

# Topology and the phases of quantum matter in two dimensions

Albanova and Nordita Colloquium,

2016 Nobel prize in physics, impact on physics research and future directions

December 8, 2016

Subir Sachdev



PERIMETER INSTITUTE  
FOR THEORETICAL PHYSICS

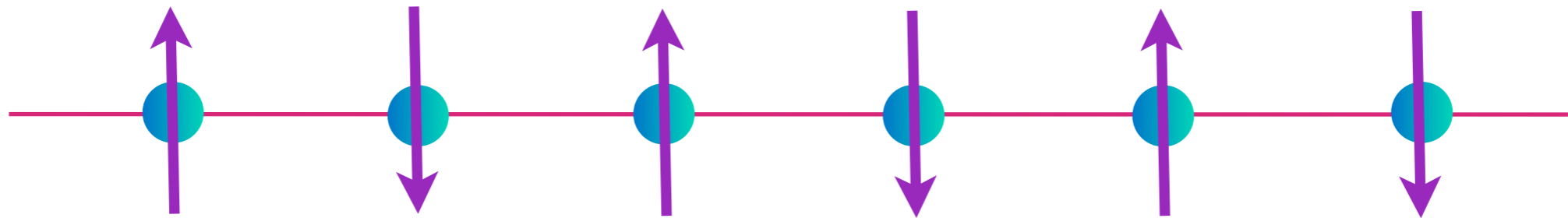
PHYSICS



HARVARD

# Quantum antiferromagnets in one dimension

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$$



Haldane: Semiclassical theory of quantum fluctuations

Berry phases of spacetime Skyrmions textures lead to sensitivity to  $2S \pmod{2}$ . For  $S=1$ , we obtain the Haldane gap SPT state.



$$\bullet \text{---} \bullet = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

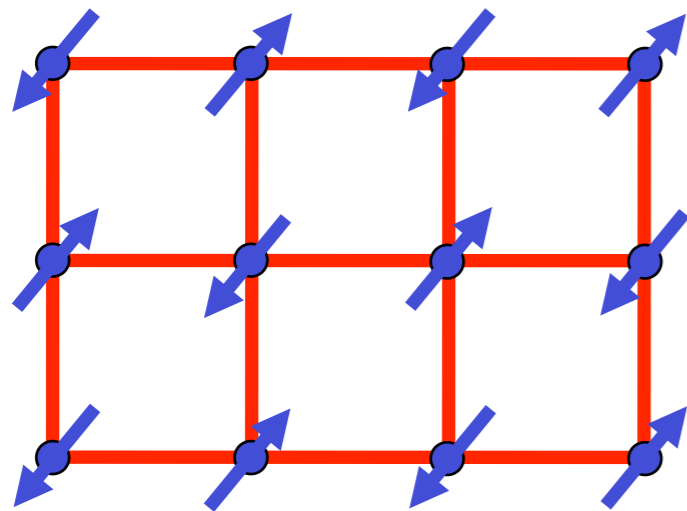
$$\bigcirc = |+\rangle\langle\uparrow\uparrow| + |0\rangle \frac{\langle\uparrow\downarrow| + \langle\downarrow\uparrow|}{\sqrt{2}} + |-\rangle\langle\downarrow\downarrow|$$

# Quantum antiferromagnets on the square lattice

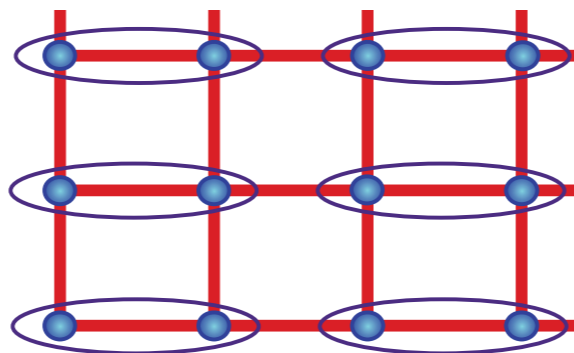
$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$$

$$= \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

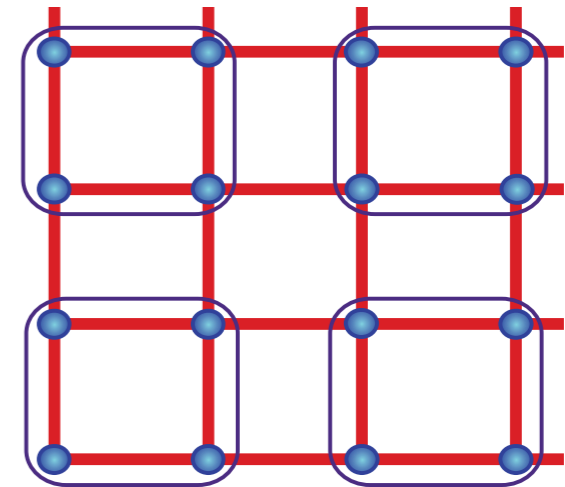
Semiclassical theory



Néel state



or



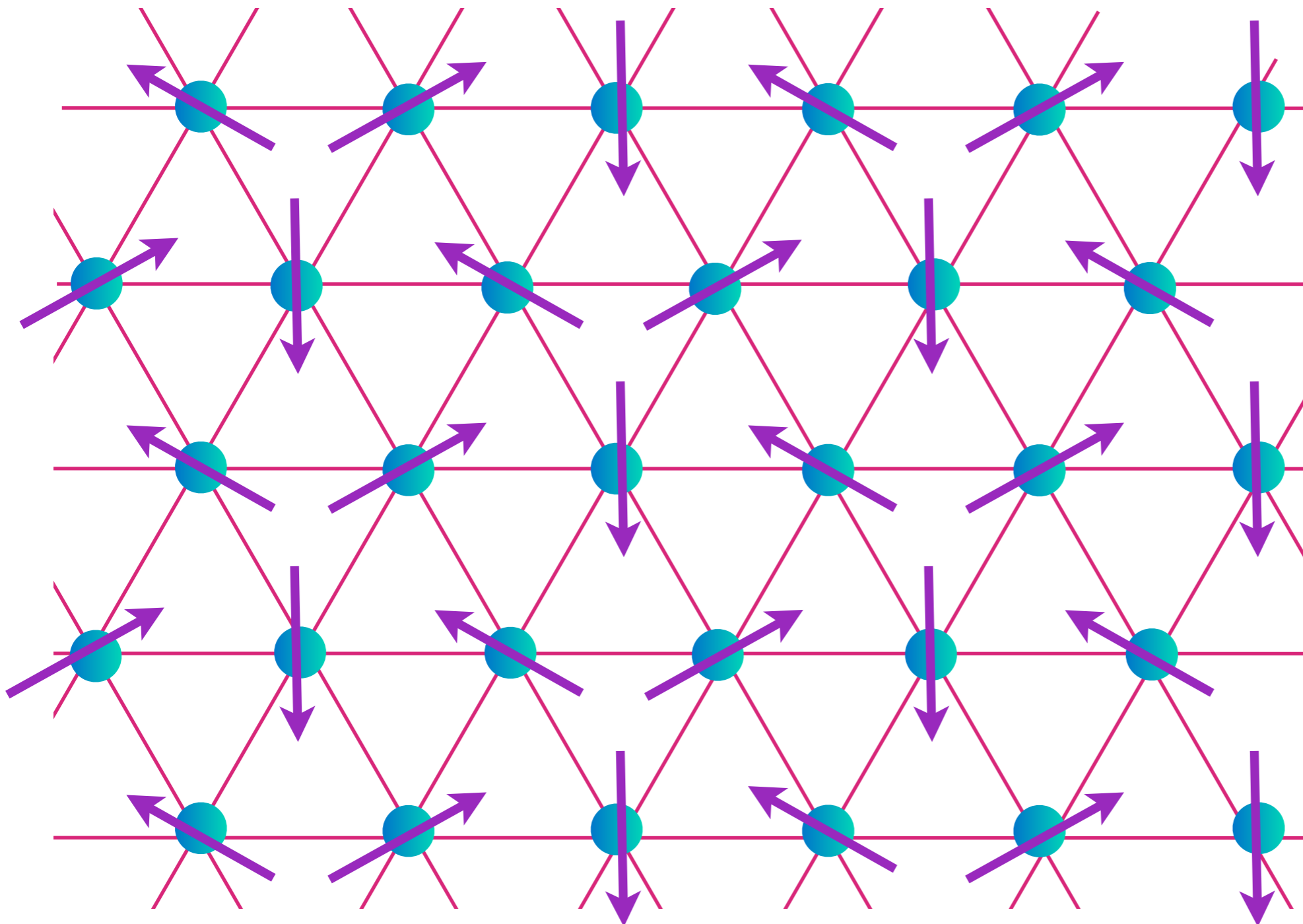
Valence bond solid (VBS) state  
with a nearly gapless, emergent “photon”

Haldane: Berry phases of spacetime instantons (“hedgehogs”) lead to sensitivity to  $2S \pmod{4}$ .

Read and S.S. : Valence bond solids break lattice translational symmetry except for  $2S \pmod{4} = 0$ .

# Quantum antiferromagnets on the triangular lattice

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$$



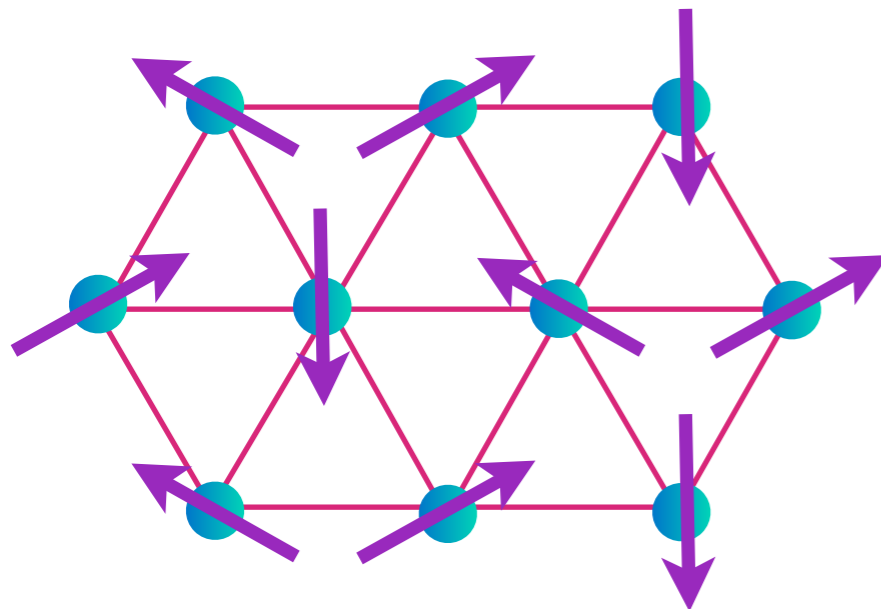
Nearest-neighbor model has non-collinear Néel order

# Quantum antiferromagnets on the triangular lattice

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$$

Semiclassical theory

(Read and S.S. (1991))



non-collinear Néel state

Semiclassical theory leads to a spin liquid state with  $Z_2$  topological order.

The topological order is associated with suppressed vortex defects in the non-collinear Neel order

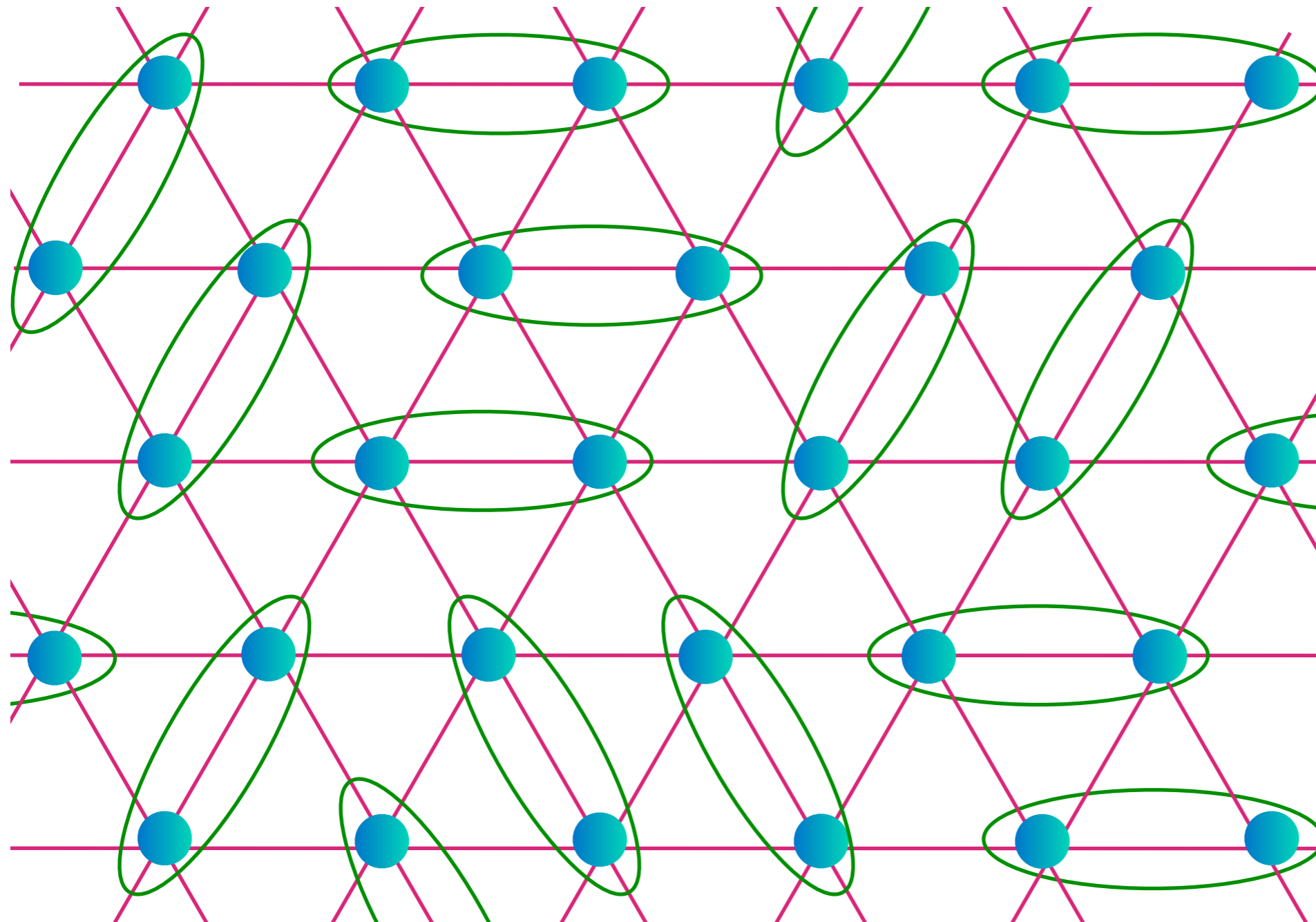
$S_c$

$S$

Vortex suppression is similar to low temperature phase of the Kosterlitz-Thouless transition. (However spin correlations decay exponentially, not as a power-law)

# Quantum antiferromagnets on the triangular lattice

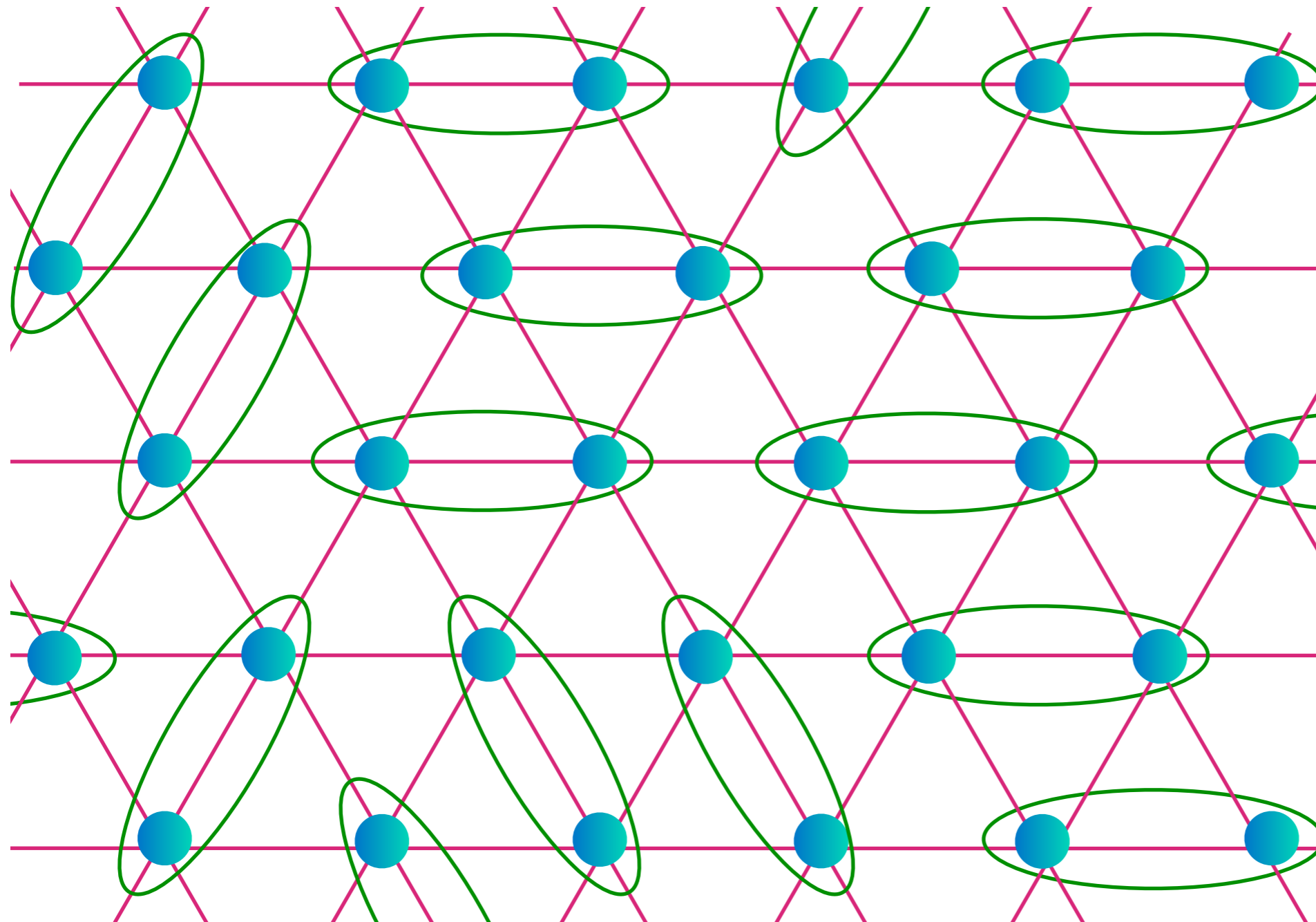
Suppressed-vortex-state is the resonating valence bond state with  $Z_2$  topological order.



$$\begin{array}{c} \text{○} \quad \text{○} \\ \hline = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \end{array}$$

# Quantum antiferromagnets on the triangular lattice

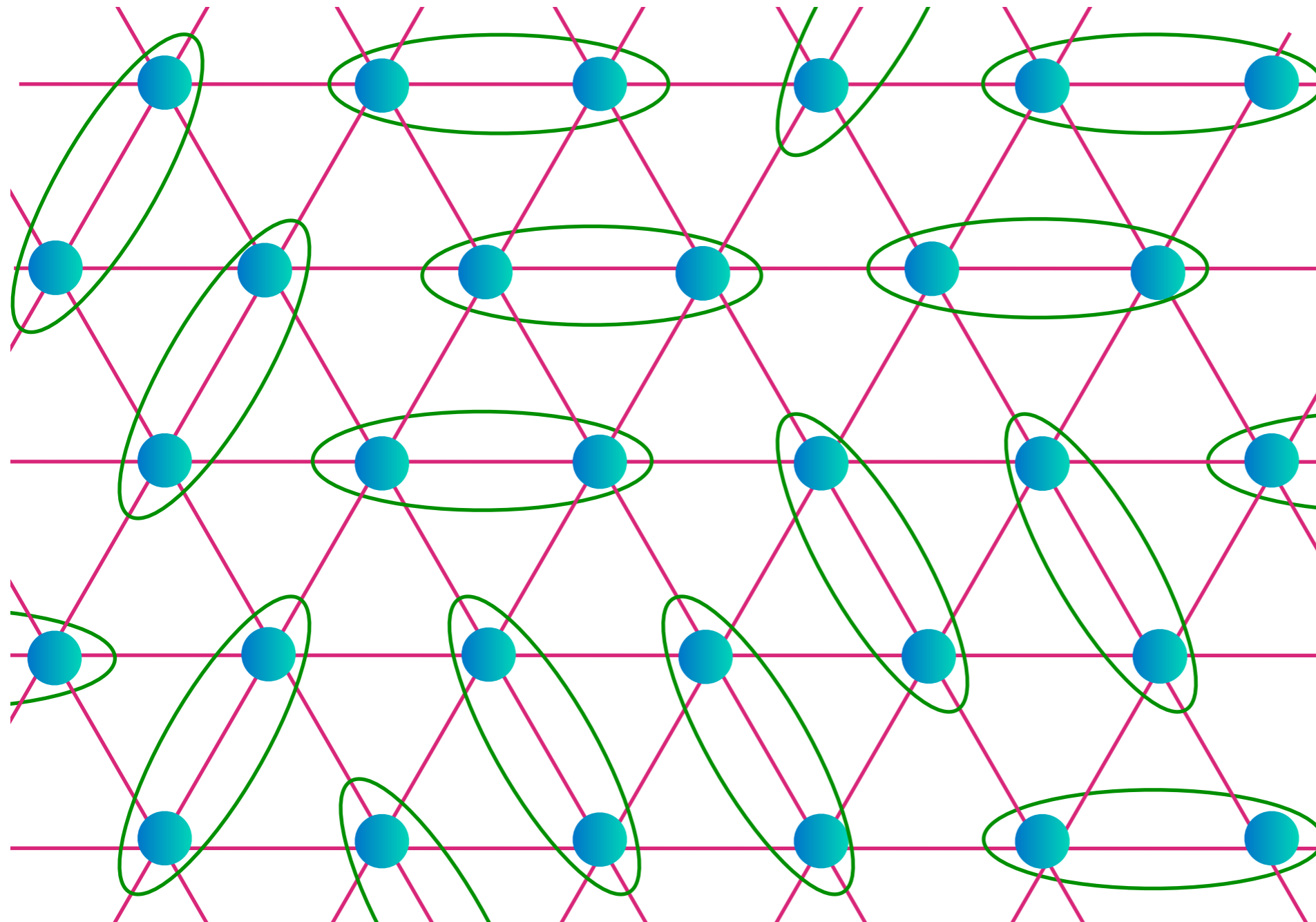
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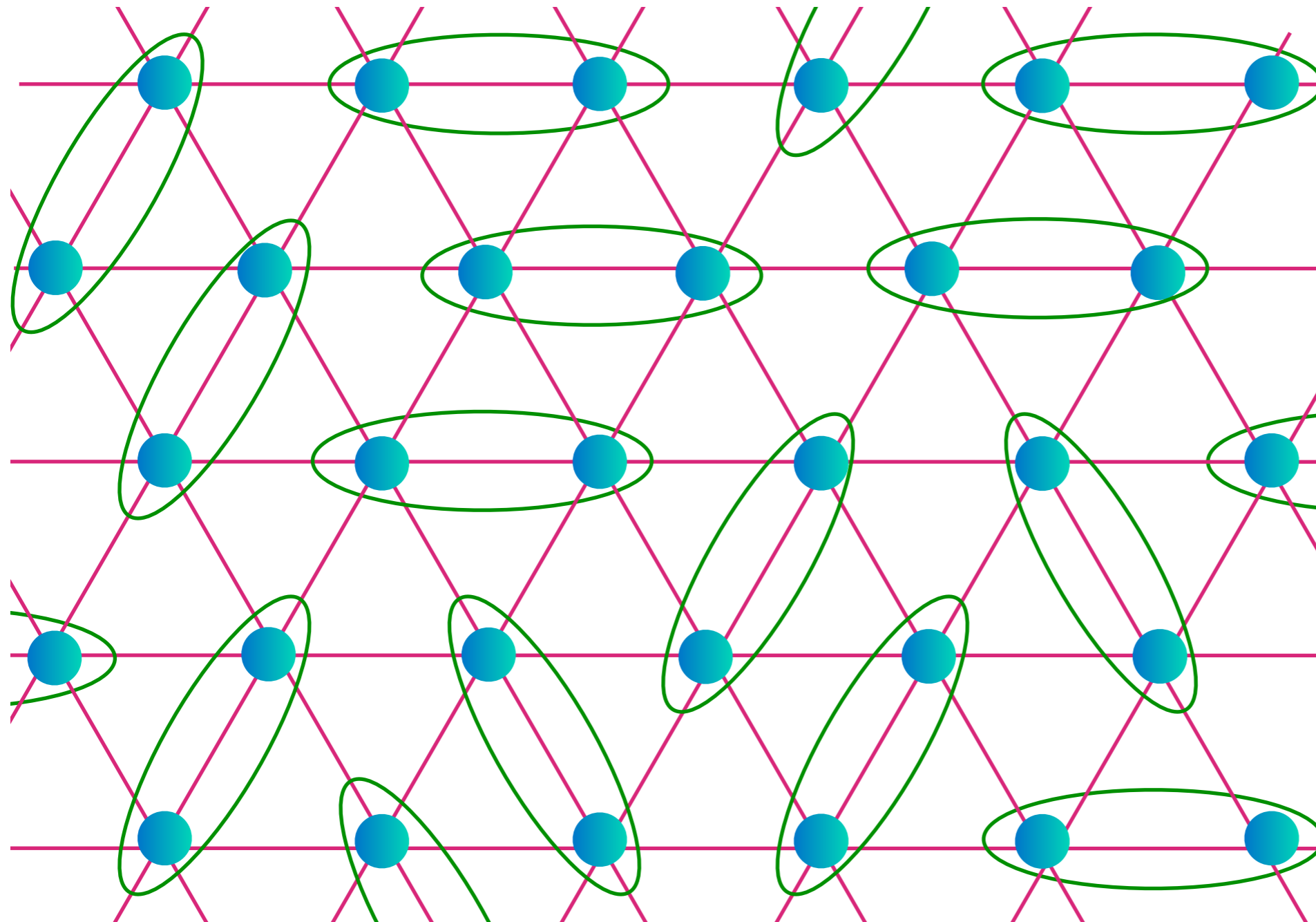


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# Quantum antiferromagnets on the triangular lattice

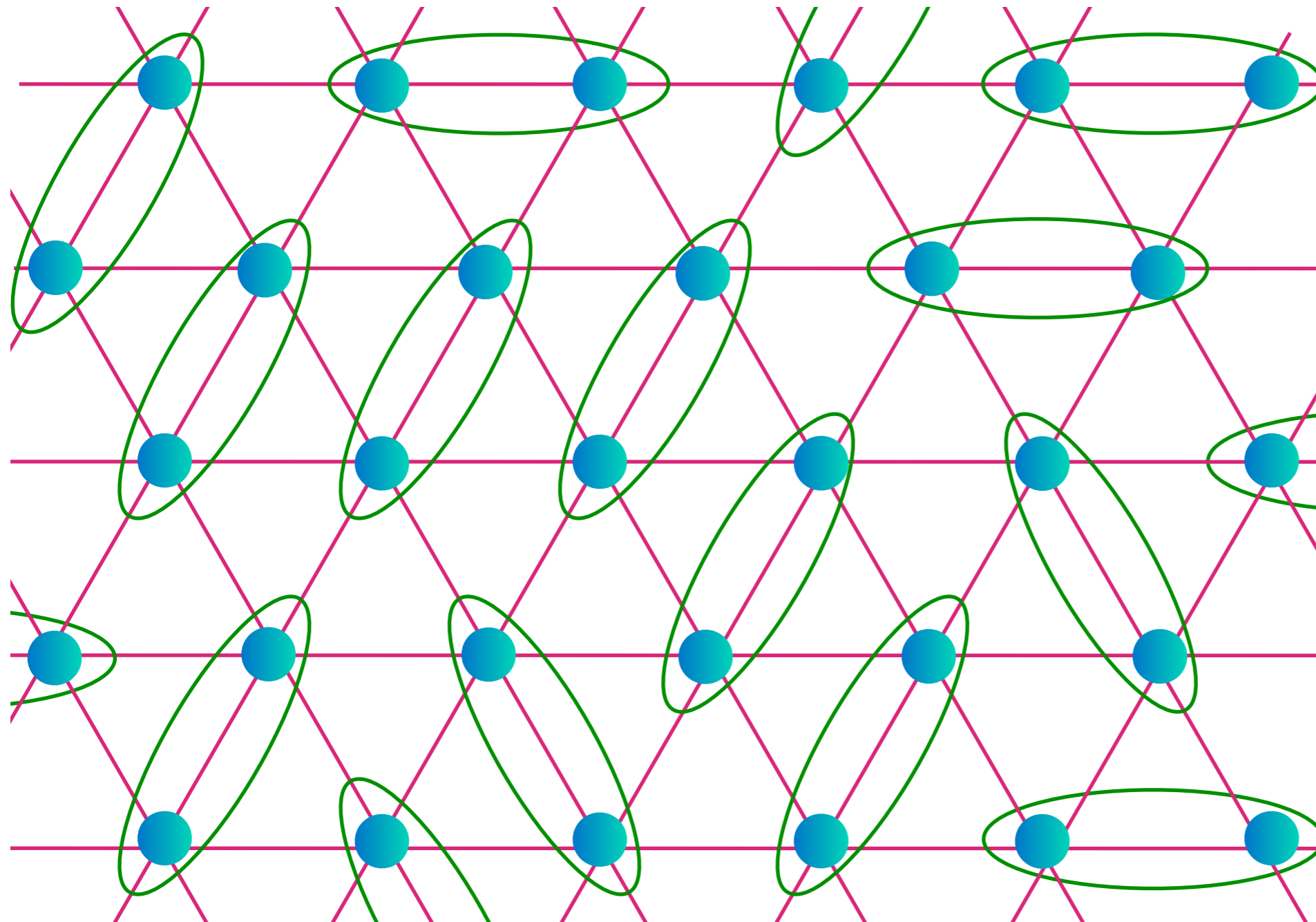
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# Quantum antiferromagnets on the triangular lattice

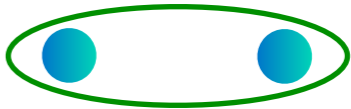
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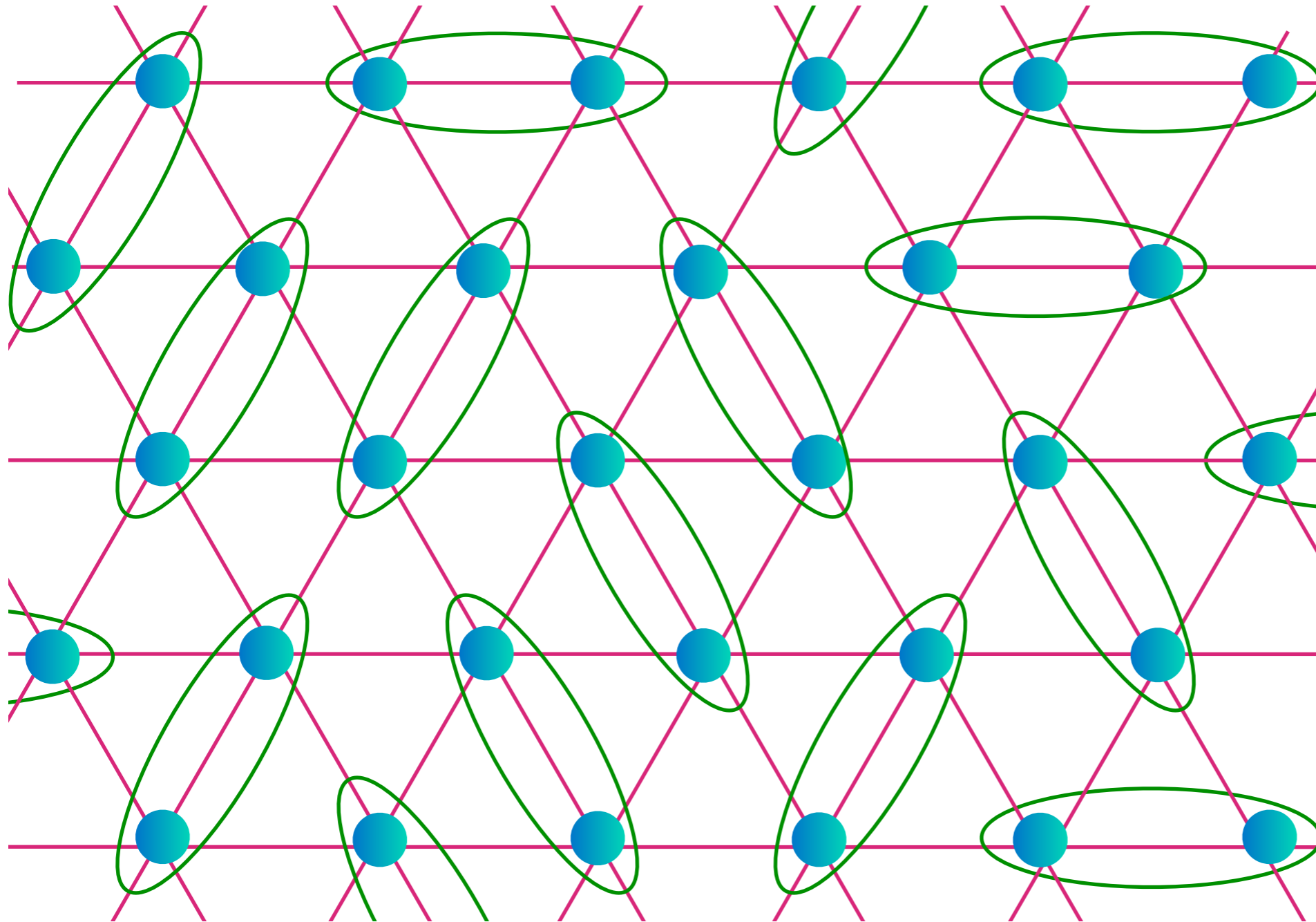


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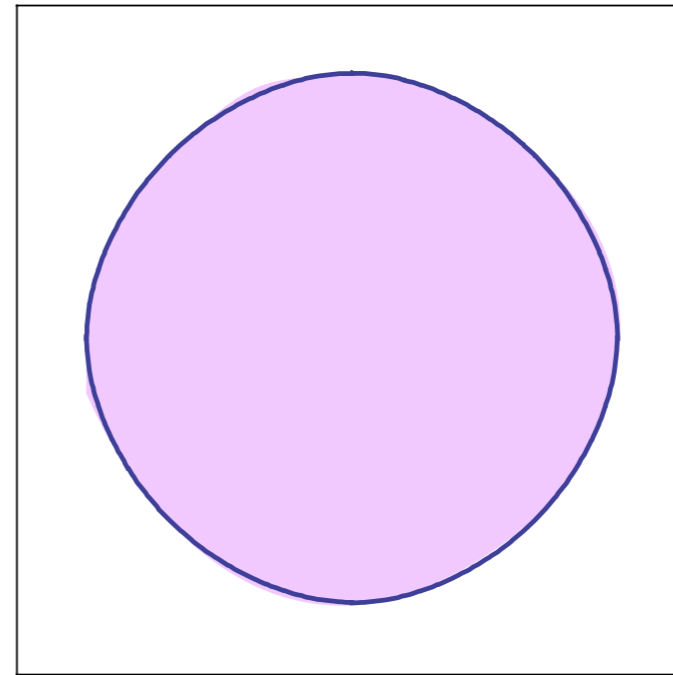
Suppressed-vortex-state is the resonating valence bond state with  $Z_2$  topological order.


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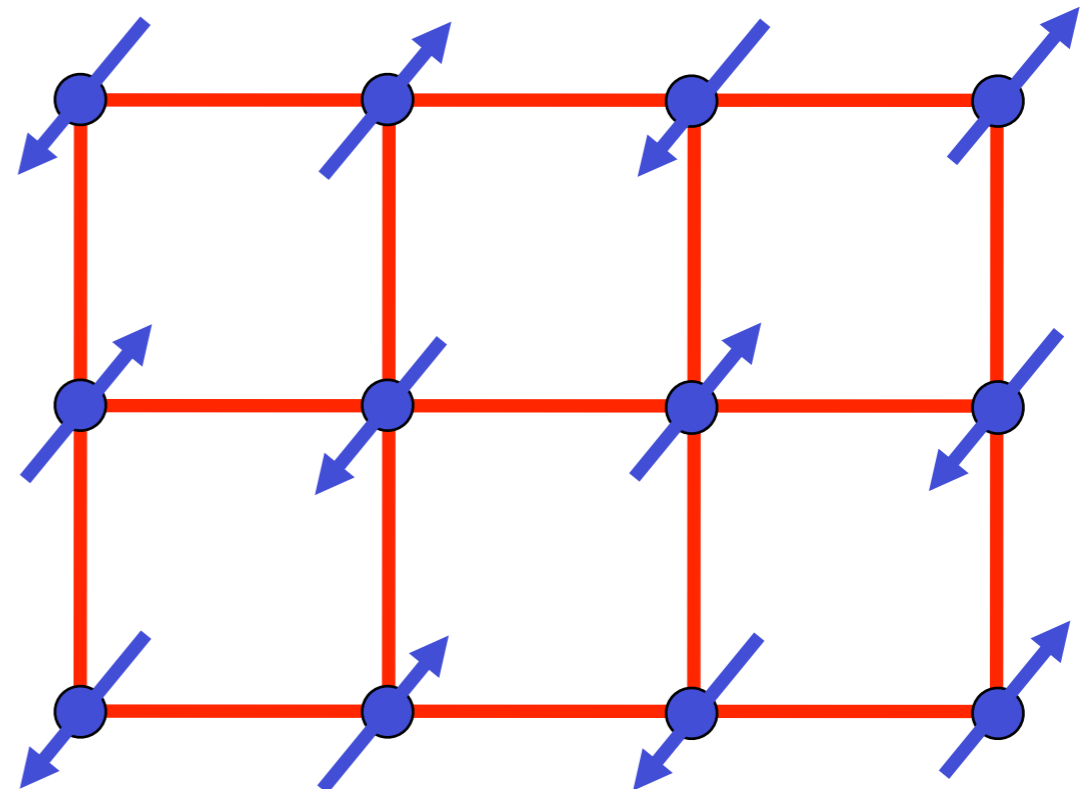


# Quantum antiferromagnetism in metals

Metal with “large”  
Fermi surface



+

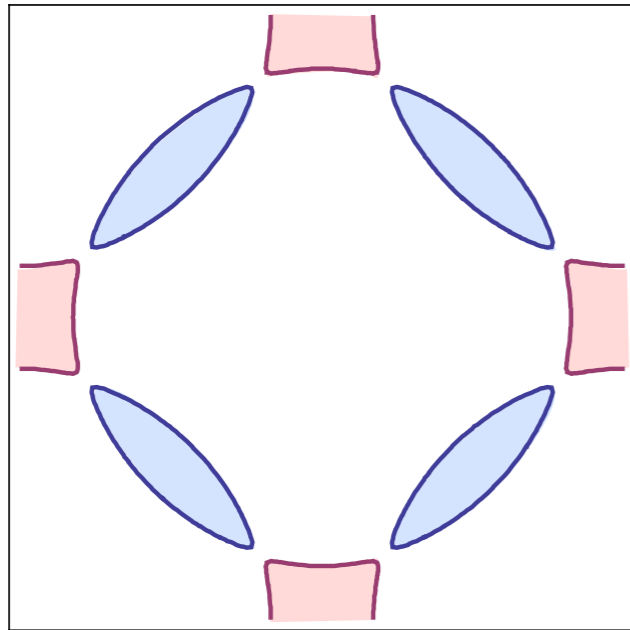


The electron spin polarization obeys

$$\langle \vec{S}(\mathbf{r}, \tau) \rangle = \vec{\varphi}(\mathbf{r}, \tau) e^{i\mathbf{K} \cdot \mathbf{r}}$$

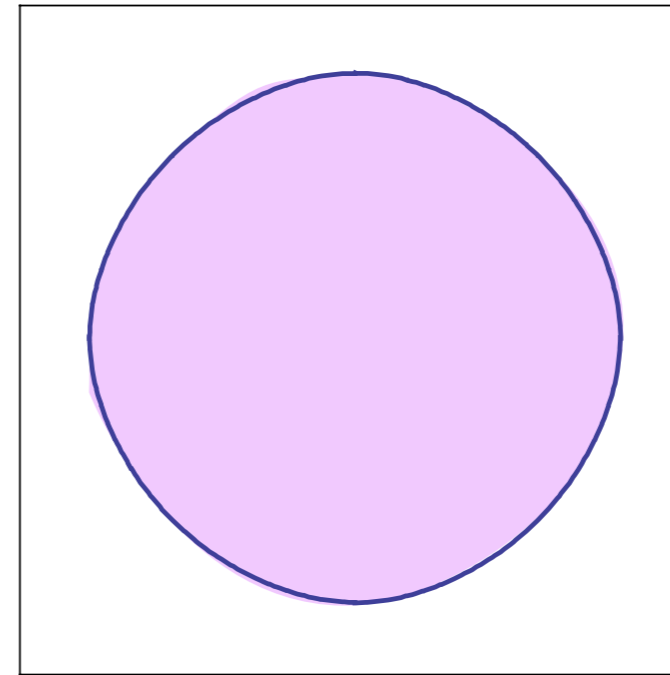
where  $\mathbf{K}$  is the ordering wavevector.

# Quantum antiferromagnetism in metals



$$\langle \vec{\varphi} \rangle \neq 0$$

Metal with  
antiferromagnetic  
order and electron  
and hole pockets



$$\langle \vec{\varphi} \rangle = 0$$

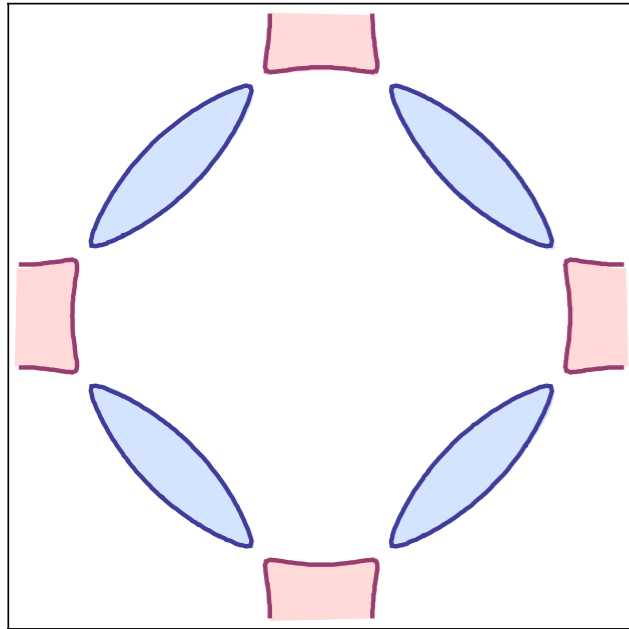
Metal with “large”  
Fermi surface

← Increasing interaction

Fermi surface reconstruction and onset of antiferromagnetism

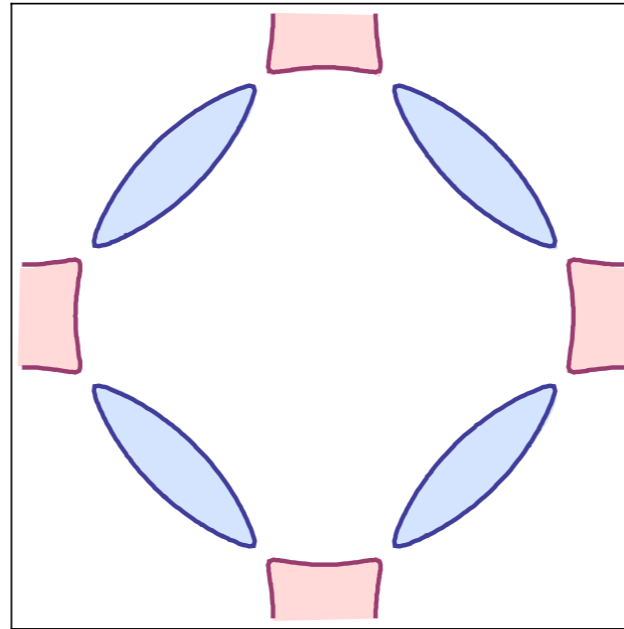
# Topological order in metals

(Senthil, S.S., Vojta (2003))



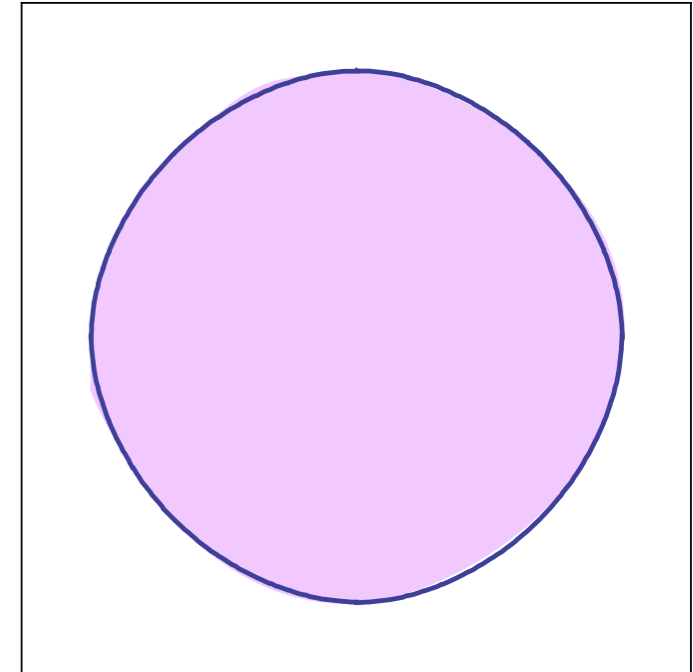
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Metal with *antiferromagnetic* order and electron and hole pockets



$$\langle \vec{\varphi} \rangle = 0$$

Metal with *topological* order and electron and hole pockets



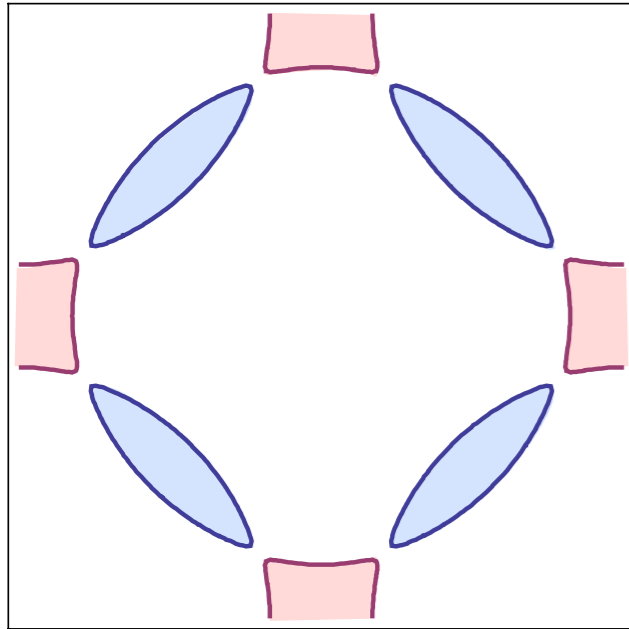
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Metal with “large” Fermi surface

Vortex suppression similar to low temperature phase of the Kosterlitz-Thouless transition.

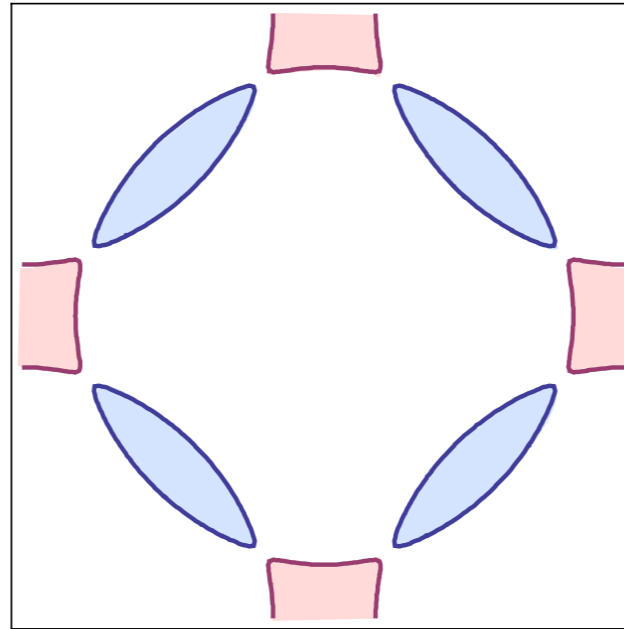
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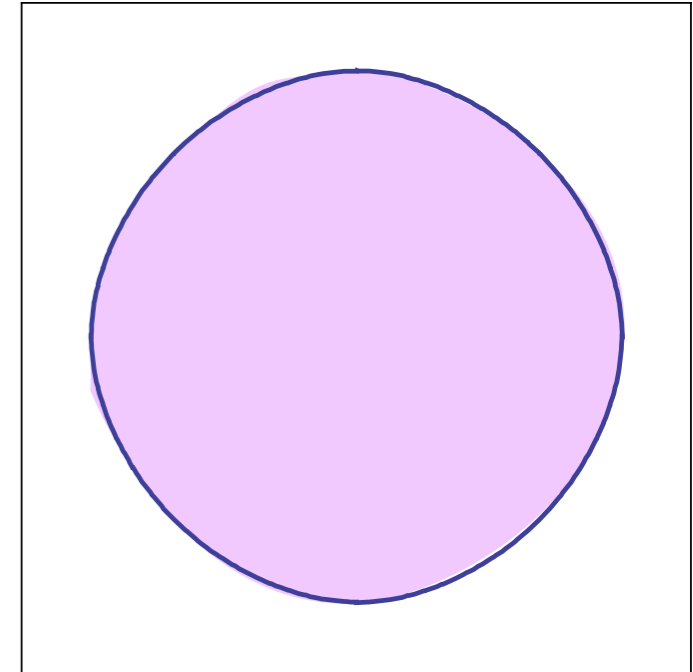
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Metal with “large”  
Fermi surface

Metal with topological order does not obey the Luttinger theorem for the volume enclosed by the Fermi surface

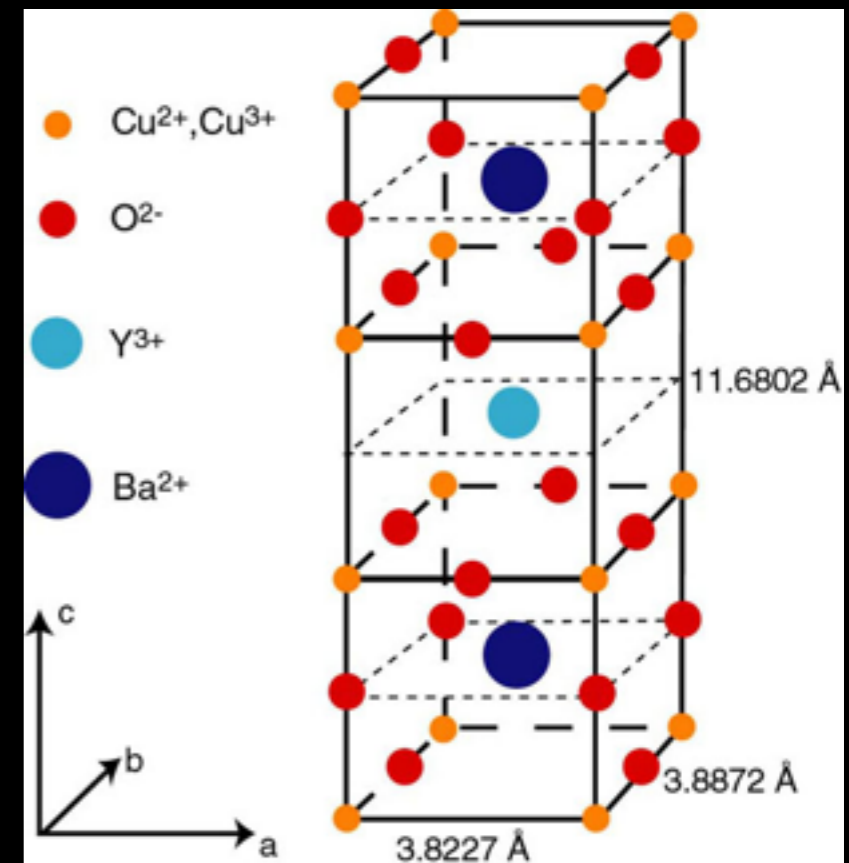
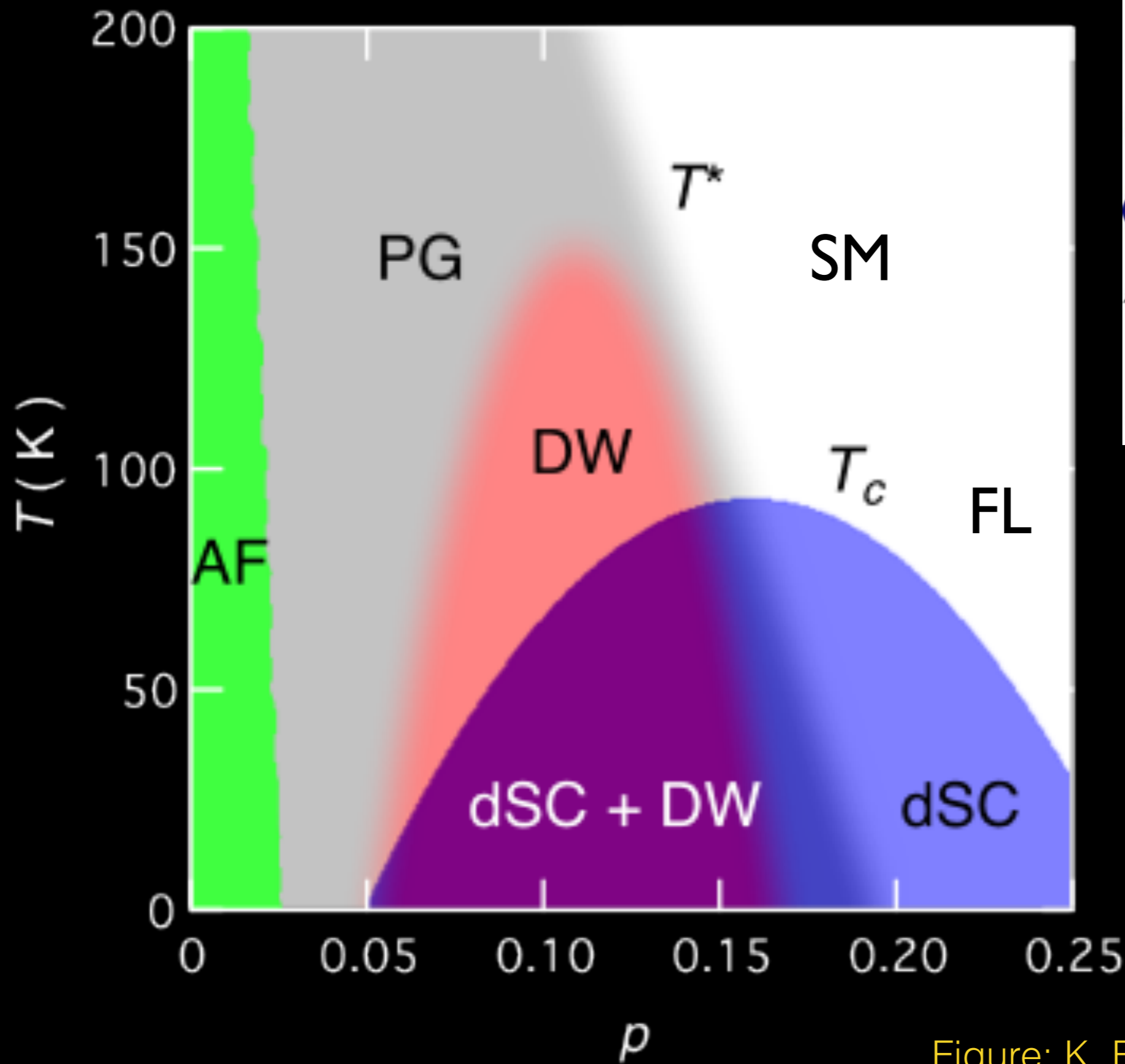
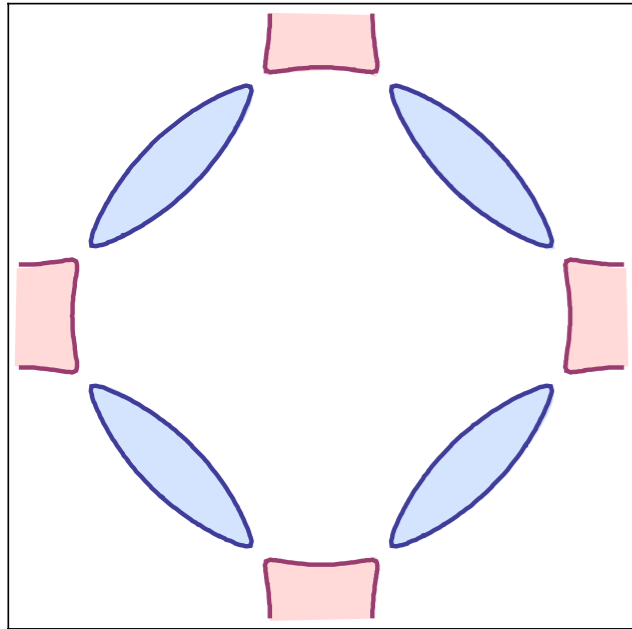


Figure: K. Fujita and J. C. Seamus Davis



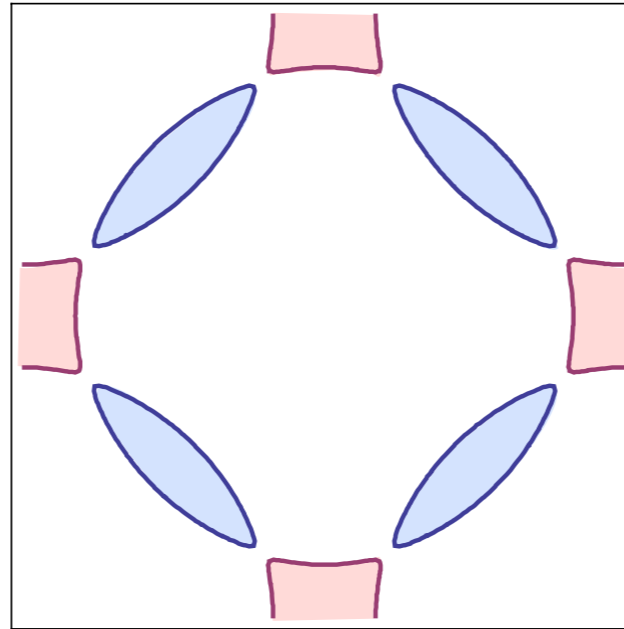
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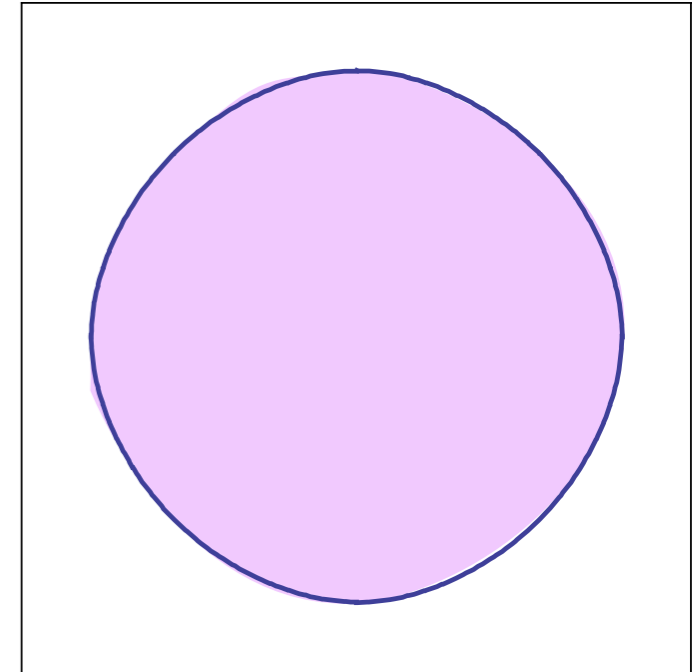
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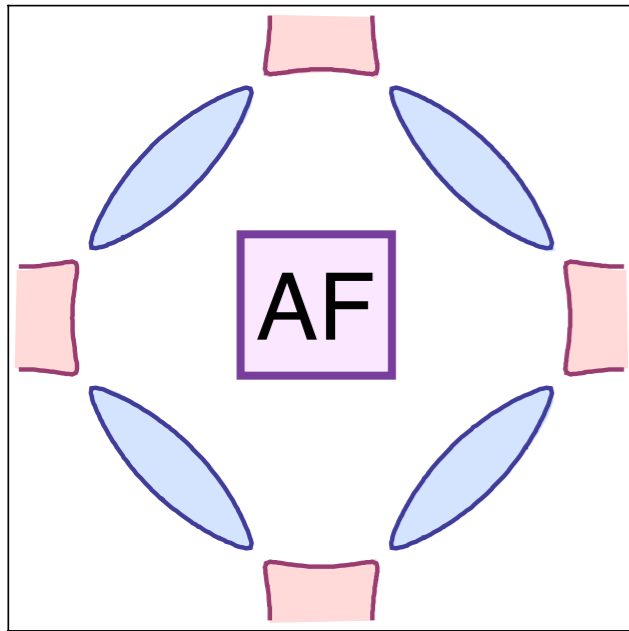
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Metal with “large”  
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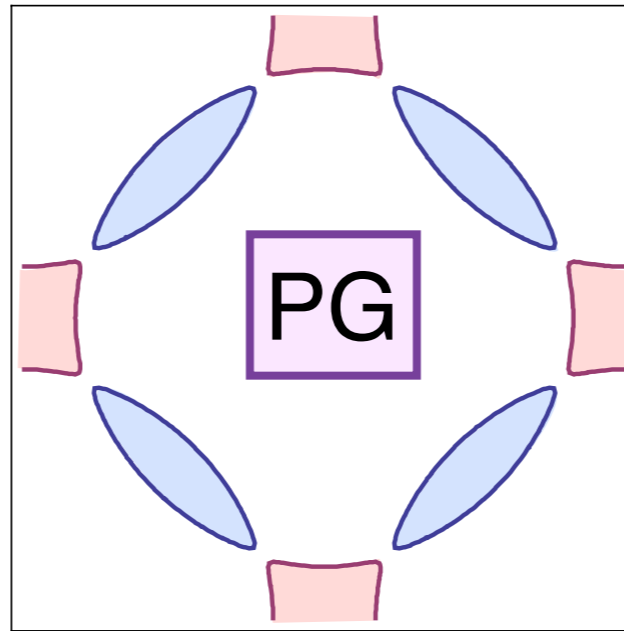
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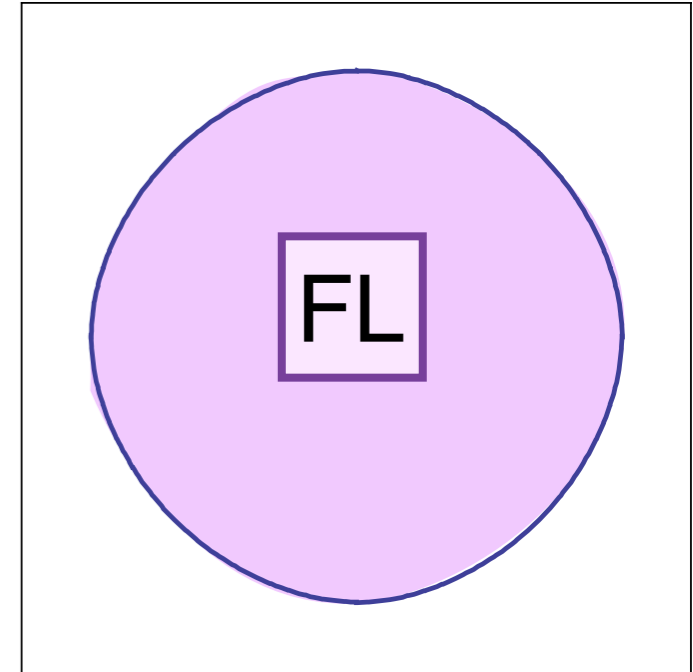
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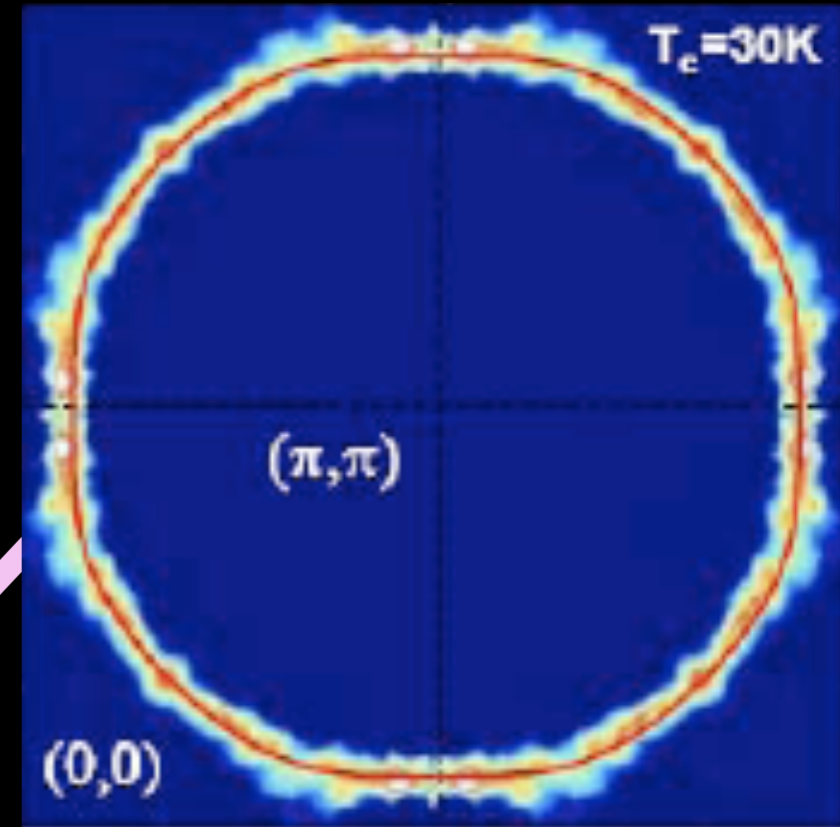
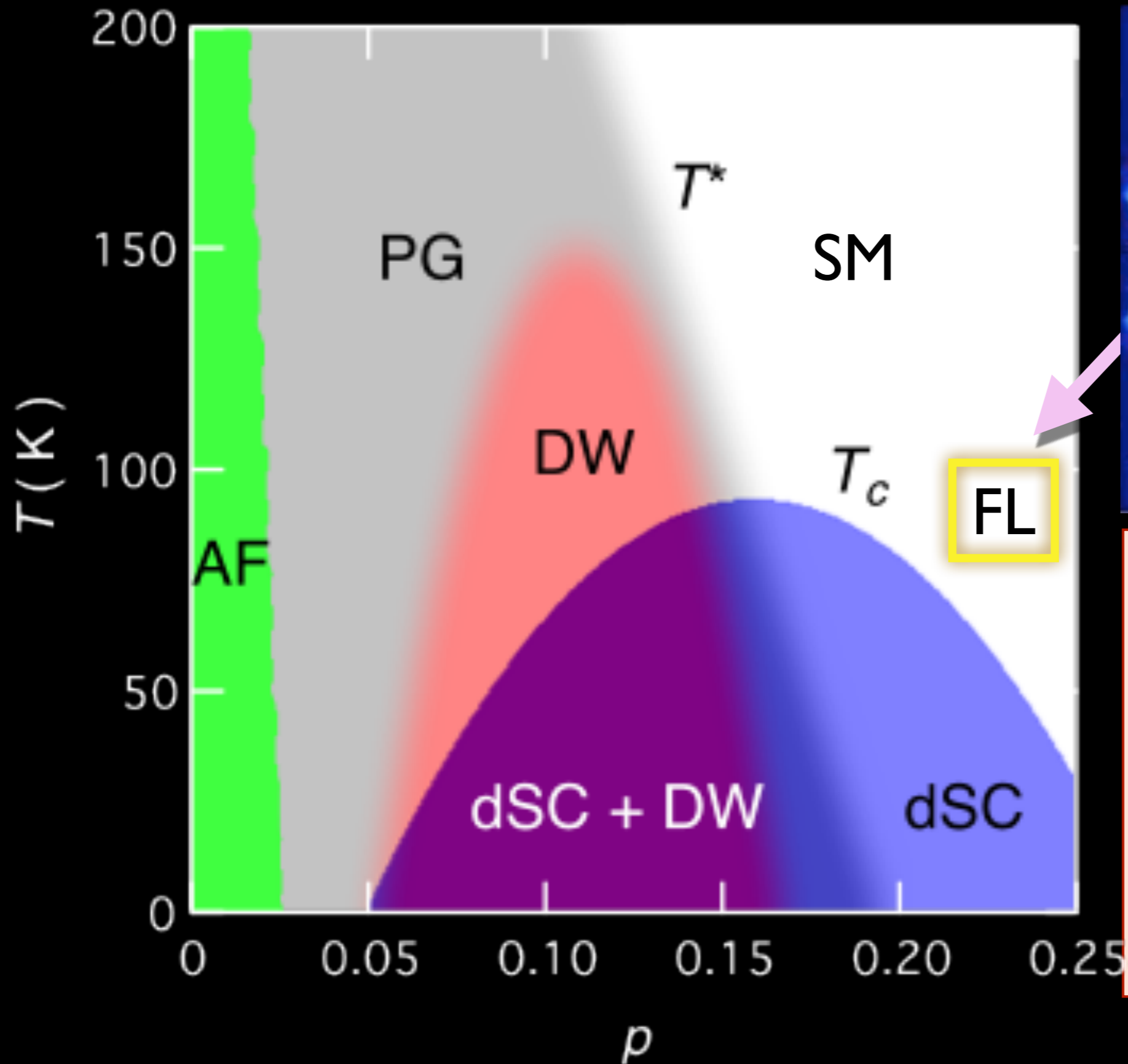


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Metal with “large”  
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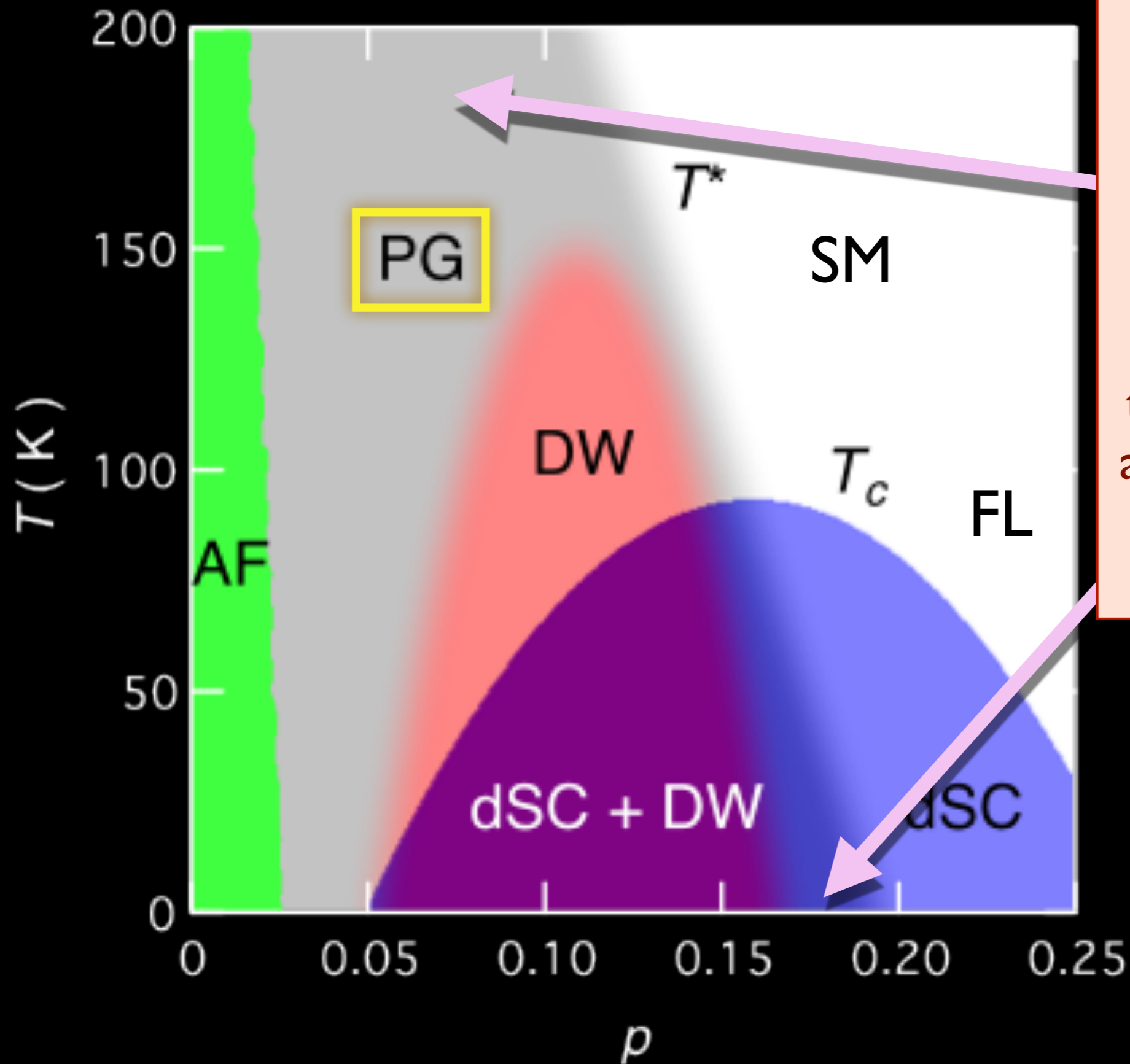


M. Platié, J. D. F. Mottershead, I. S. Elfimov, D. C. Peets, Ruixing Liang, D. A. Bonn, W. N. Hardy, S. Chiuzbaian, M. Falub, M. Shi, L. Patthey, and A. Damascelli, Phys. Rev. Lett. **95**, 077001 (2005)



A conventional metal:  
the Fermi liquid  
with Fermi  
surface of size  
 $1+p$

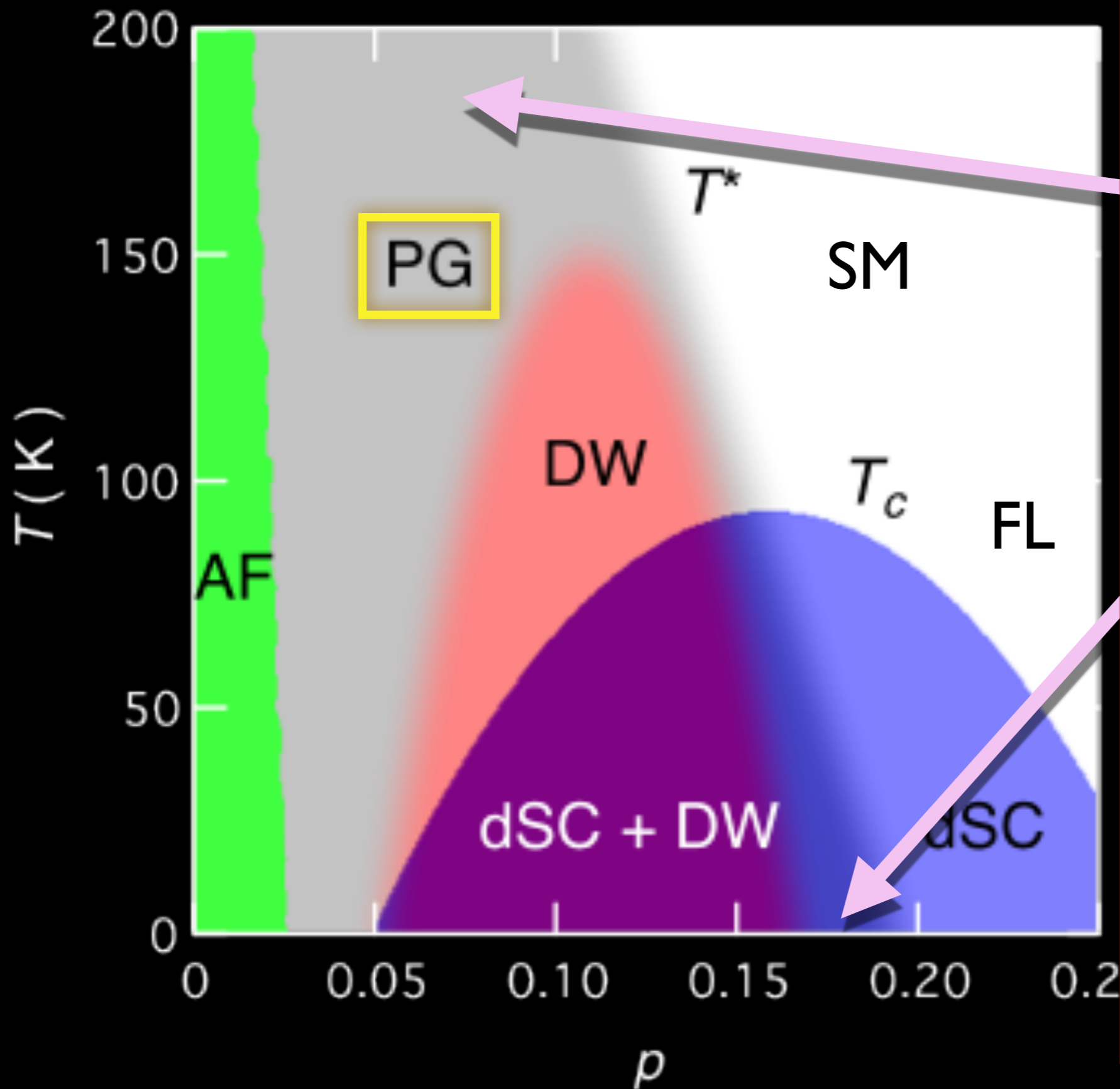
S. Badoux, W. Tabis, F. Laliberté, G. Grissonnanche, B. Vignolle, D. Vignolles, J. Béard, D.A. Bonn, W.N. Hardy, R. Liang, N. Doiron-Leyraud, L. Taillefer, and C. Proust, Nature **531**, 210 (2016).



Pseudogap  
metal

at low  $p$

Many indications that this metal behaves like a Fermi liquid, but with Fermi surface size  $p$  and *not*  $1+p$ .

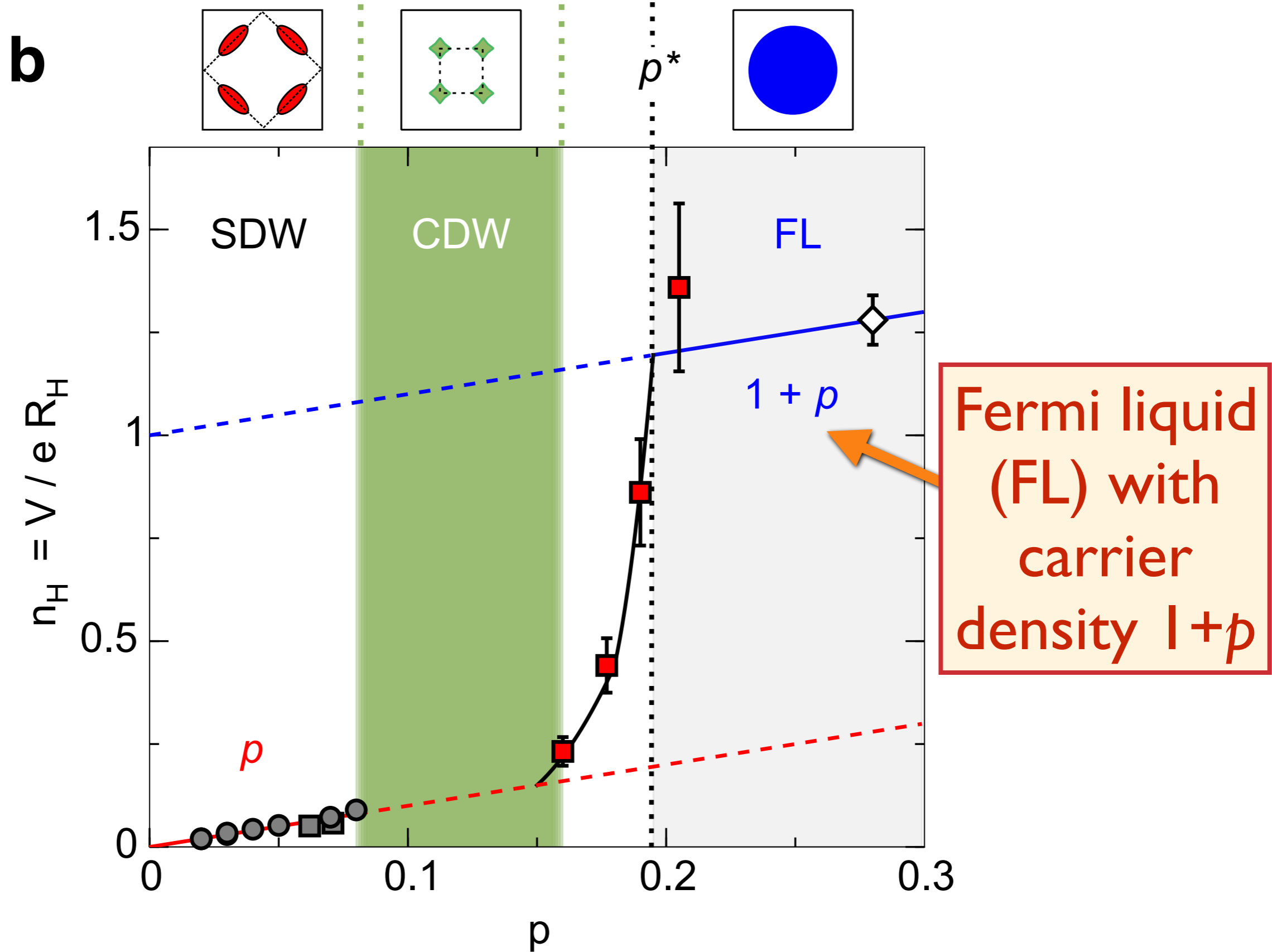


## Pseudogap metal at low $p$

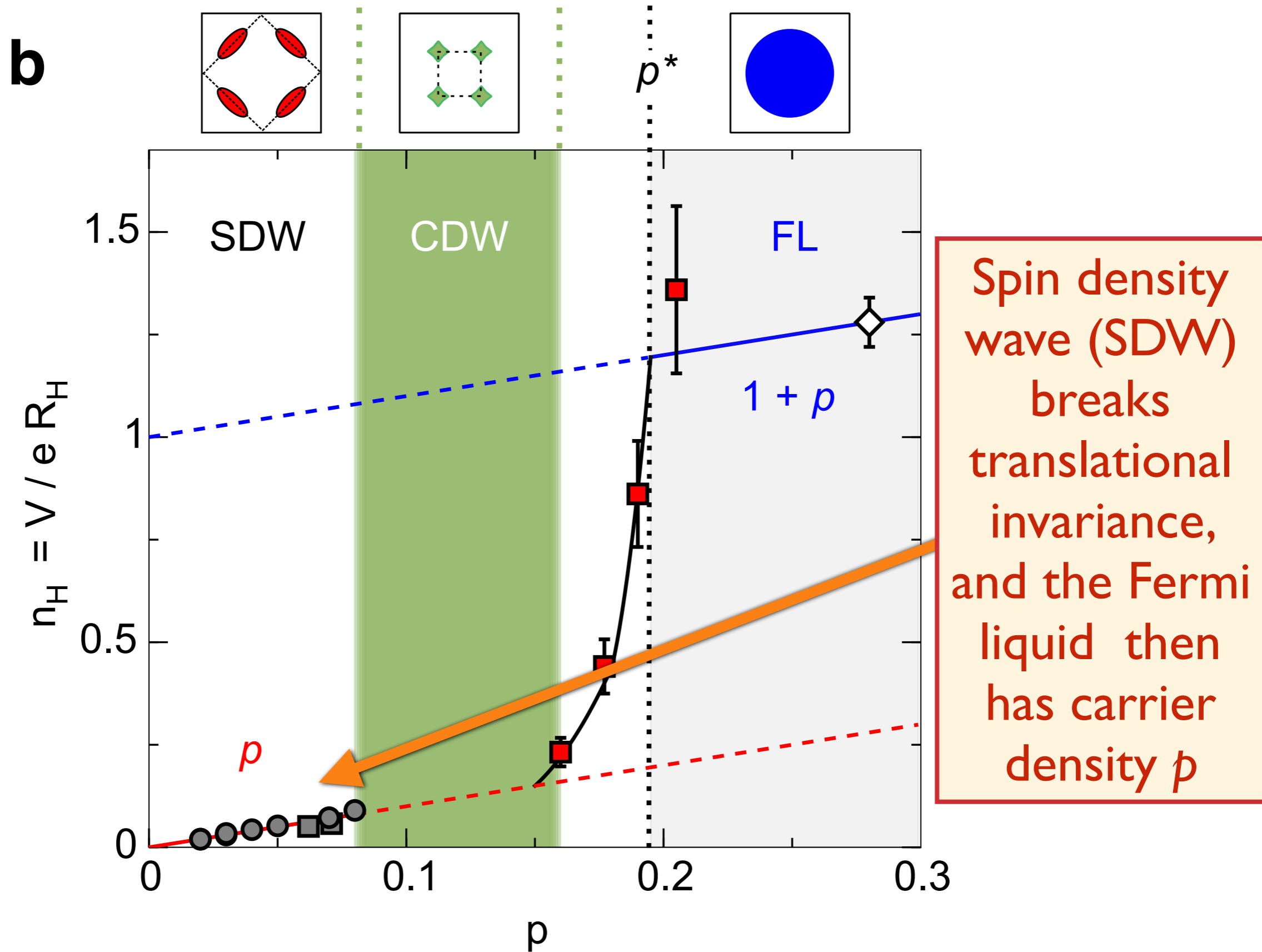
Many indications that this metal behaves like a Fermi liquid, but with Fermi surface size  $p$  and *not*  $1+p$ .

If present at  $T=0$ , a metal with a size  $p$  Fermi surface (and translational symmetry preserved) must have topological order

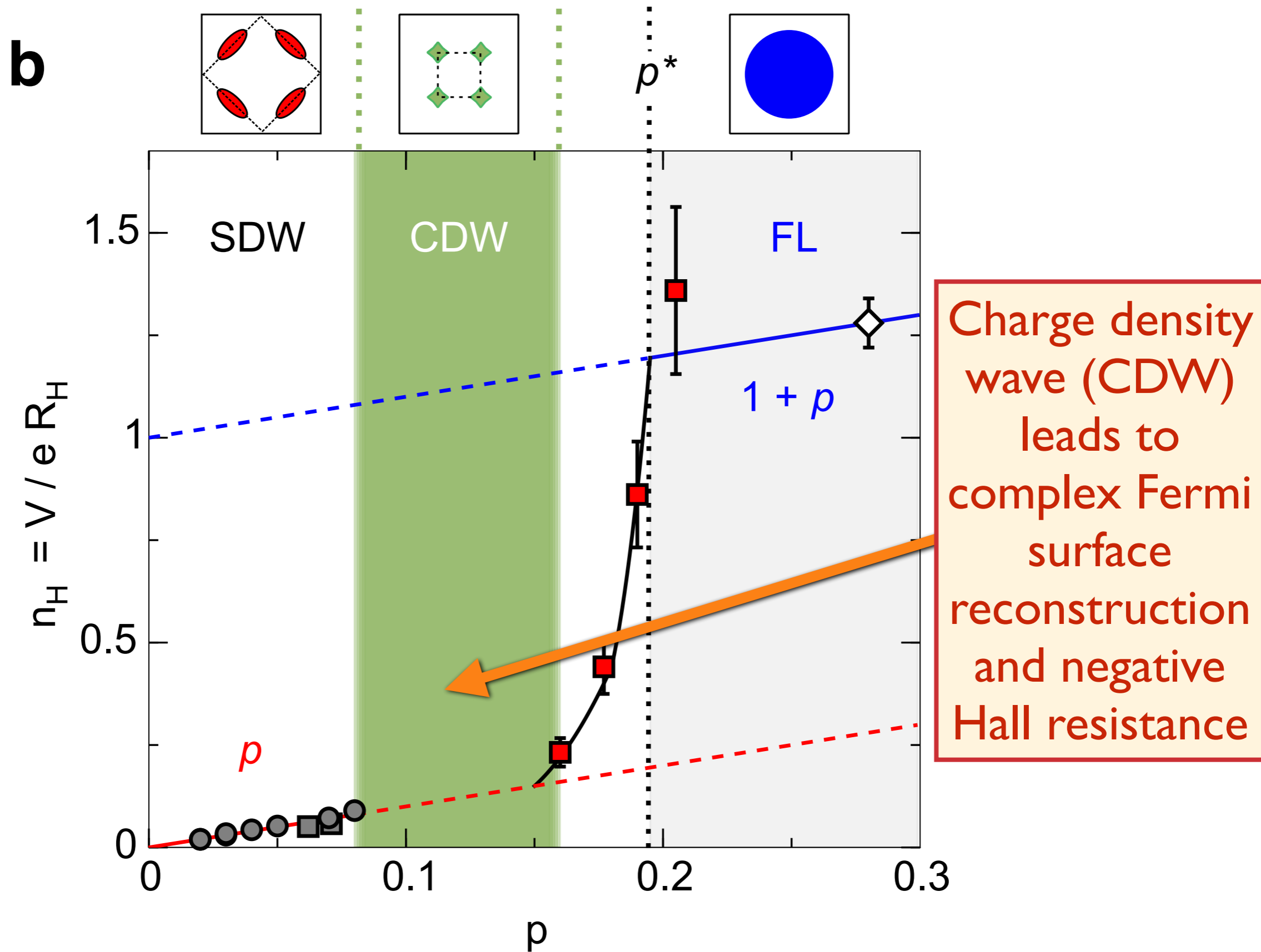
# Hall effect measurements in YBCO



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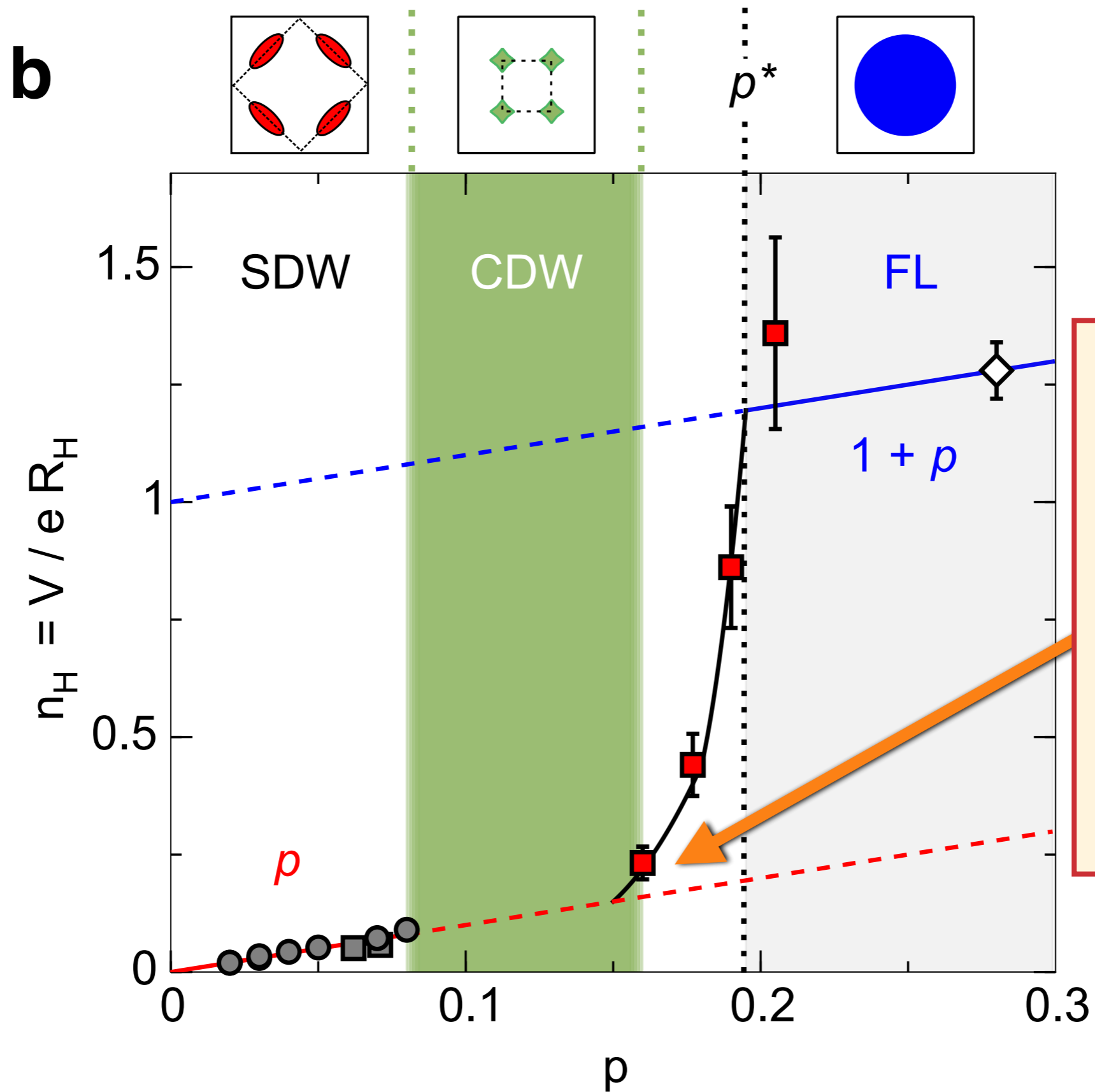


# Hall effect measurements in YBCO

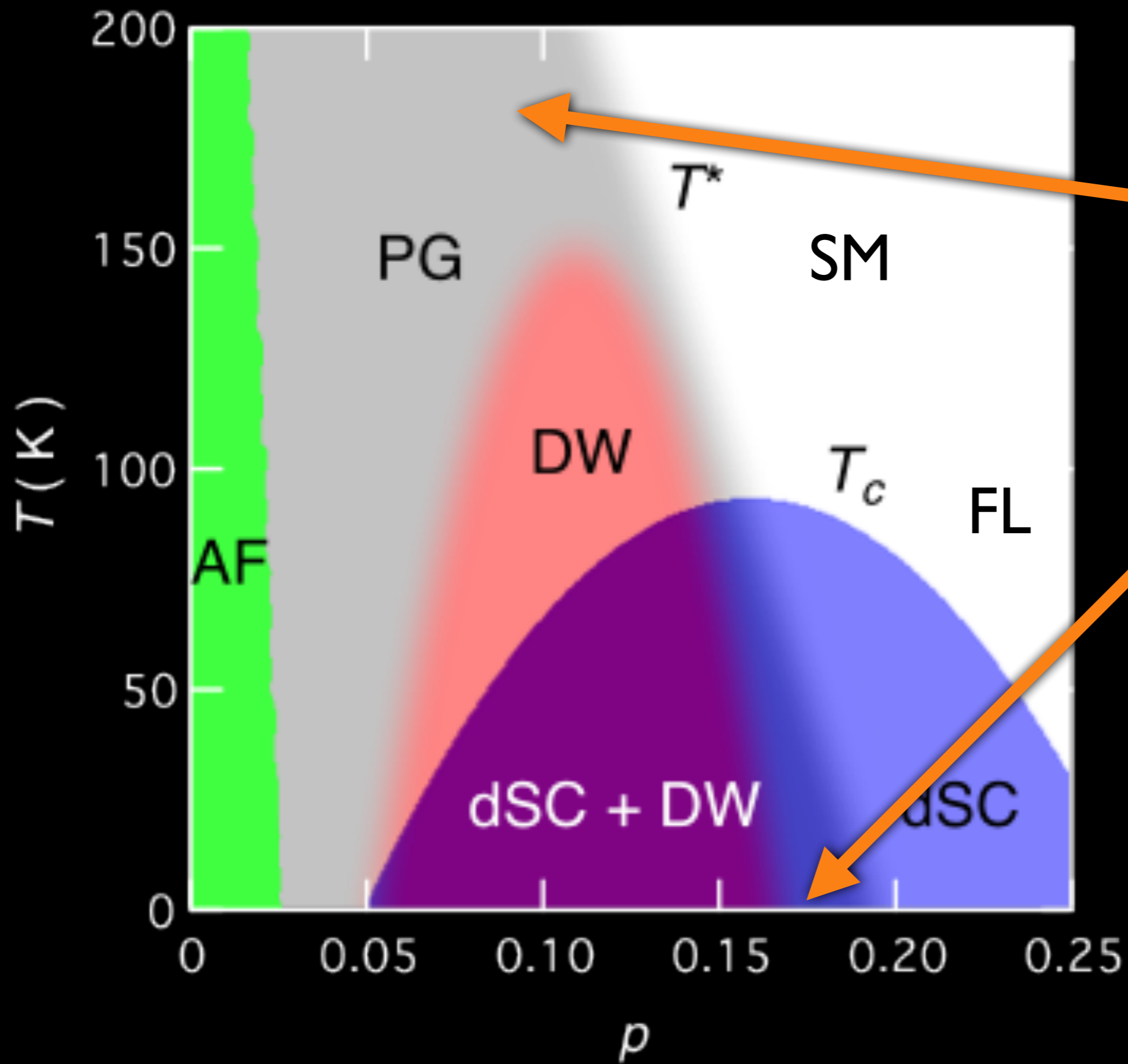




# Hall effect measurements in YBCO

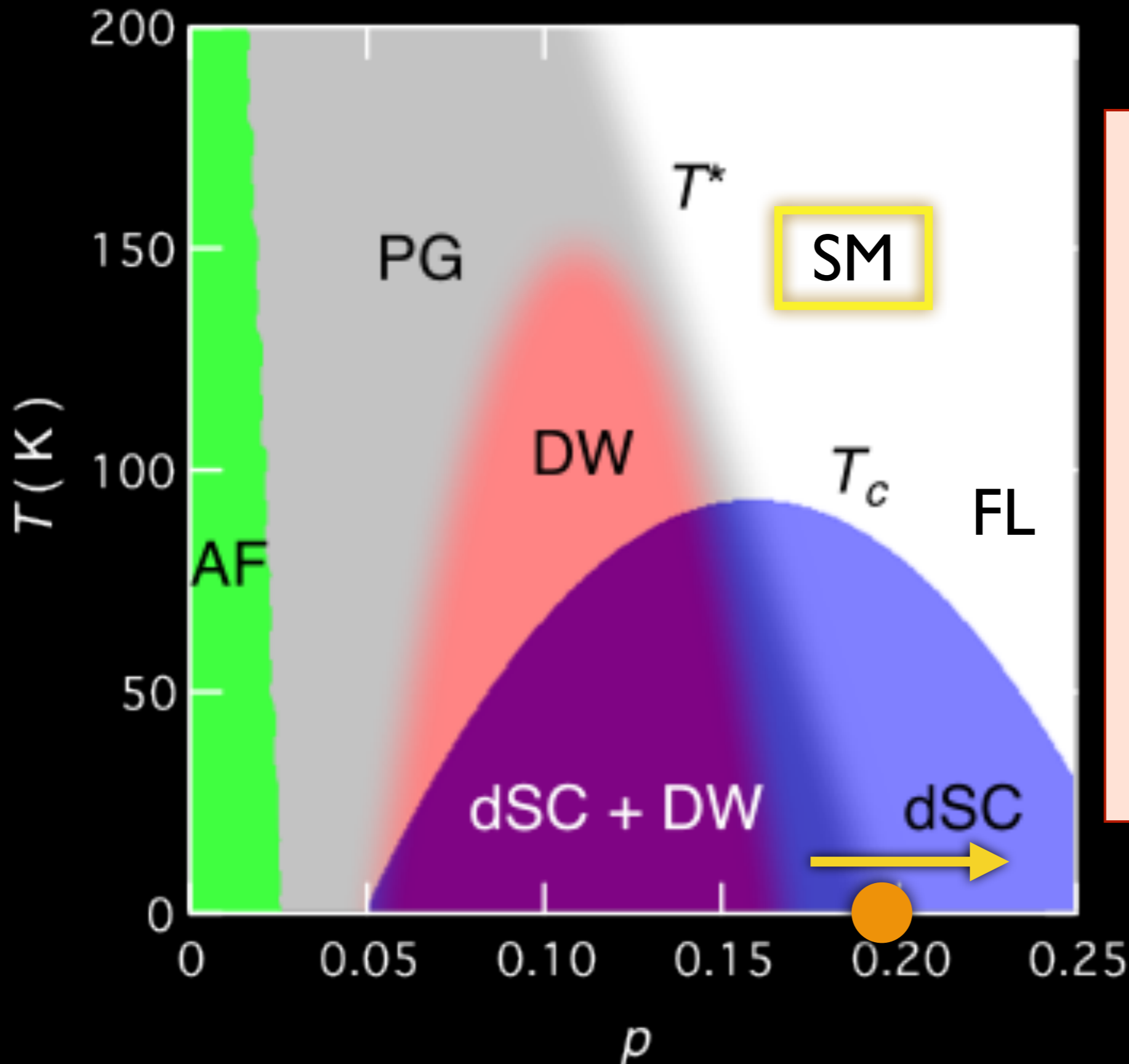


Evidence for a metal with topological order: Fermi surface of size  $p$  !



**Metal with topological order?**

S. Sachdev, M. A. Metlitski, Y. Qi, and C. Xu, PRB **80**, 155129 (2009); D. Chowdhury and S. Sachdev, PRB **91**, 115123 (2015); S. Sachdev and D. Chowdhury, arXiv:1605.03579.



Gauge theory  
for a  
topological  
phase  
transition,  
and  
for the strange  
metal (SM)