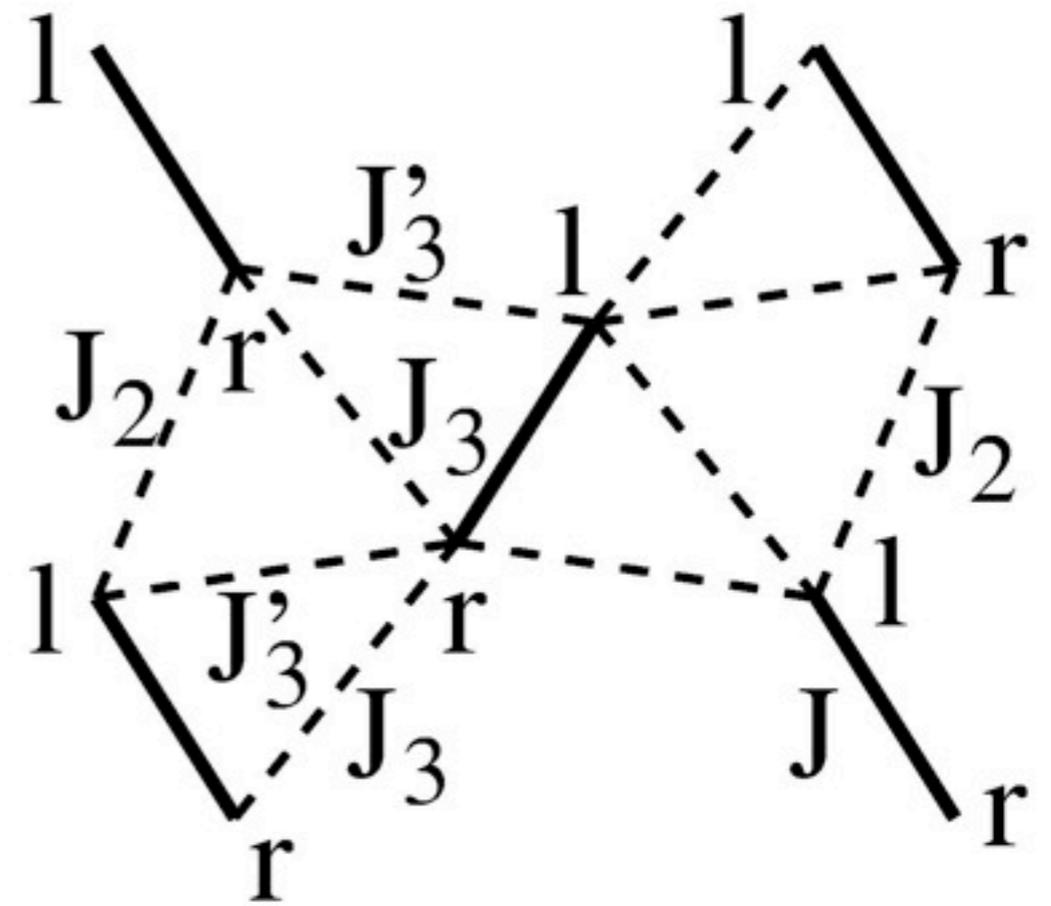
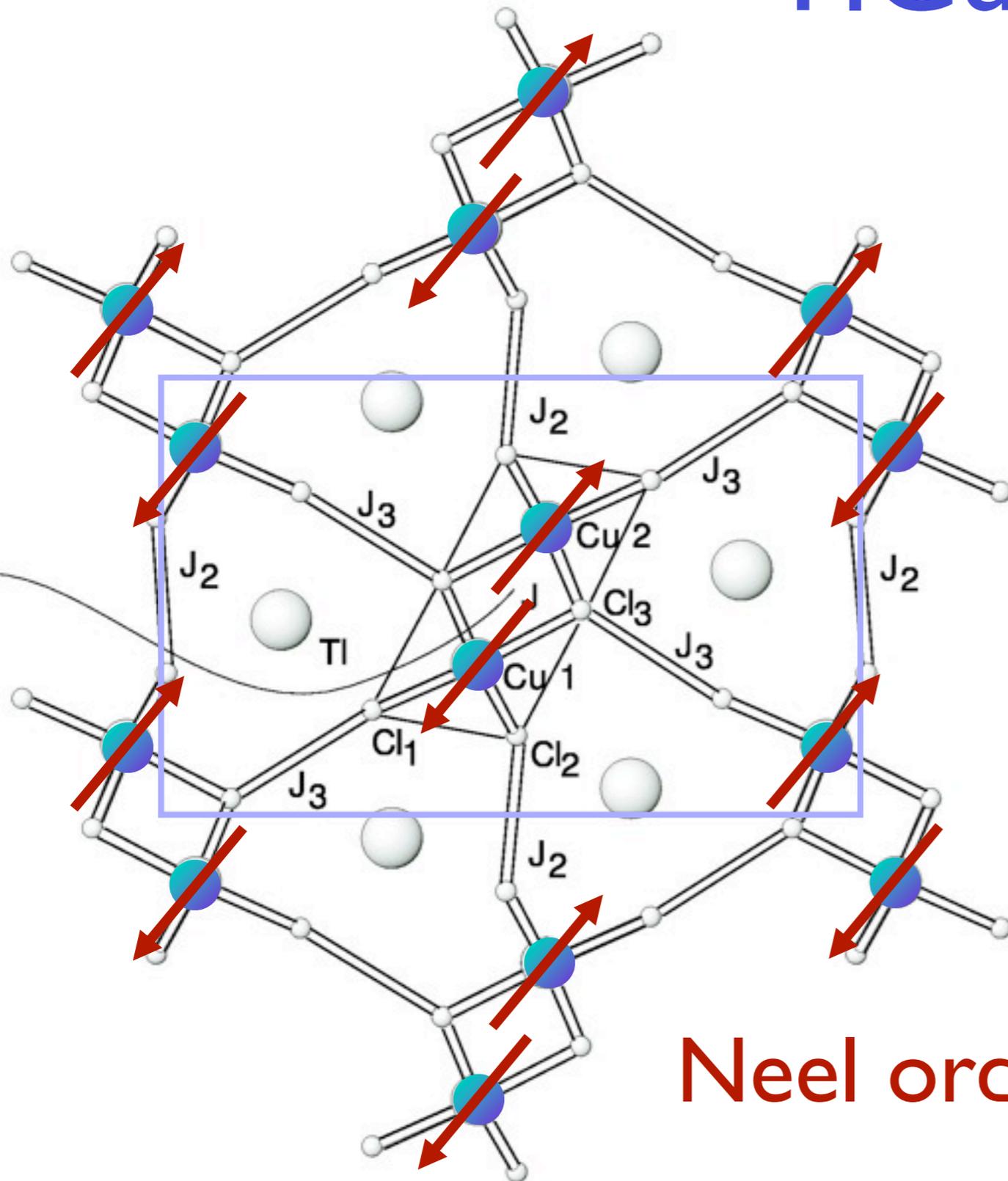
An insulator whose spin susceptibility vanishes exponentially as the temperature  $T$  tends to zero.

# TlCuCl<sub>3</sub>



Quantum paramagnet at  
ambient pressure

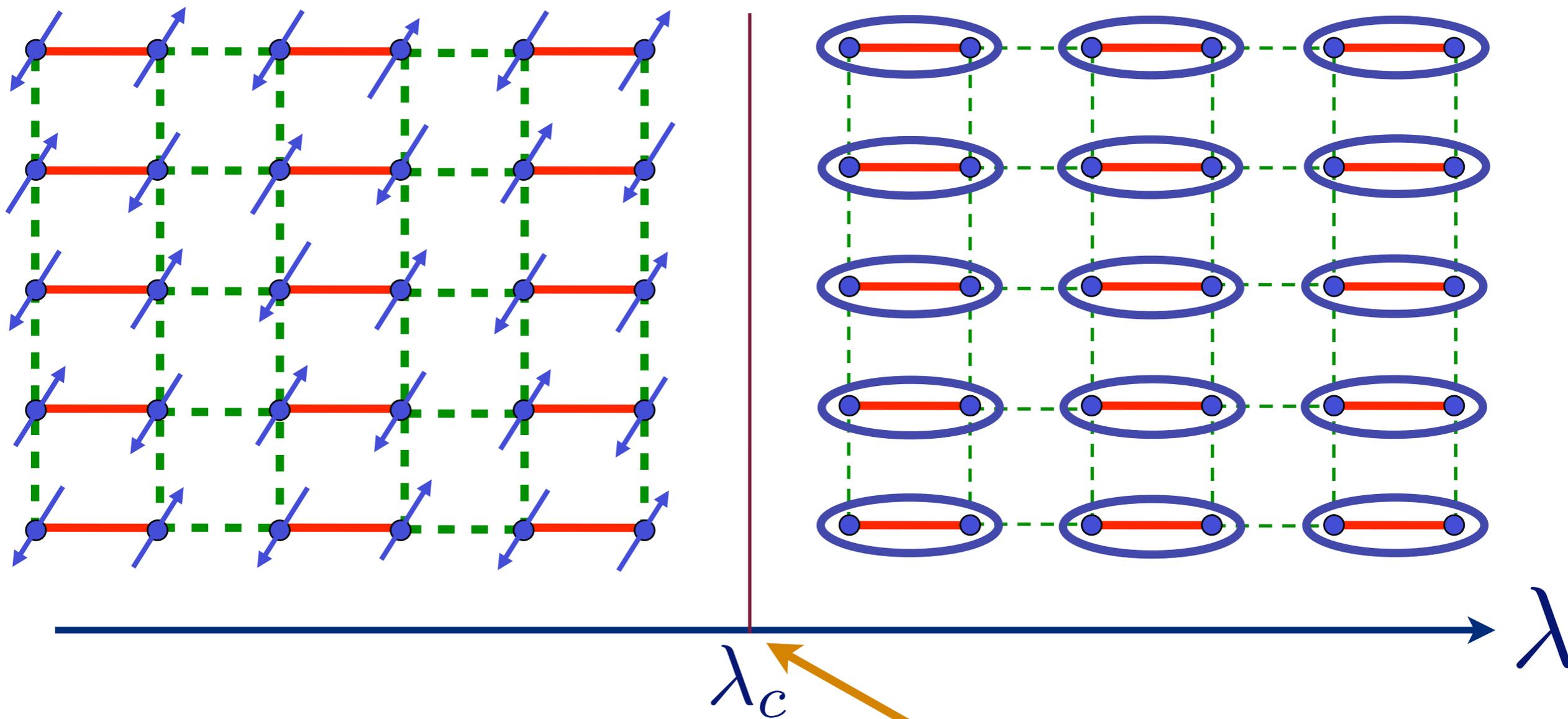
# TlCuCl<sub>3</sub>



Neel order under pressure

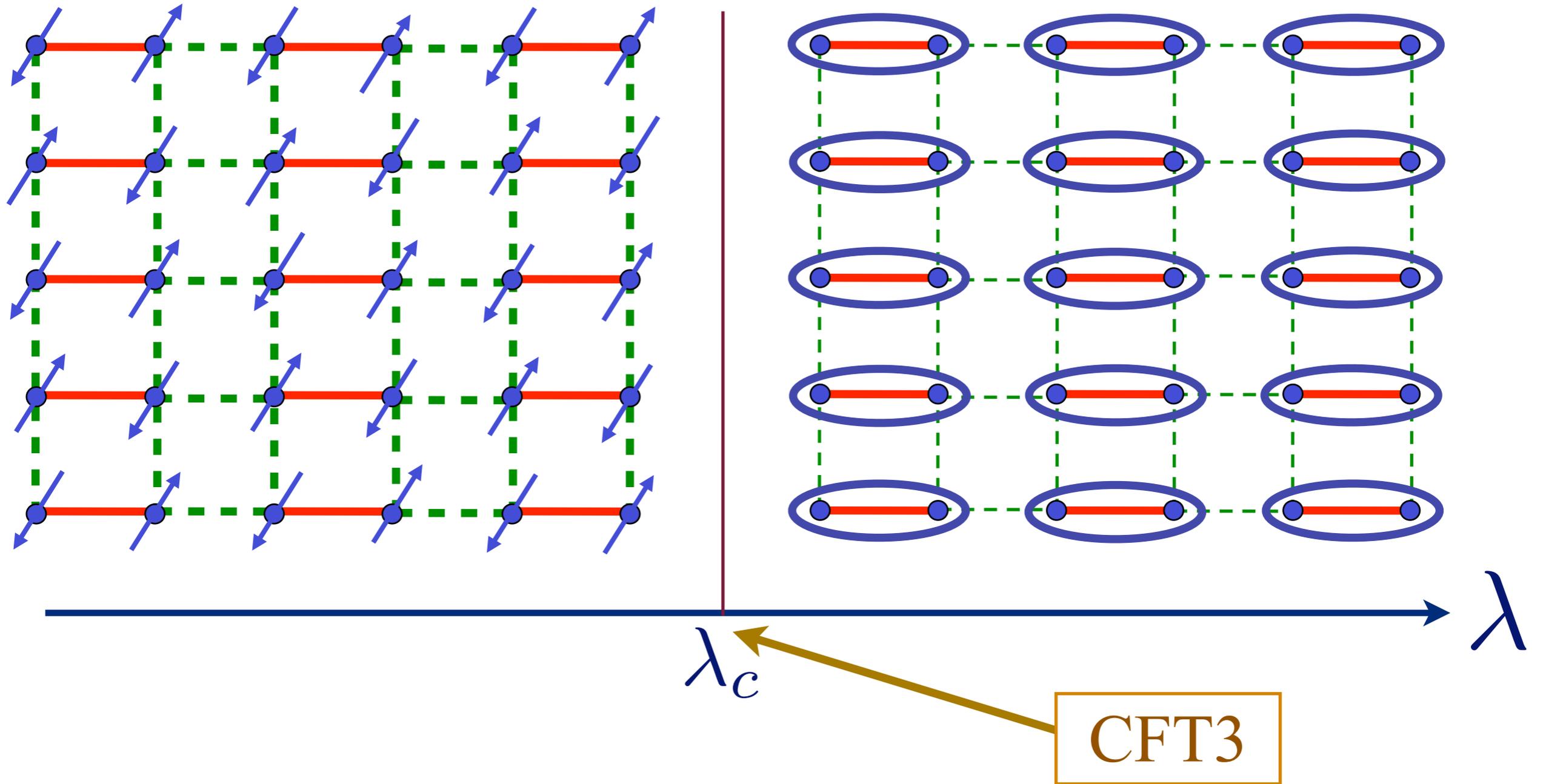
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$$\text{[Diagram of two blue dots connected by a red line, enclosed in a blue oval]} = \frac{1}{\sqrt{2}} \left( |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$



Quantum critical point with non-local entanglement in spin wavefunction

# Description using Landau-Ginzburg field theory

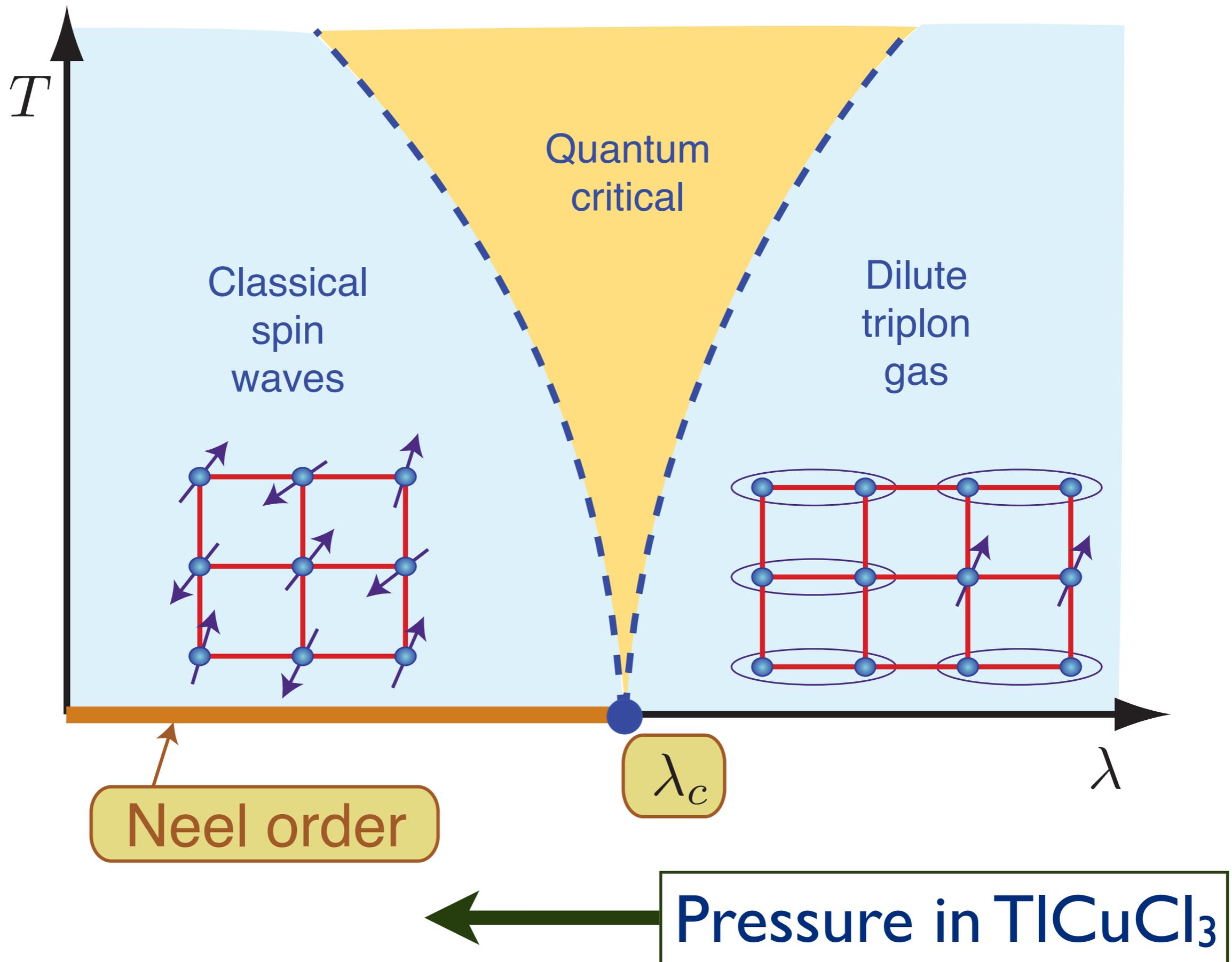


$O(3)$  order parameter  $\vec{\varphi}$

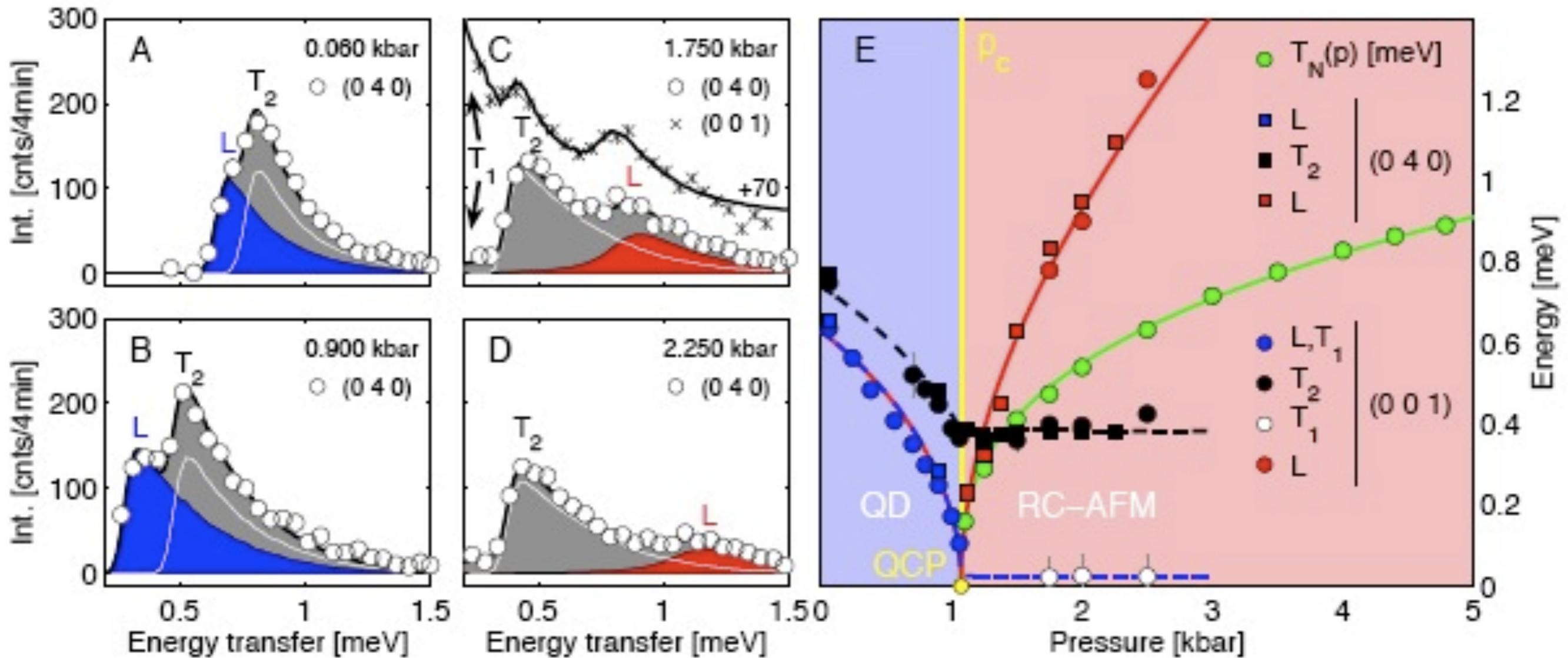
$$\mathcal{S} = \int d^2 r d\tau \left[ (\partial_\tau \vec{\varphi})^2 + c^2 (\nabla_r \vec{\varphi})^2 + (\lambda - \lambda_c) \vec{\varphi}^2 + u (\vec{\varphi}^2)^2 \right]$$



S. Sachdev and J. Ye, *Phys. Rev. Lett.* **69**, 2411 (1992).  
A. V. Chubukov, S. Sachdev, and J. Ye, *Phys. Rev. B* **49**, 11919 (1994).



# TiCuCl<sub>3</sub> with varying pressure



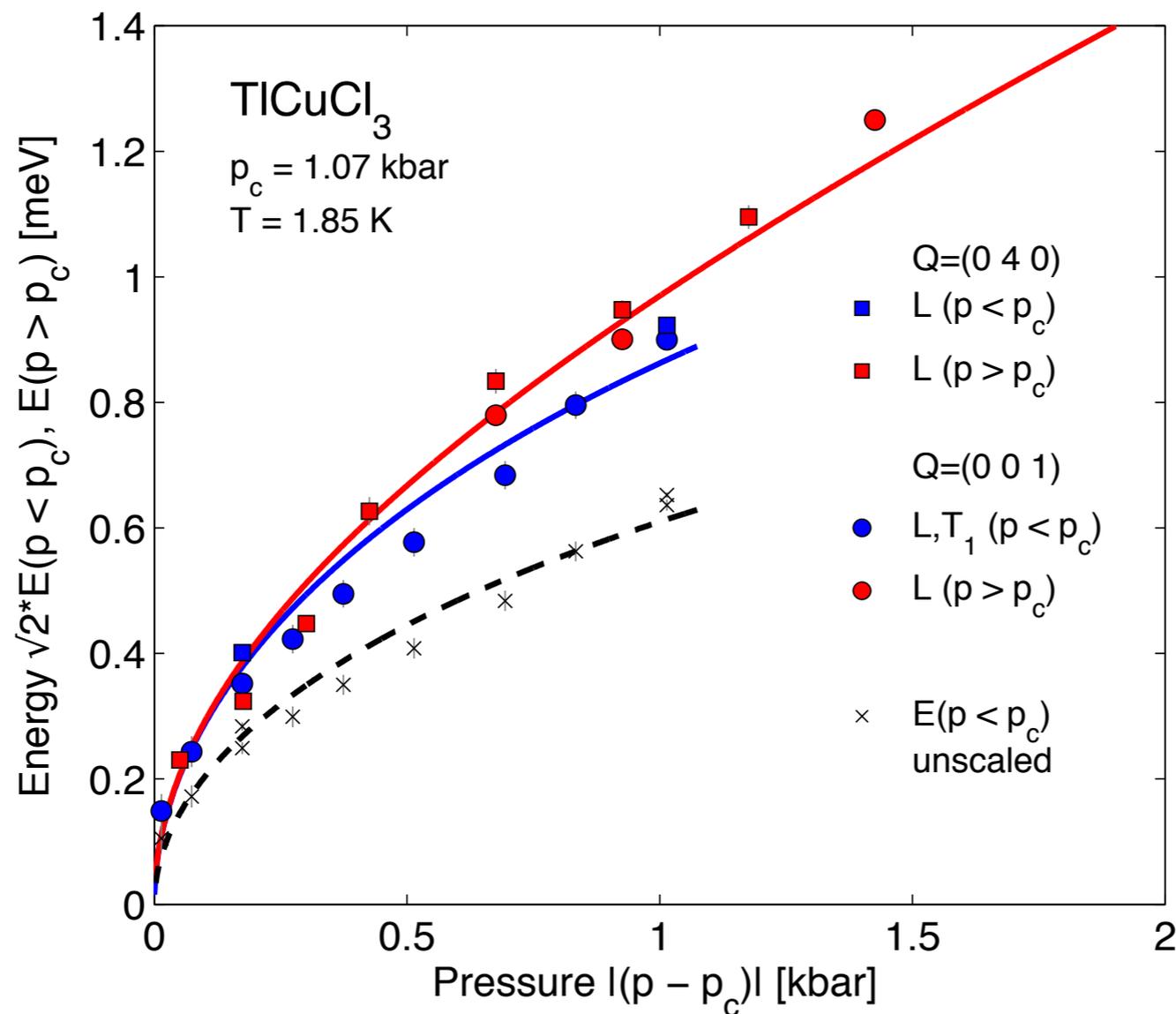
Observation of 3 → 2 low energy modes, emergence of new Higgs particle in the Néel phase, and vanishing of Néel temperature at the quantum critical point

Christian Ruegg, Bruce Normand, Masashige Matsumoto, Albert Furrer, Desmond McMorrow, Karl Kramer, Hans-Ulrich Gudel, Severian Gvasaliya, Hannu Mutka, and Martin Boehm, *Phys. Rev. Lett.* **100**, 205701 (2008)

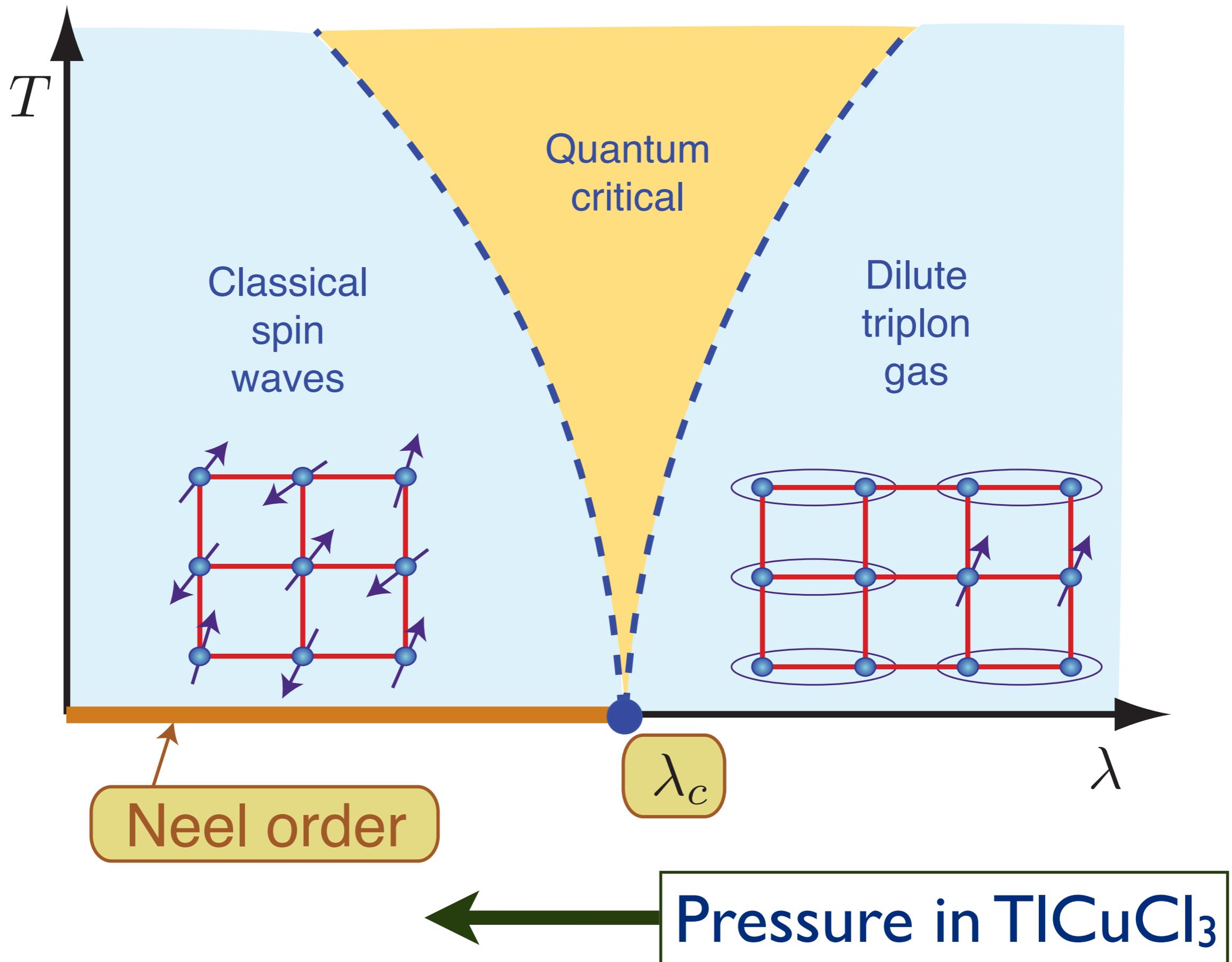
# Prediction of quantum field theory

$$\frac{\text{Energy of Higgs particle}}{\text{Energy of triplon}} = \sqrt{2}$$

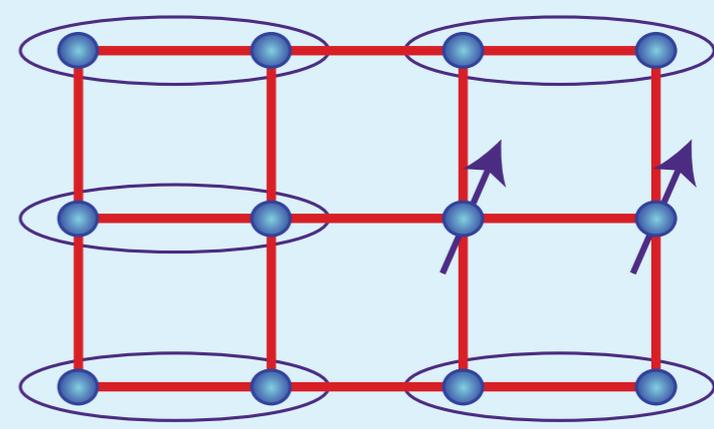
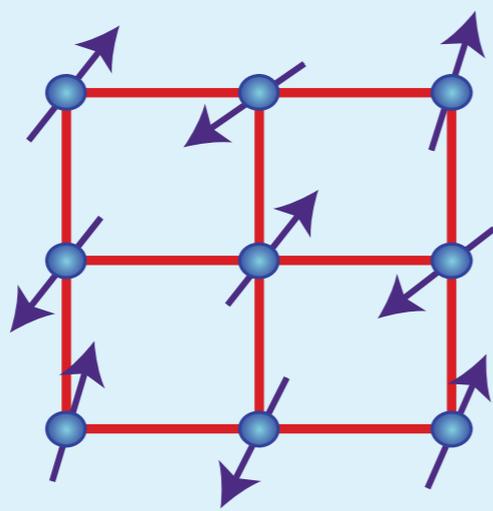
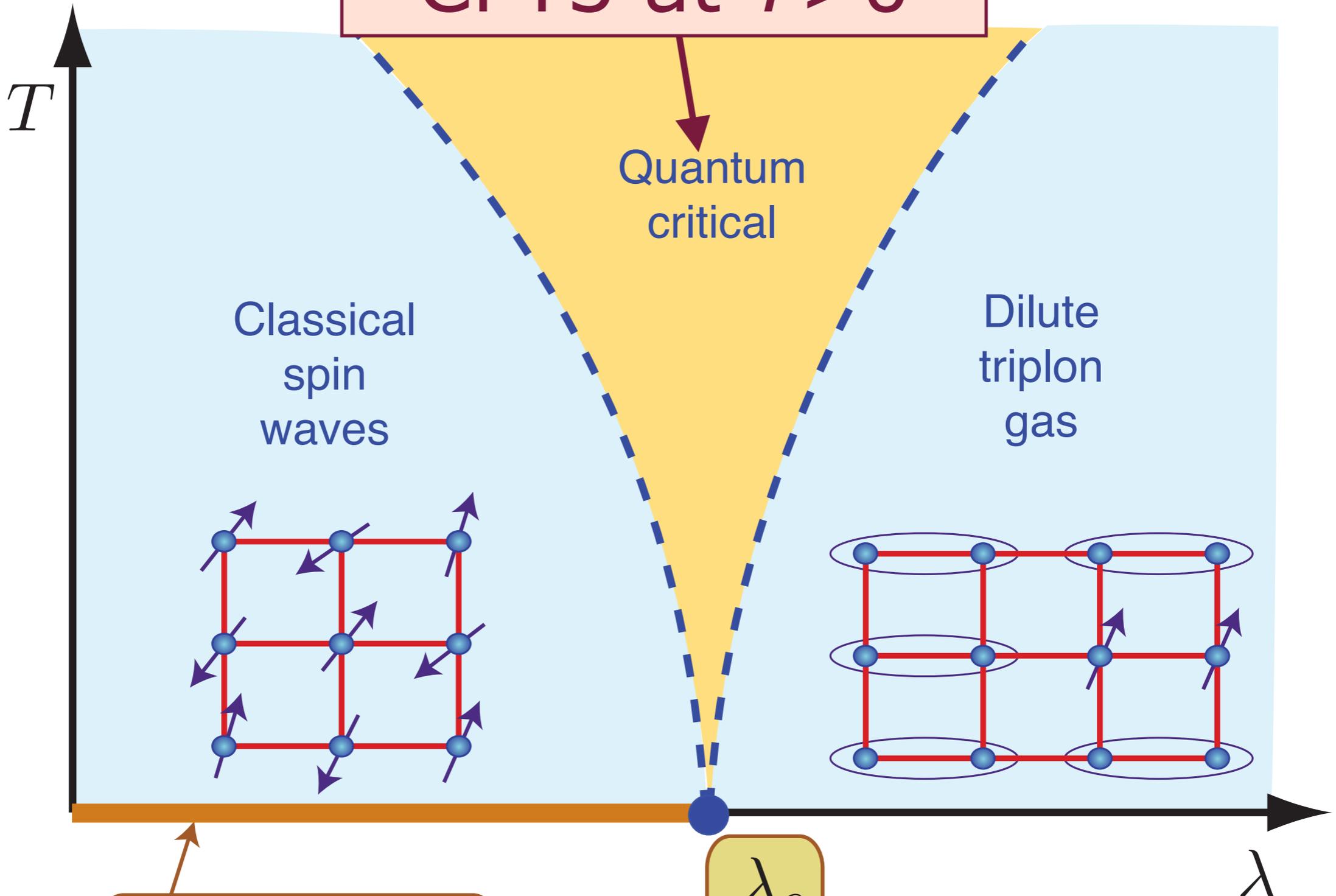
$$V(\vec{\varphi}) = (\lambda - \lambda_c)\vec{\varphi}^2 + u(\vec{\varphi}^2)^2$$



S. Sachdev and J. Ye, *Phys. Rev. Lett.* **69**, 2411 (1992).  
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# CFT3 at $T > 0$



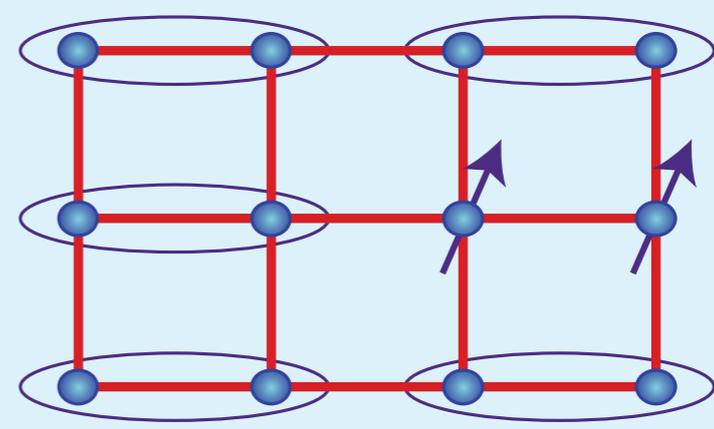
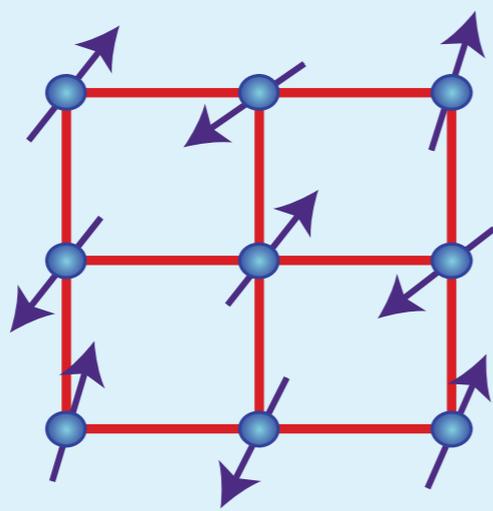
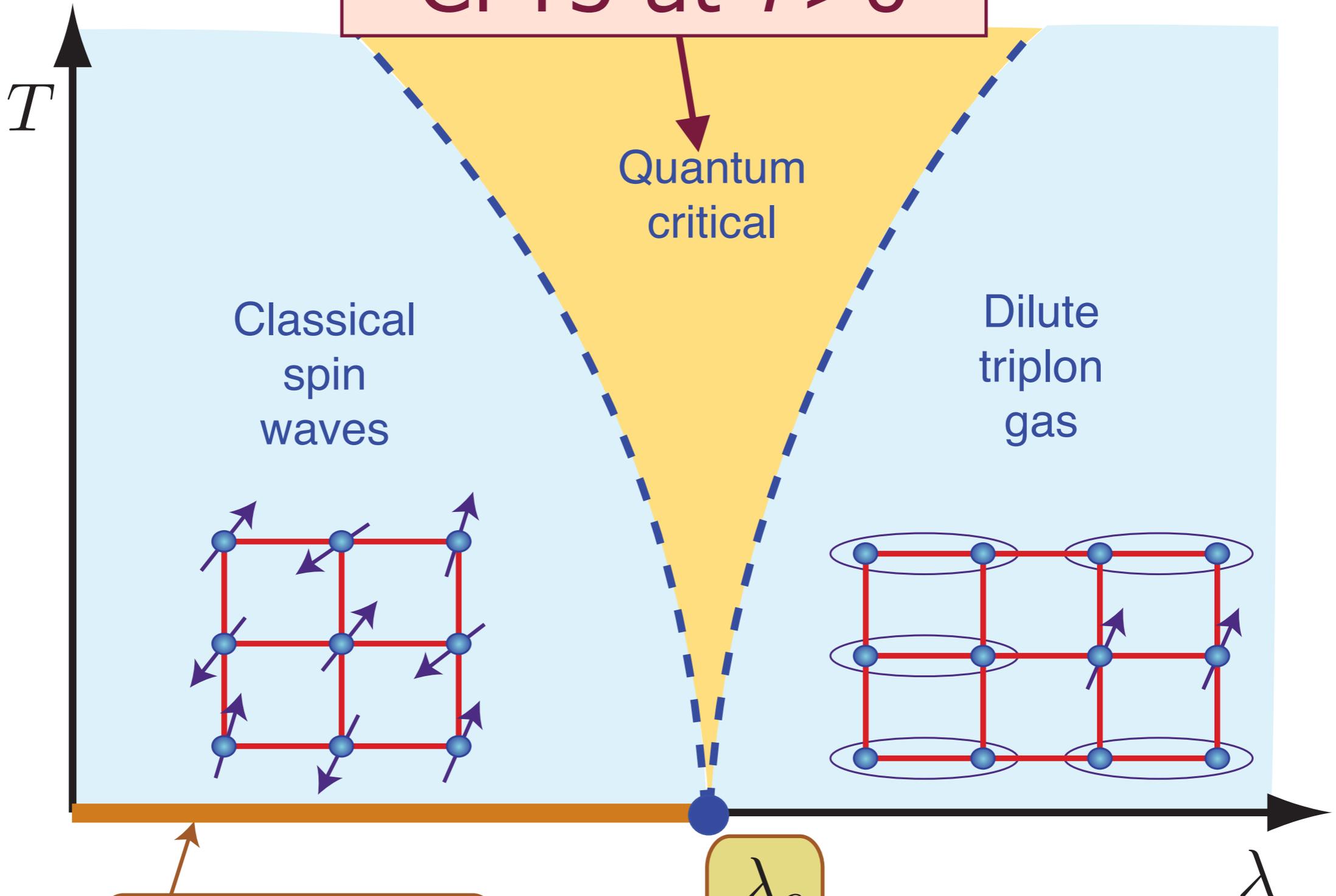
Neel order

$\lambda_c$

Pressure in  $\text{TlCuCl}_3$



# CFT3 at $T > 0$



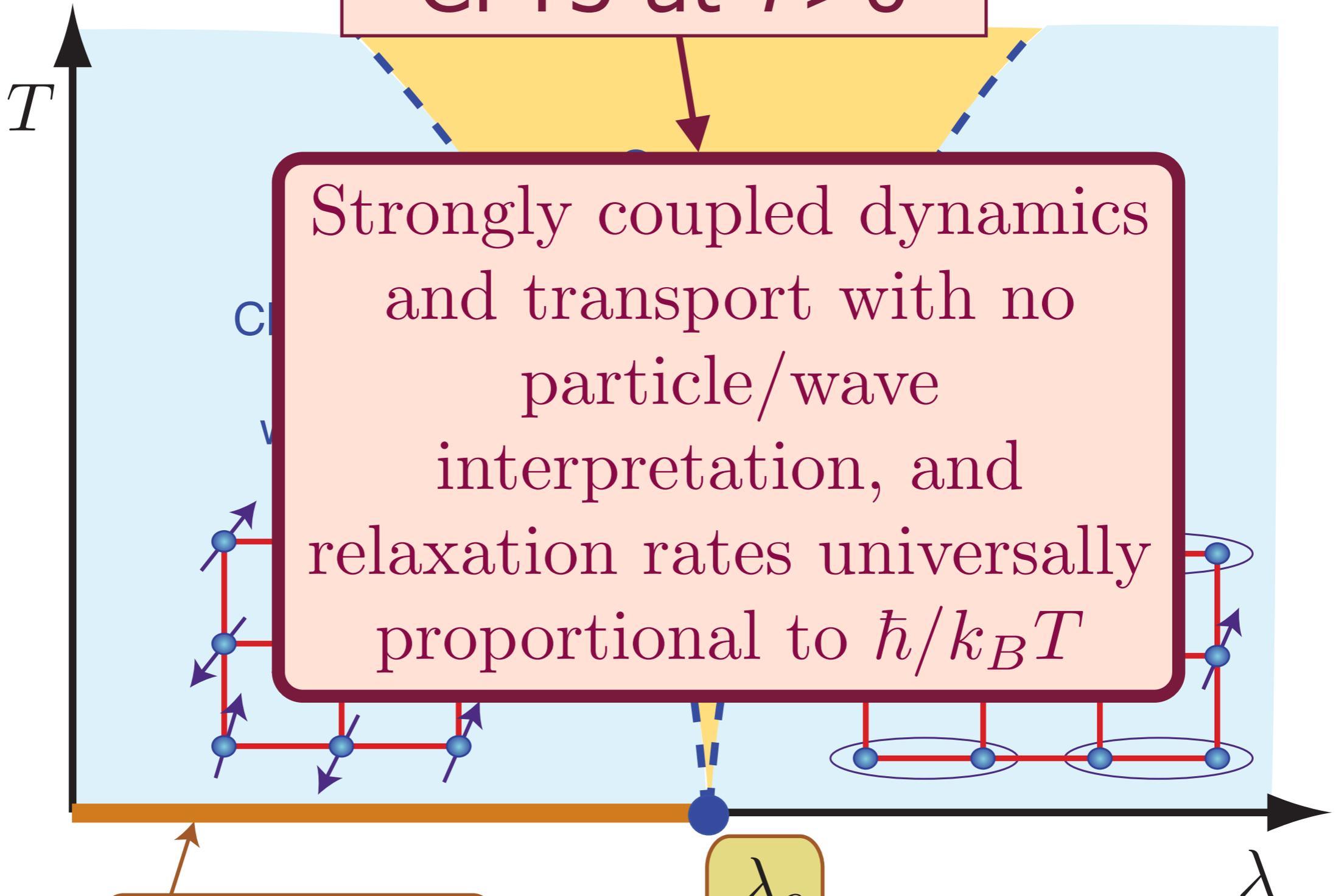
Neel order

$\lambda_c$

Pressure in  $\text{TlCuCl}_3$



CFT3 at  $T > 0$



Strongly coupled dynamics and transport with no particle/wave interpretation, and relaxation rates universally proportional to  $\hbar/k_B T$

Neel order

$\lambda_c$

Pressure in  $TlCuCl_3$



# Outline

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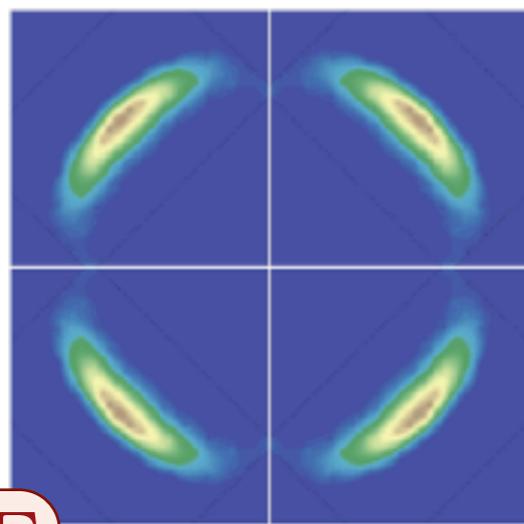
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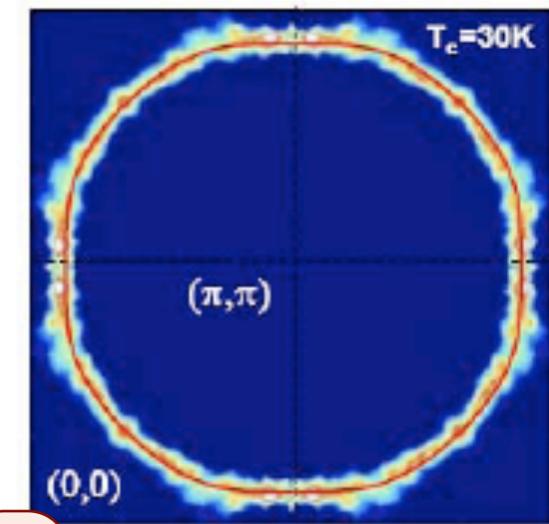
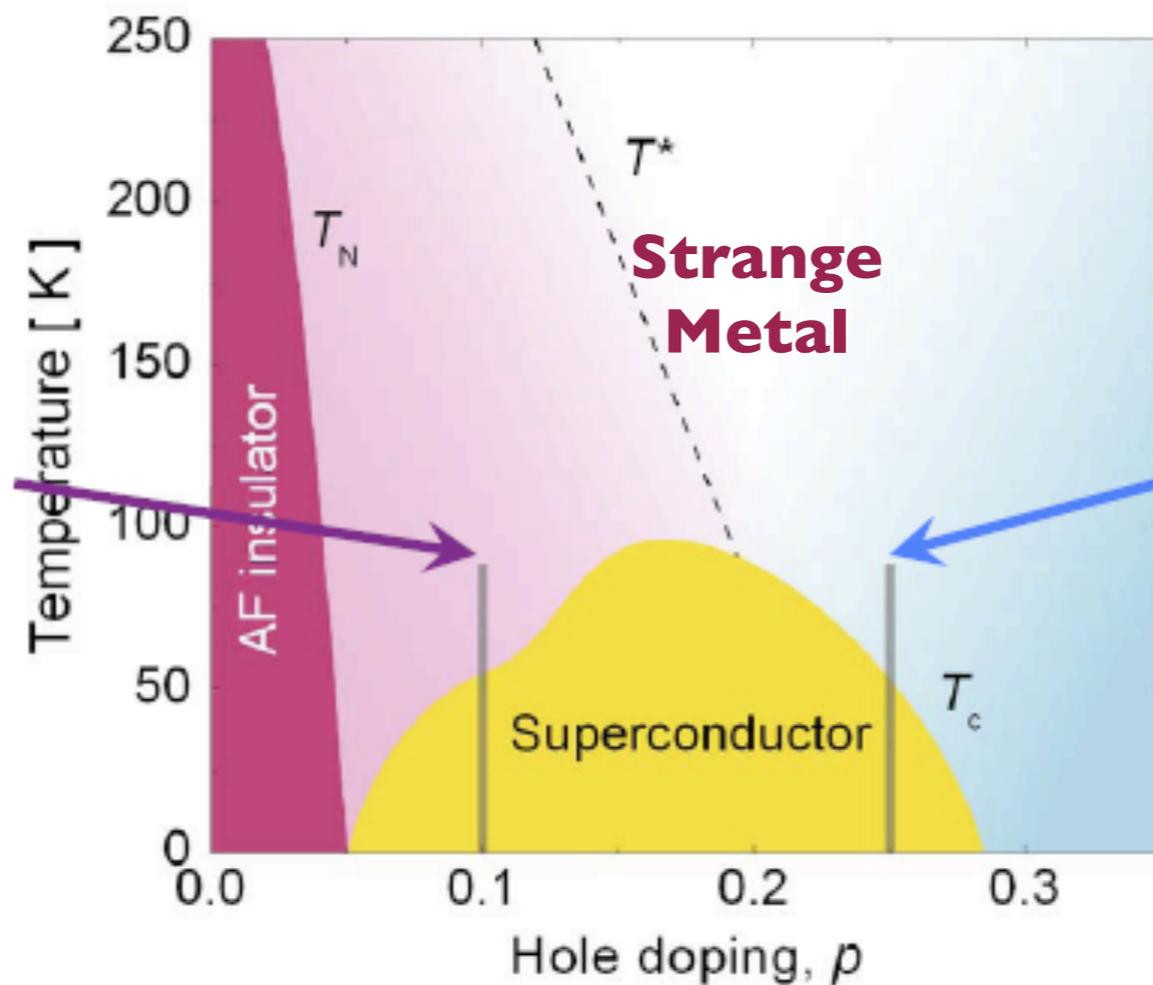
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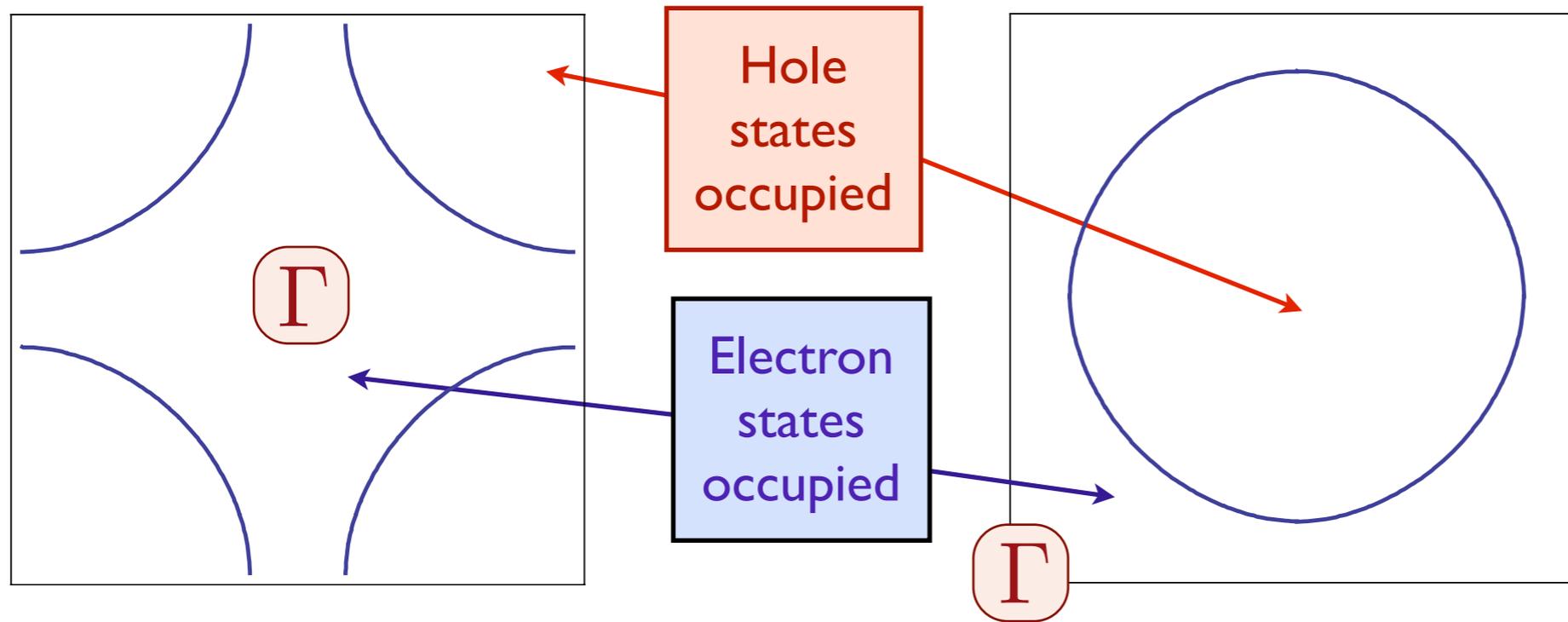
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Smaller hole  
Fermi-pockets

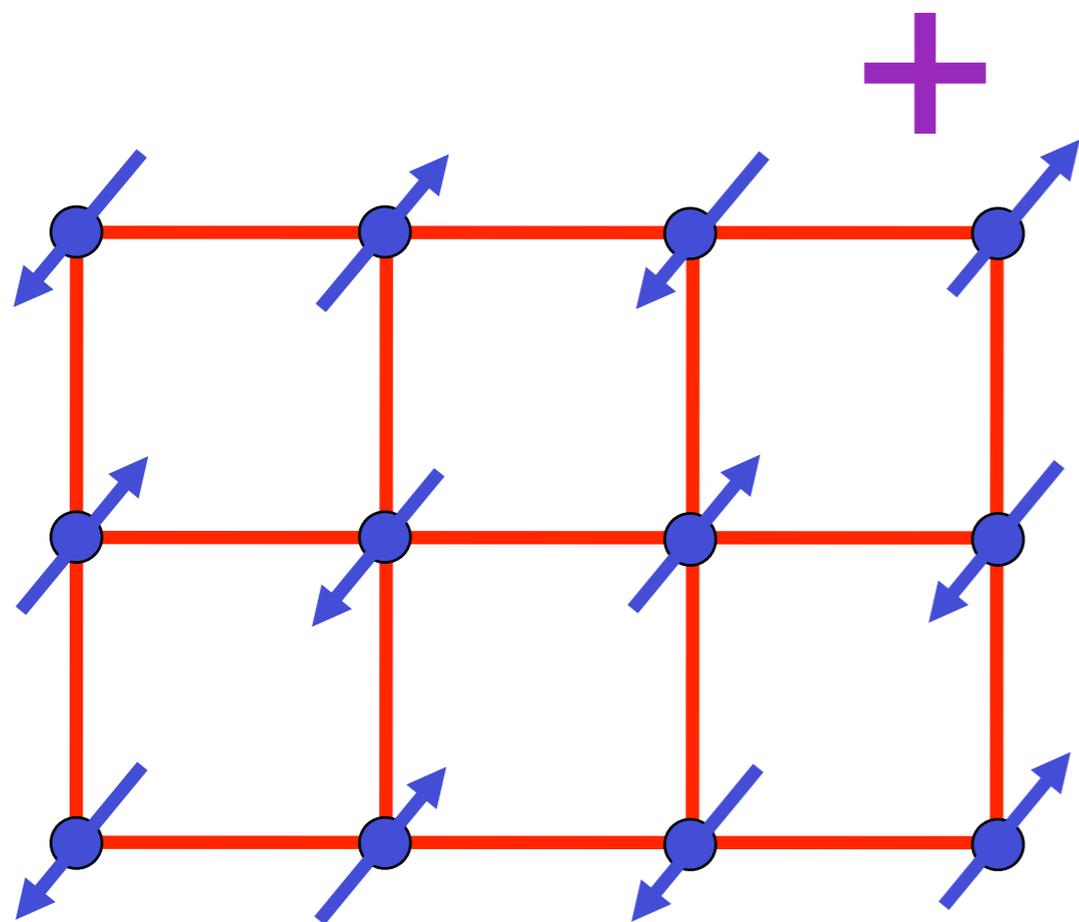
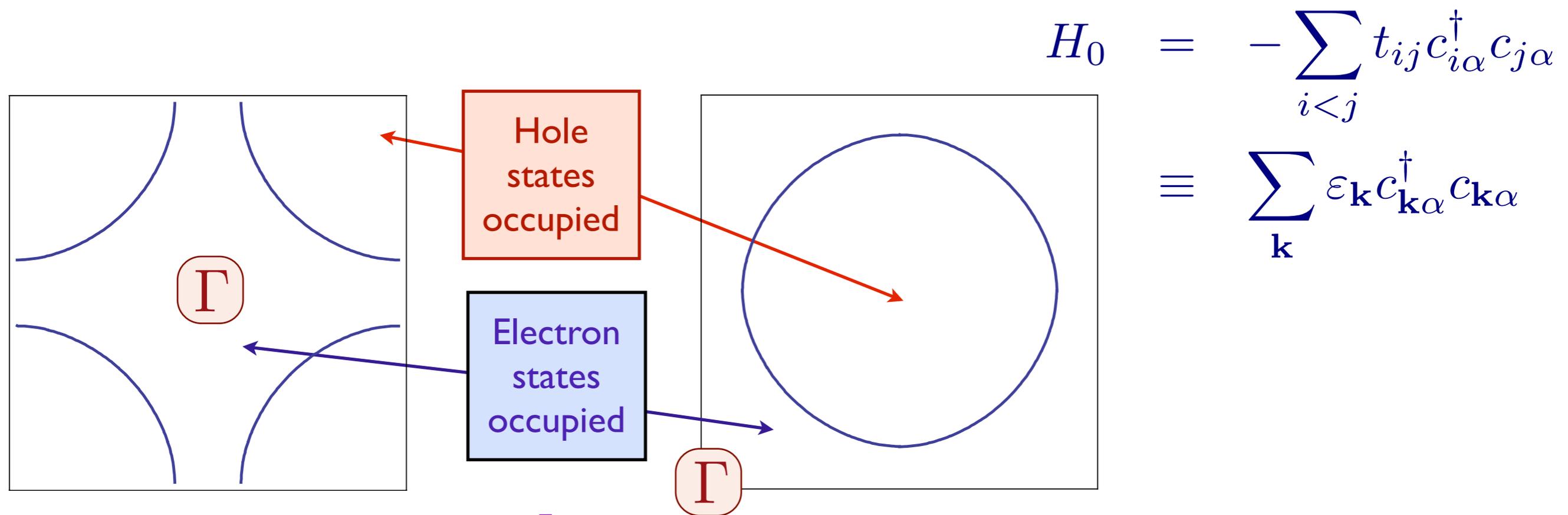
Large hole  
Fermi surface

# Fermi surface+antiferromagnetism



$$H_0 = - \sum_{i < j} t_{ij} c_{i\alpha}^\dagger c_{j\alpha}$$
$$\equiv \sum_{\mathbf{k}} \varepsilon_{\mathbf{k}} c_{\mathbf{k}\alpha}^\dagger c_{\mathbf{k}\alpha}$$

# Fermi surface+antiferromagnetism

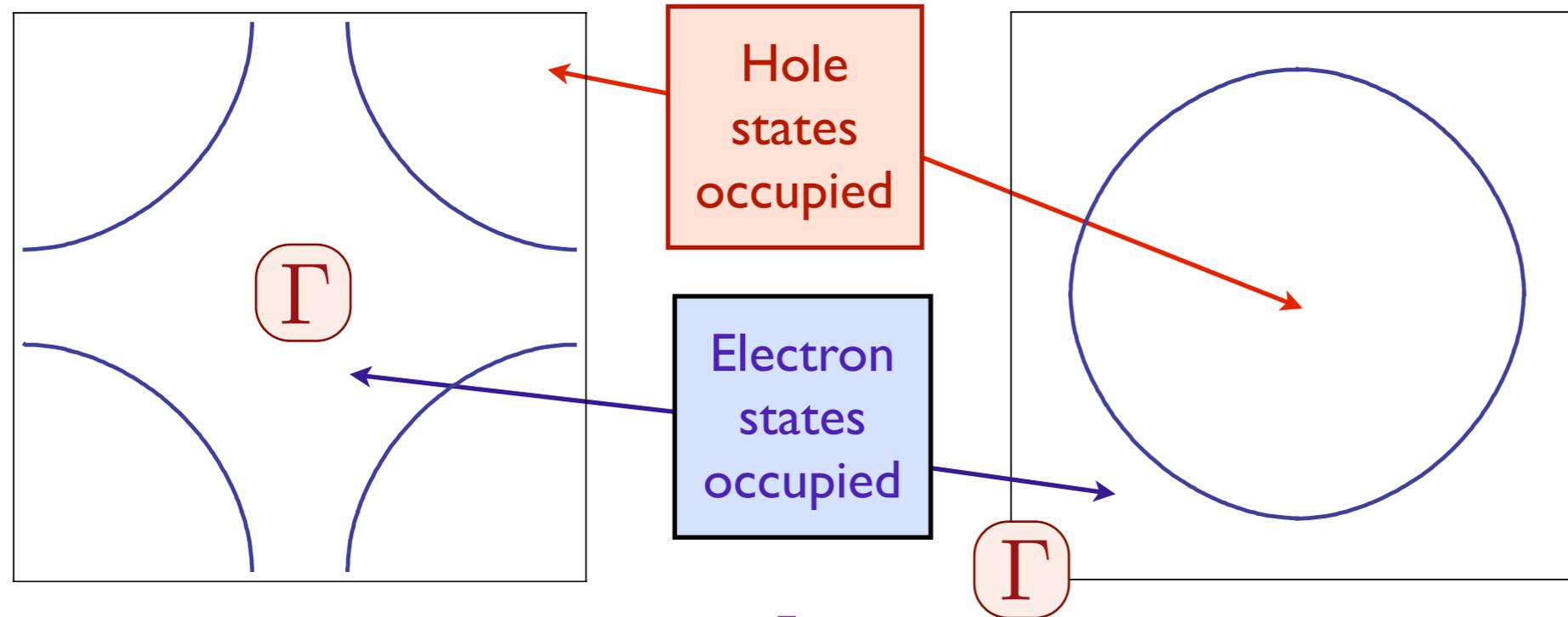


The electron spin polarization obeys

$$\langle \vec{S}(\mathbf{r}, \tau) \rangle = \vec{\varphi}(\mathbf{r}, \tau) e^{i\mathbf{K} \cdot \mathbf{r}}$$

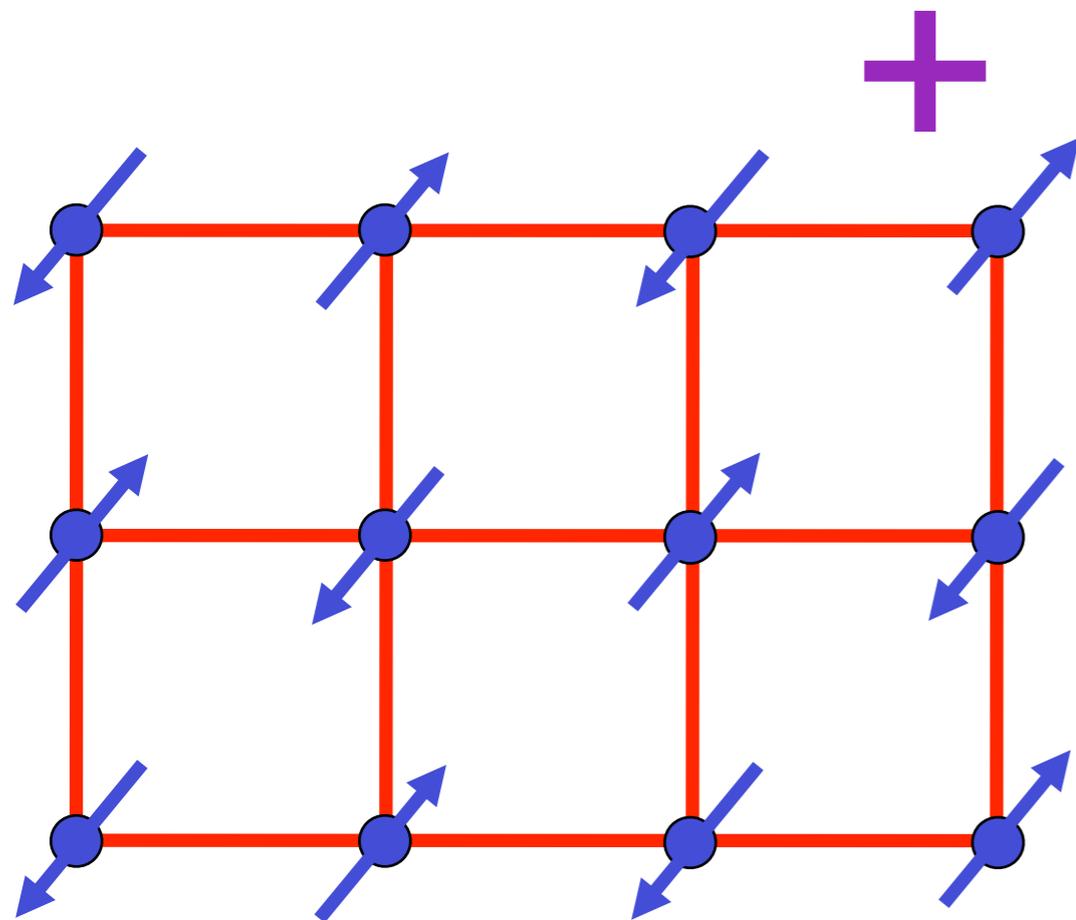
where  $\mathbf{K}$  is the ordering wavevector.

# Fermi surface+antiferromagnetism



$$H_0 = - \sum_{i < j} t_{ij} c_{i\alpha}^\dagger c_{j\alpha}$$

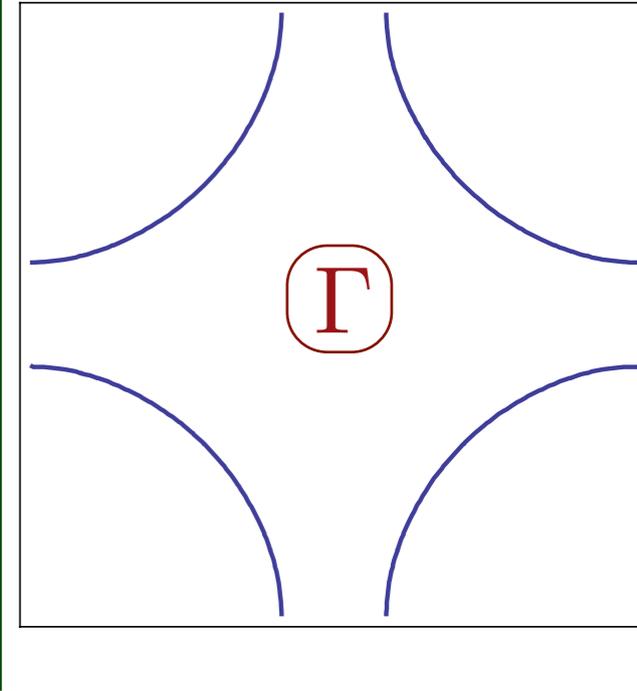
$$\equiv \sum_{\mathbf{k}} \varepsilon_{\mathbf{k}} c_{\mathbf{k}\alpha}^\dagger c_{\mathbf{k}\alpha}$$



$$H_{\text{sdw}} = \sum_{\mathbf{k}, \mathbf{q}, \alpha, \beta} \vec{\varphi}_{\mathbf{q}} \cdot c_{\mathbf{k}+\mathbf{q}, \alpha}^\dagger \vec{\sigma}_{\alpha\beta} c_{\mathbf{k}+\mathbf{K}, \beta}$$

# Hole-doped cuprates

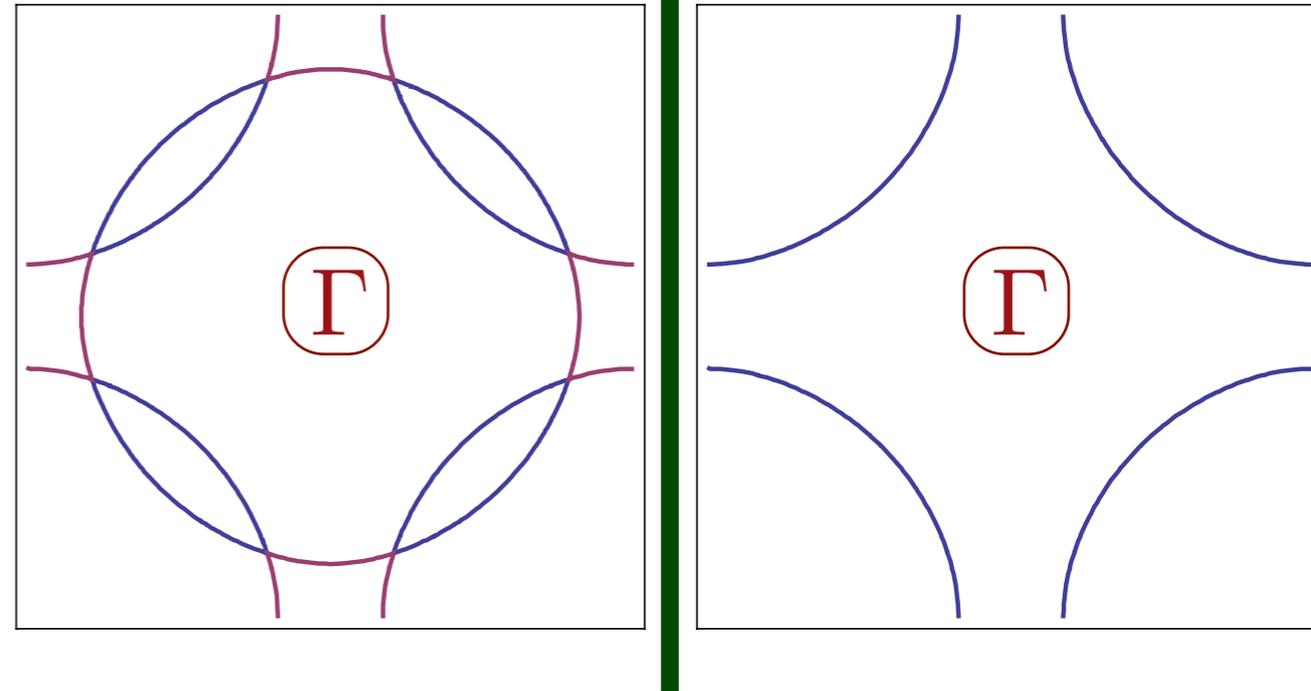
← Increasing SDW order →



S. Sachdev, A. V. Chubukov, and A. Sokol, *Phys. Rev. B* **51**, 14874 (1995).  
A. V. Chubukov and D. K. Morr, *Physics Reports* **288**, 355 (1997).

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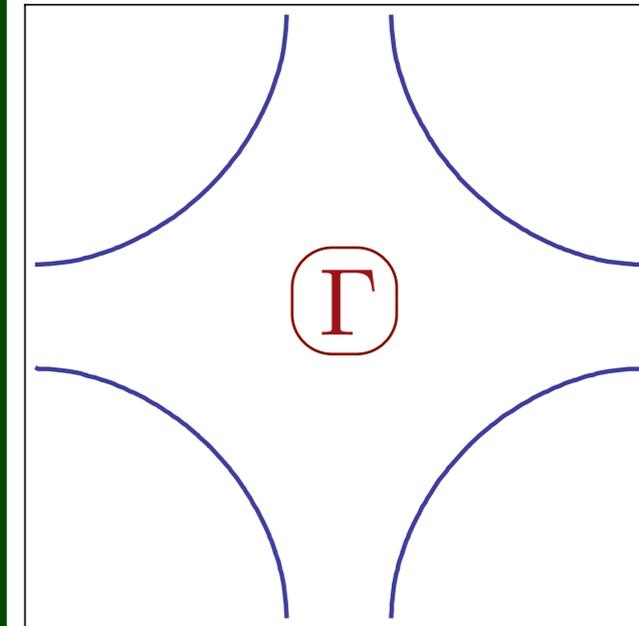
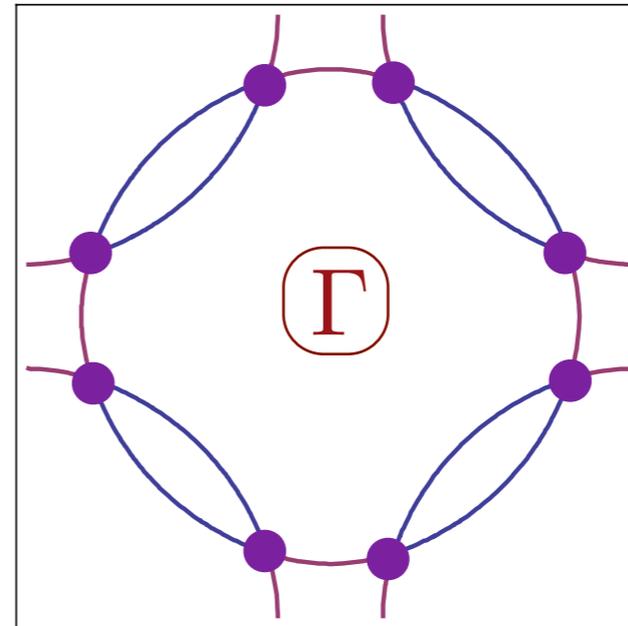


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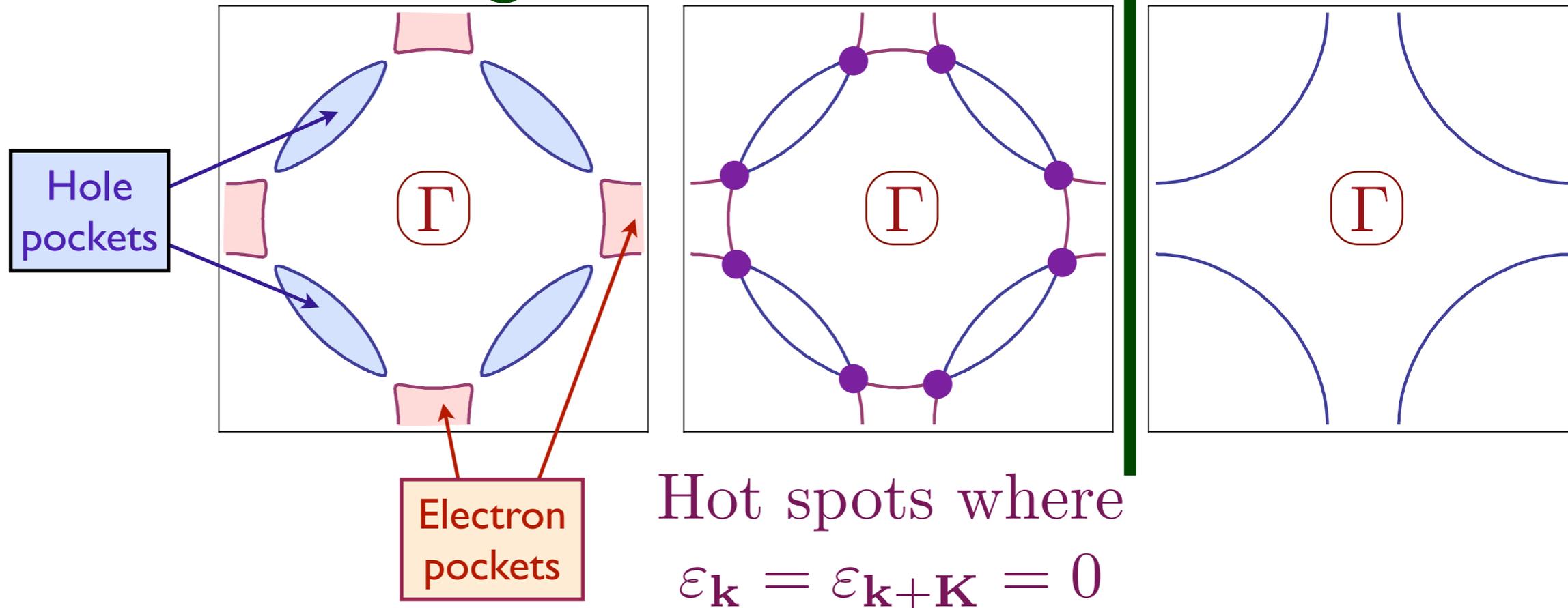


Hot spots where  
 $\varepsilon_{\mathbf{k}} = \varepsilon_{\mathbf{k}+\mathbf{K}} = 0$



# Hole-doped cuprates

← Increasing SDW order →

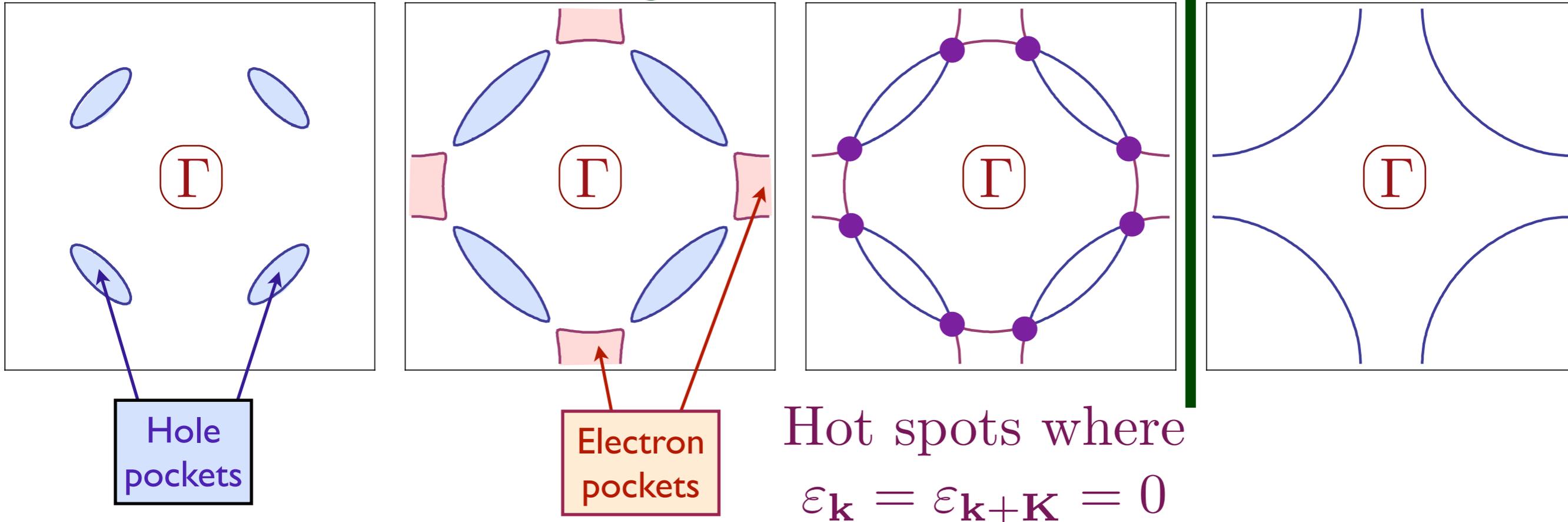


Fermi surface breaks up at hot spots  
into electron and hole “pockets”

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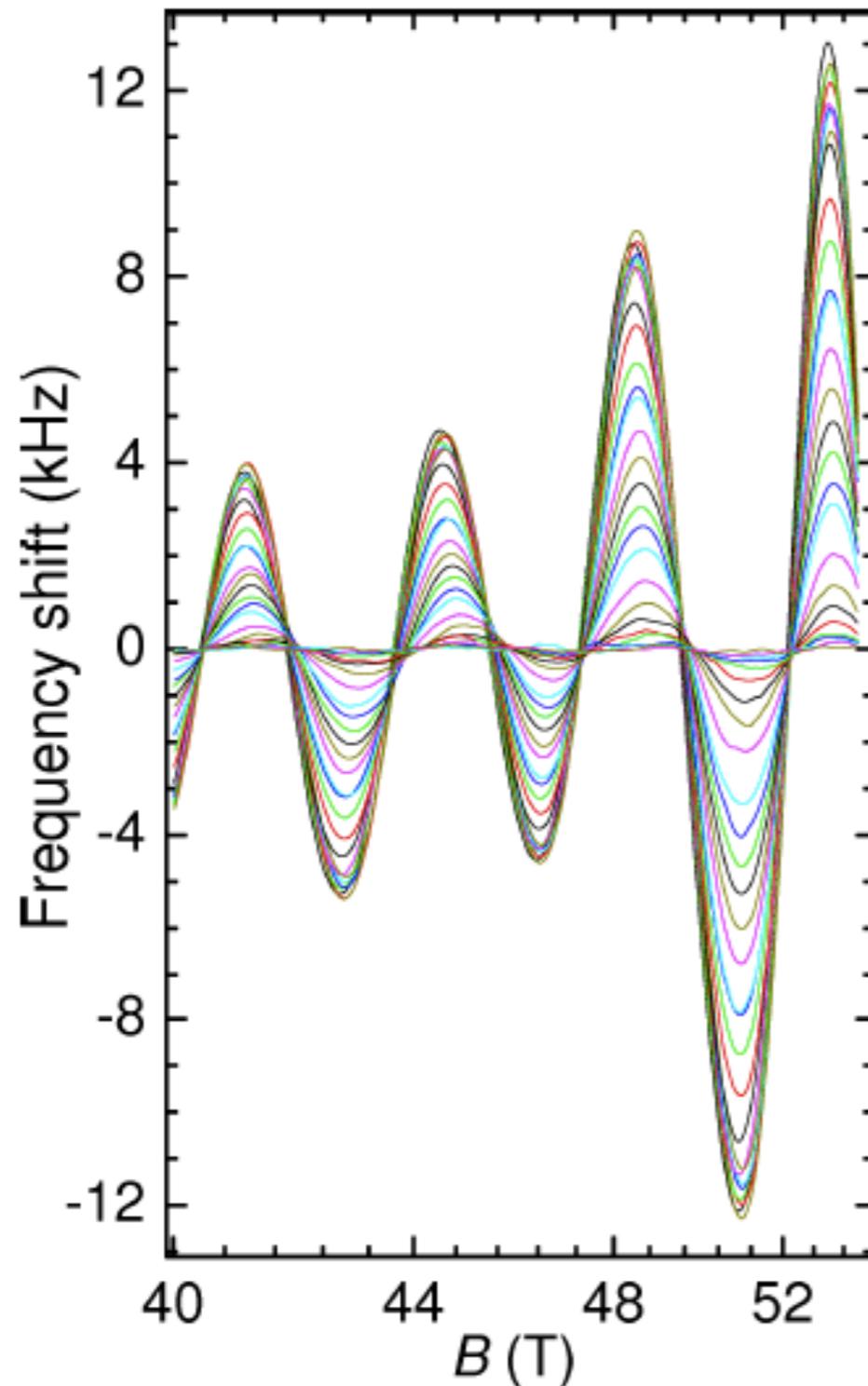
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A.V. Chubukov and D. K. Morr, *Physics Reports* **288**, 355 (1997).

# Evidence for small Fermi pockets

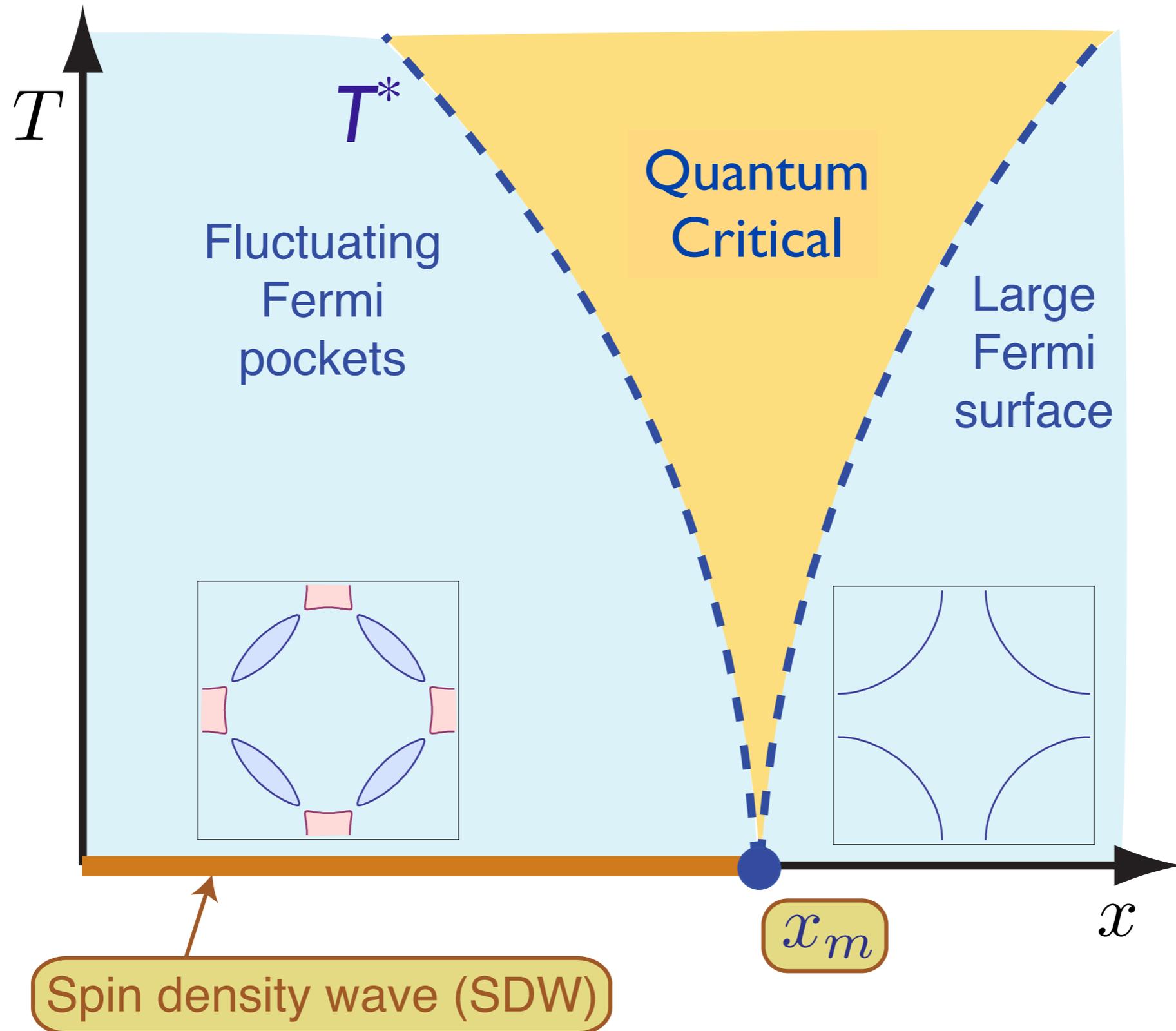


Suchitra E. Sebastian, N. Harrison,  
M. M. Altarawneh, Ruixing Liang, D.A. Bonn,  
W. N. Hardy, and G. G. Lonzarich  
*Physical Review B* **81**, 140505(R) (2010)

Original observation:  
N. Doiron-Leyraud, C. Proust,  
D. LeBoeuf, J. Levallois,  
J.-B. Bonnemaïson, R. Liang,  
D.A. Bonn, W. N. Hardy,  
and L. Taillefer,  
*Nature* **447**, 565 (2007)

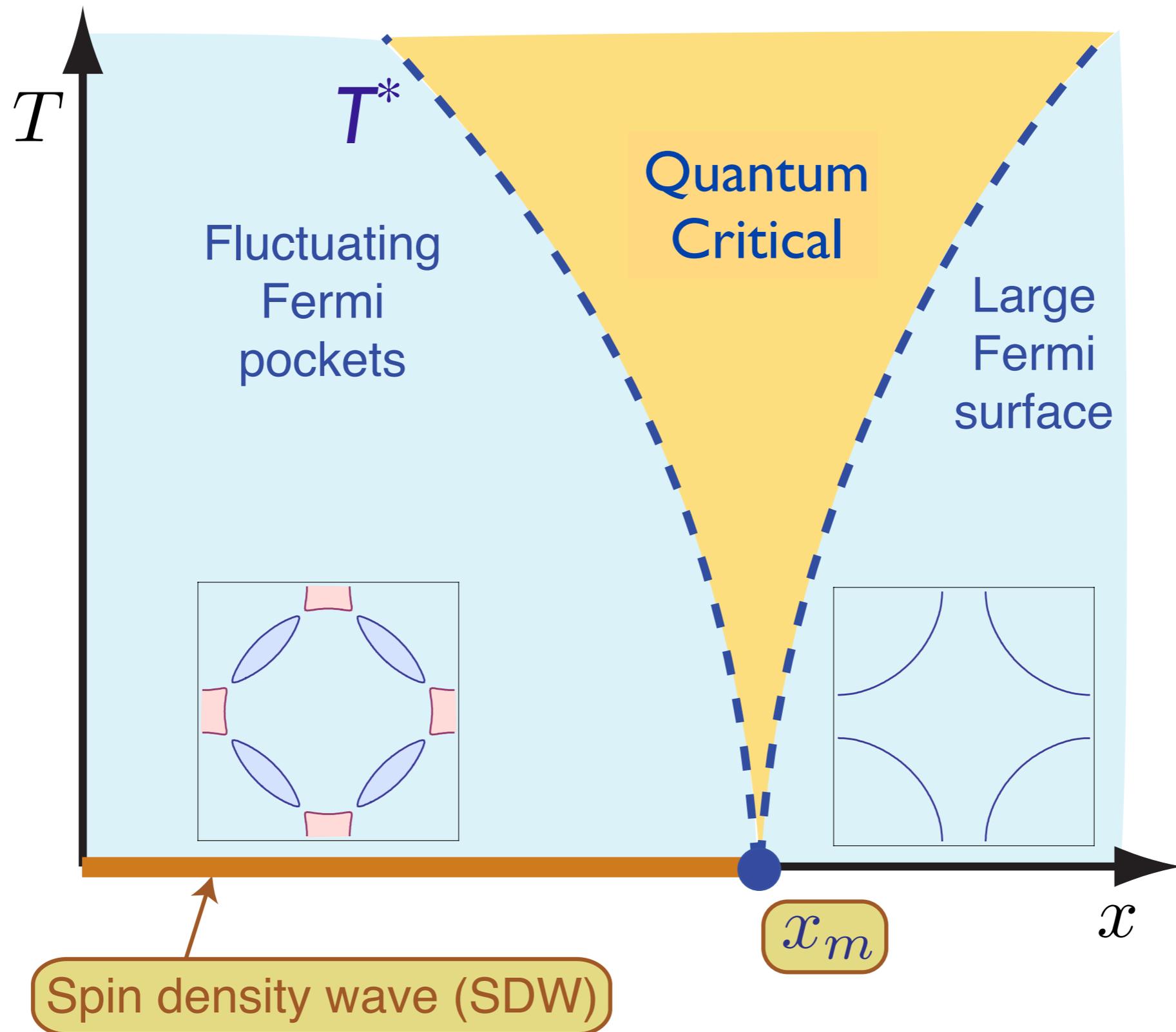
FIG. 2: Magnetic quantum oscillations measured in  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  with  $x \approx 0.56$  (after background polynomial subtraction). This restricted interval in  $B = |\mathbf{B}|$  furnishes a dynamic range of  $\sim 50$  dB between  $T = 1$  and 18 K. The actual  $T$  values are provided in Fig. 3.

# Theory of quantum criticality in the cuprates



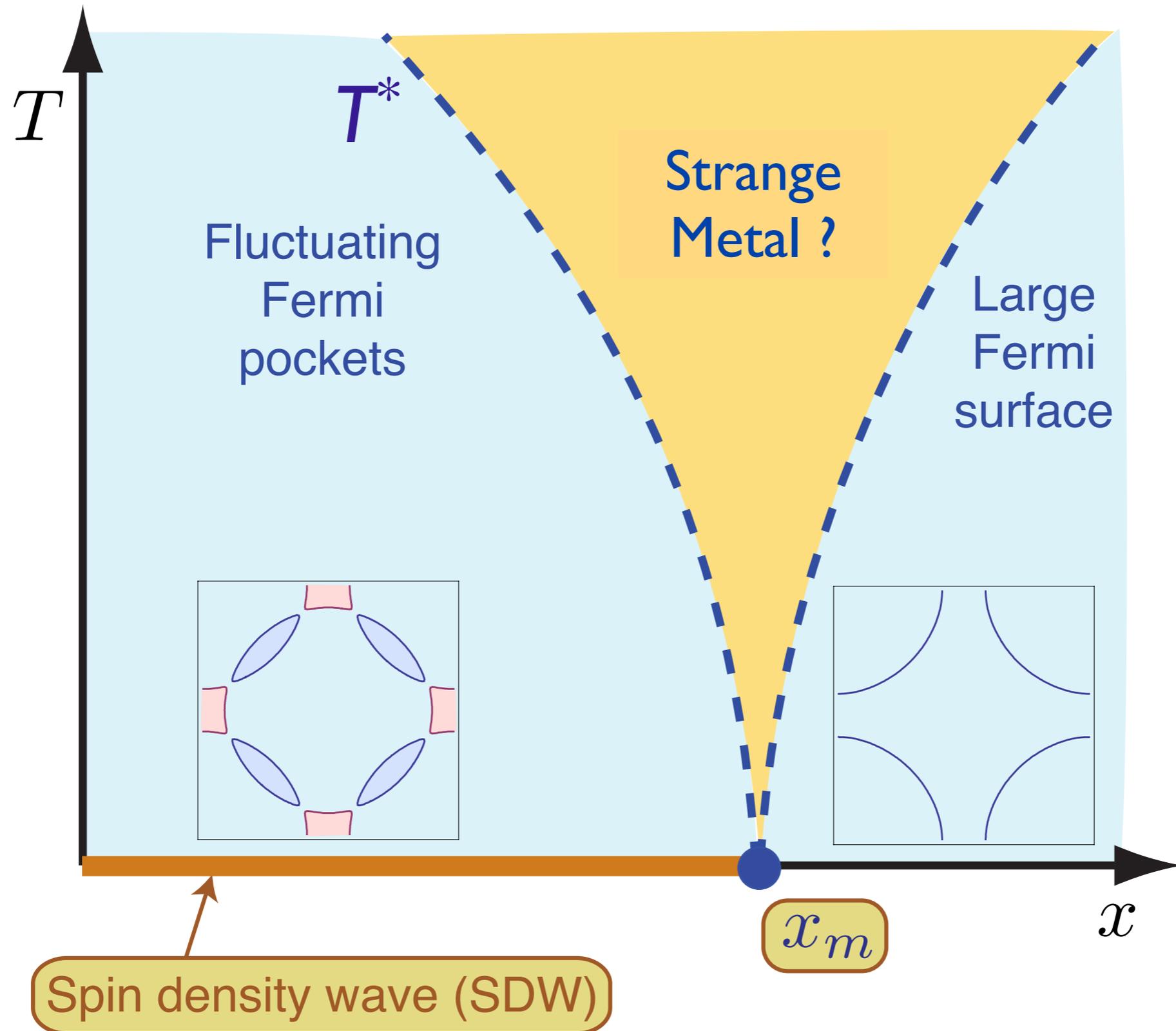
Underlying SDW ordering quantum critical point  
in metal at  $x = x_m$

# Theory of quantum criticality in the cuprates



Underlying SDW ordering quantum critical point  
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## ***d*-wave pairing near a spin-density-wave instability**

D. J. Scalapino, E. Loh, Jr.,\* and J. E. Hirsch†

*Institute for Theoretical Physics, University of California, Santa Barbara, California 93106*

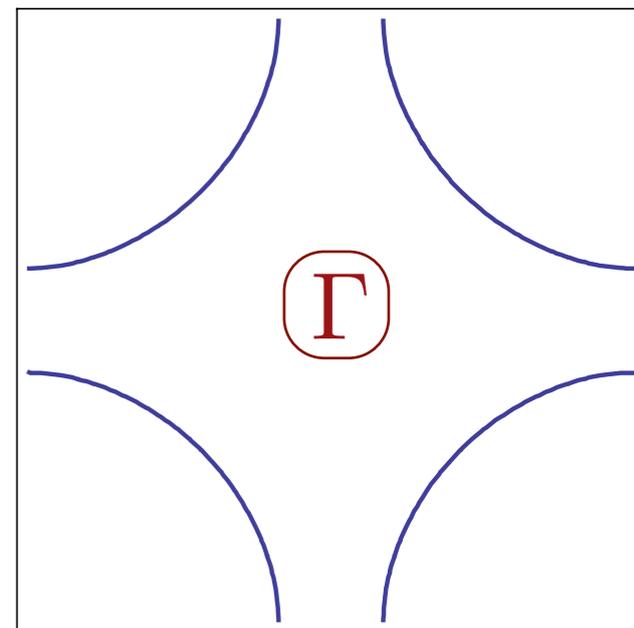
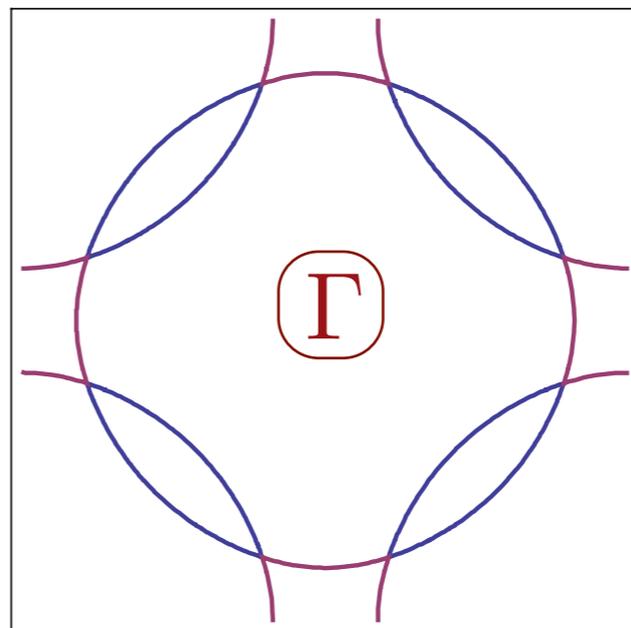
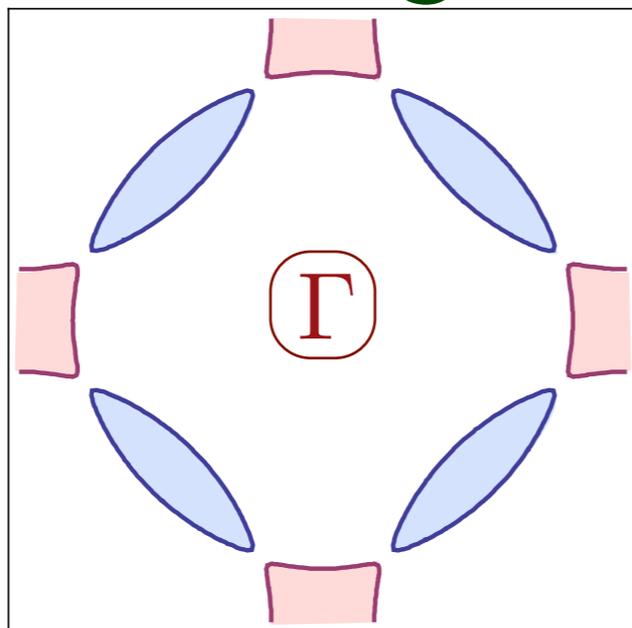
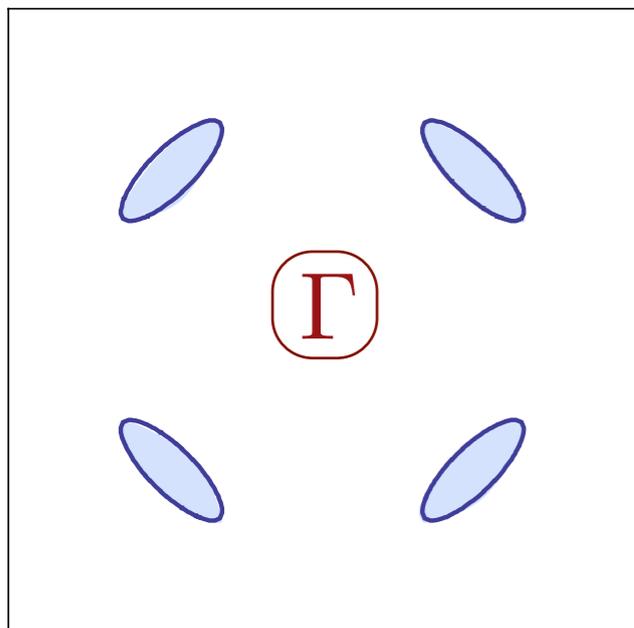
(Received 23 June 1986)

We investigate the three-dimensional Hubbard model and show that paramagnon exchange near a spin-density-wave instability gives rise to a strong singlet *d*-wave pairing interaction. For a cubic band the singlet ( $d_{x^2-y^2}$  and  $d_{3z^2-r^2}$ ) channels are enhanced while the singlet ( $d_{xy}, d_{xz}, d_{yz}$ ) and triplet *p*-wave channels are suppressed. A unique feature of this pairing mechanism is its sensitivity to band structure and band filling.

Physical Review B **34**, 8190 (1986)

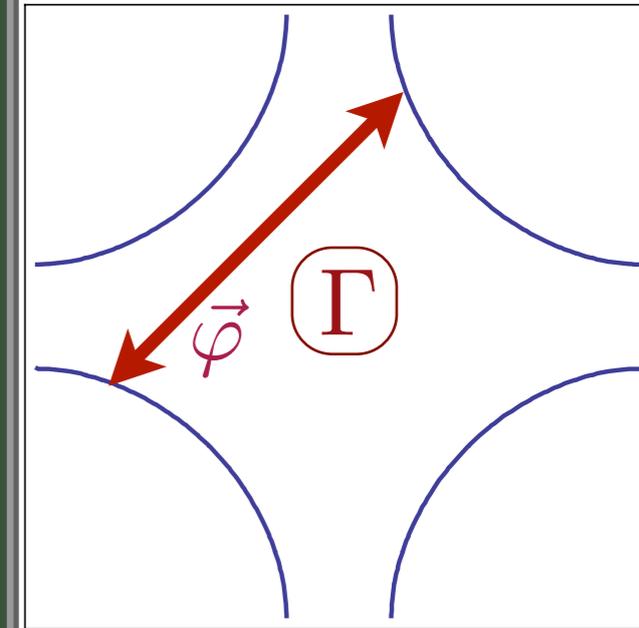
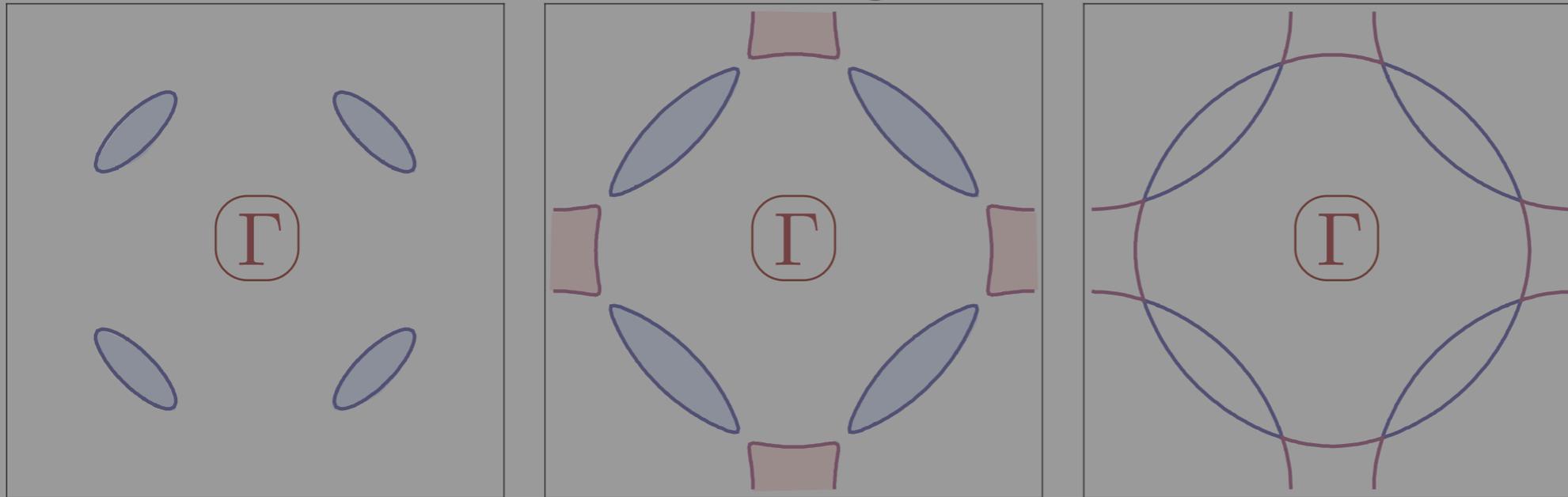
# Spin density wave theory in hole-doped cuprates

← Increasing SDW order →



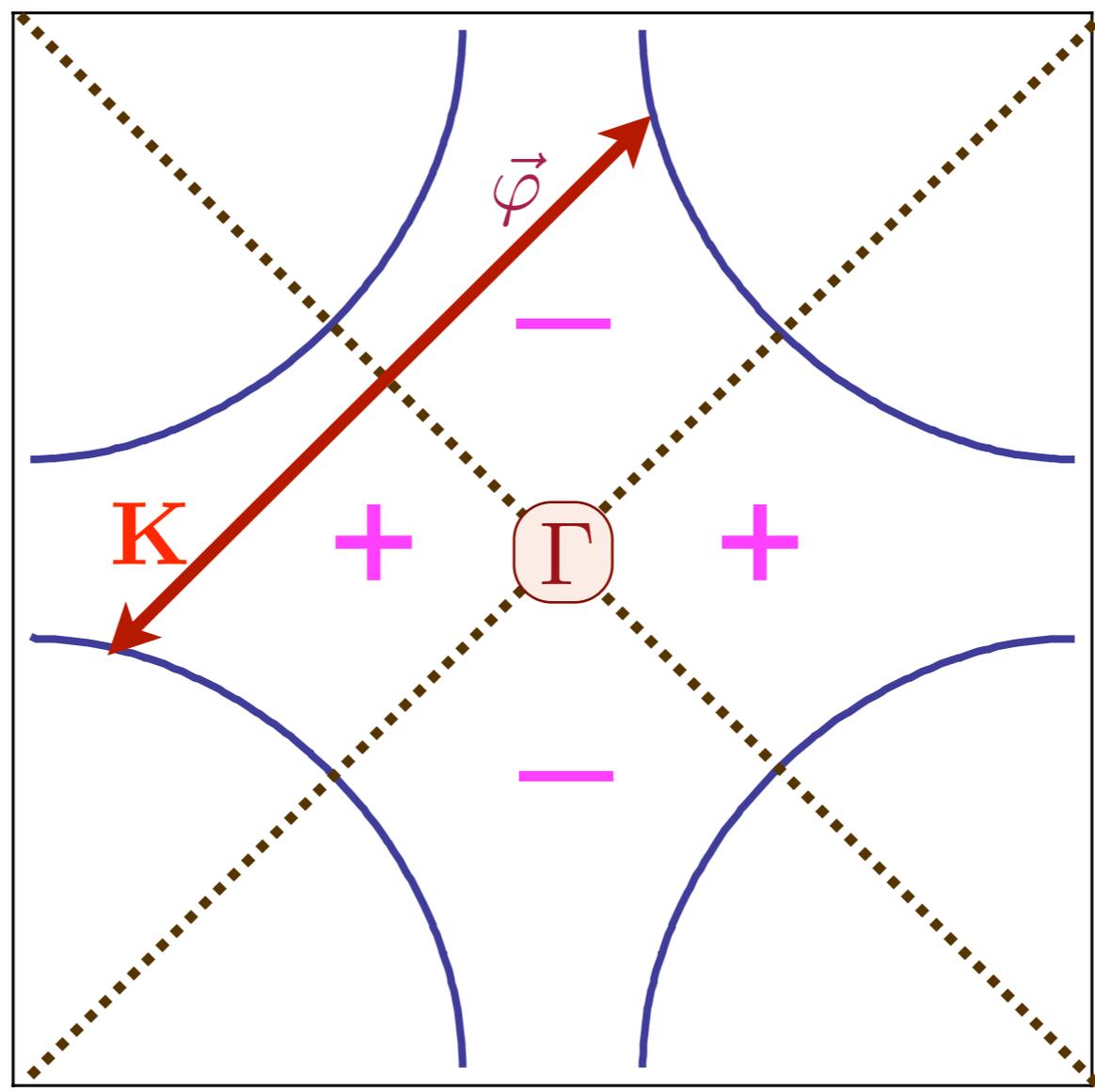
# Spin-fluctuation exchange theory of d-wave superconductivity in the cuprates

← Increasing SDW order →



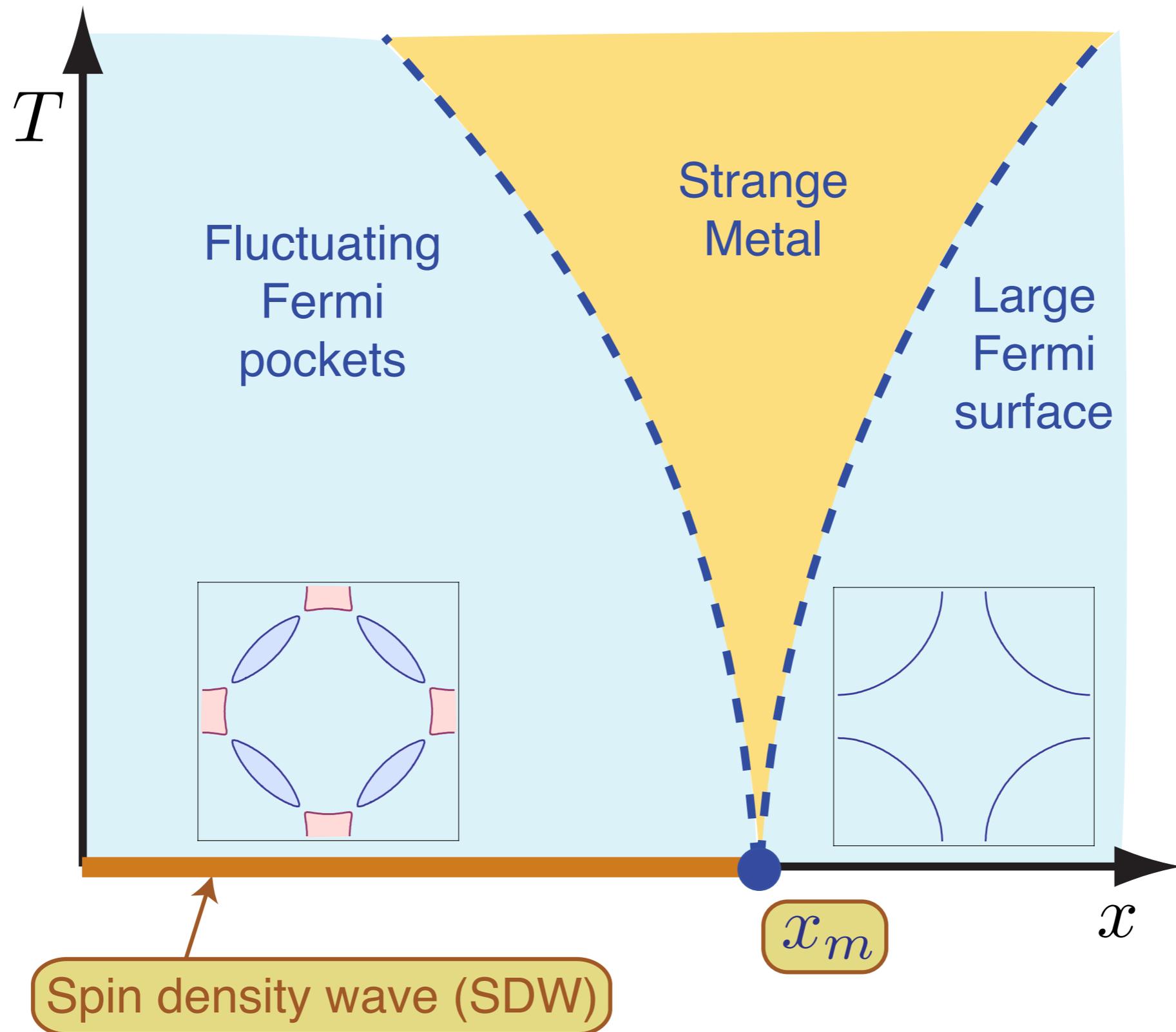
Fermions at the *large* Fermi surface exchange fluctuations of the SDW order parameter  $\vec{\varphi}$ .

# $d$ -wave pairing of the large Fermi surface



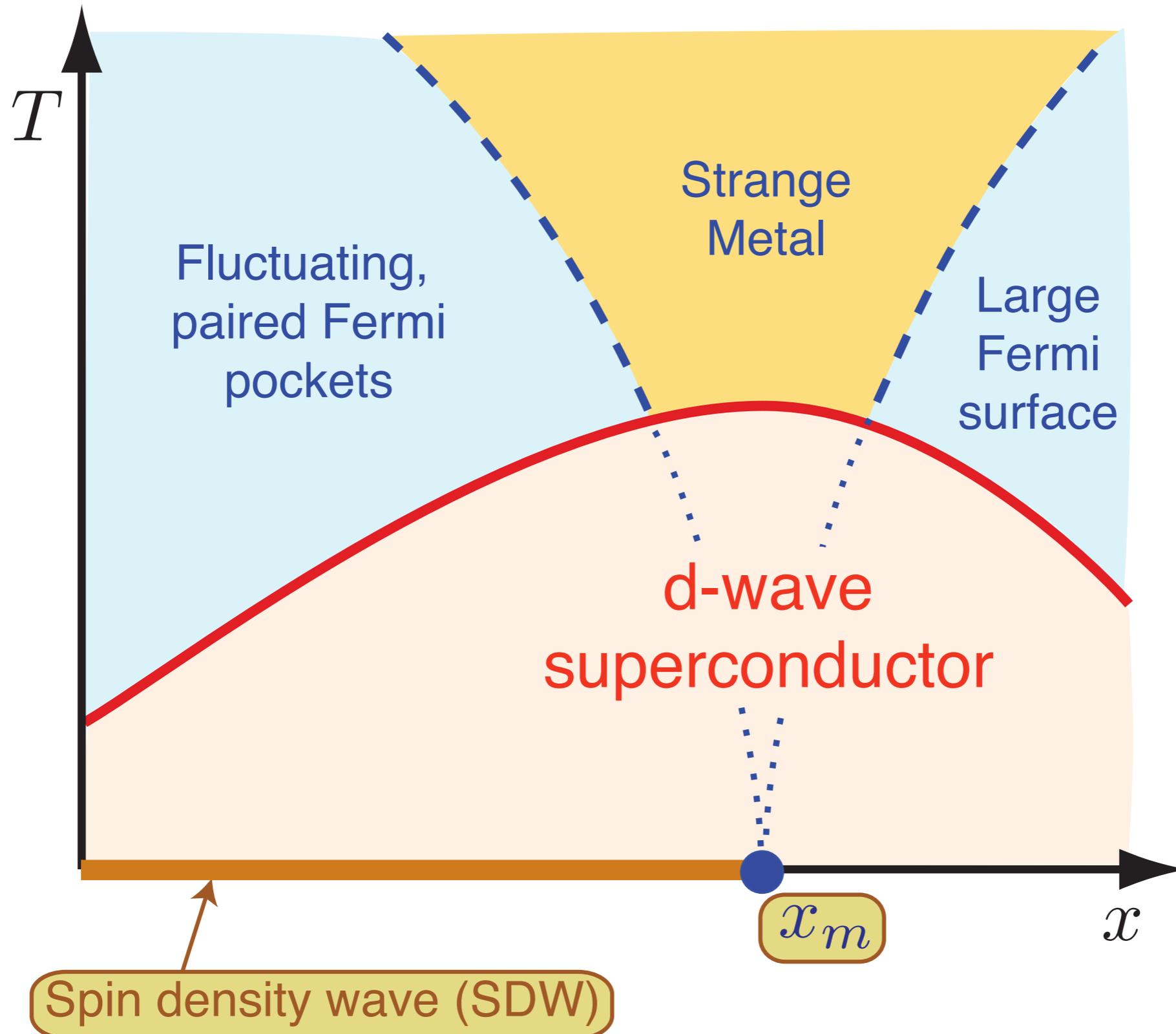
$$\langle c_{\mathbf{k}\uparrow} c_{-\mathbf{k}\downarrow} \rangle \propto \Delta_{\mathbf{k}} = \Delta_0 (\cos(k_x) - \cos(k_y))$$

# Theory of quantum criticality in the cuprates



Underlying SDW ordering quantum critical point  
in metal at  $x = x_m$

# Theory of quantum criticality in the cuprates



Ar. Abanov,  
A.V. Chubukov, and  
J. Schmalian,  
*Advances in Physics*  
**52**, 119 (2003).

Onset of  $d$ -wave superconductivity  
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# Phenomenological quantum theory of competing orders

Competition between superconductivity (SC) and spin-density wave (SDW) order

Begin with the Landau-Ginzburg field theory for quantum fluctuations of the antiferromagnetism ( $\vec{\varphi}$ ).

$$\mathcal{S} = \int d^2r d\tau \left[ \frac{1}{2} (\partial_\tau \vec{\varphi})^2 + \frac{c^2}{2} (\nabla_x \vec{\varphi})^2 + \frac{s}{2} \vec{\varphi}^2 + \frac{u}{4} (\vec{\varphi}^2)^2 \right]$$

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Competition between superconductivity (SC) and spin-density wave (SDW) order

Begin with the Landau-Ginzburg field theory for quantum fluctuations of the antiferromagnetism ( $\vec{\varphi}$ ). Include the Landau-Ginzburg mean-field action for superconductivity in an applied magnetic field  $H = \nabla \times \mathcal{A}$ :

$$\mathcal{S} = \int d^2r d\tau \left[ \frac{1}{2} (\partial_\tau \vec{\varphi})^2 + \frac{c^2}{2} (\nabla_x \vec{\varphi})^2 + \frac{s}{2} \vec{\varphi}^2 + \frac{u}{4} (\vec{\varphi}^2)^2 \right. \\ \left. + \int d^2r \left[ |(\nabla_x - i(2e/\hbar c)\mathcal{A})\Delta|^2 - |\Delta|^2 + \frac{|\Delta|^4}{2} \right] \right]$$

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Include the simplest allowed coupling between the two orders,  $\kappa > 0$ , with a positive sign implying repulsion or competition between them.

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Competition between superconductivity (SC) and spin-density wave (SDW) order

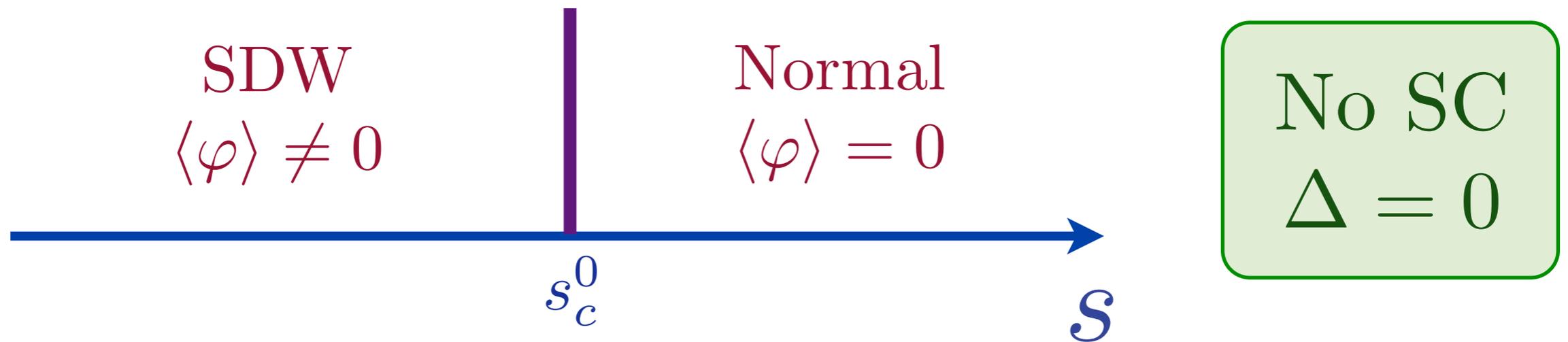
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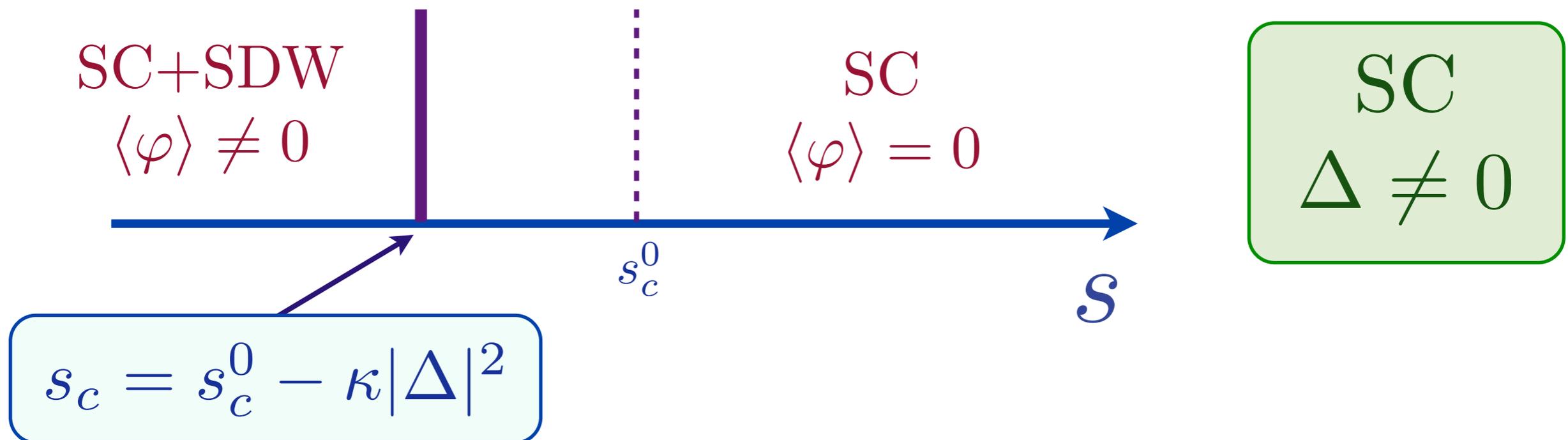
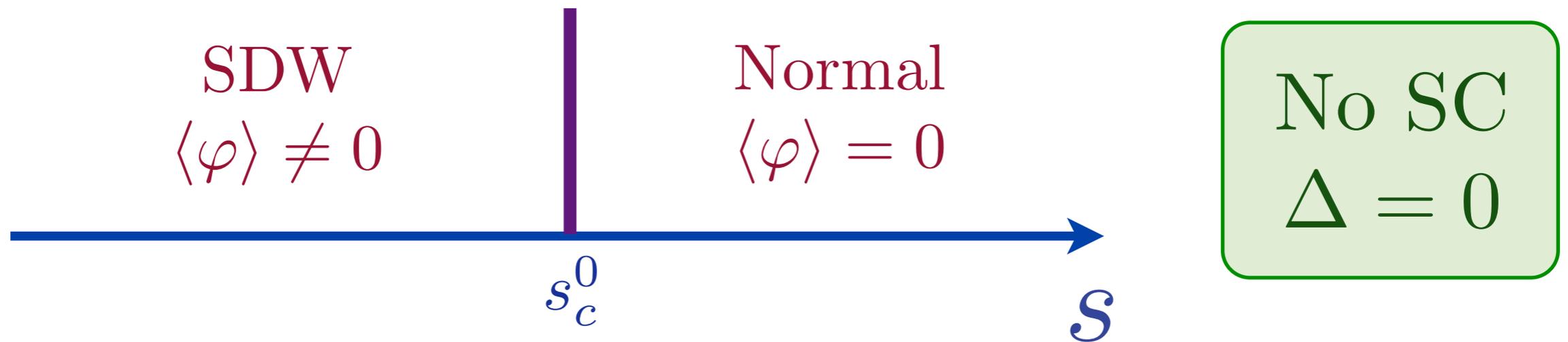
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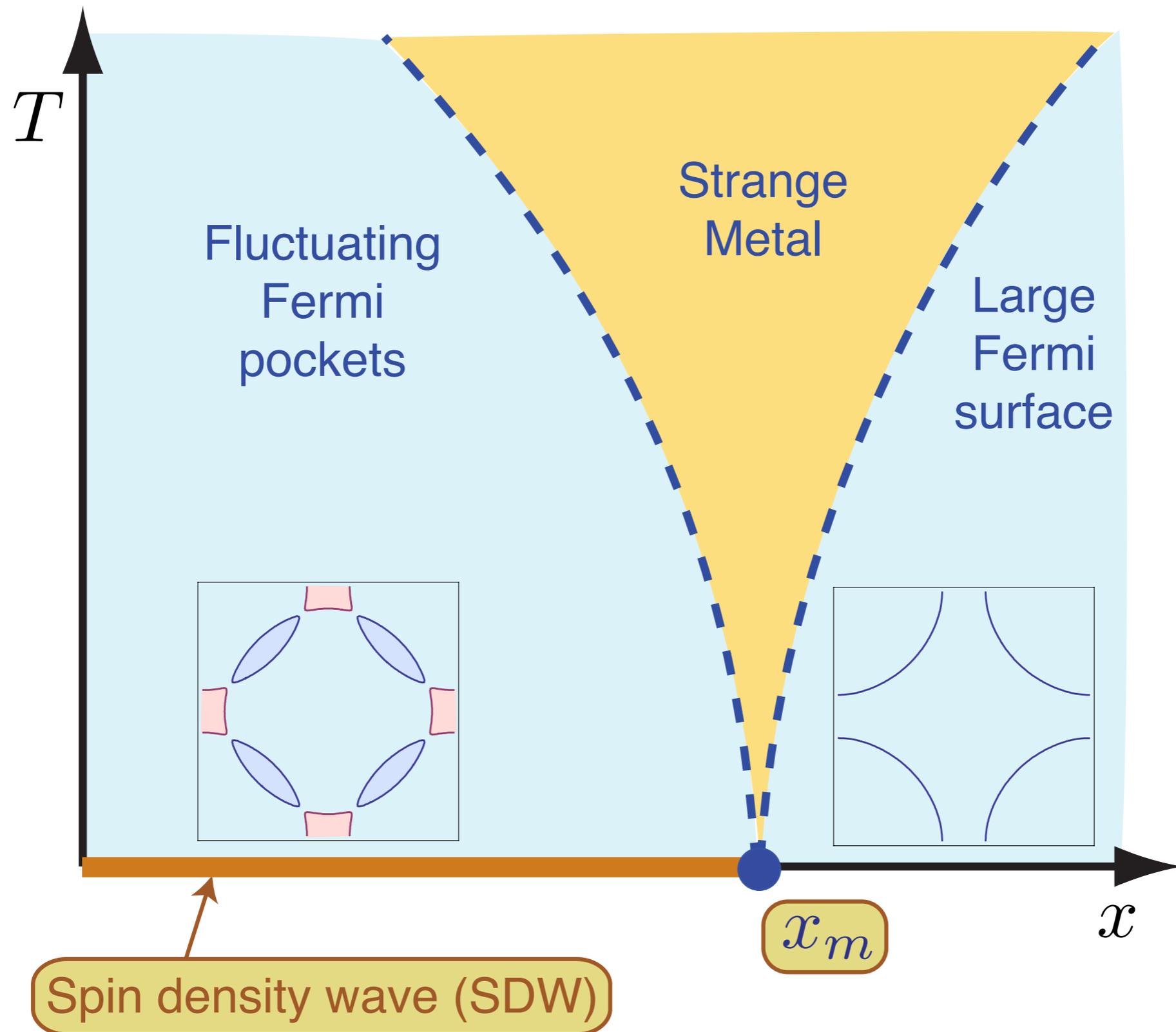


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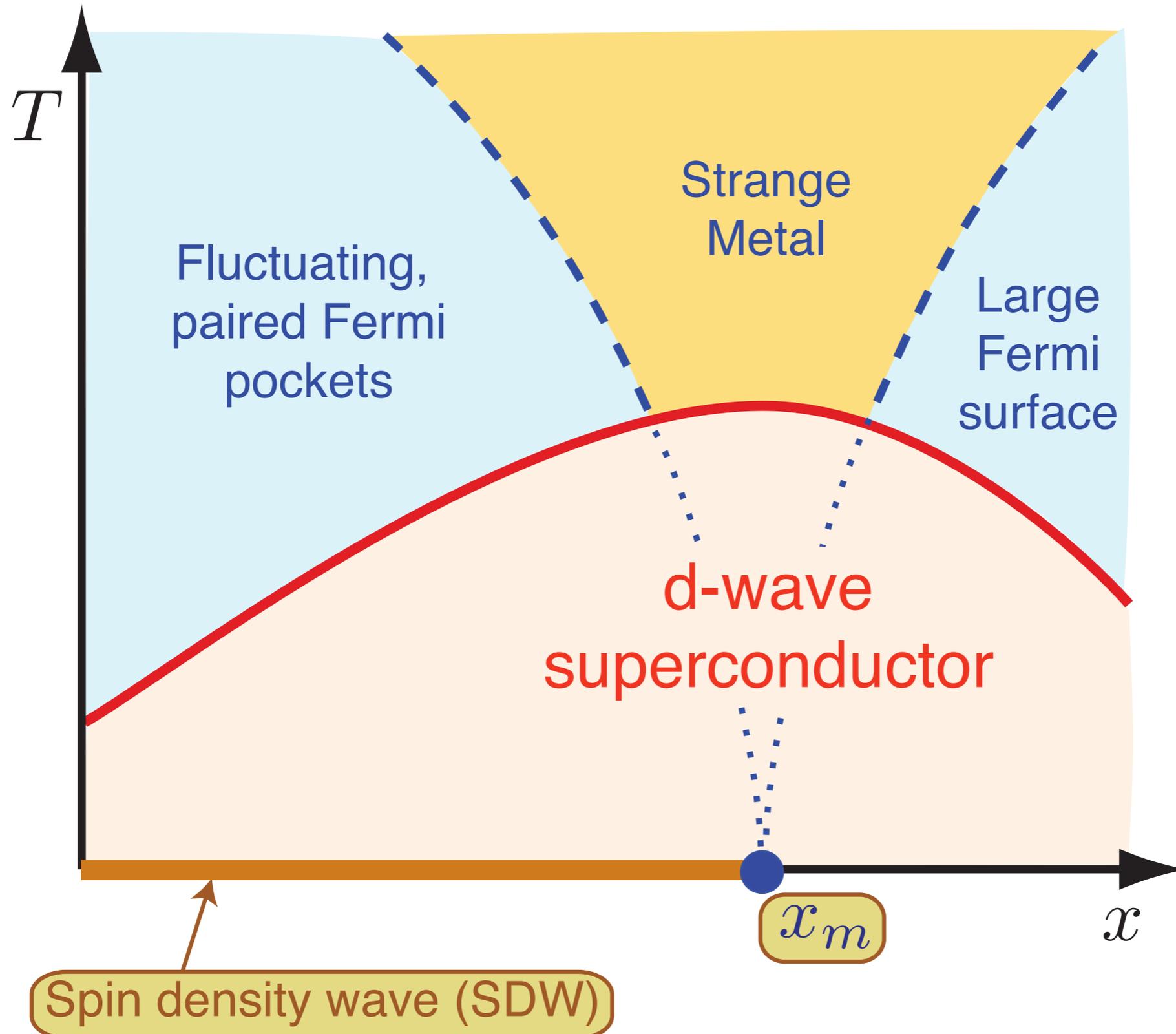


# Theory of quantum criticality in the cuprates



Underlying SDW ordering quantum critical point  
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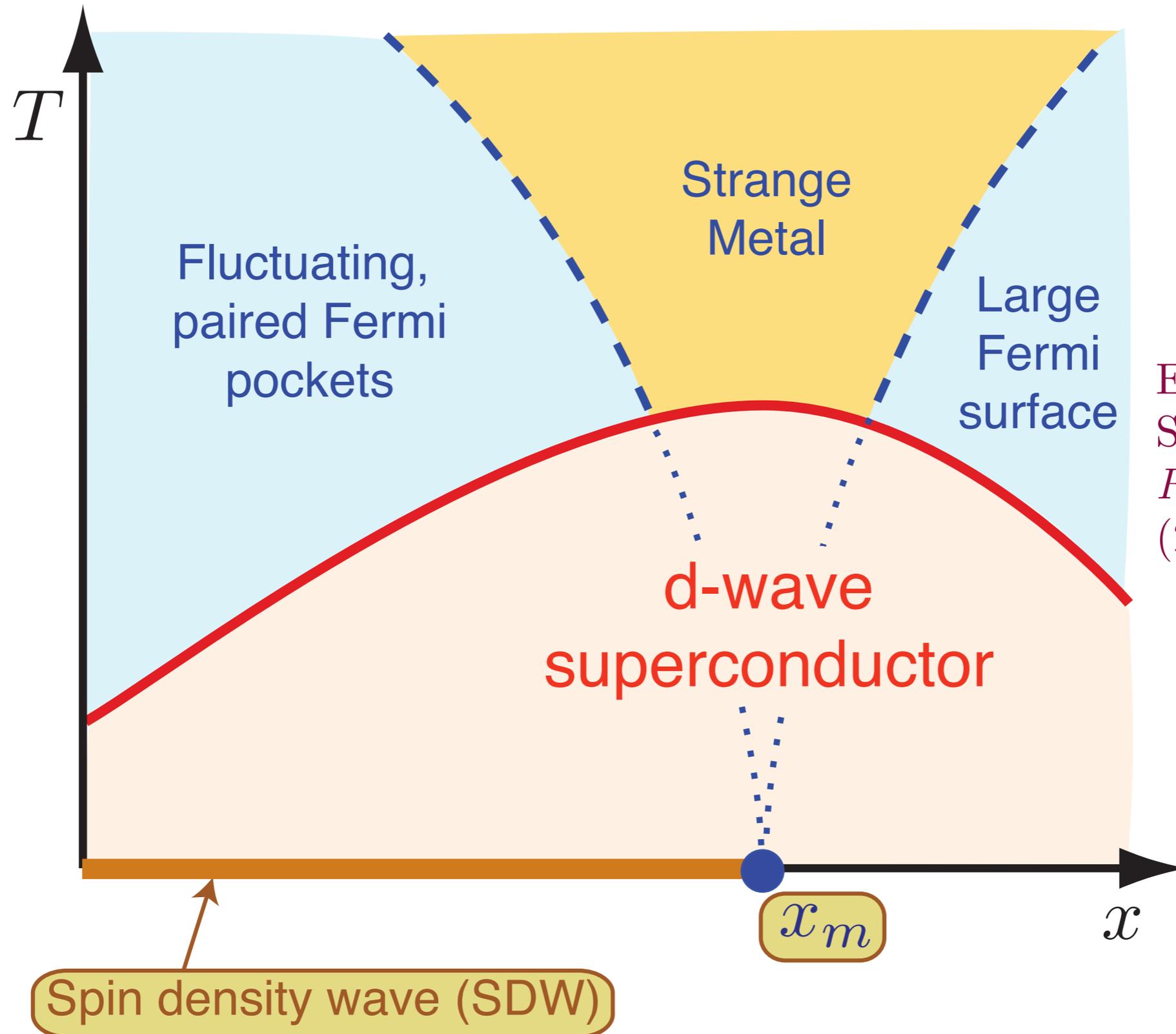


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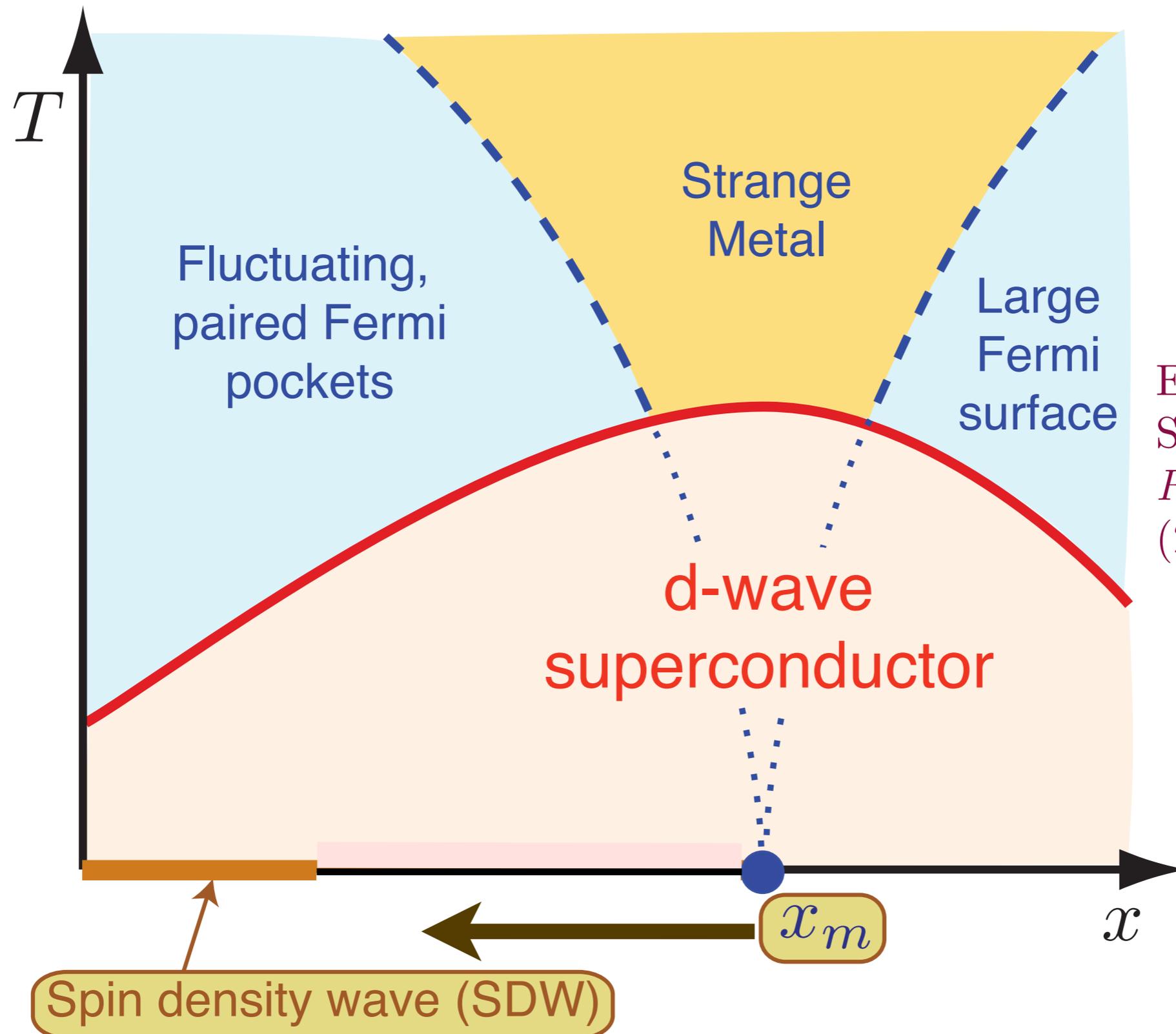
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E. G. Moon and S. Sachdev, *Phy. Rev. B* **80**, 035117 (2009)

Competition between SDW order and superconductivity moves the actual quantum critical point to  $x = x_s < x_m$ .

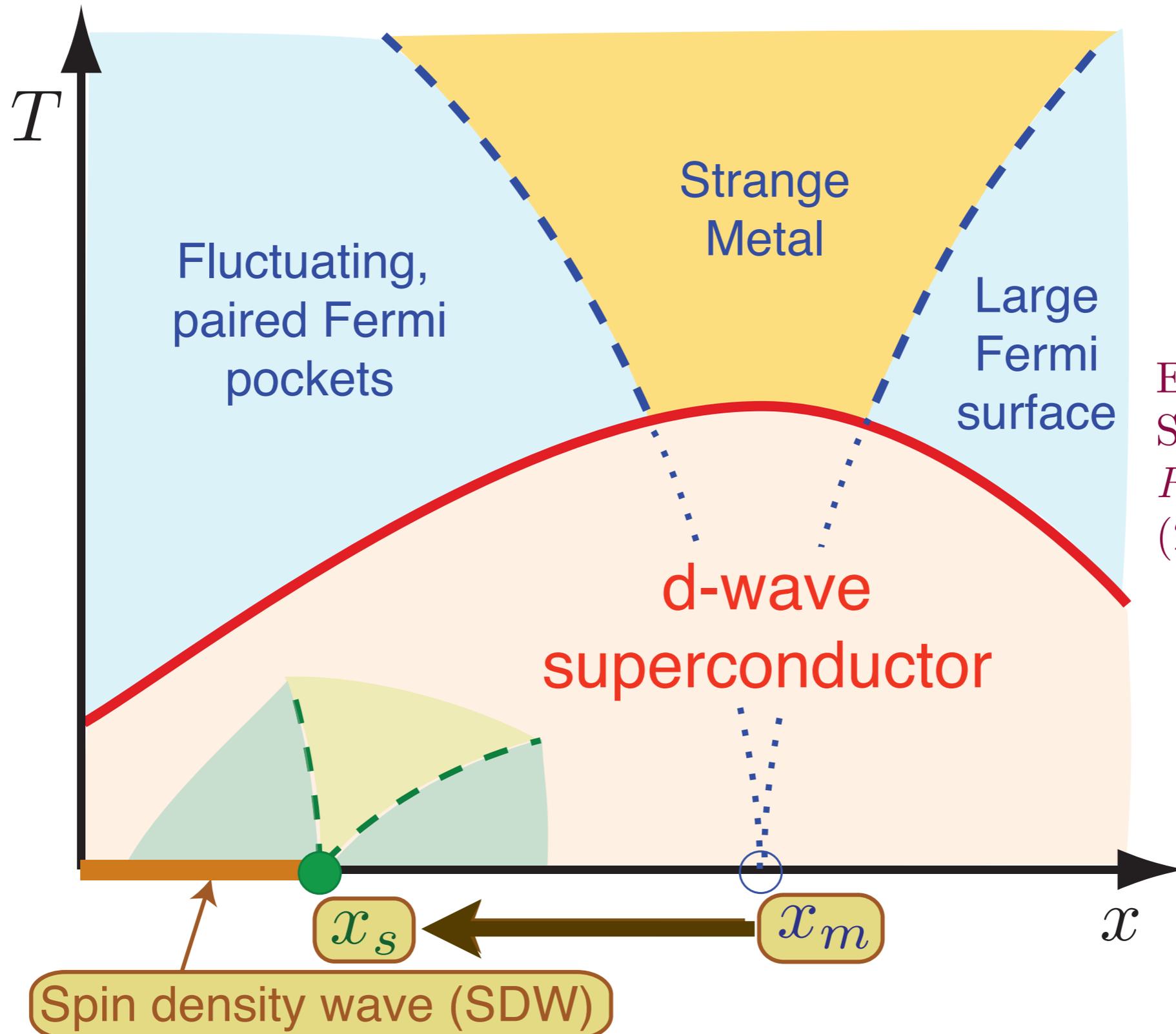
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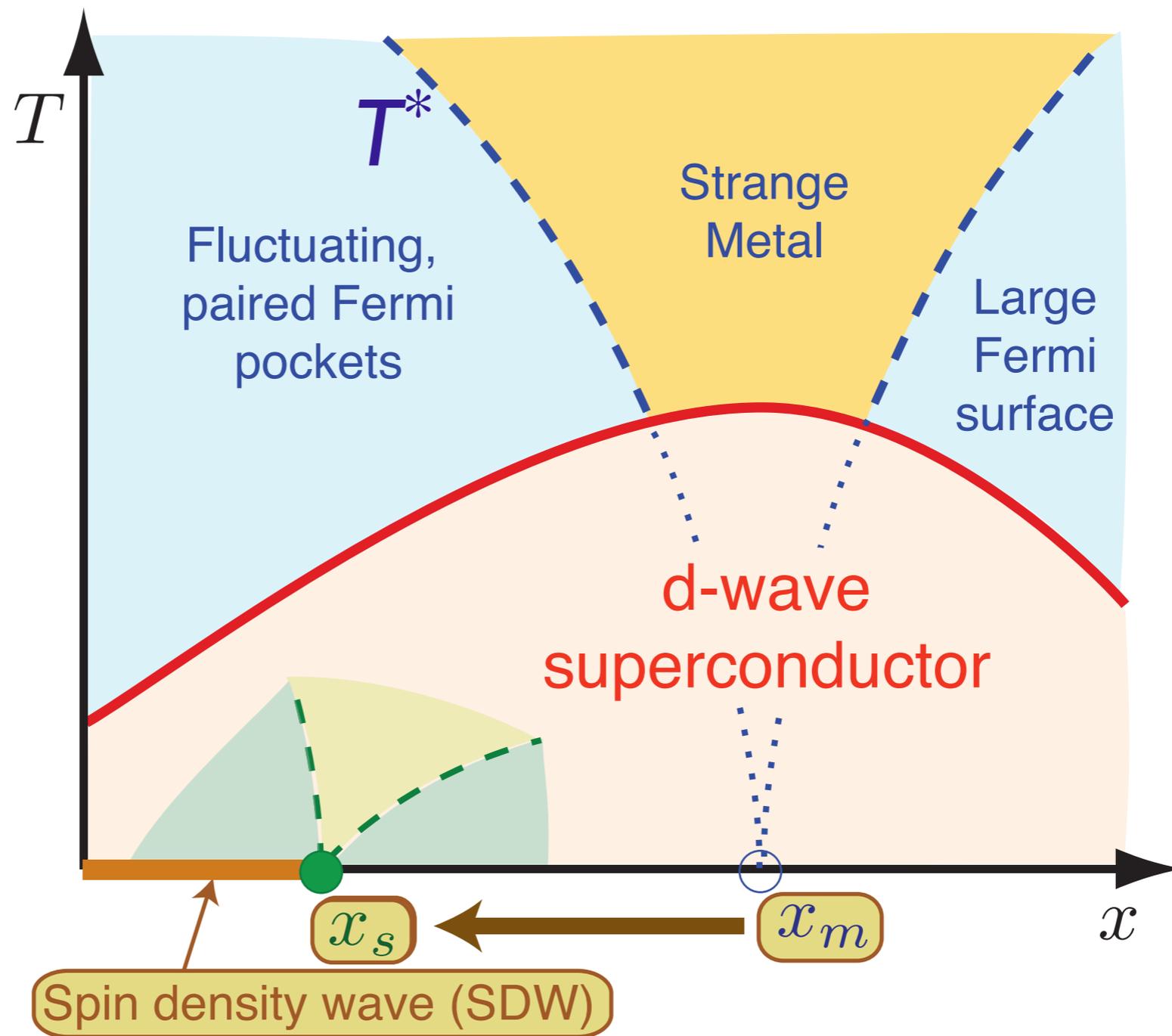
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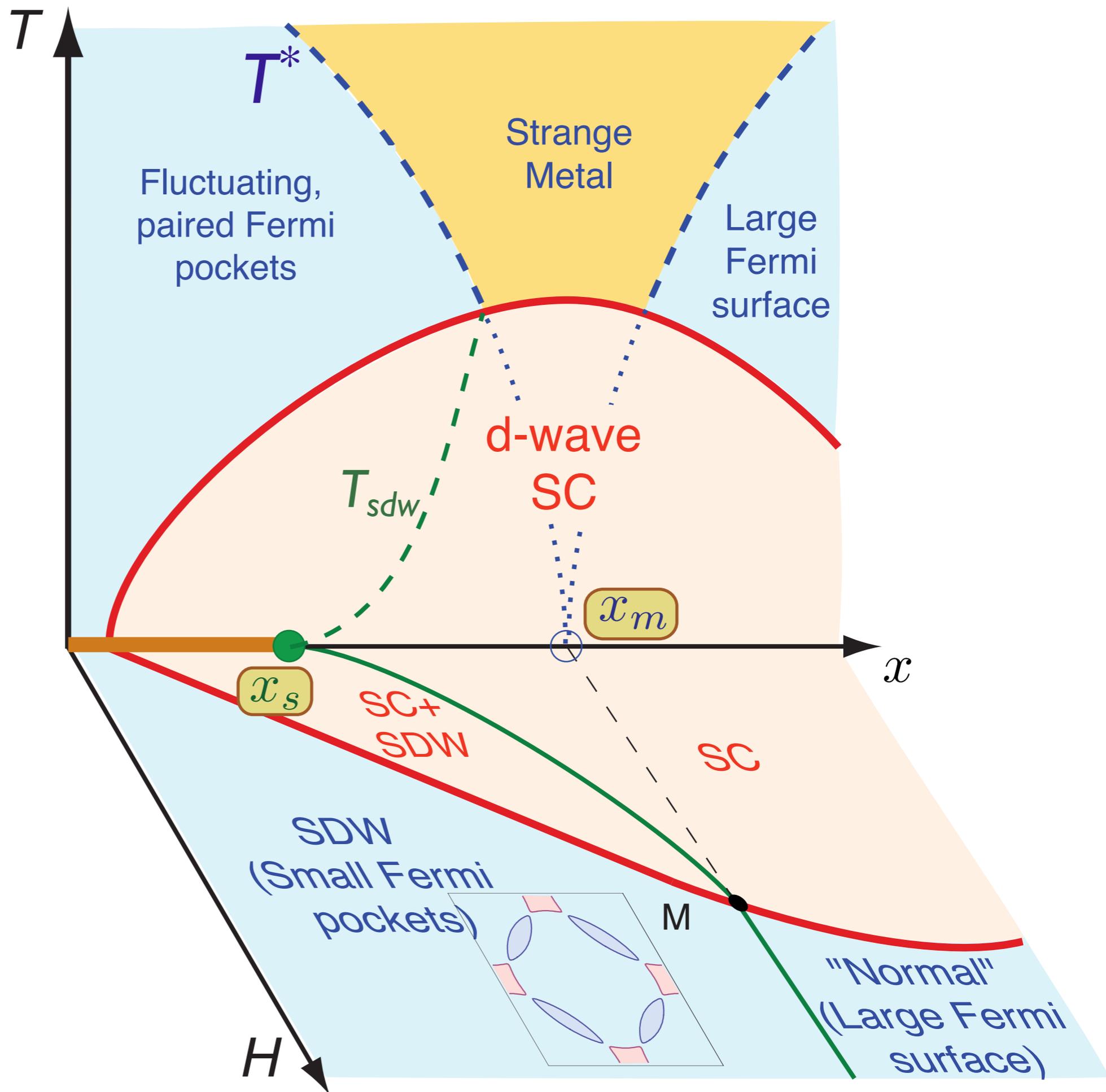


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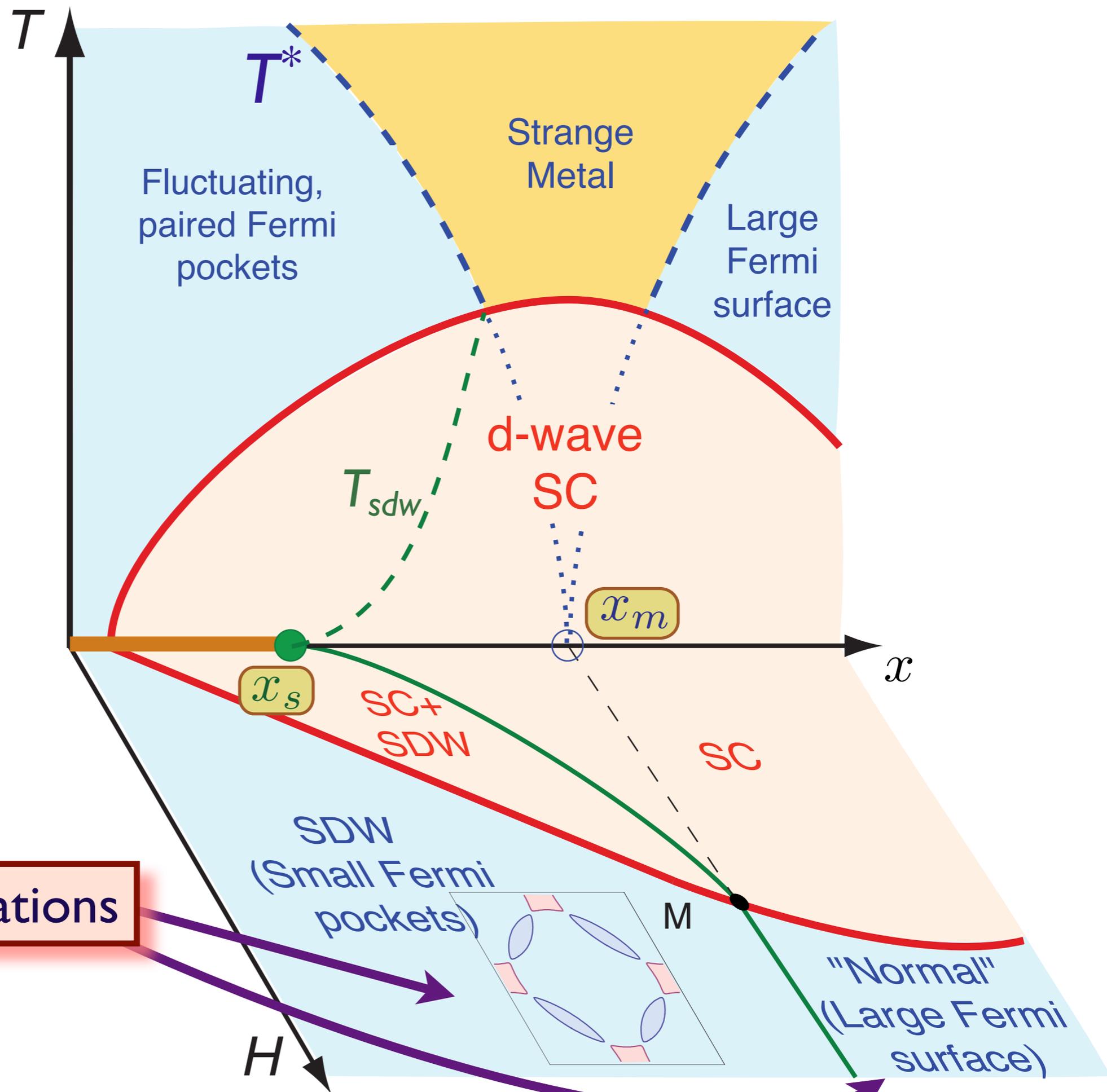
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E. Demler, S. Sachdev  
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Rev. Lett.* **87**,  
067202 (2001).

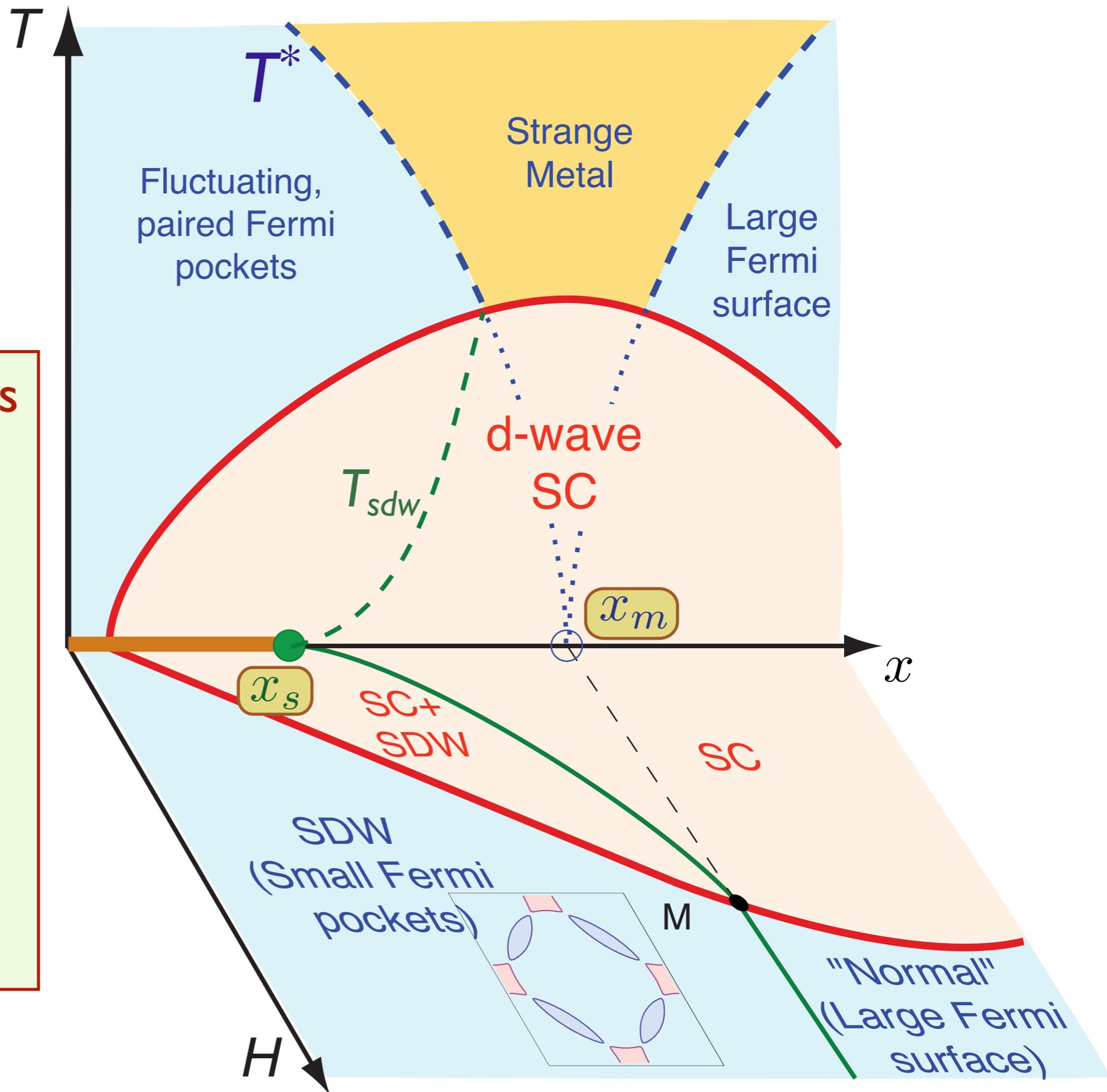


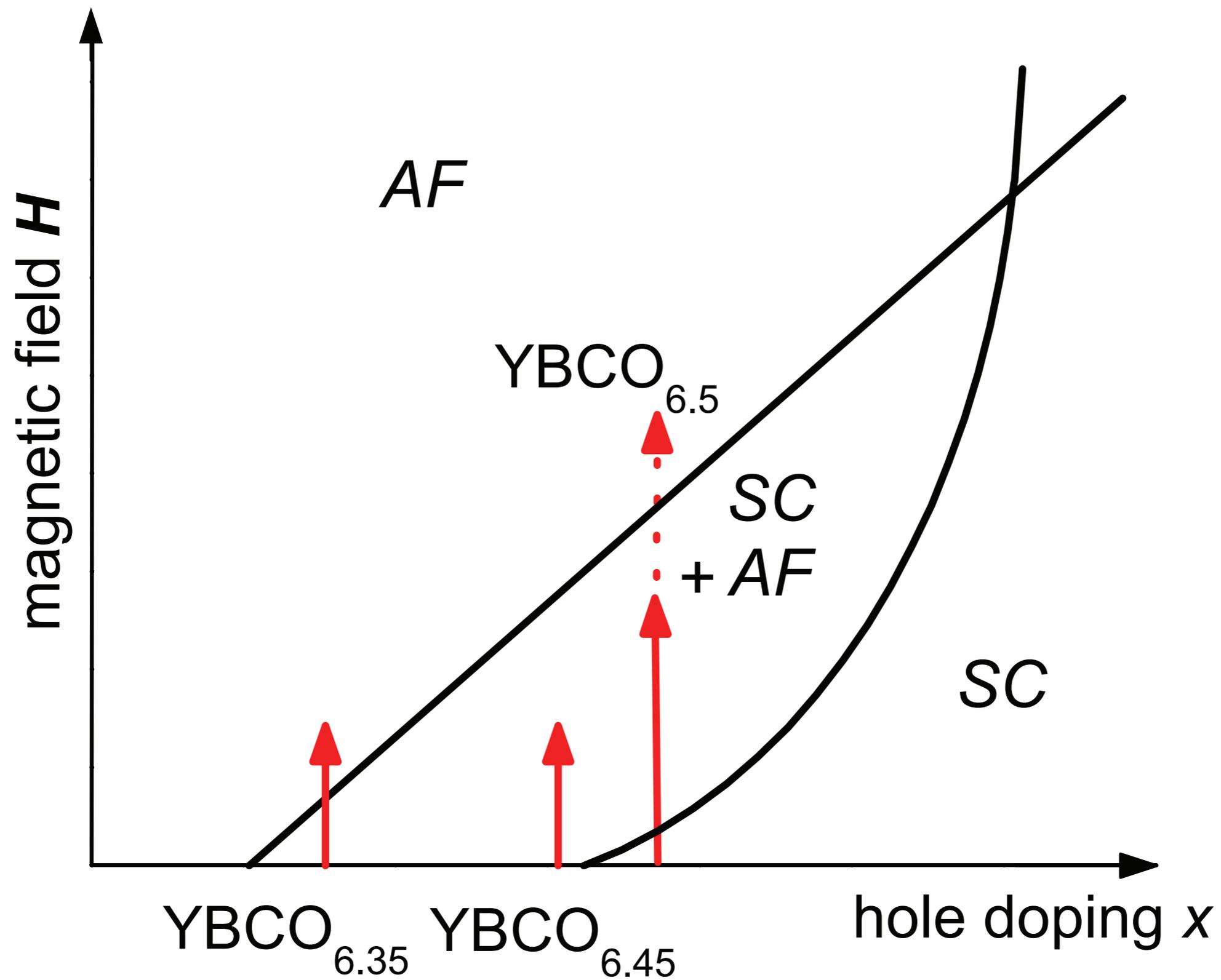
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Many experiments  
have presented  
evidence for the  
predicted  
green quantum  
phase transition  
line  
from SC to SC  
+SDW in a  
magnetic field

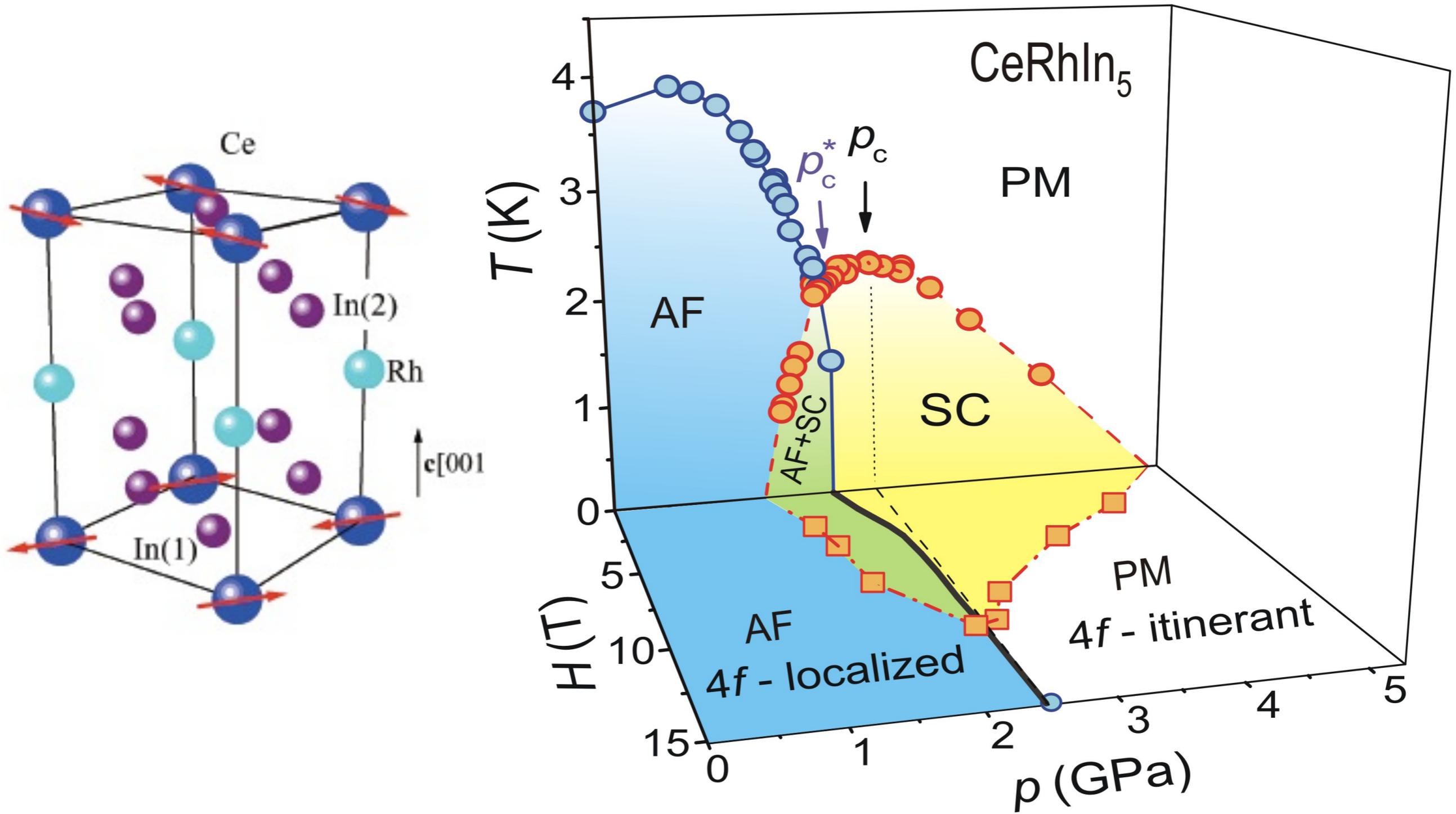




D. Haug, V. Hinkov, Y. Sidis, P. Bourges, N. B. Christensen, A. Ivanov, T. Keller, C. T. Lin, and B. Keimer, arXiv:1008.4298



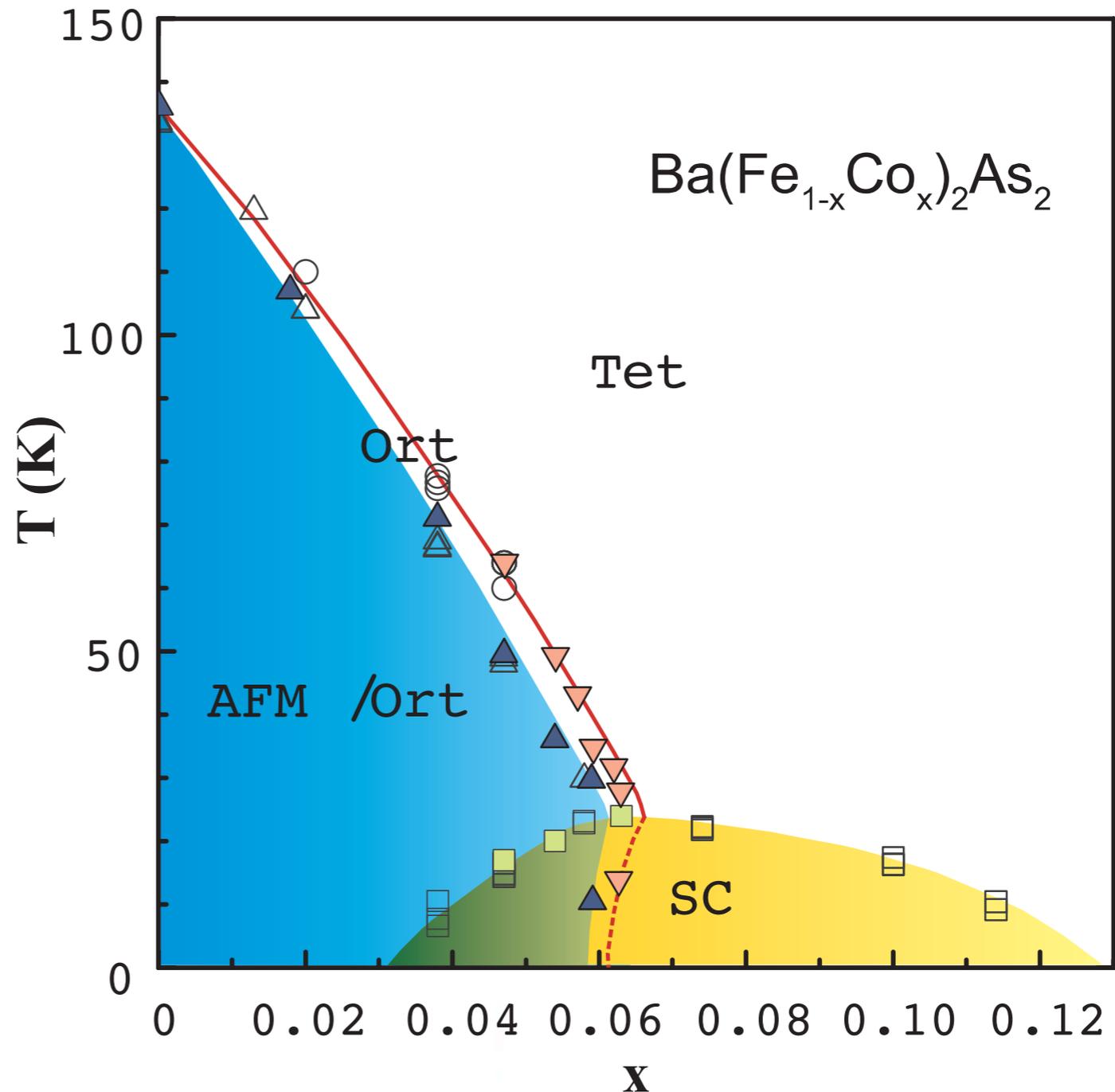
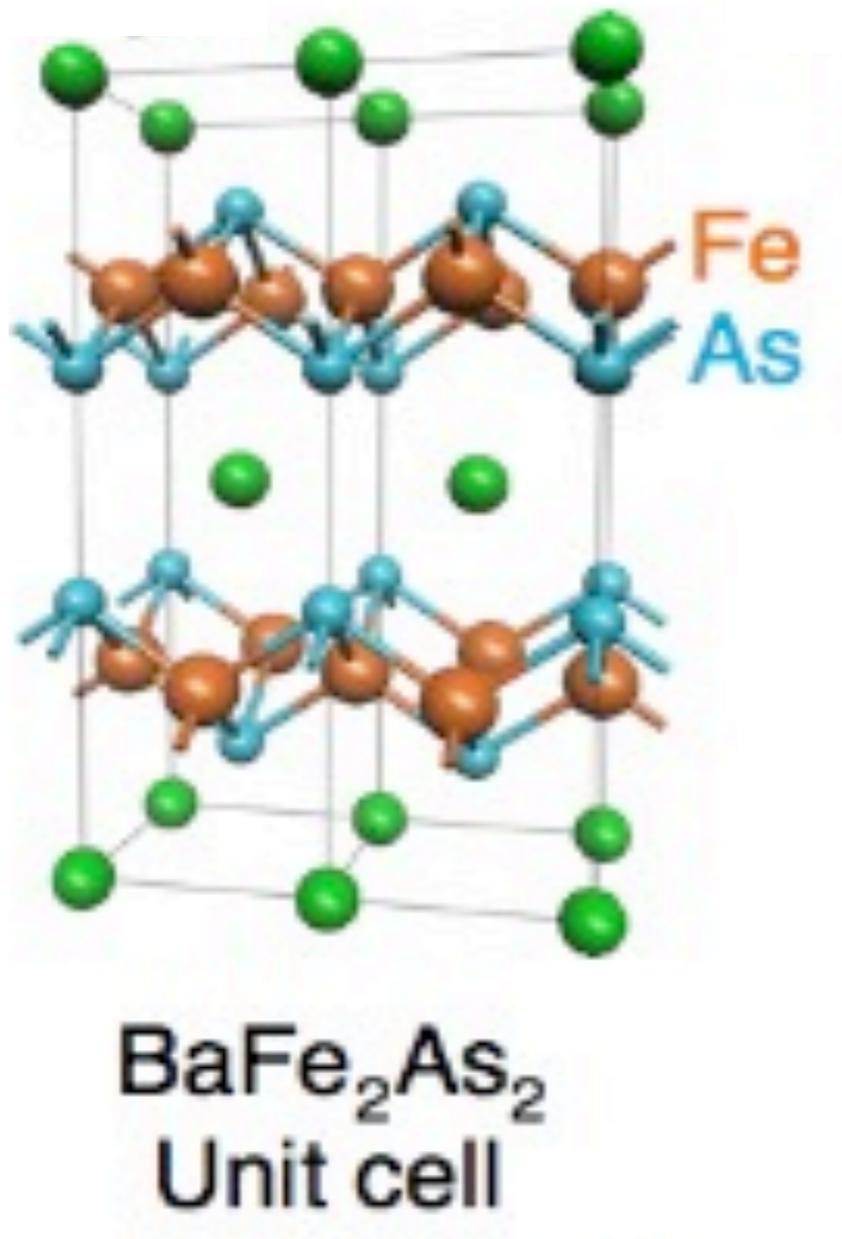
# Similar phase diagram for CeRhIn<sub>5</sub>



G. Knebel, D. Aoki, and J. Flouquet, arXiv:0911.5223

# Iron pnictides:

a new class of high temperature superconductors



S. Nandi, M. G. Kim, A. Kreyssig, R. M. Fernandes, D. K. Pratt, A. Thaler, N. Ni,  
S. L. Bud'ko, P. C. Canfield, J. Schmalian, R. J. McQueeney, A. I. Goldman,  
*Physical Review Letters* **104**, 057006 (2010).

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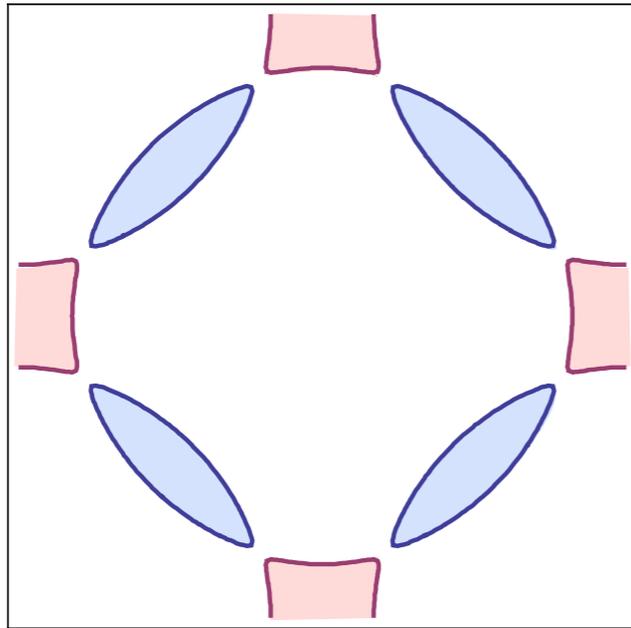
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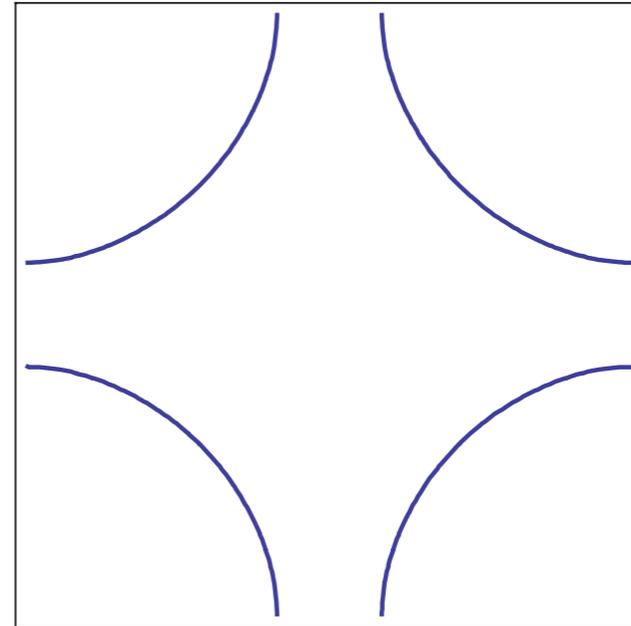
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$$\langle \vec{\varphi} \rangle \neq 0$$



Metal with electron  
and hole pockets

$$\langle \vec{\varphi} \rangle = 0$$

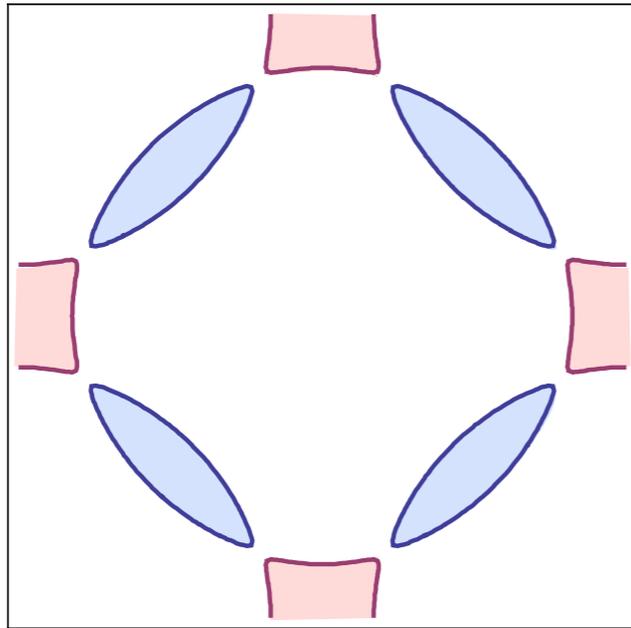


Metal with "large"  
Fermi surface

$S$

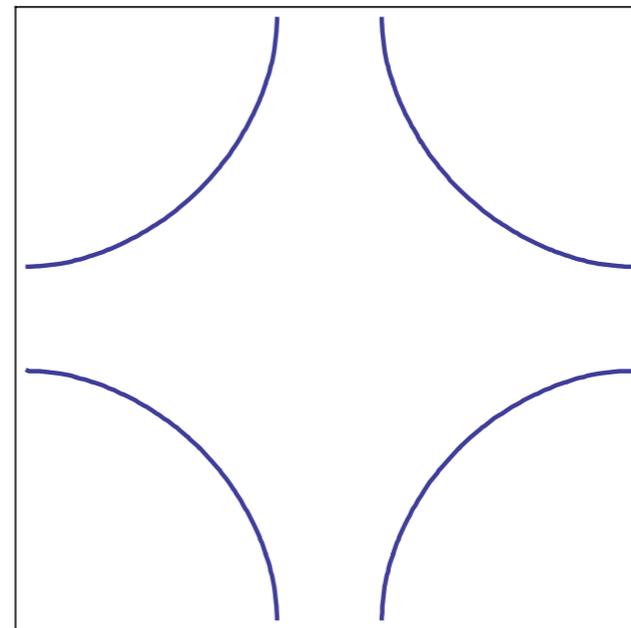
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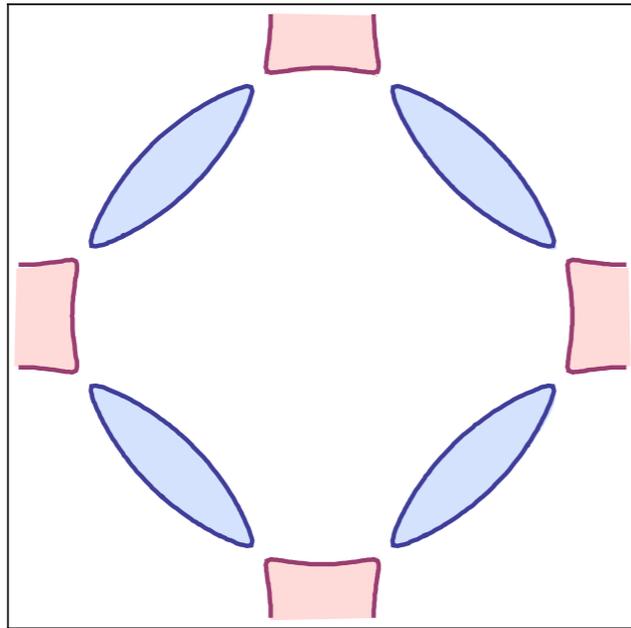


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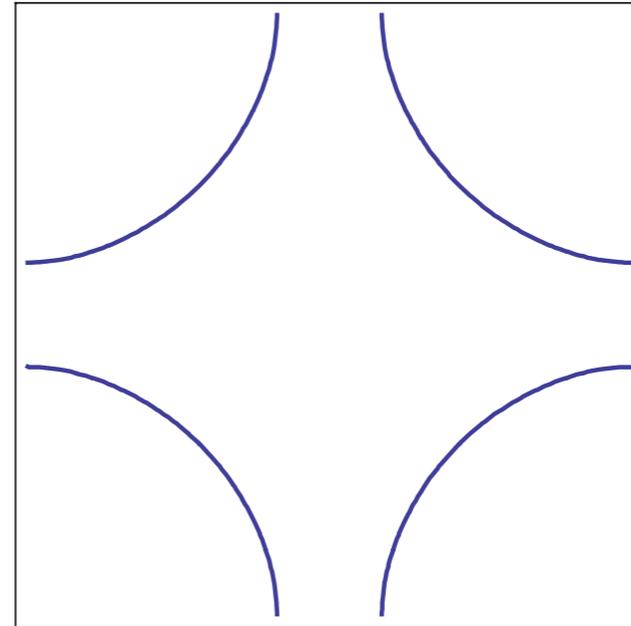
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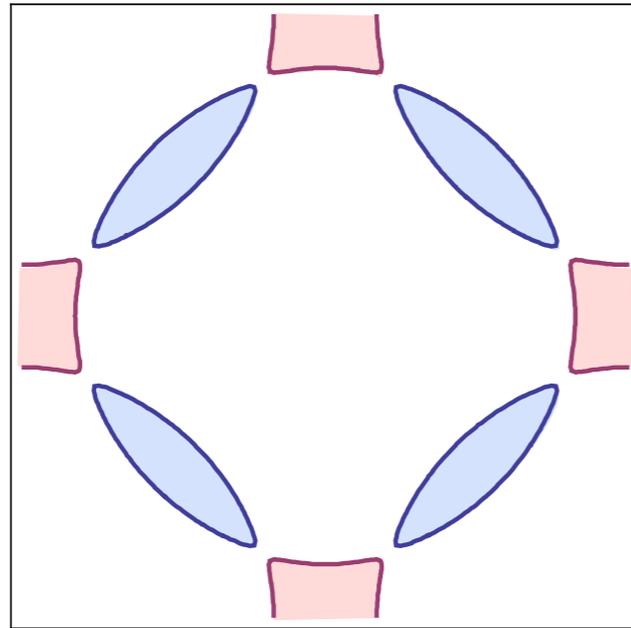


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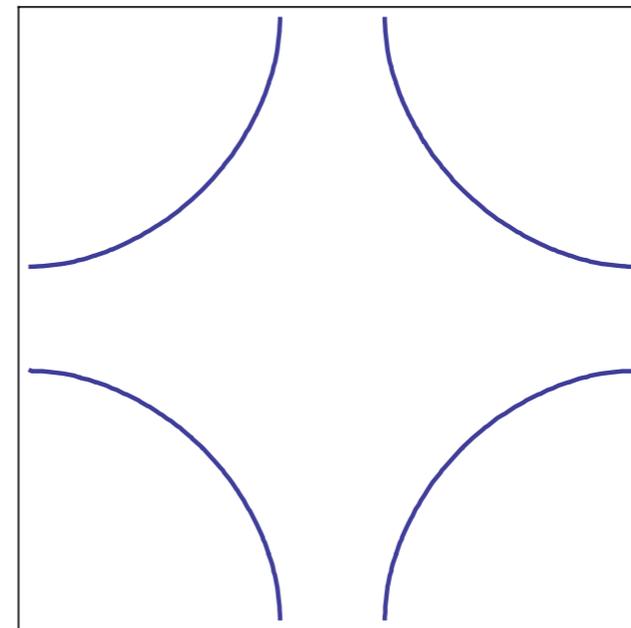
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Fermi surface

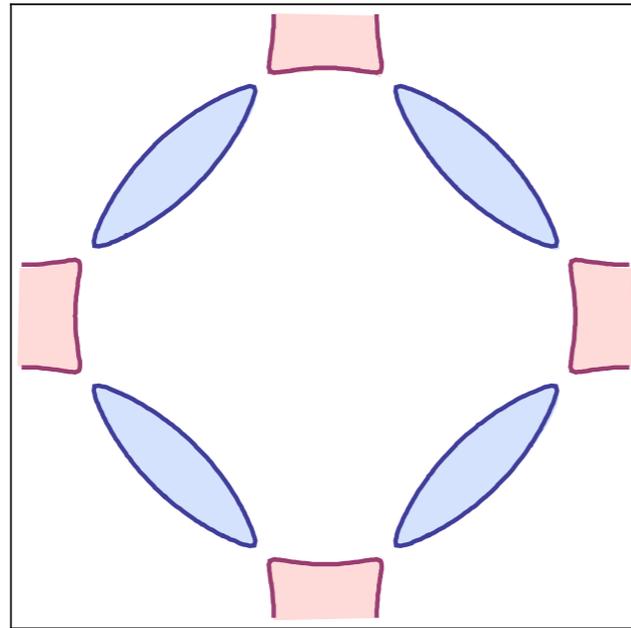
$S$

- Quantum critical theory is strongly-coupled in two (but not higher) spatial dimensions (but not a CFT).



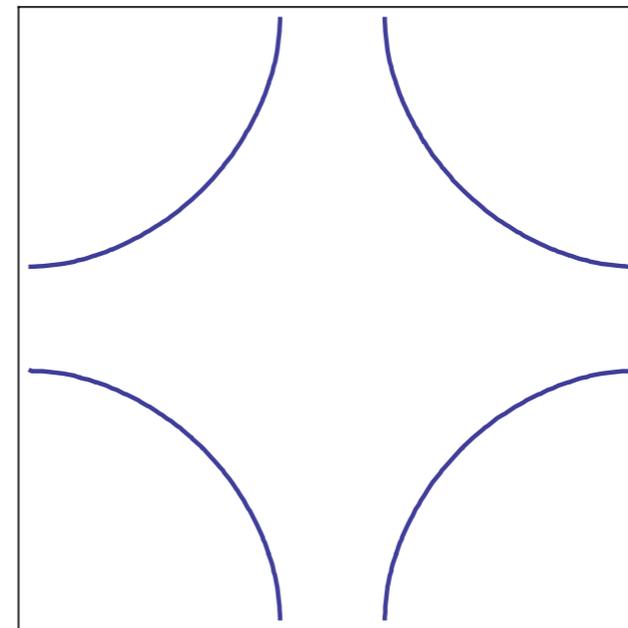
# Quantum criticality of the onset of antiferromagnetism in a metal

$$\langle \vec{\varphi} \rangle \neq 0$$



Metal with electron and hole pockets

$$\langle \vec{\varphi} \rangle = 0$$



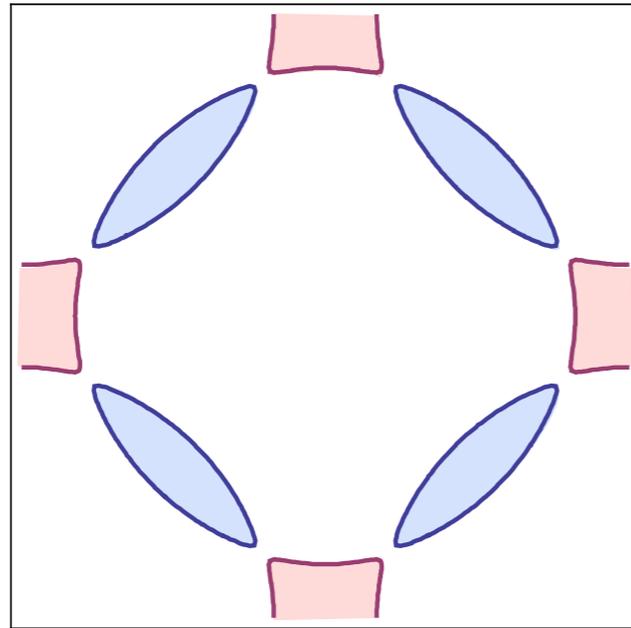
Metal with "large" Fermi surface

$\rightarrow S$

- Theory has strong (log-squared) instability already at the Fermi energy to superconductivity with sign-changing pairing amplitude near quantum criticality.

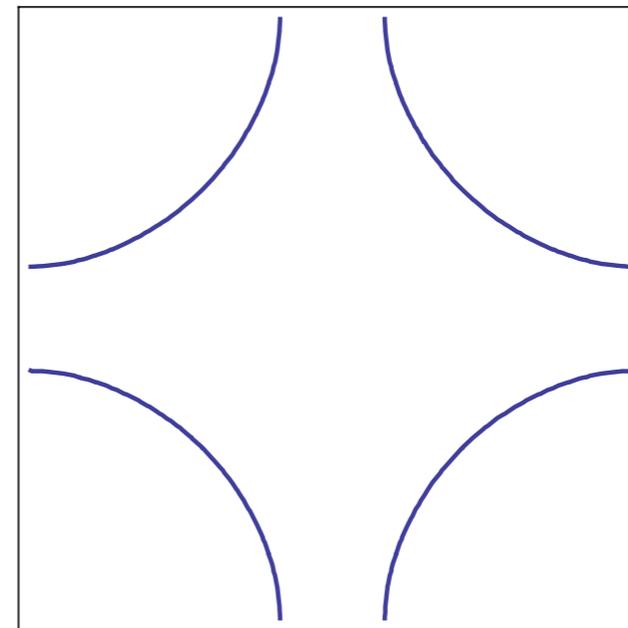
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$S$

There is non-Fermi liquid behavior at the QCP not only at hotspots, but on entire Fermi surface.

## Questions

- *Can quantum fluctuations near the loss of antiferromagnetism induce higher temperature superconductivity ?*
- *If so, why is there no antiferromagnetism in the cuprates near the point where the superconductivity is strongest ?*
- *What is the physics of the strange metal ?*

## Questions and answers

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**Yes**

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Competition between antiferromagnetism and superconductivity has shifted the antiferromagnetic quantum-critical point (QCP), and shrunk the region of antiferromagnetism. This QCP shift is largest in the cuprates

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● *What is the physics of the strange metal ?*

Proposal: strongly-coupled quantum criticality of fluctuating antiferromagnetism in a metal