

A simple model of many-particle entanglement: how it describes black holes and superconductors

New Horizons in Physics-IPA50
Commemorating 50 years of Indian Physics Association
APS-IPA Joint Lecture
February 27, 2021

Subir Sachdev

Talk online: sachdev.physics.harvard.edu

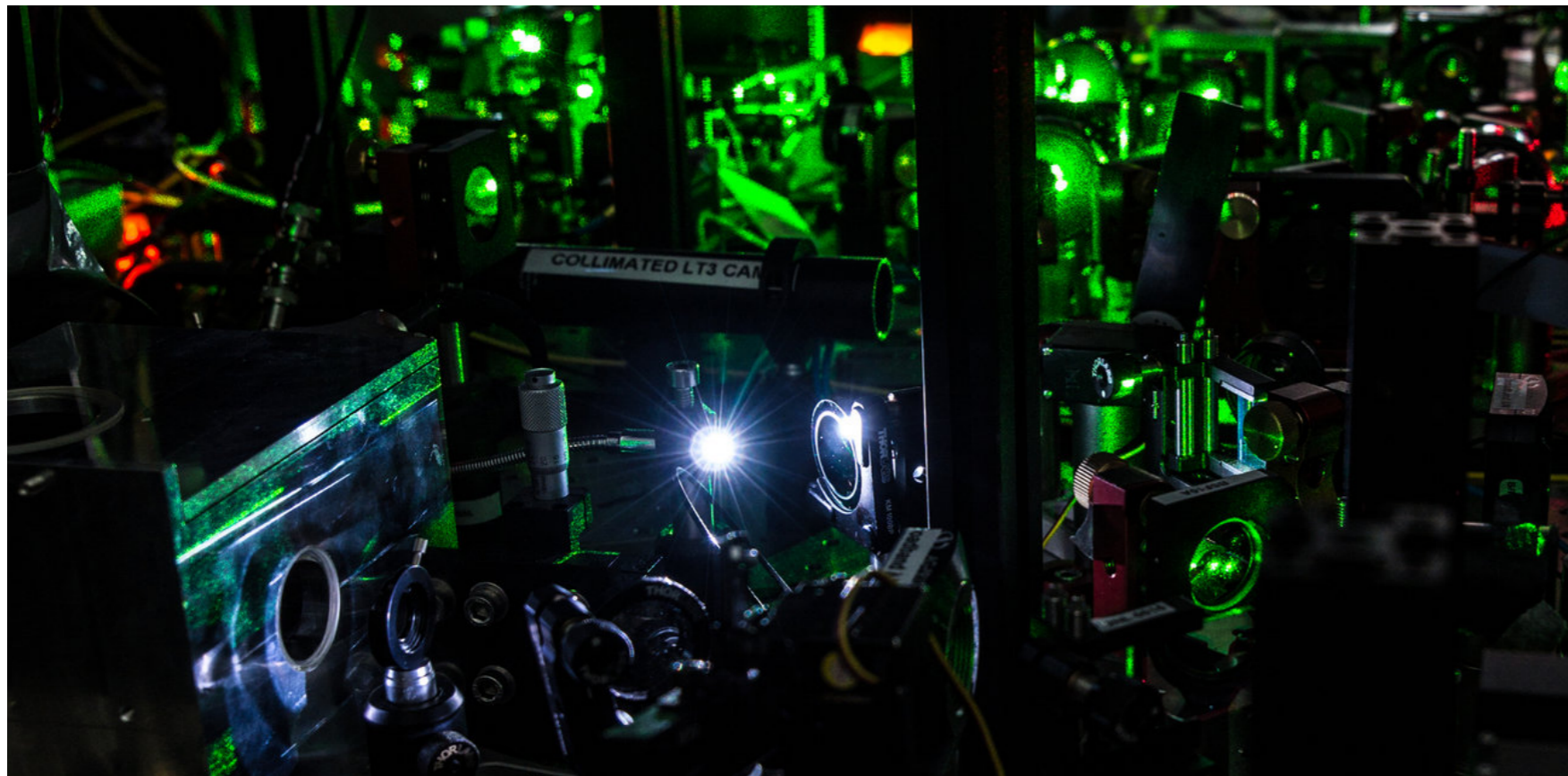


The New York Times

Sorry, Einstein. Quantum Study Suggests ‘Spooky Action’ Is Real.

By **JOHN MARKOFF** OCT. 21, 2015

In a landmark study, scientists at Delft University of Technology in the Netherlands reported that they had conducted an experiment that they say proved one of the most fundamental claims of quantum theory — that objects separated by great distance can instantaneously affect each other’s behavior.



Part of the laboratory setup for an experiment at Delft University of Technology, in which two diamonds were set 1.3 kilometers apart, entangled and then shared information.

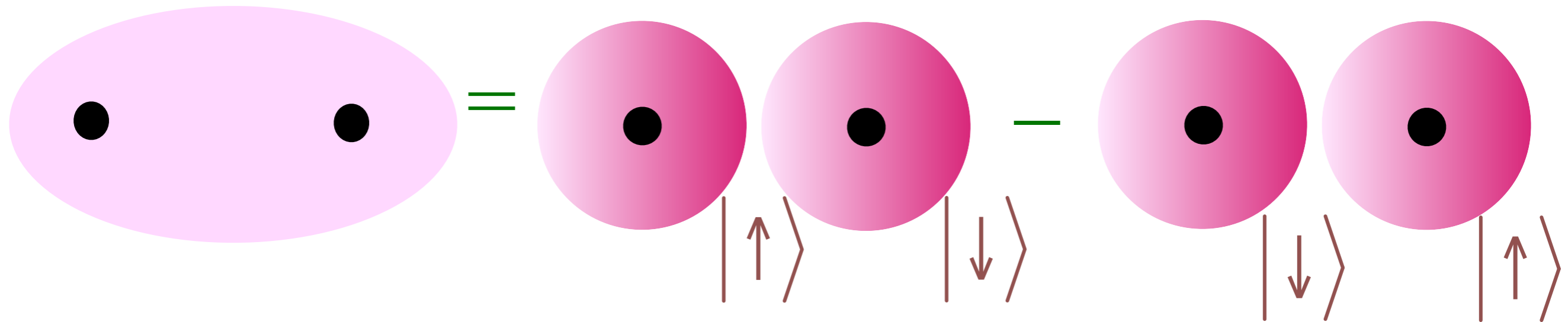
Quantum entanglement

Principles of Quantum Mechanics: II. Quantum Entanglement

Quantum Entanglement: quantum superposition with more than one particle



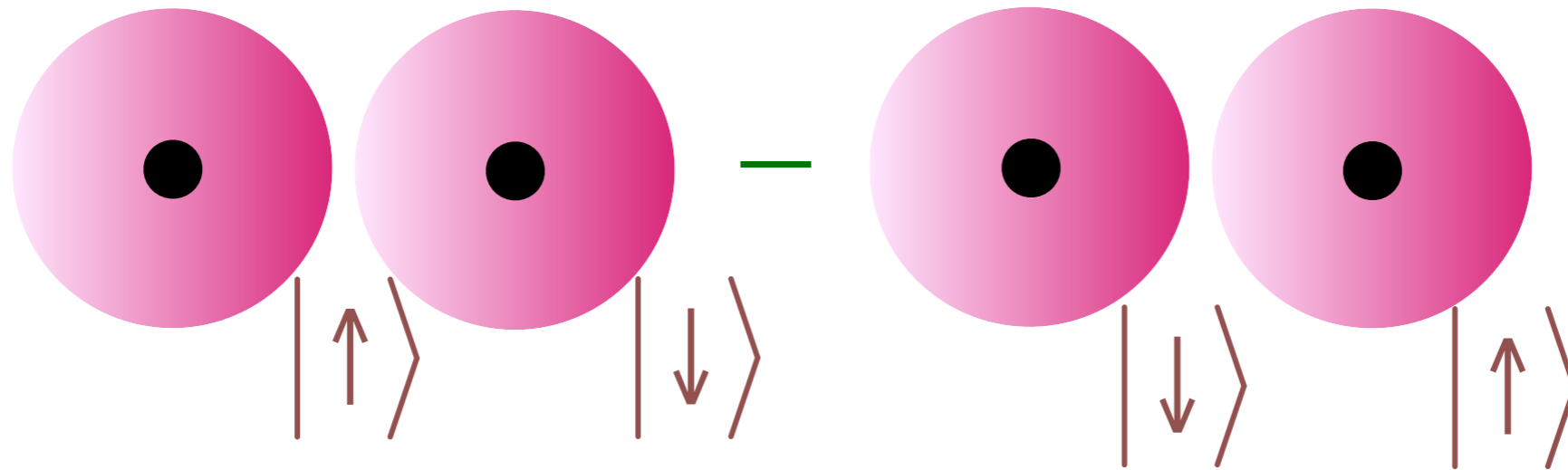
Hydrogen molecule:



$$= \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

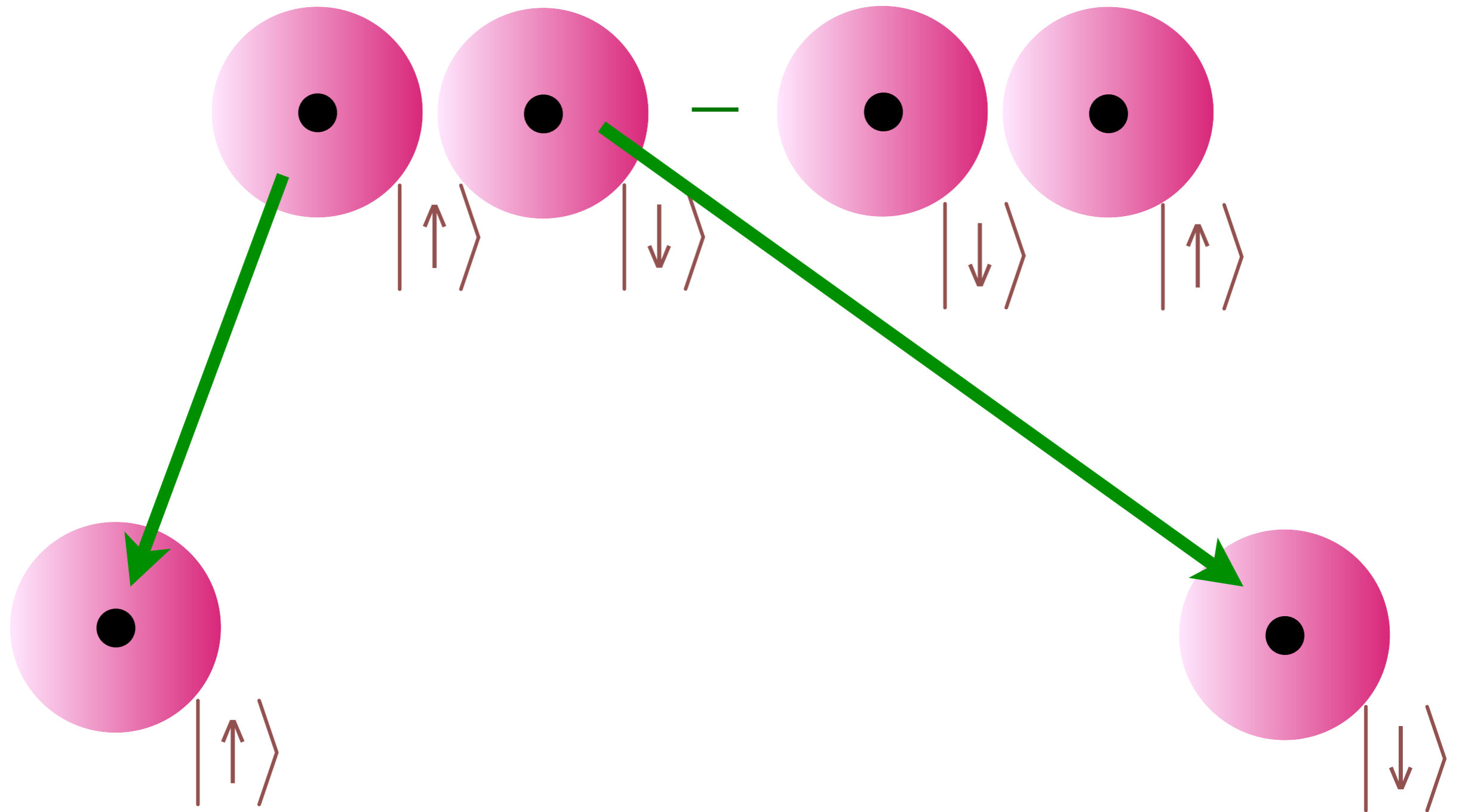
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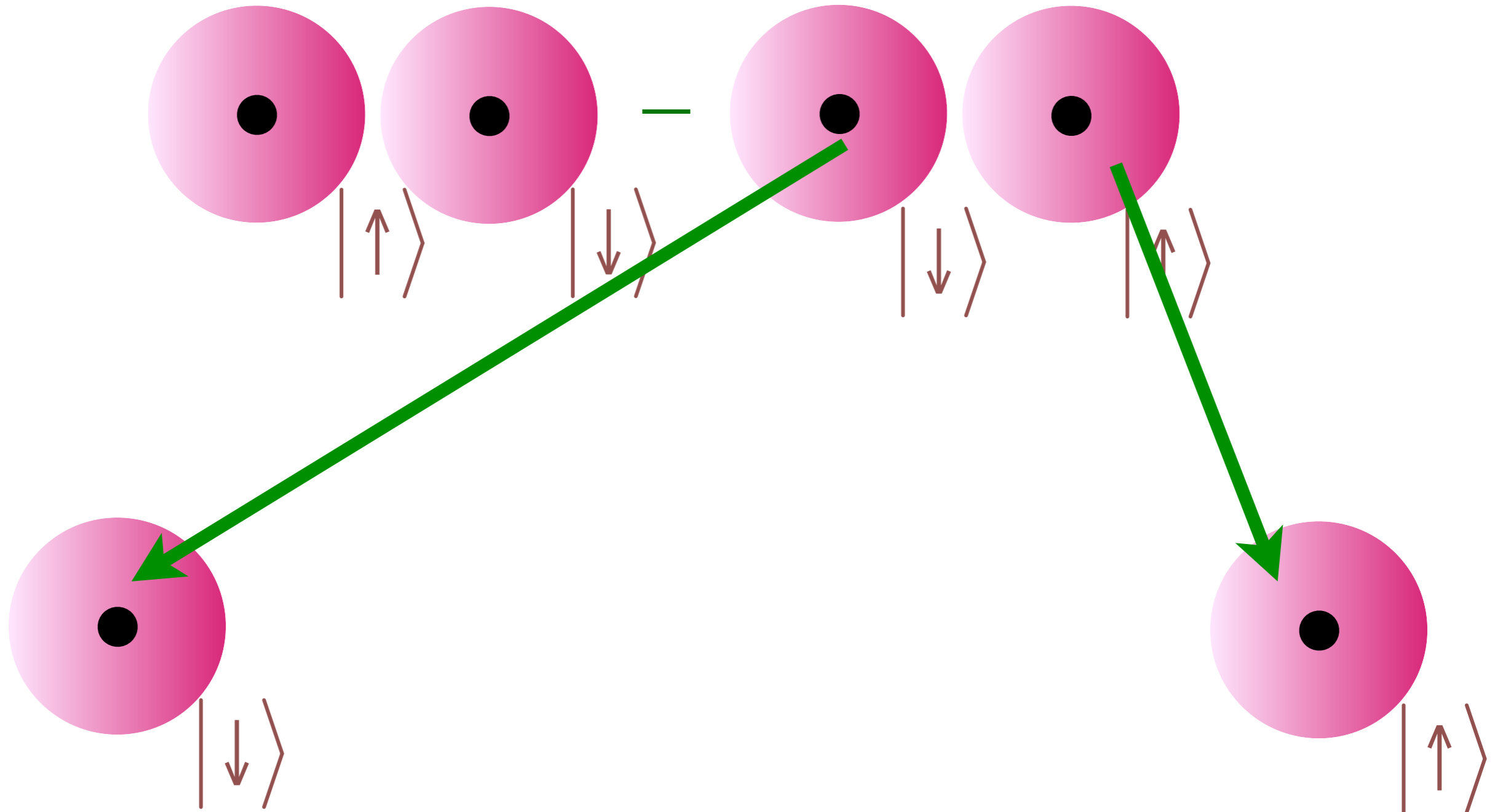
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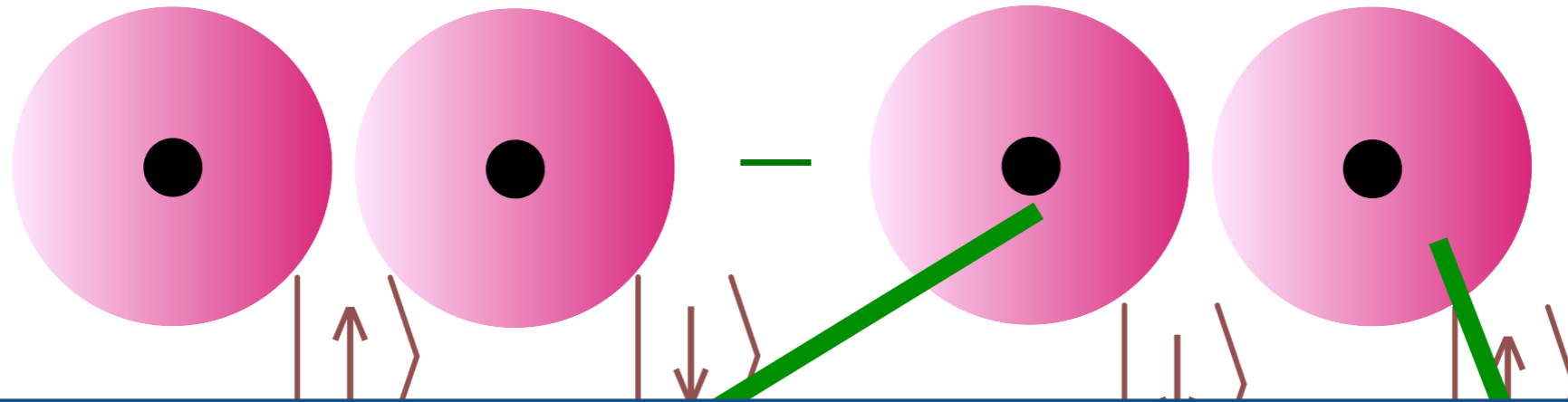
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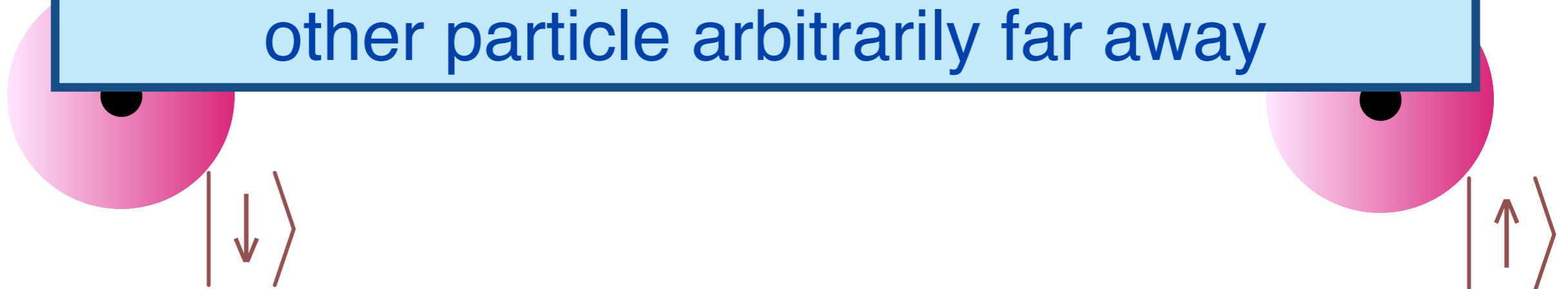


Principles of Quantum Mechanics: II. Quantum Entanglement

Quantum Entanglement: quantum superposition with more than one particle



Einstein-Podolsky-Rosen “paradox” (1935):
Measurement of one particle
instantaneously determines the state of the
other particle arbitrarily far away

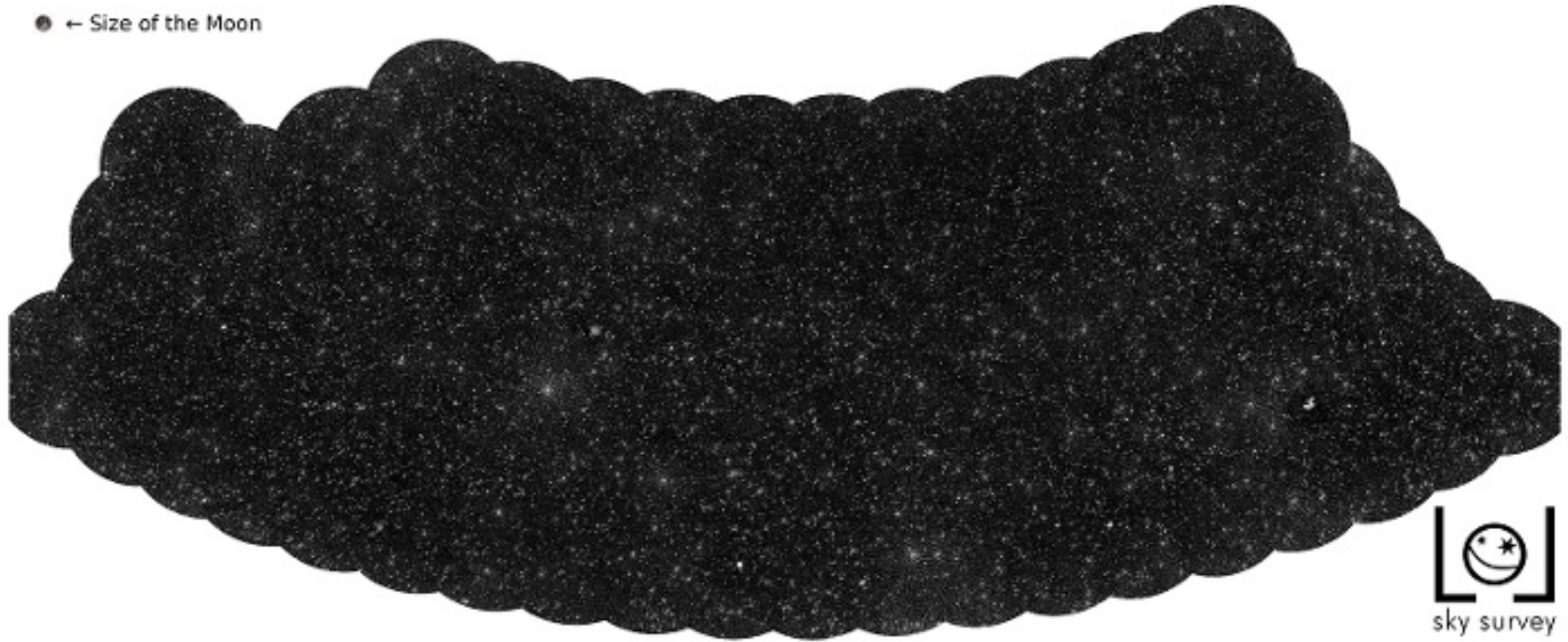


Quantum entanglement

**Quantum
entanglement**

**Black
holes**

● ← Size of the Moon



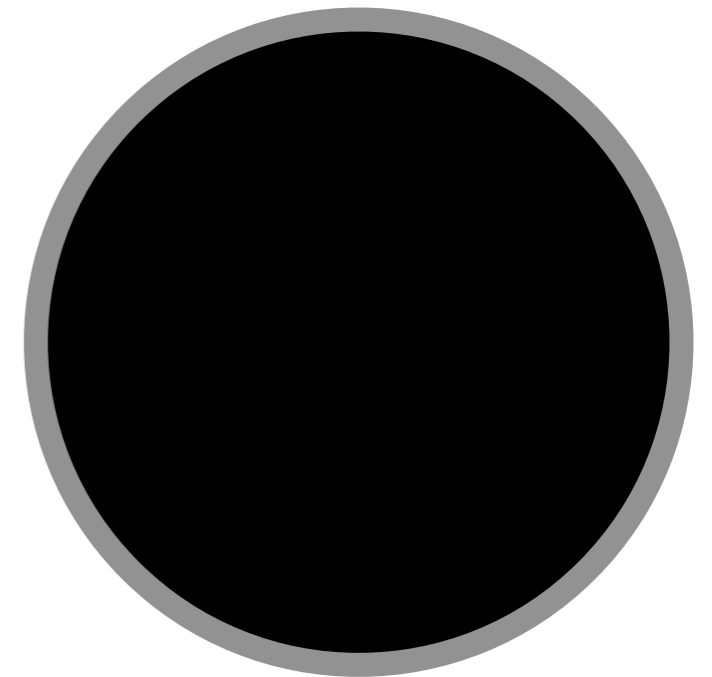
**LOFAR LBA Sky Survey showing 25000 supermassive
black holes on 4% of the northern sky.
Obtained by 52 radio telescopes across Europe**

de Gasperin et al. (2021)

Black Holes

Objects so dense that light is gravitationally bound to them.

In Einstein's theory, the region inside the black hole **horizon** is disconnected from the rest of the universe.



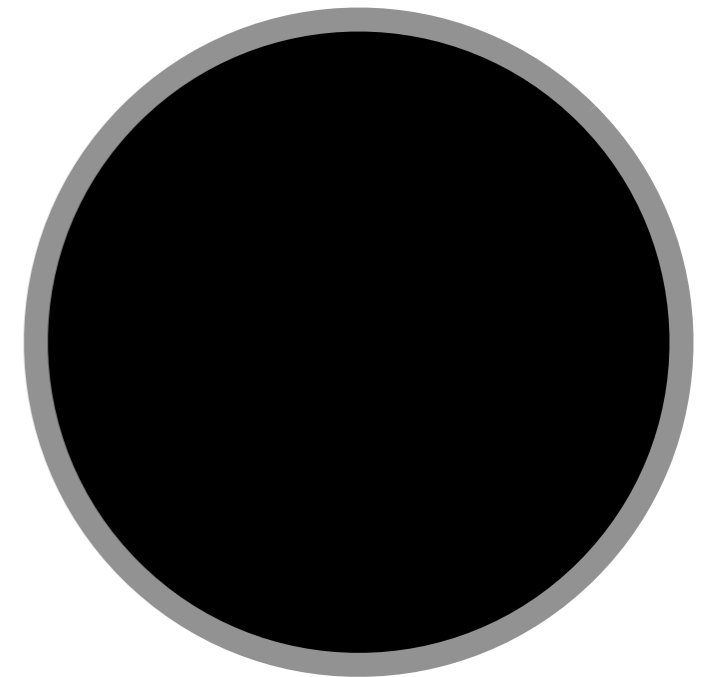
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G Newton's constant, c velocity of light, M mass of black hole

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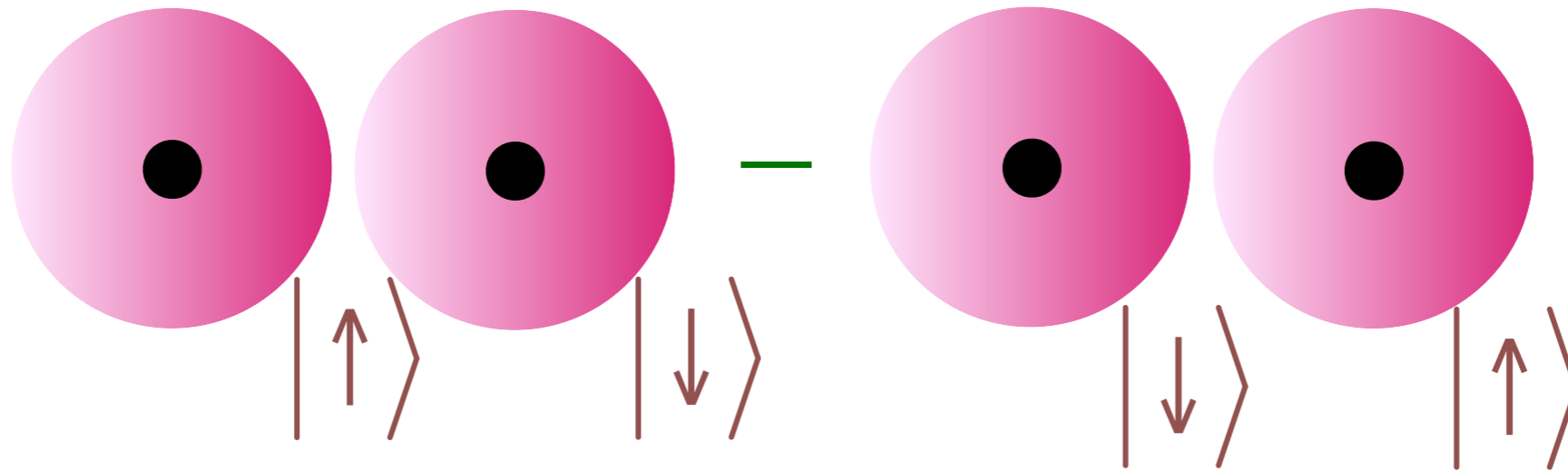
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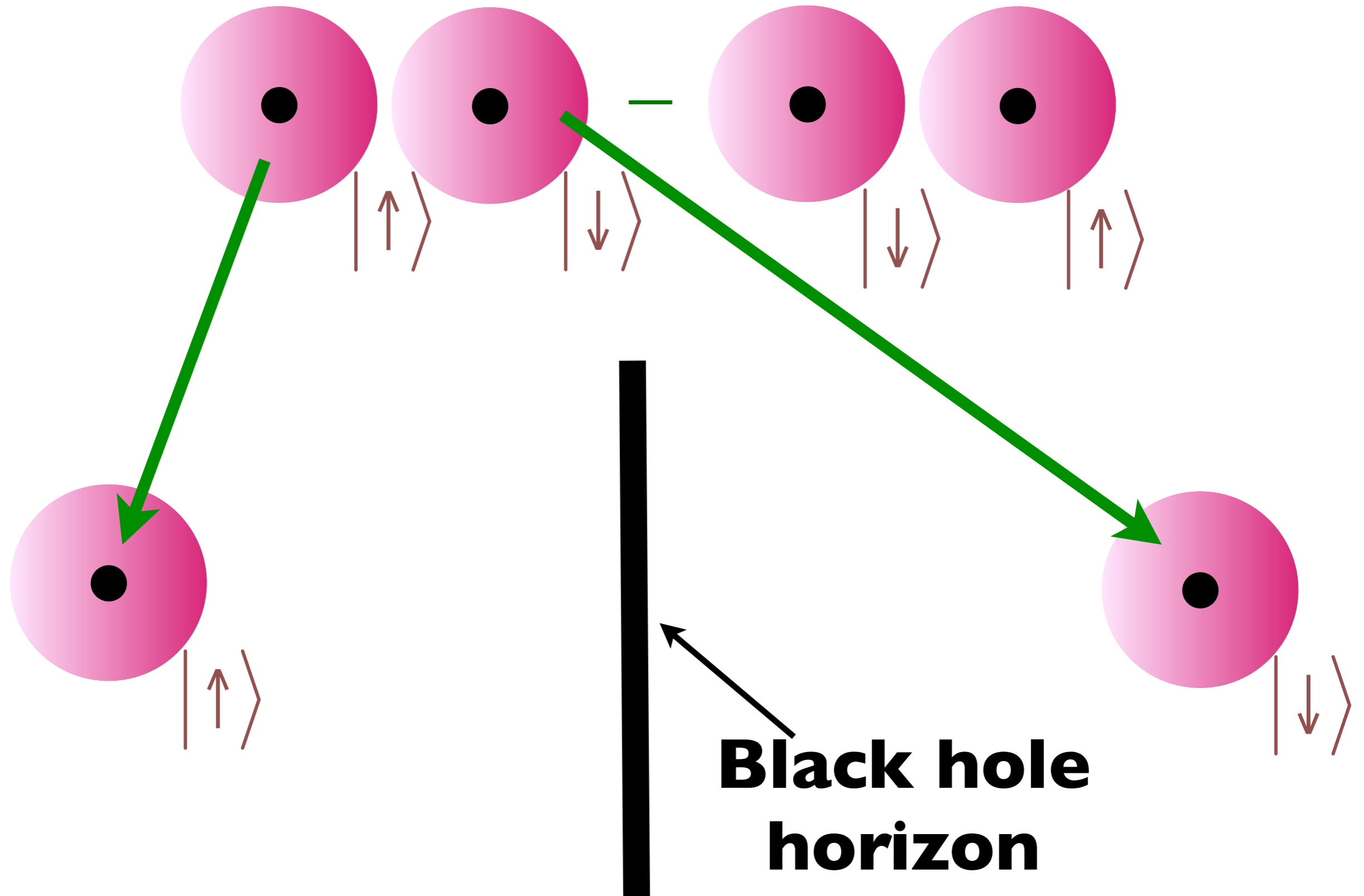
$$\text{Horizon radius } R = \frac{2GM}{c^2}$$

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For $M = \text{earth's mass}$, $R \approx 9 \text{ mm!}$

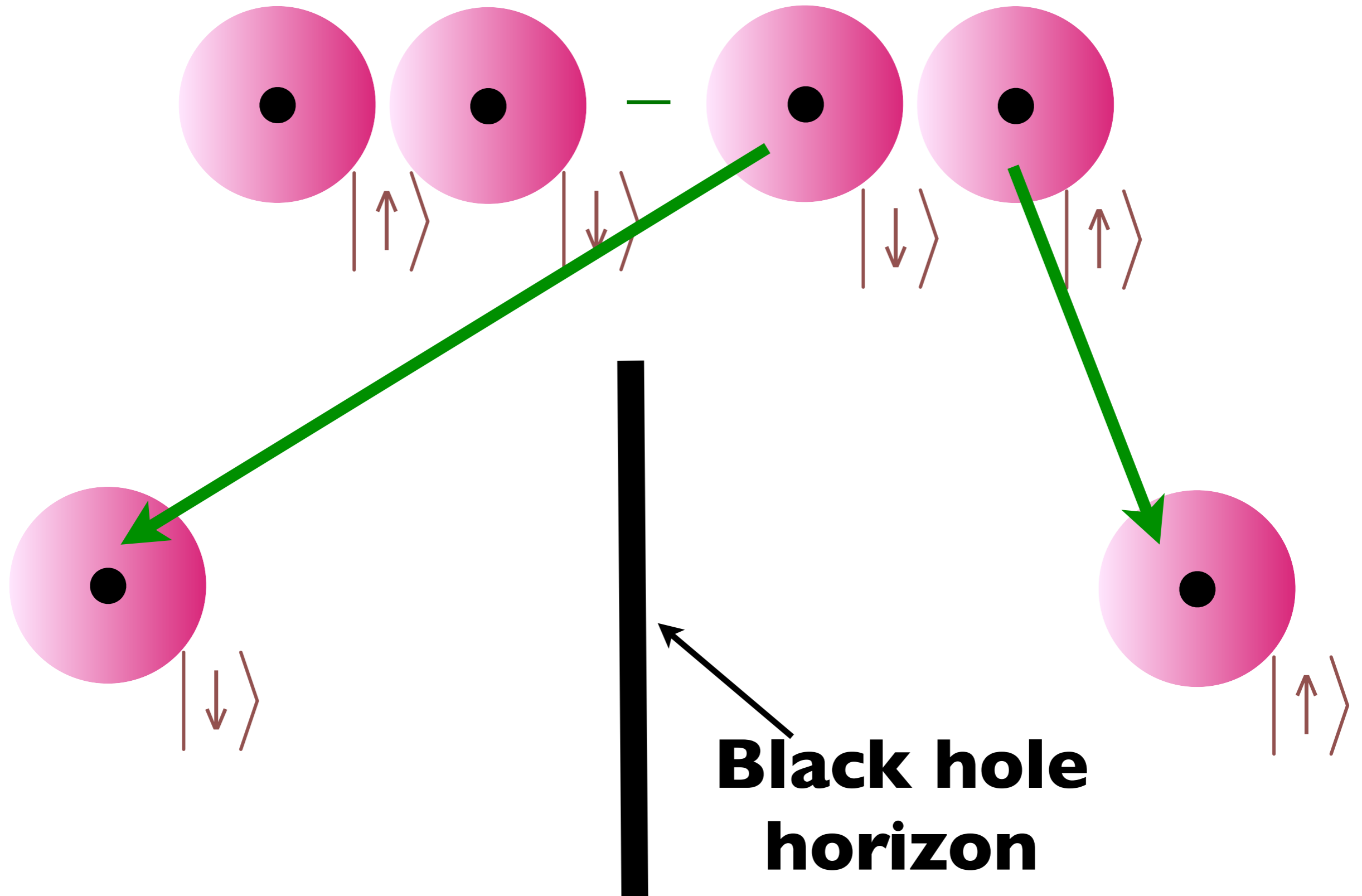
Quantum Entanglement across a black hole horizon



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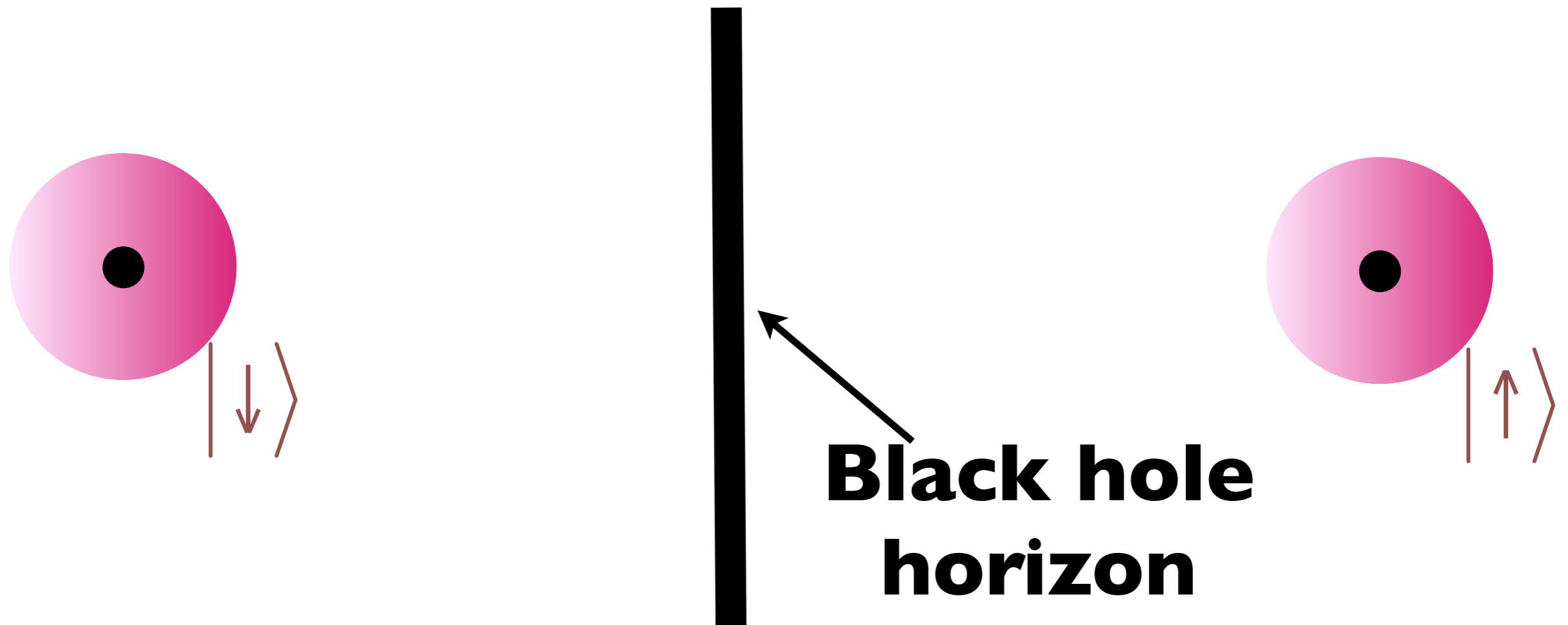


Quantum Entanglement across a black hole horizon



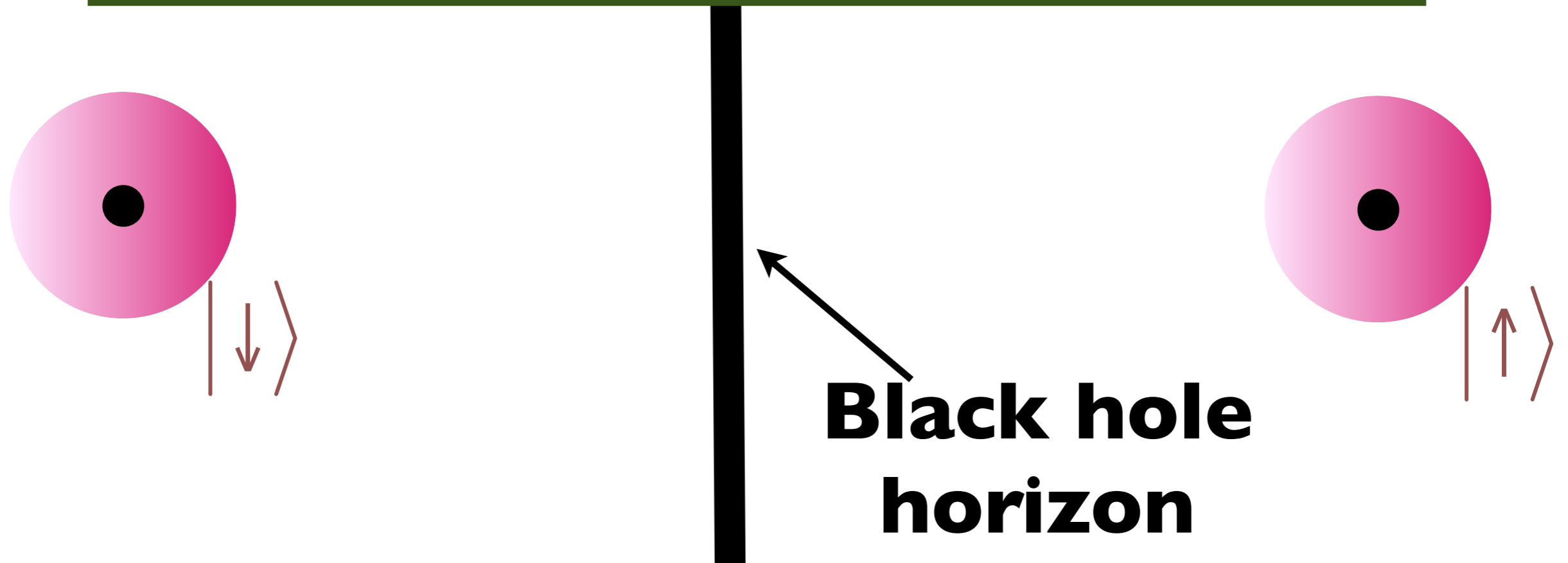
Quantum Entanglement across a black hole horizon

There is quantum entanglement between the inside and outside of a black hole



Quantum Entanglement across a black hole horizon

Hawking used this to show that black hole horizons have an entropy and a temperature
(because to an outside observer, the state of the electron inside the black hole is an unknown)



Quantum Black holes

- Black holes have an entropy and a temperature, T_H .
- The entropy, S_{BH} is proportional to their surface area.

J. D. Bekenstein, PRD **7**, 2333 (1973)
S.W. Hawking, Nature **248**, 30 (1974)

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All many-body quantum systems
(without quantum gravity)
have an entropy
proportional to their volume !?!?

Thermodynamics of quantum black holes:

$$\int \mathcal{D}g_{\mu\nu} \exp \left(-\frac{1}{\hbar} \mathcal{S}_{\text{Einstein gravity}}^{(d)}[g_{\mu\nu}] \right)$$

Metric of
spacetime



Quantum gravity: a summation over all possible configurations of spacetime, each weighted by a factor which is the exponential of (the ‘action’ of Einstein gravity)/(Planck’s constant)

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Gibbons, Hawking (1977)

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Holography and duality

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$$= \exp(S_{BH}) \times \left(\text{Many body quantum theory} \right)$$

in $d - 1$ dimensions *without* gravity

Metric of spacetime

Black holes are represented as a 'hologram' by a quantum many-body system in one lower dimension.

Duality: a 'change of variables' between the many-particle configurations and the metric of spacetime

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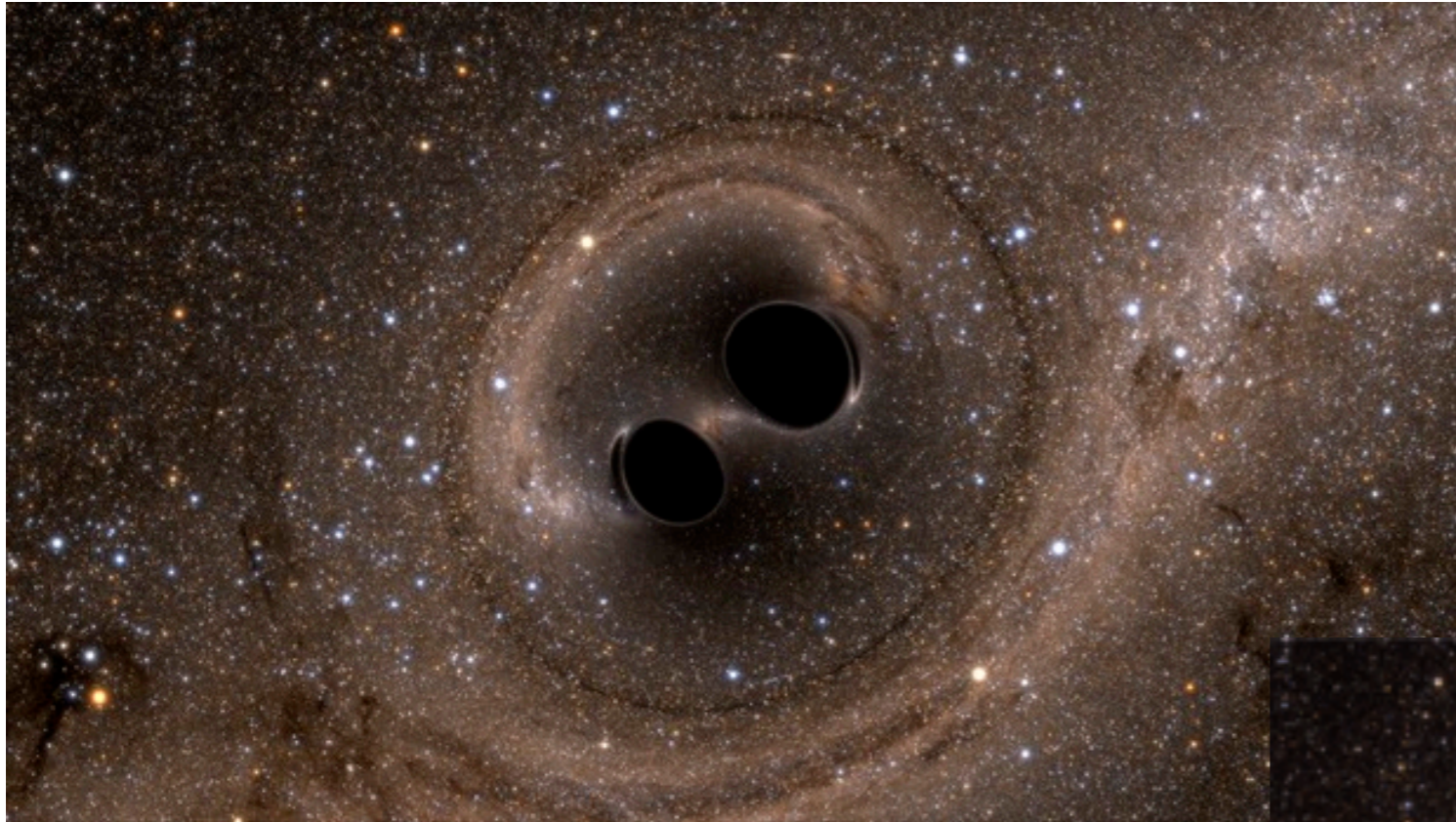
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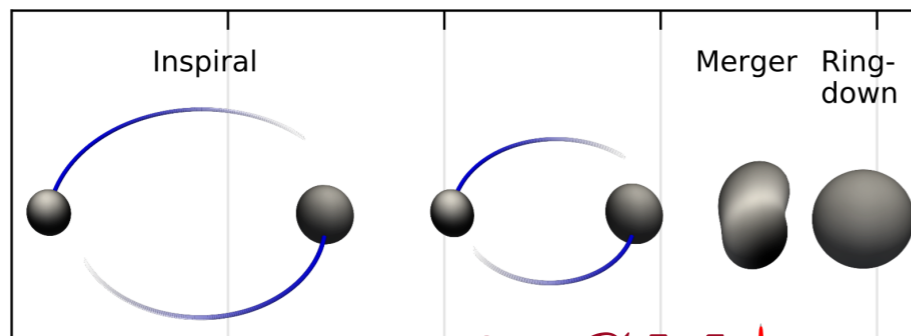
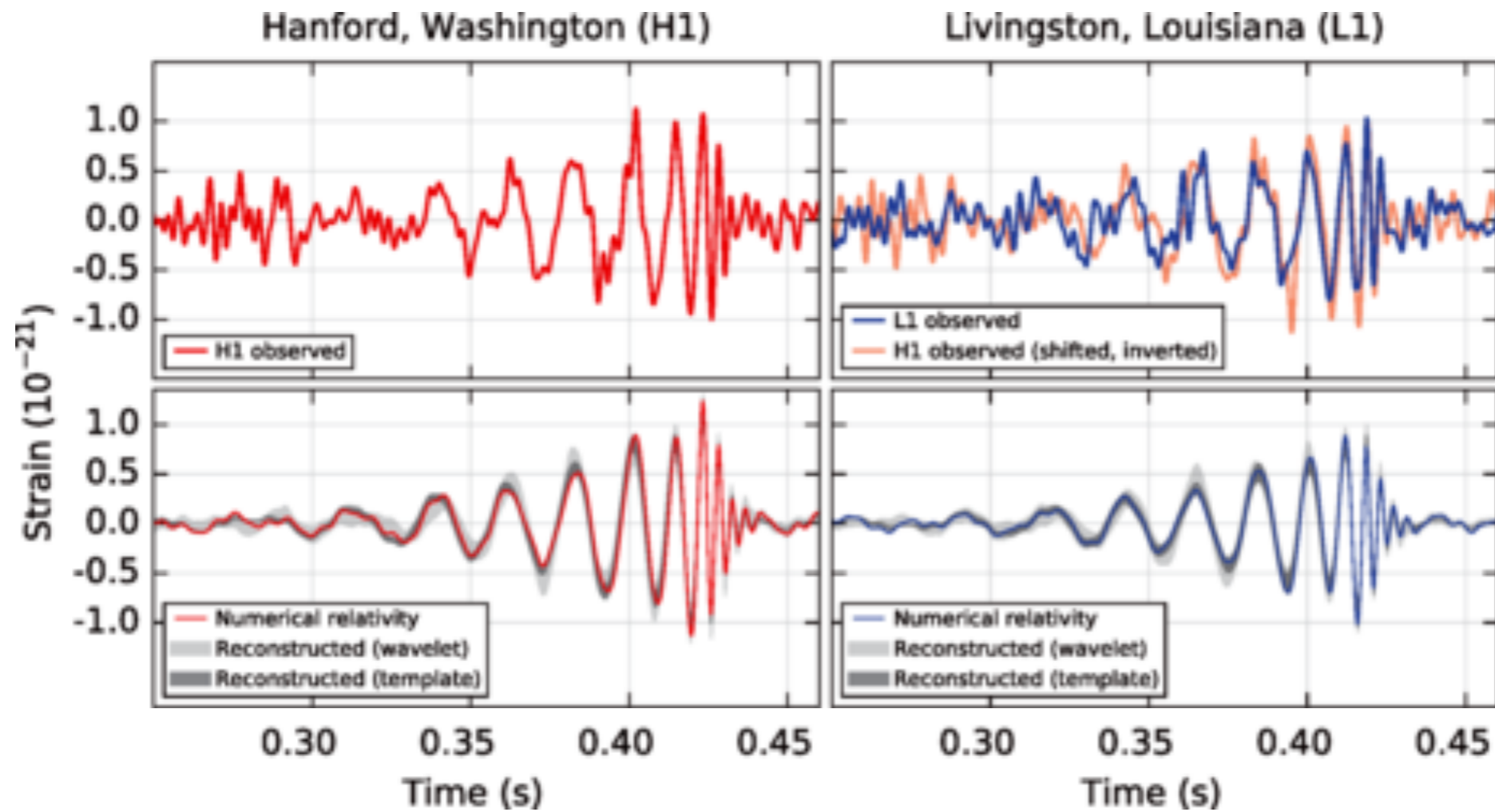
Susskind, Maldacena.....

On September 14, 2015, LIGO detected the merger of two black holes, each weighing about 30 solar masses, with radii of about 100 km, 1.3 billion light years away



0.1 seconds later !

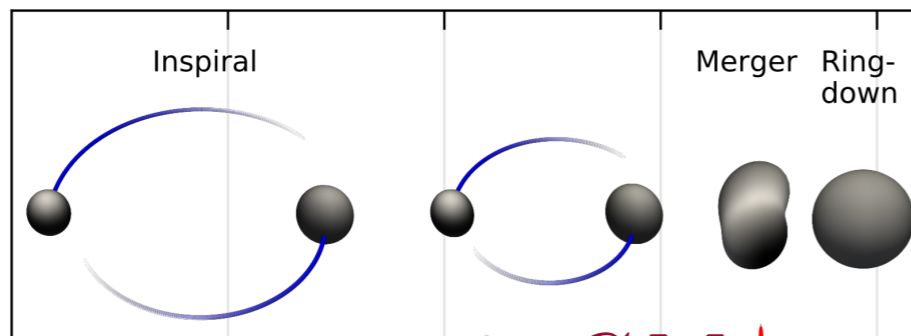
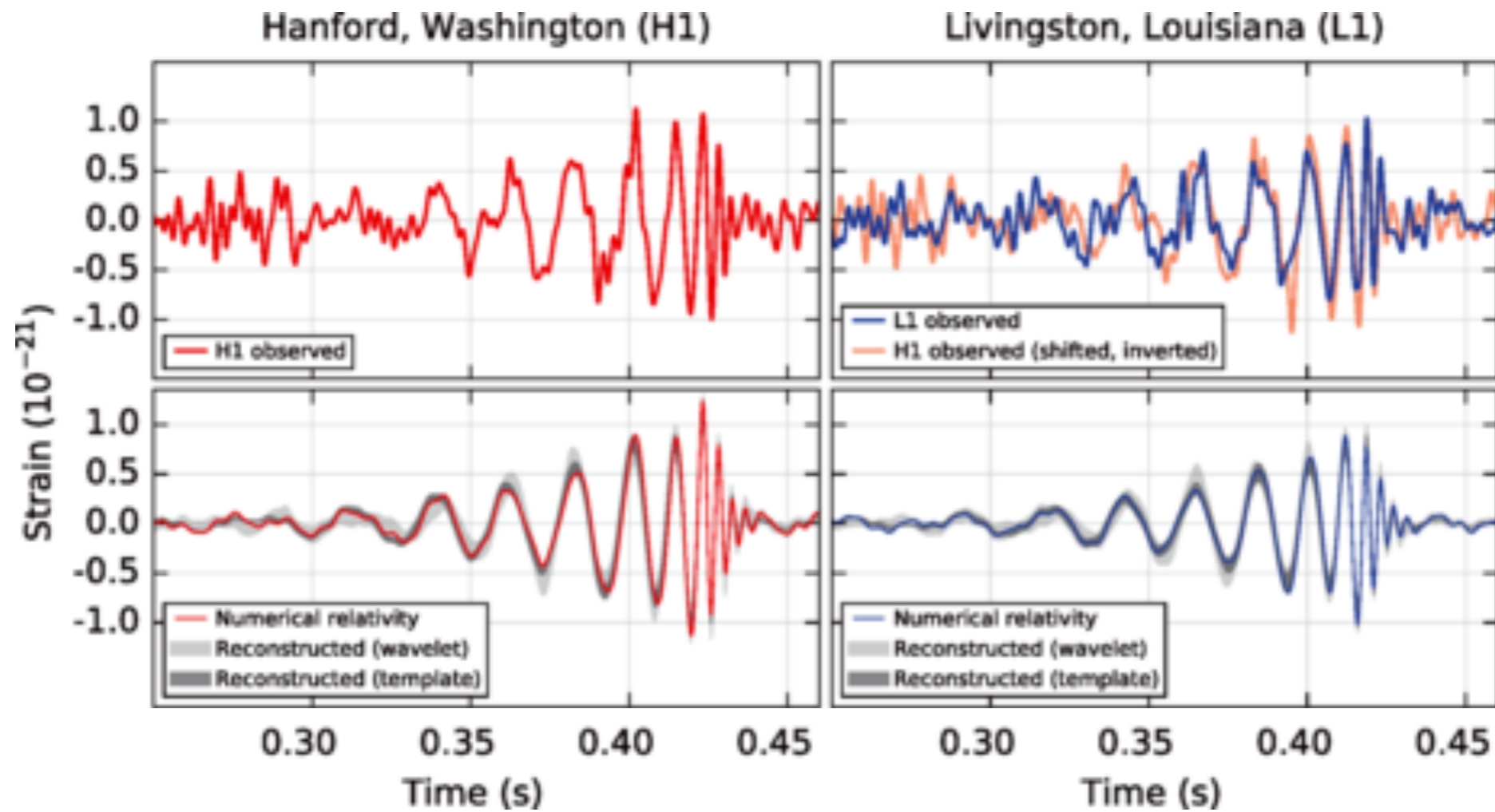




LIGO
September 14, 2015

- The ring-down time $\frac{8\pi GM}{c^3} \sim 8$ milliseconds. Curiously, for essentially all types of black holes, the ring-down time equals

$$\frac{\hbar}{k_B T_H}, \quad \hbar \text{ Planck's constant, } k_B \text{ Boltzmann's constant}$$



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The hologram of a black hole
in d dimensions
is a quantum many-particle system
in $(d - 1)$ dimensions
which relaxes to thermal equilibrium
in a Planckian time $\sim \hbar/(k_B T)$

**Quantum
entanglement**

**Black
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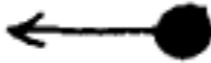
**A simple
many-particle
(SYK) model**

Ordinary metals




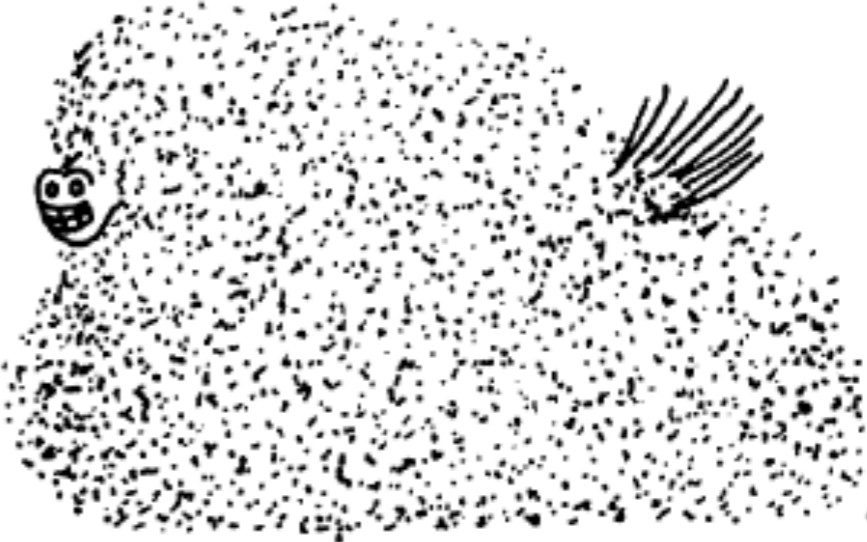
Ordinary metals are shiny, and they conduct heat and electricity efficiently. Each atom donates electrons which are delocalized throughout the entire crystal

Almost all many-electron systems are described by the quasiparticle concept: a quasiparticle is an “excited lump” in the many-electron state which responds just like an ordinary particle. The existence of quasiparticles implies limited many-particle entanglement


real particle


quasi particle

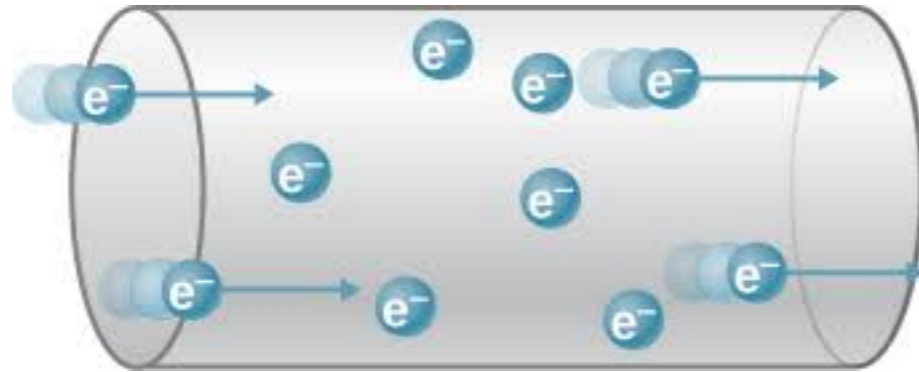

real horse


quasi horse

RDM

R.D. Mattuck

Current flow with quasiparticles



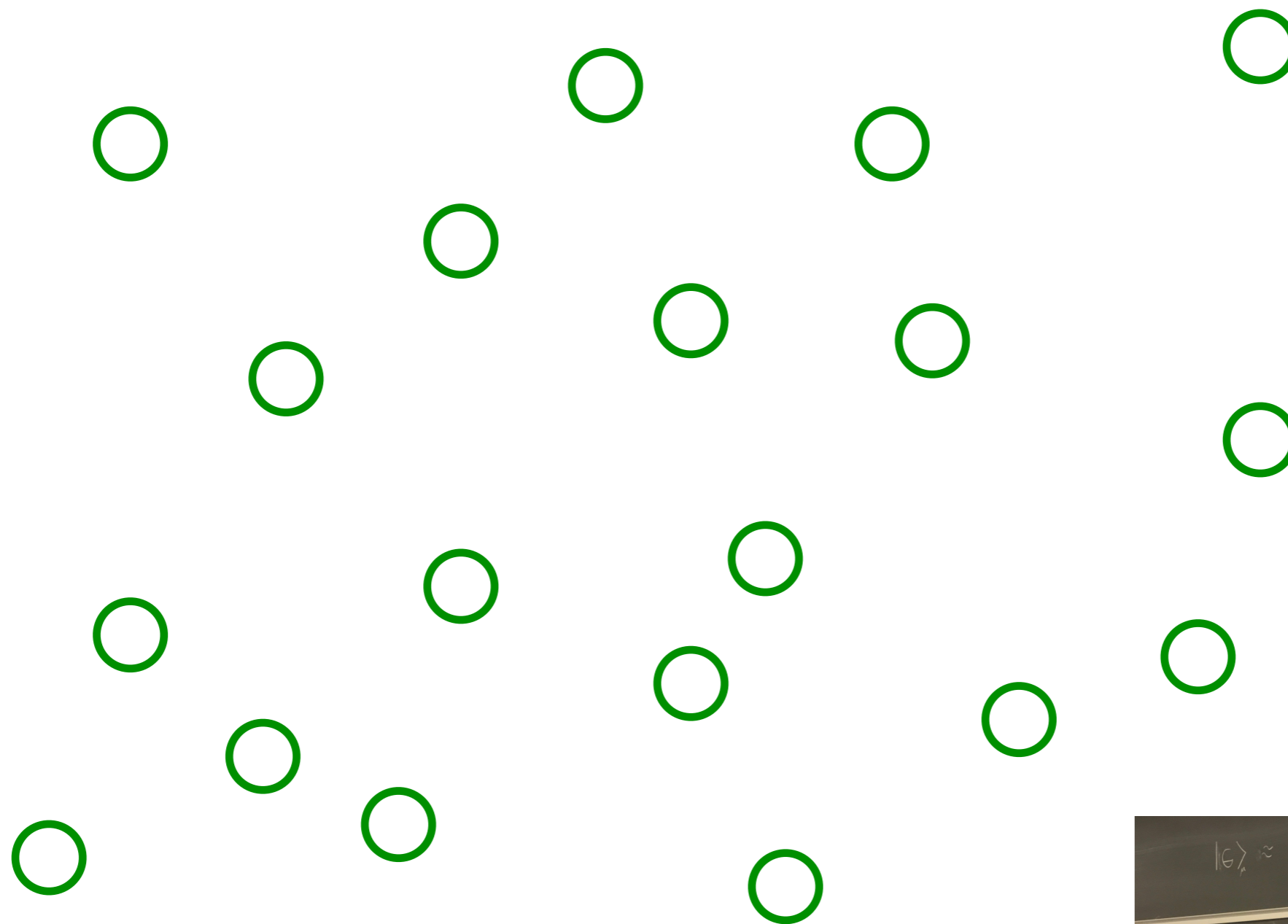
Flowing quasiparticles scatter off each other in a typical scattering time τ

This time is much longer than a limiting
'Planckian time' $\frac{\hbar}{k_B T}$.

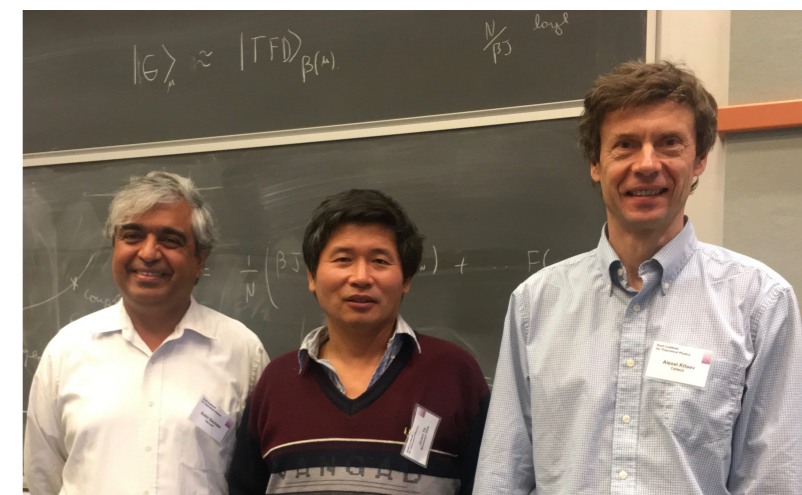
The long scattering time implies that quasiparticles are well-defined.

The Sachdev-Ye-Kitaev (SYK) model

Sachdev, Ye (1993); Kitaev (2015)

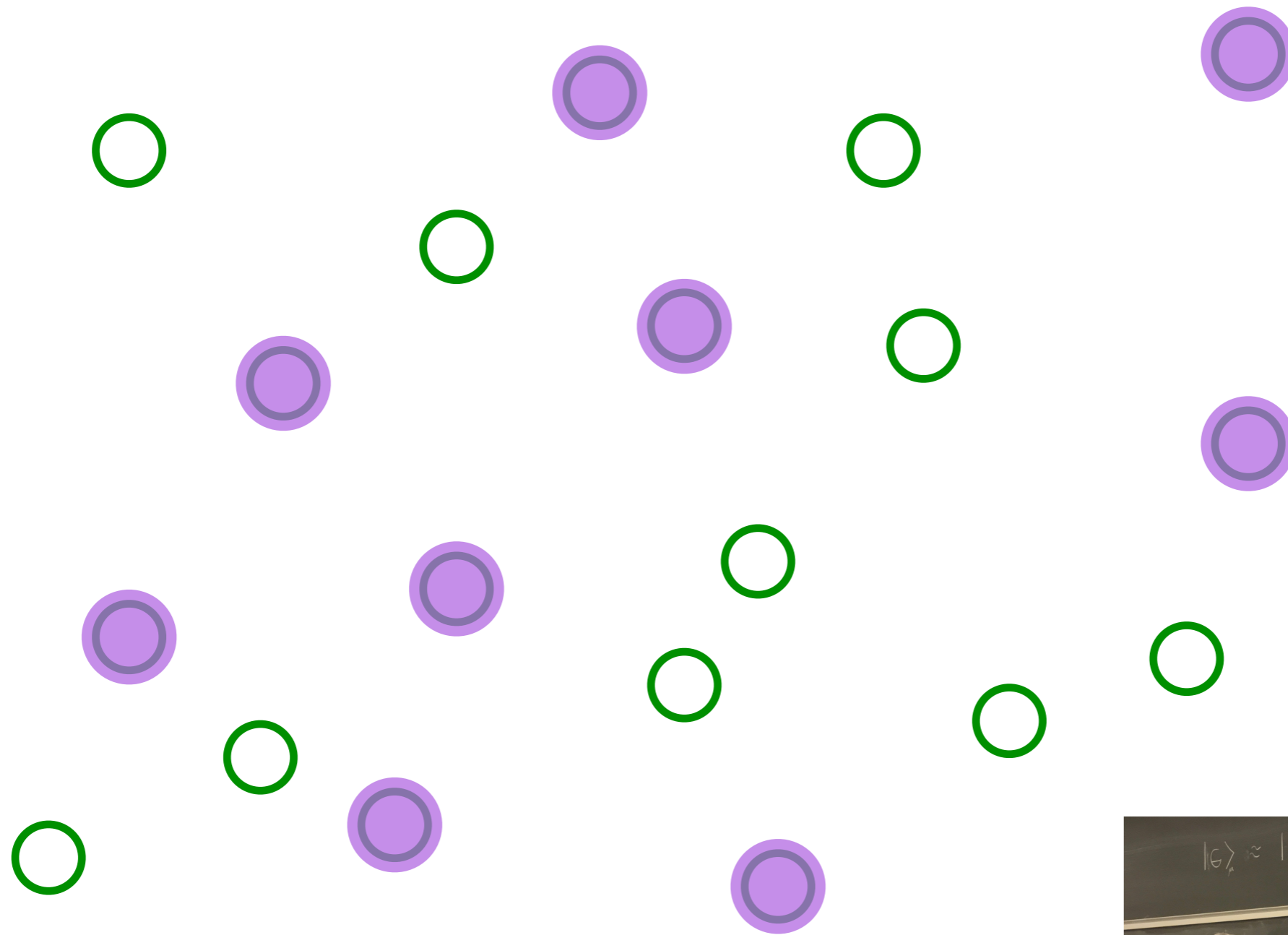


Pick a set of random positions

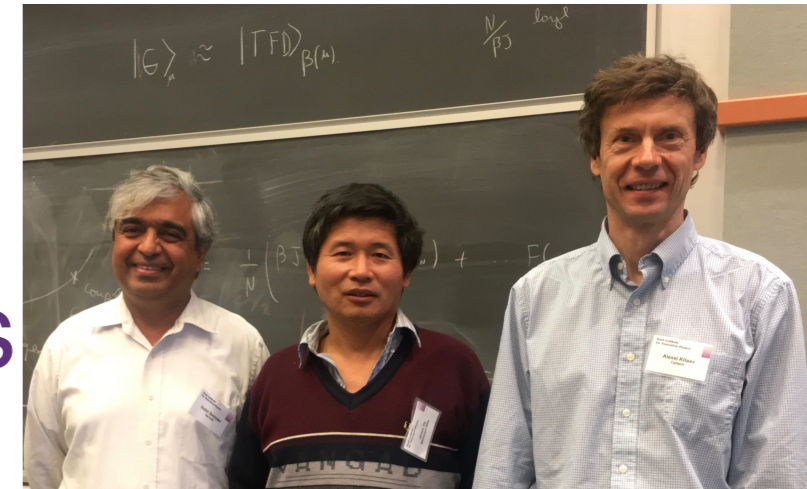


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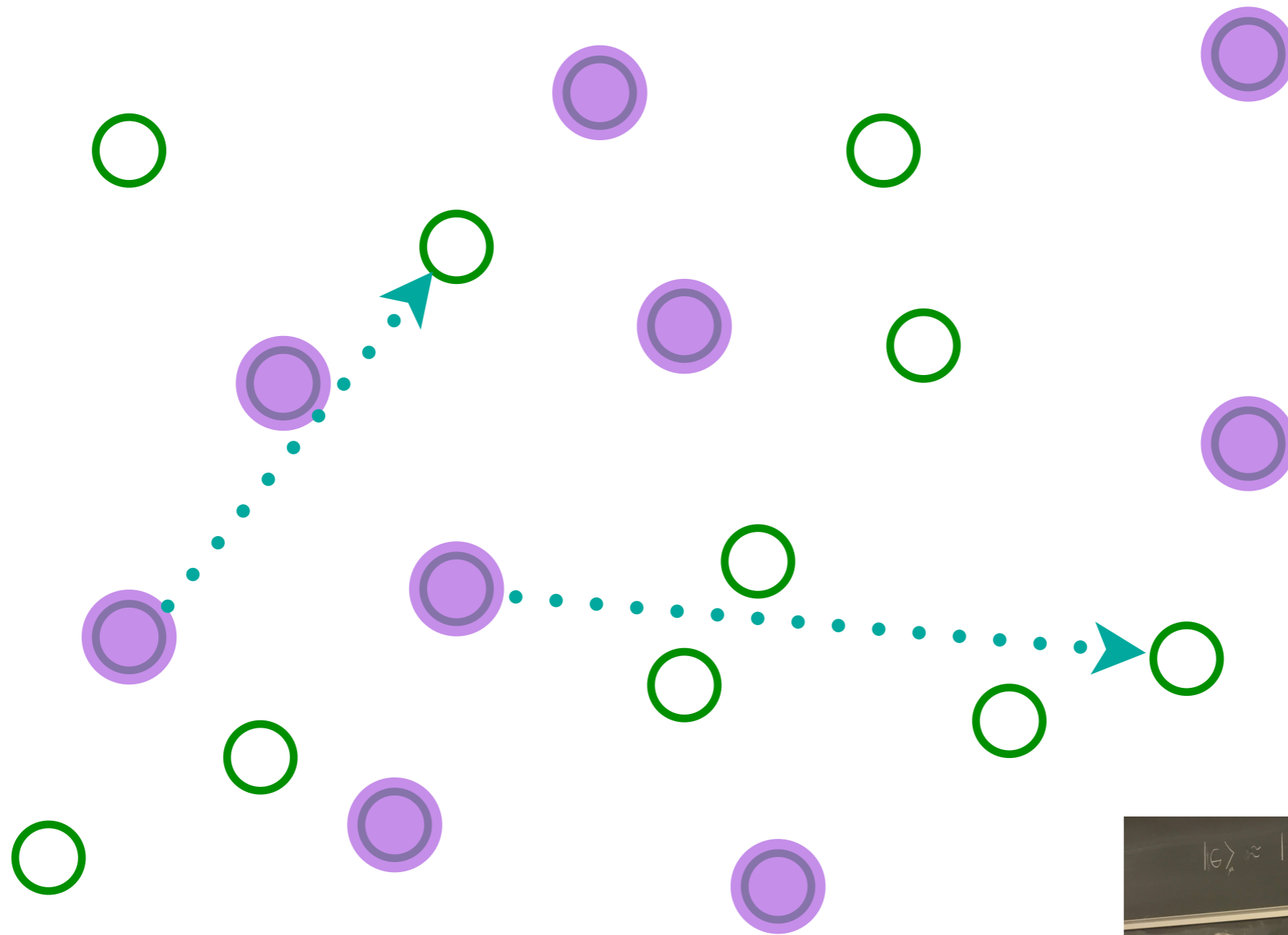


Place electrons randomly on some sites

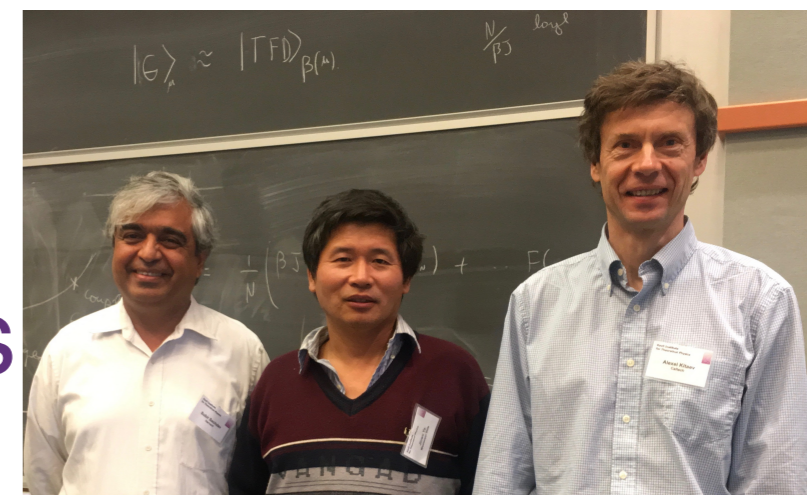


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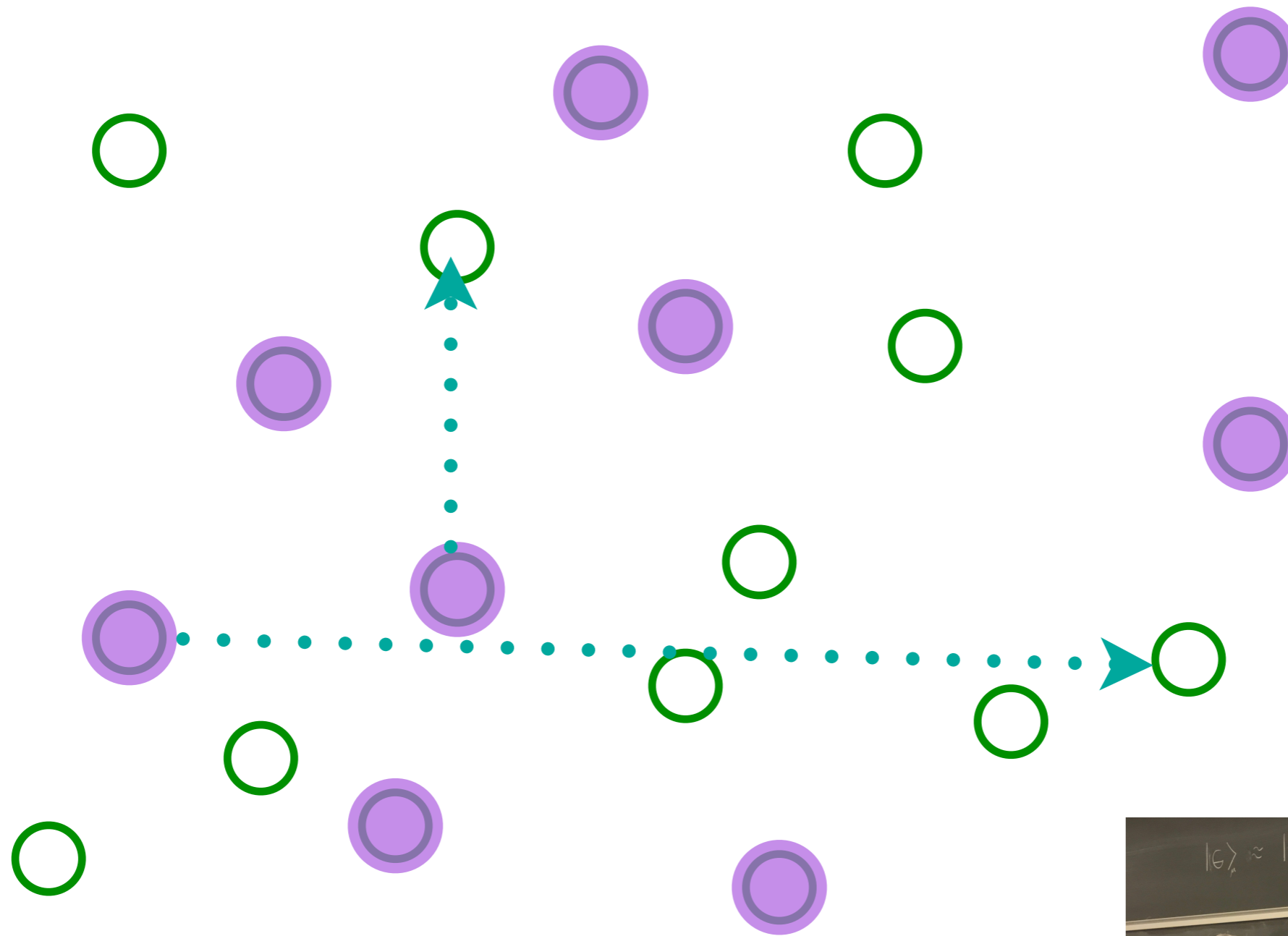


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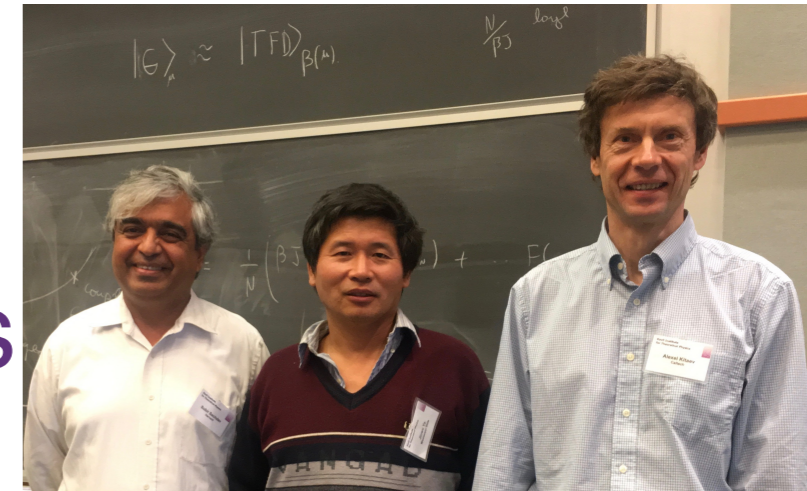


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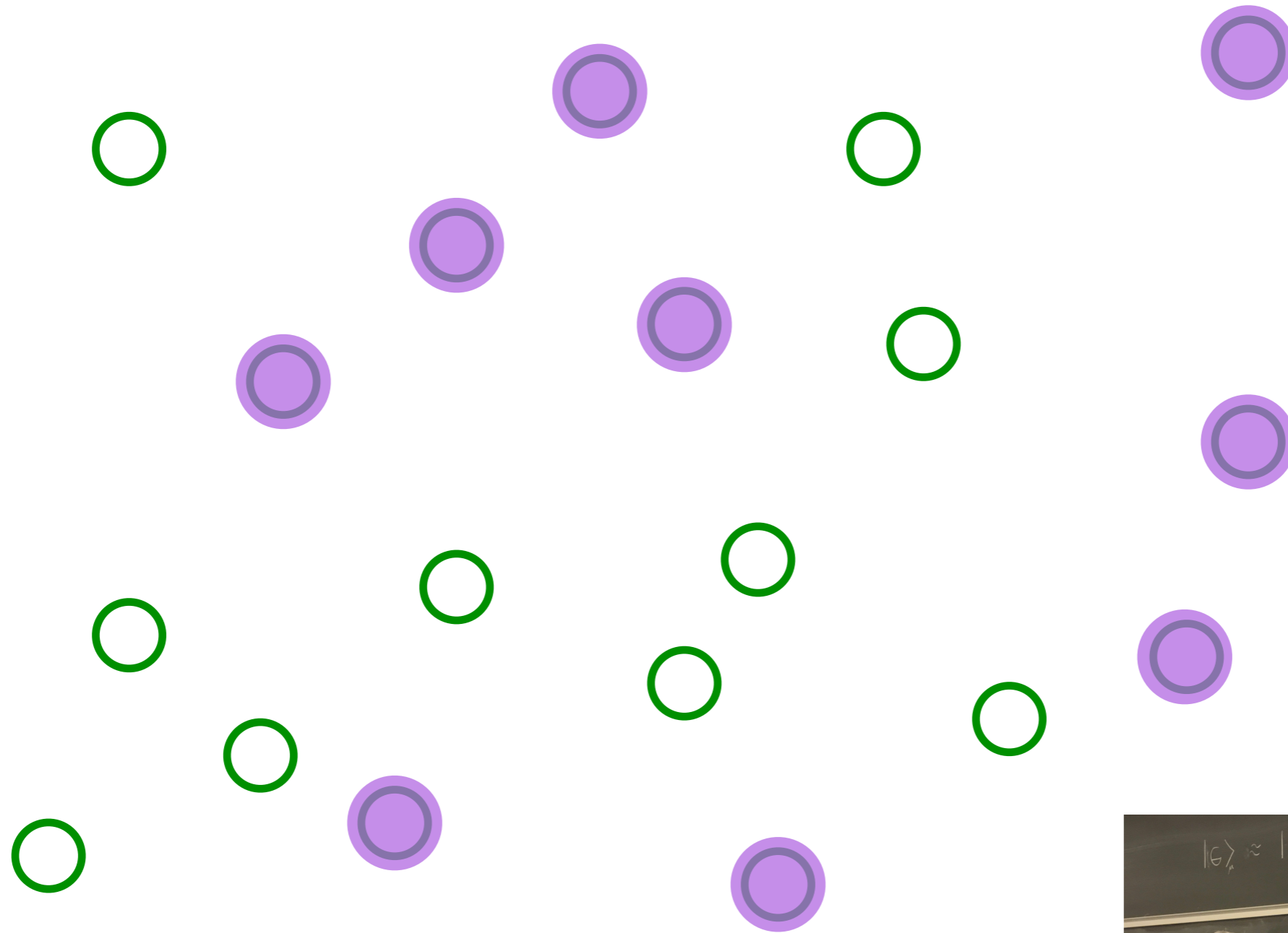


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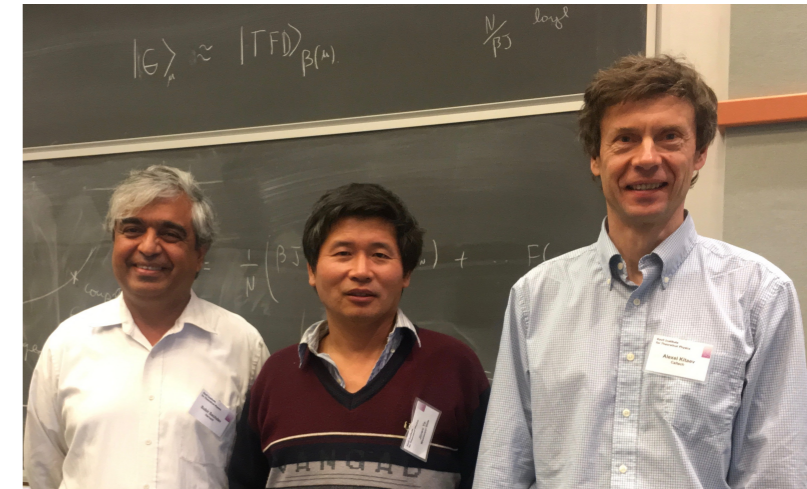


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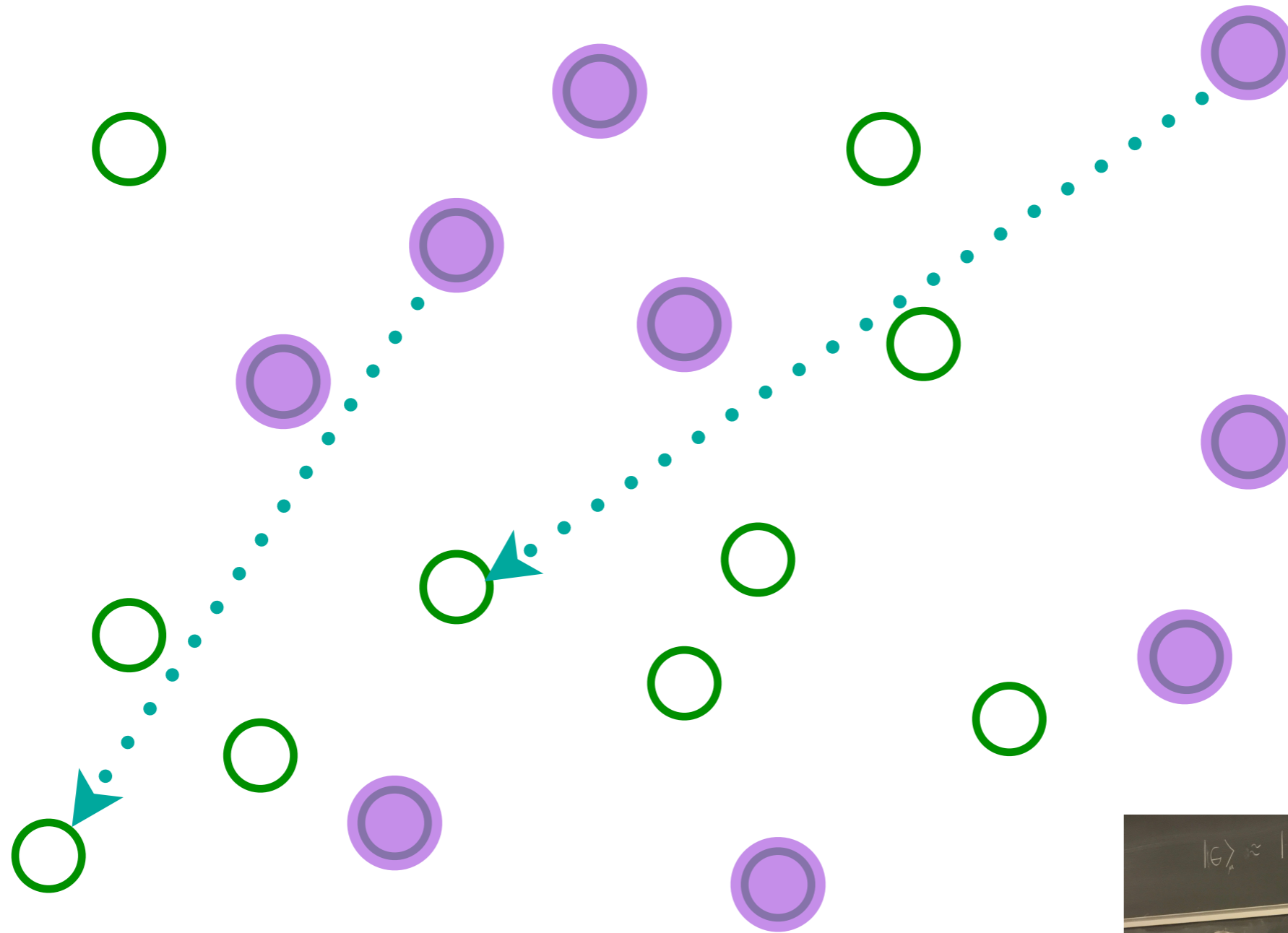


Entangle electrons pairwise randomly

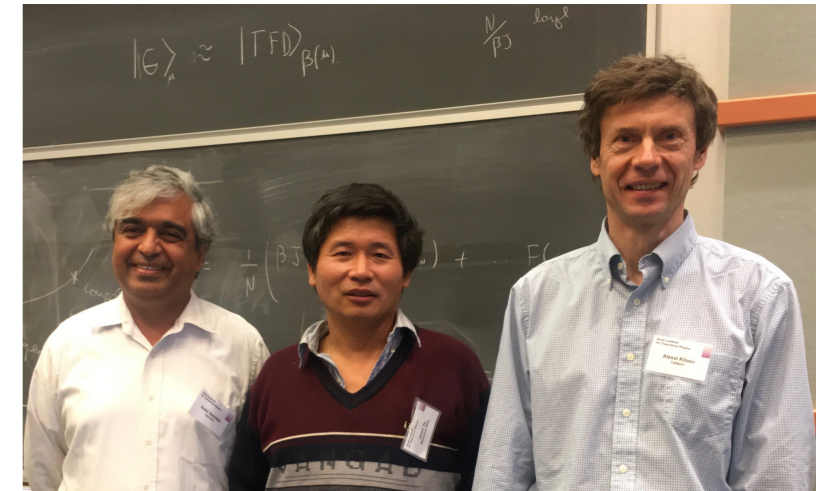


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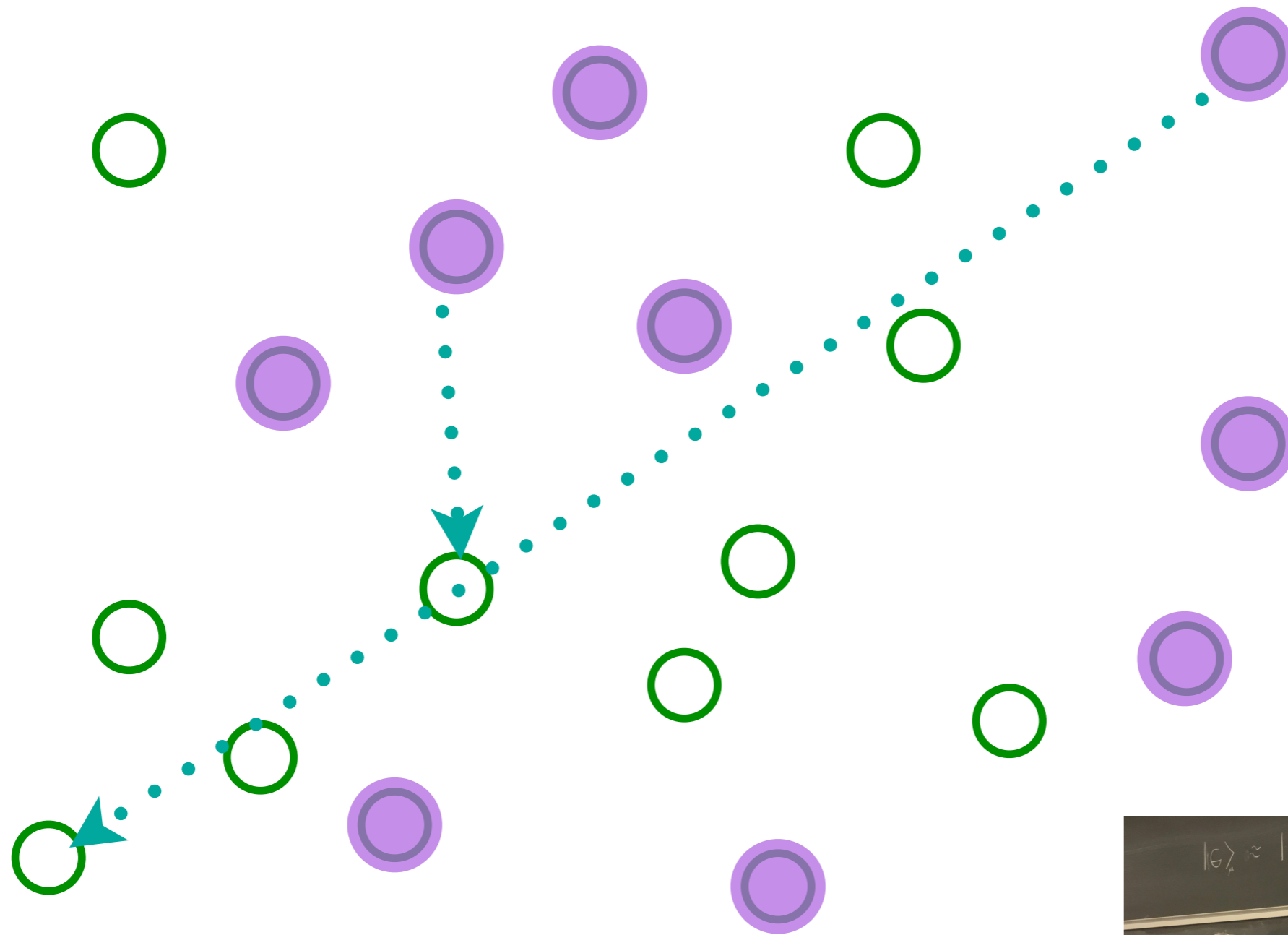


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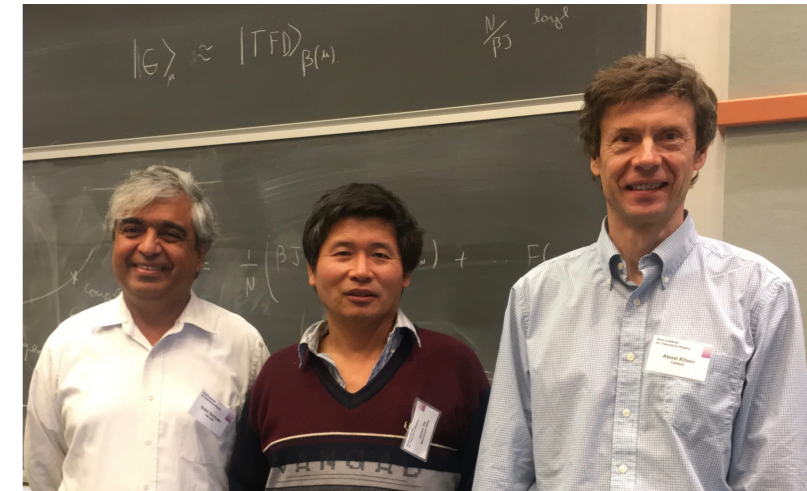


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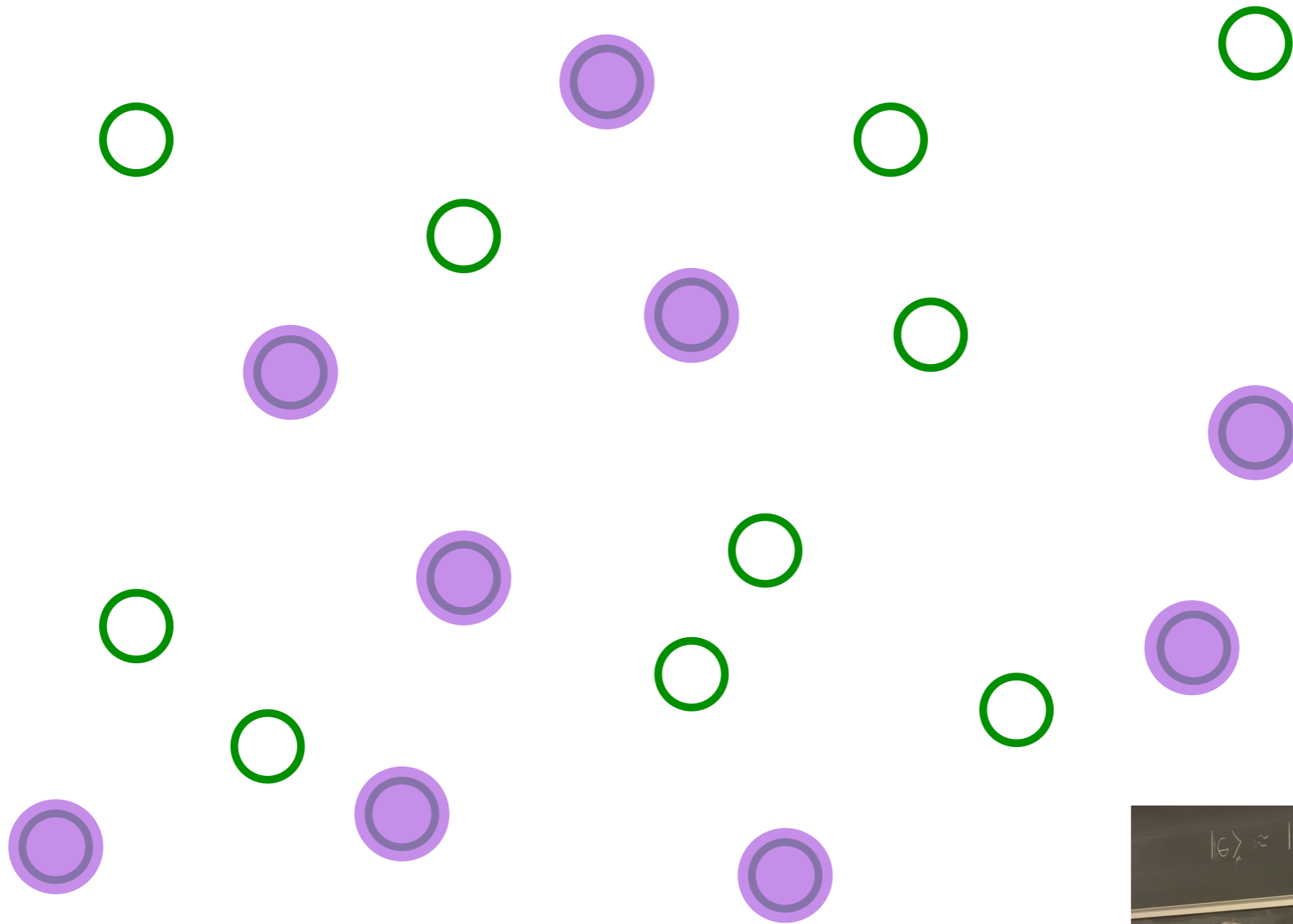


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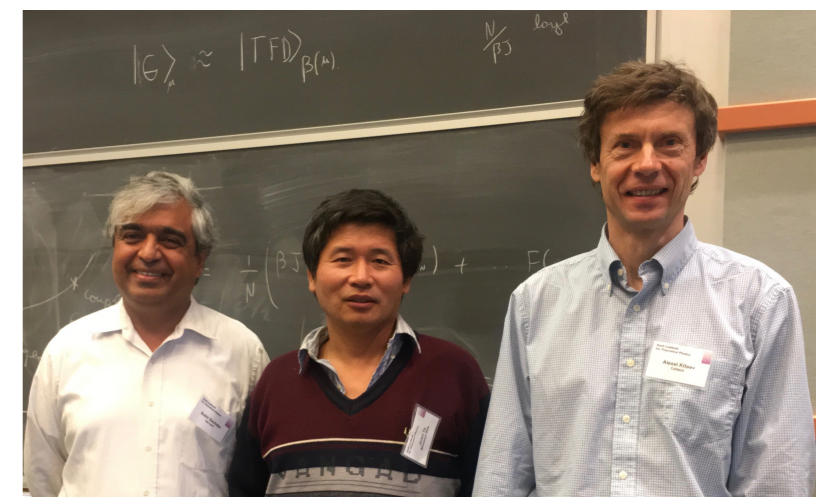


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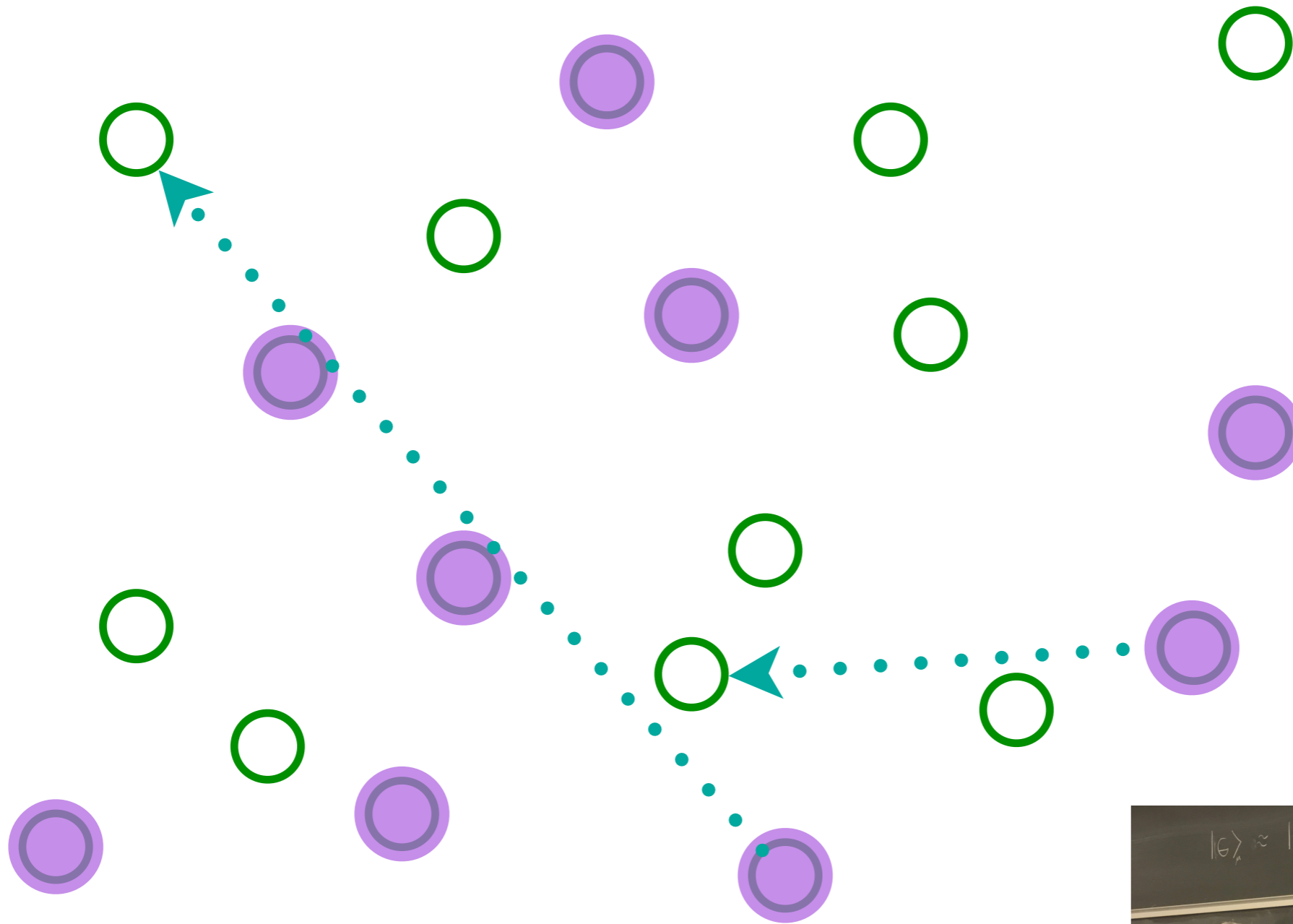


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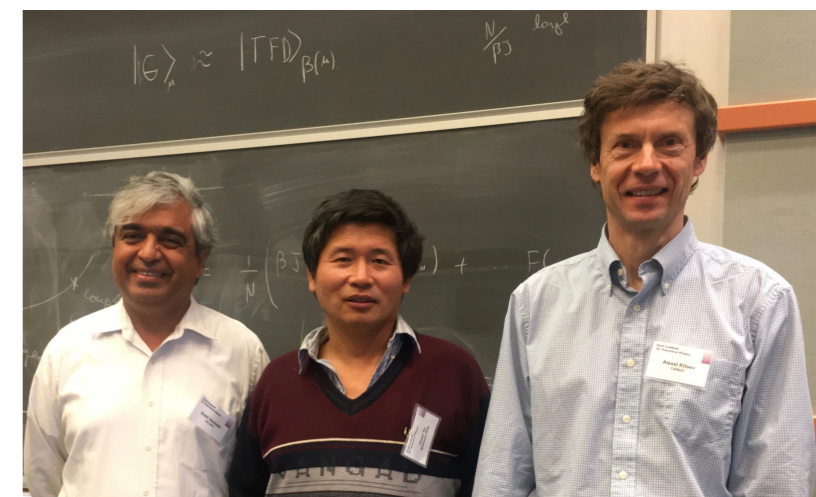


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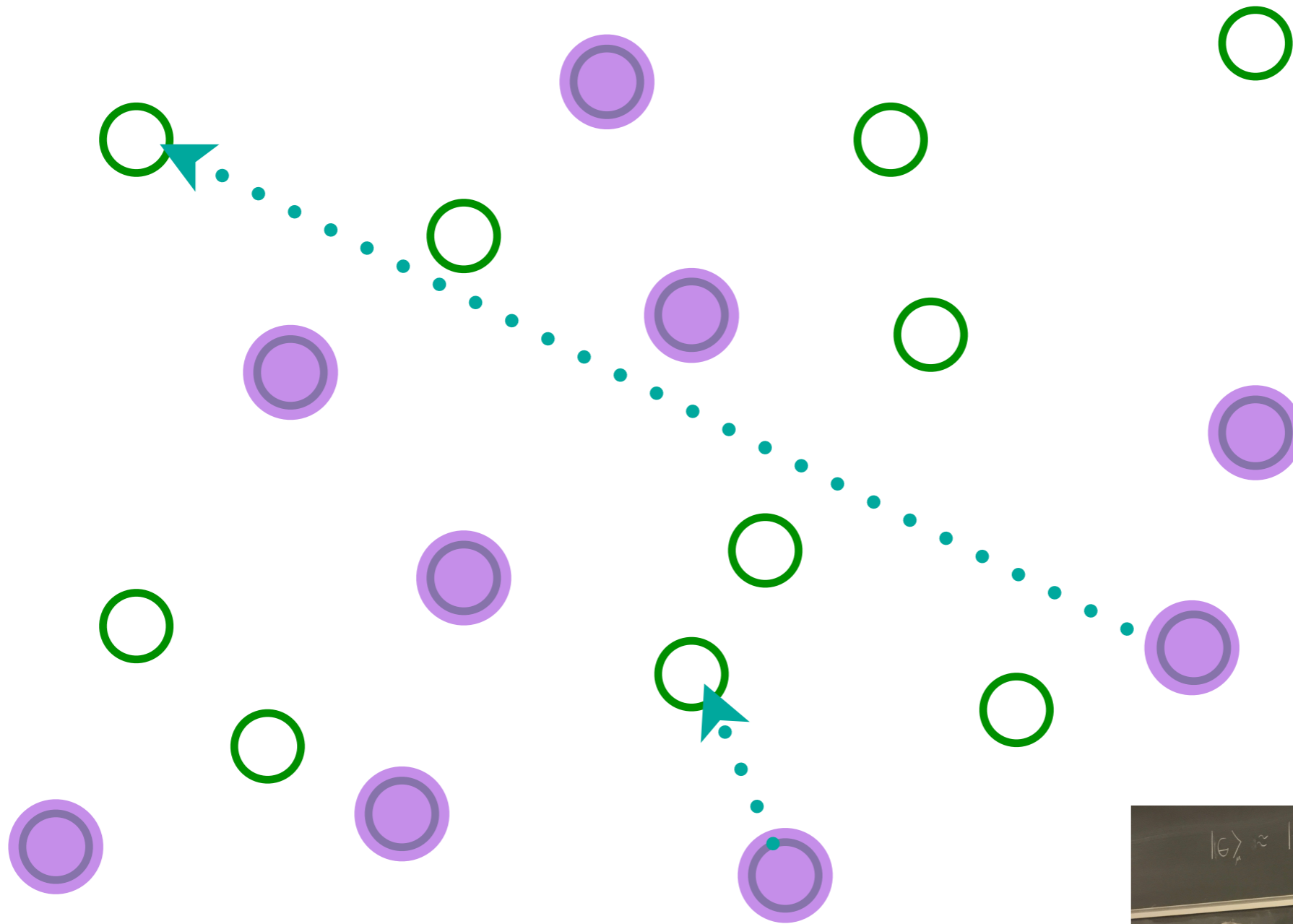


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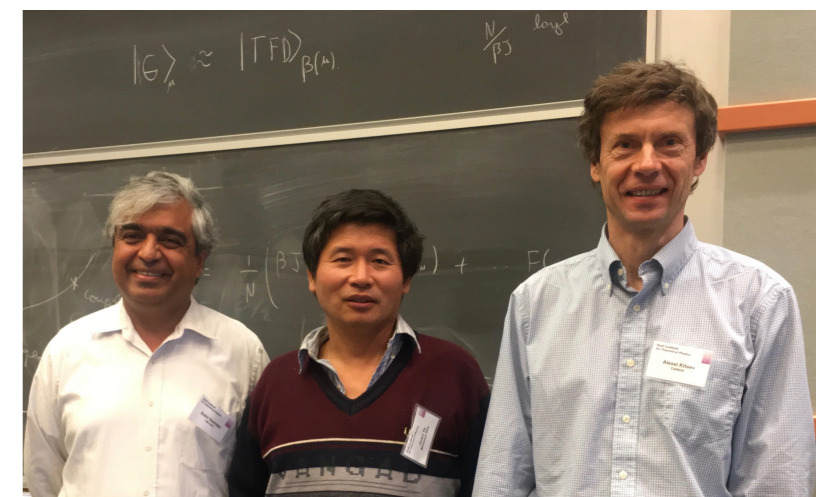


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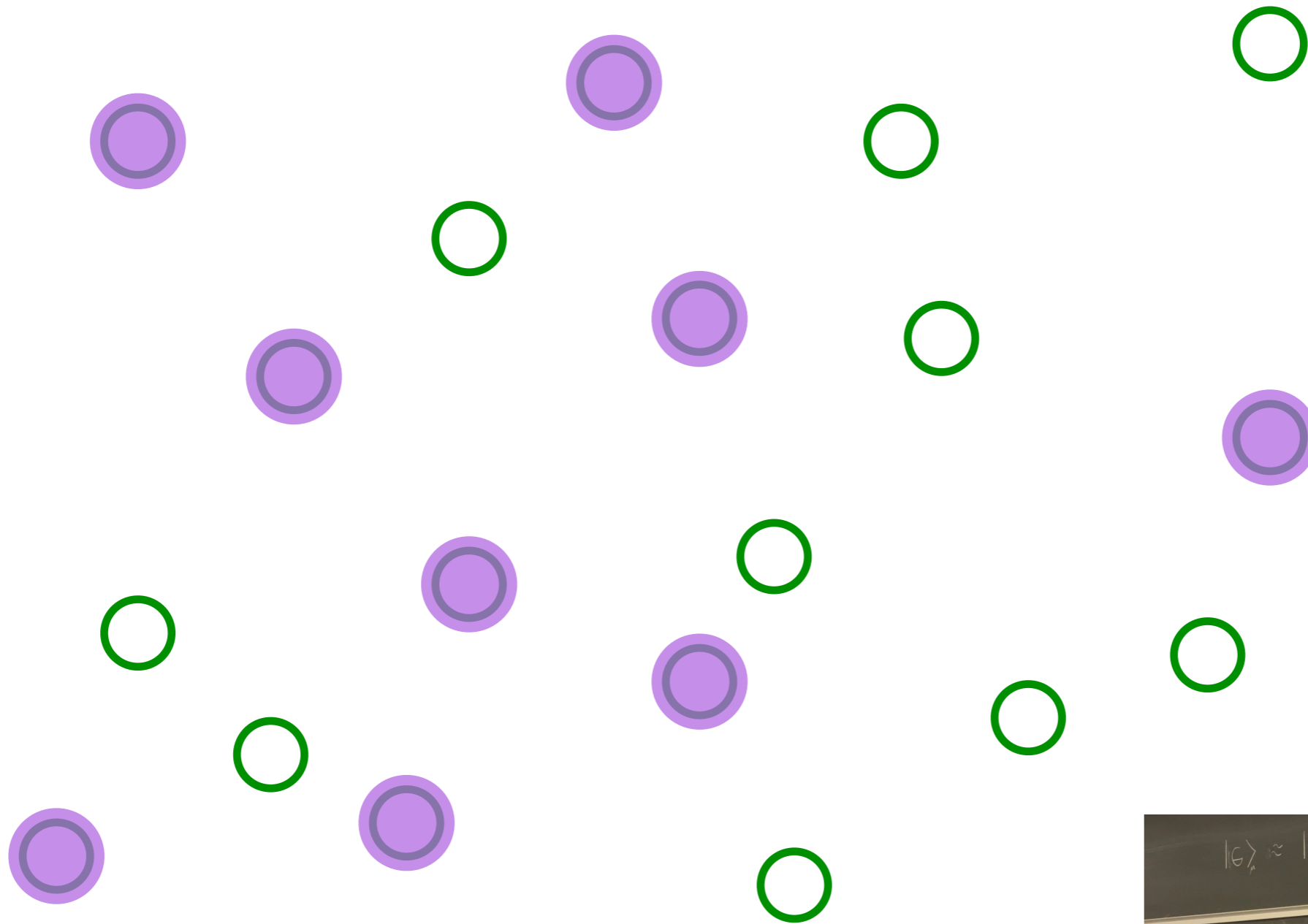


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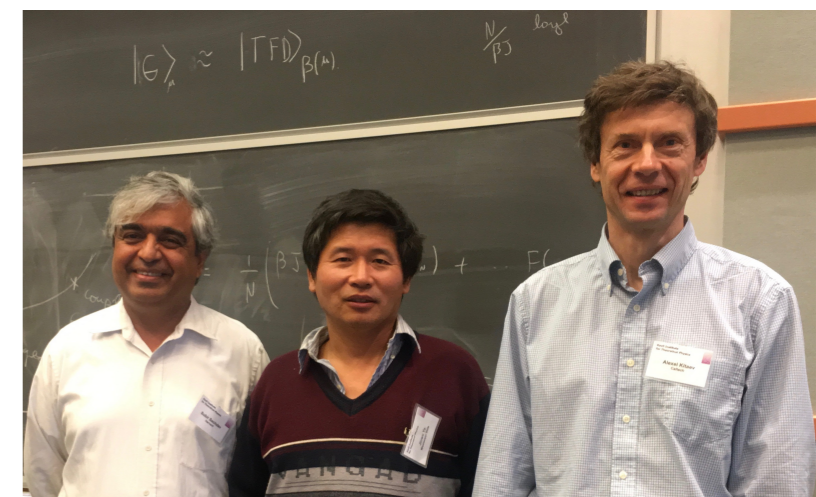


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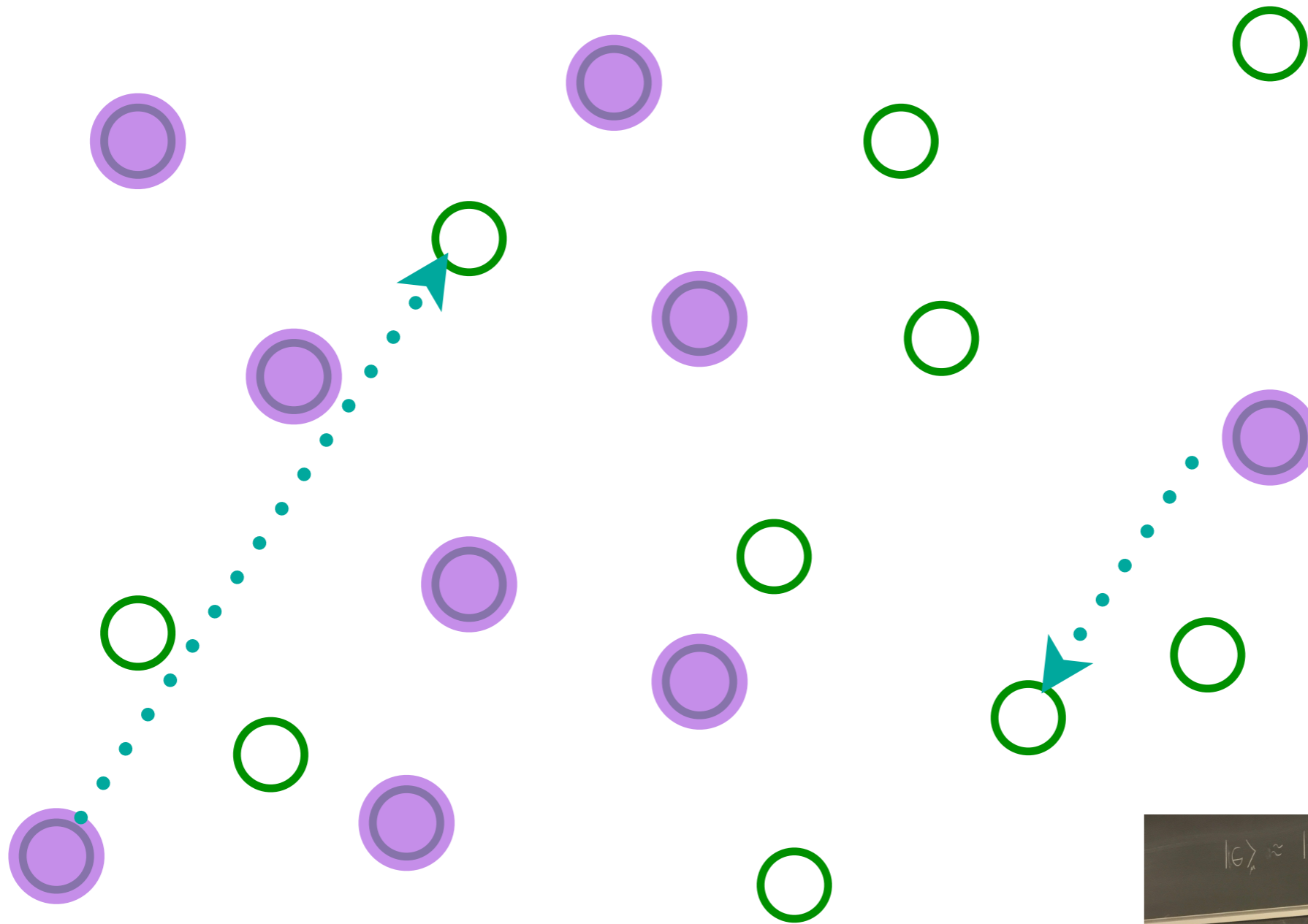


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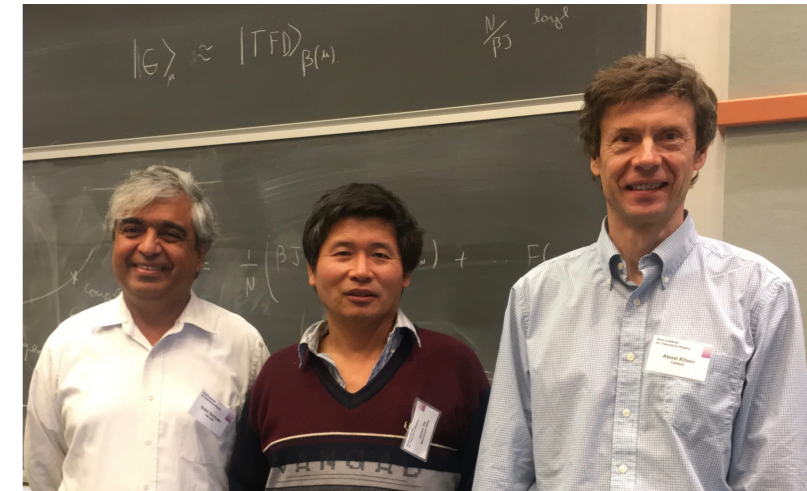


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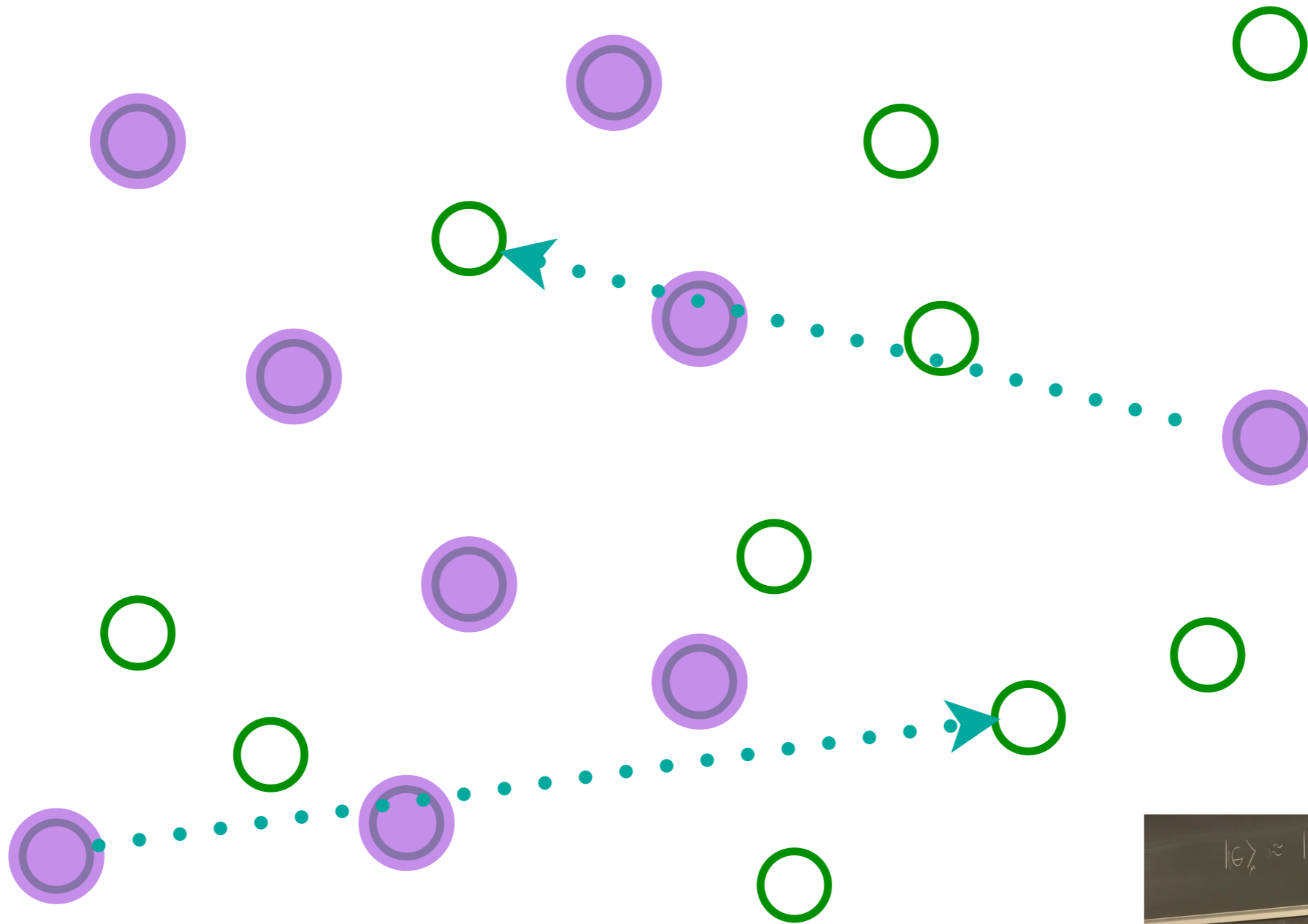


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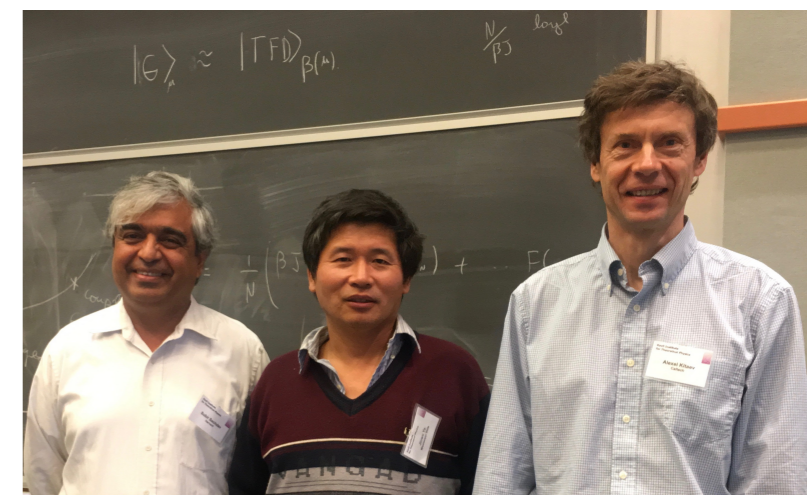


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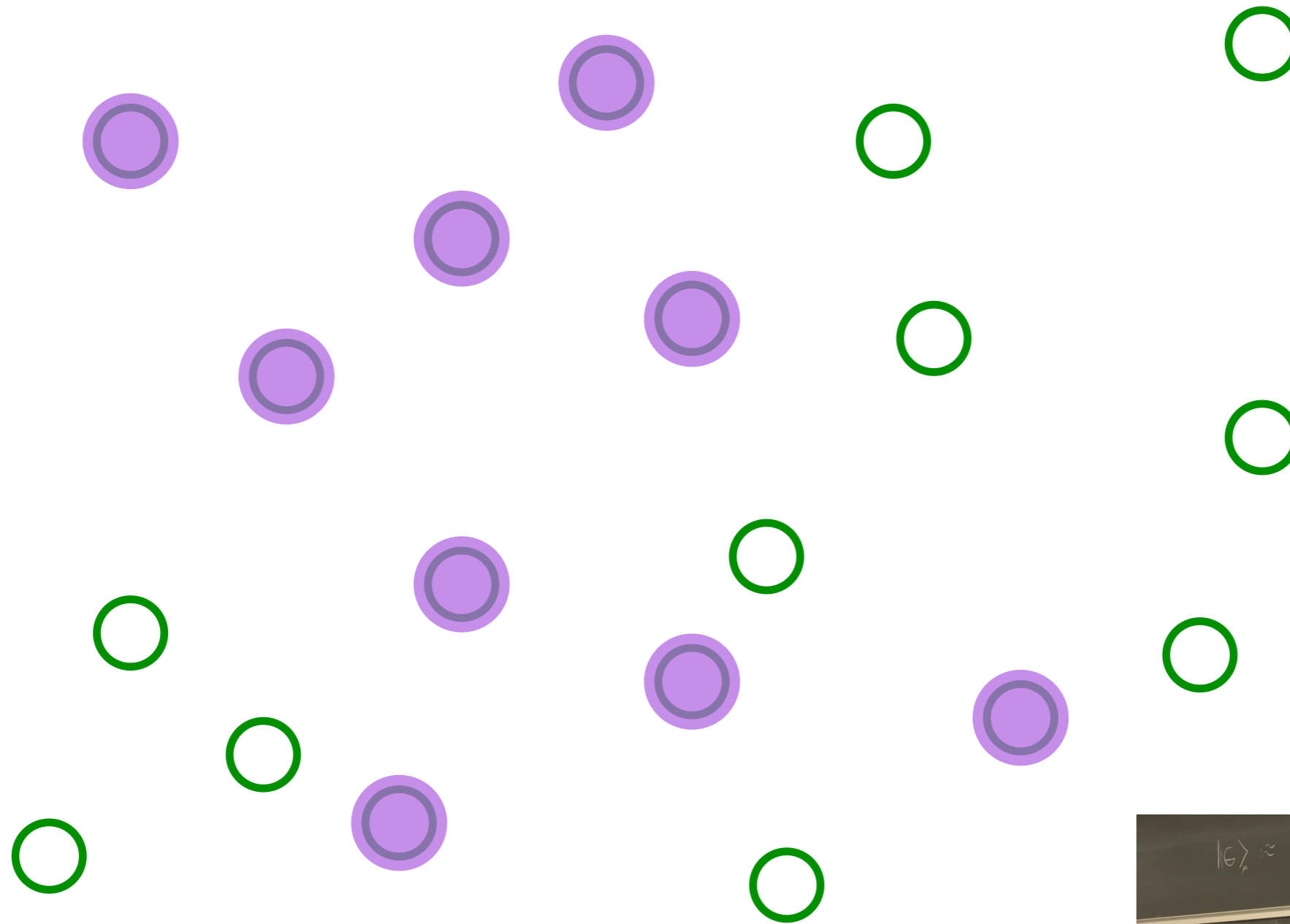


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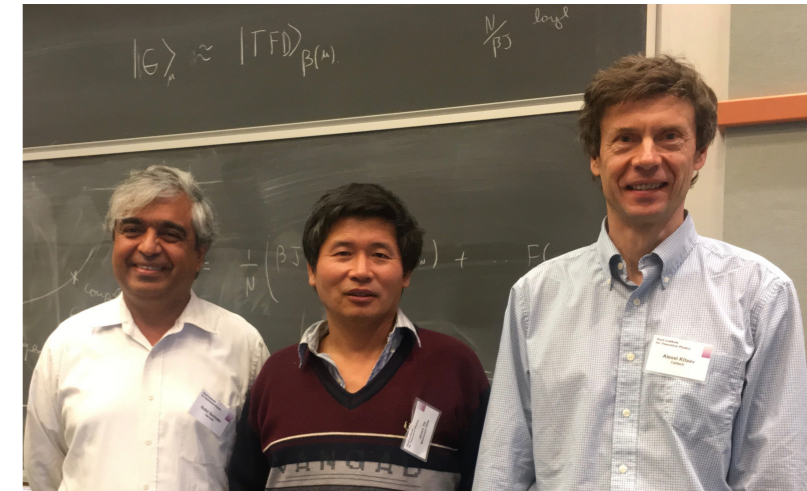


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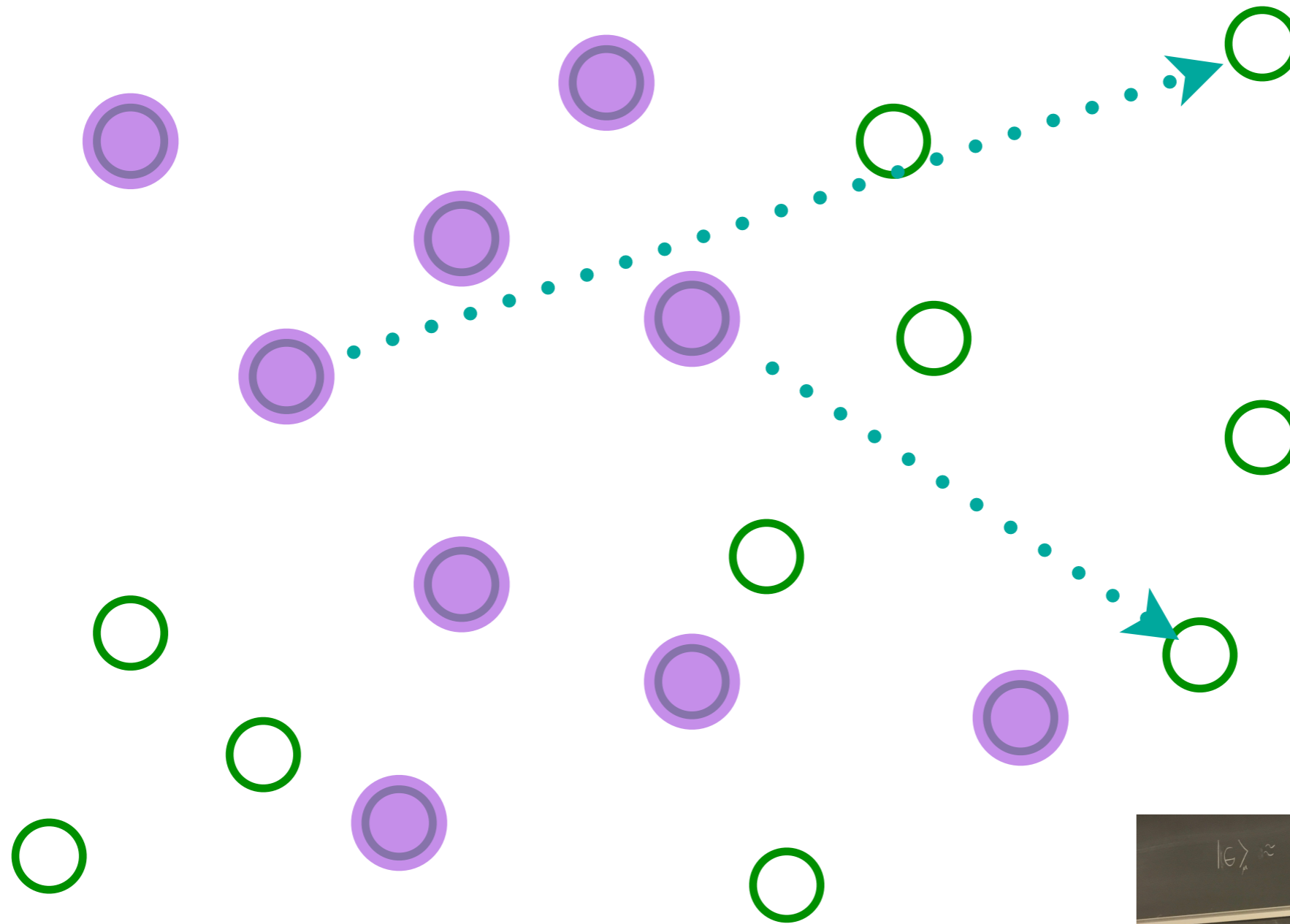


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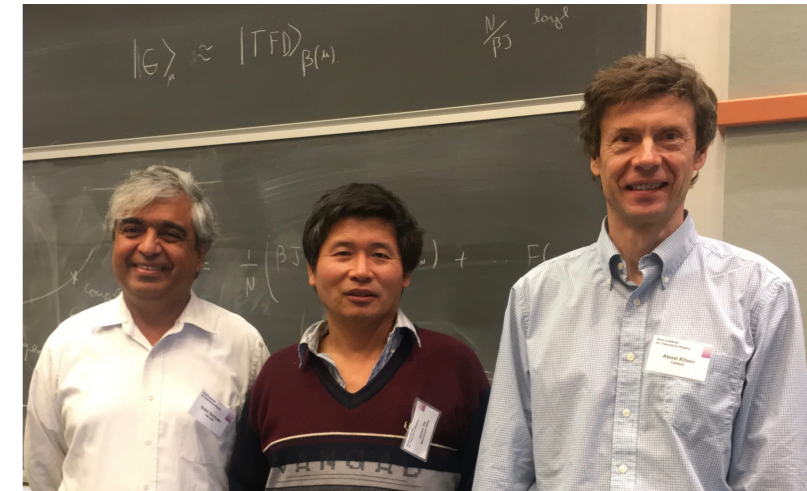


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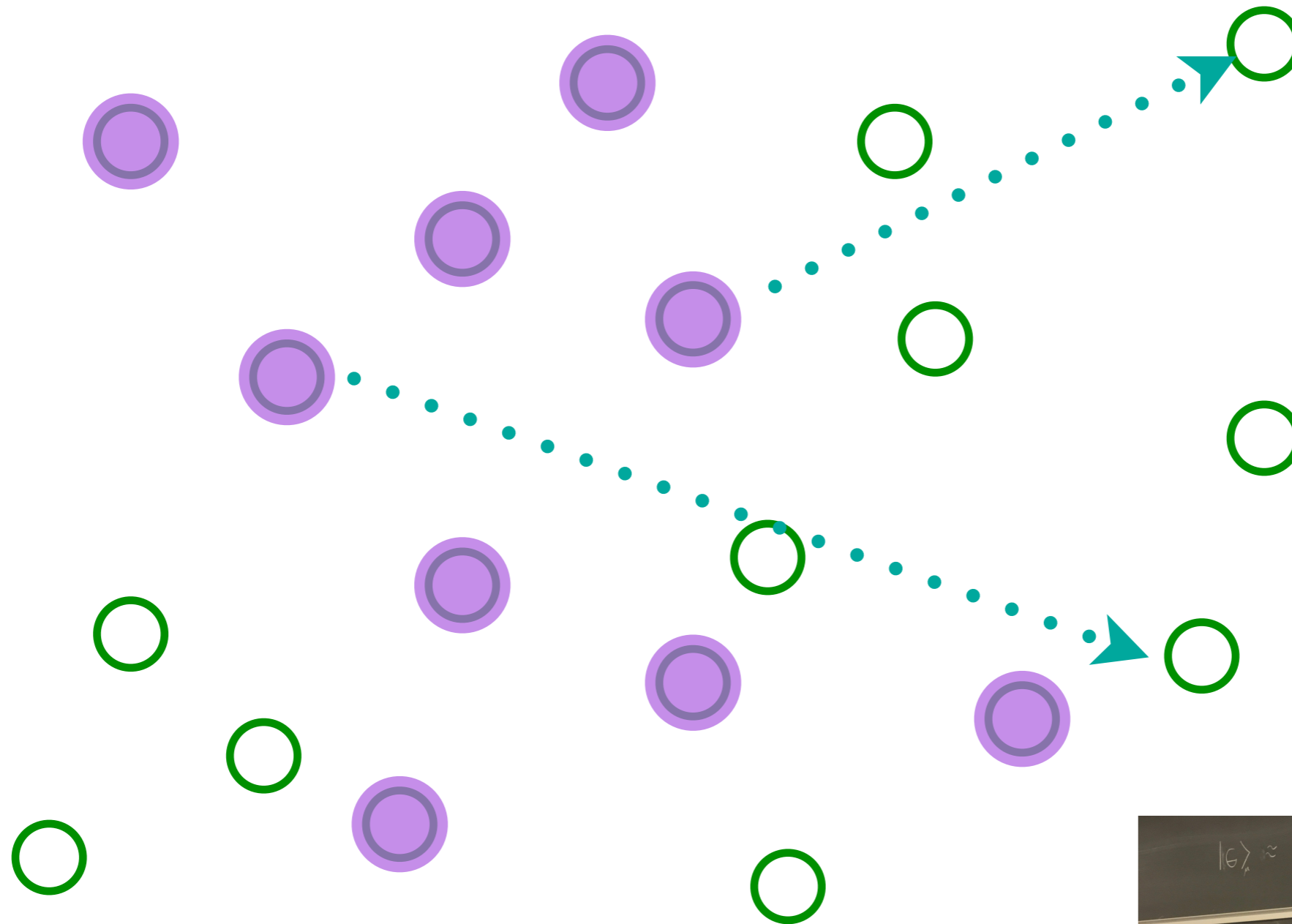


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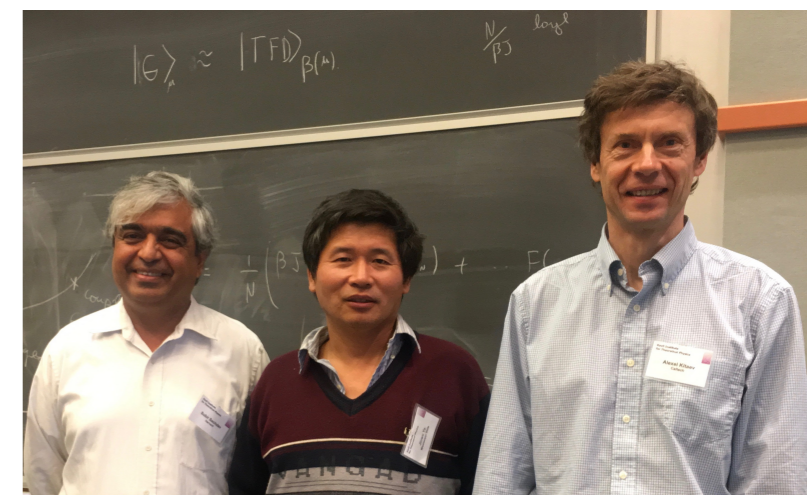


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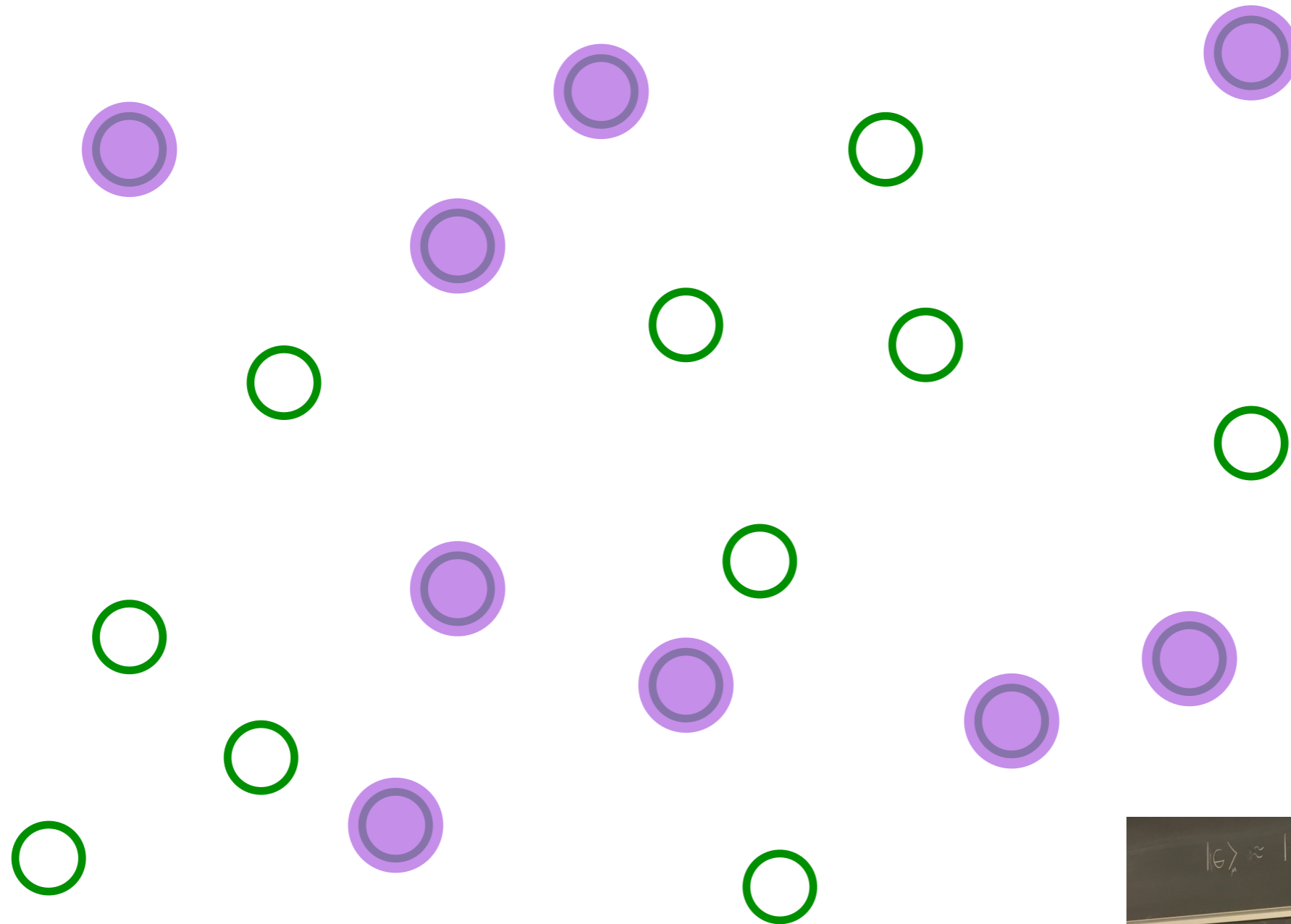


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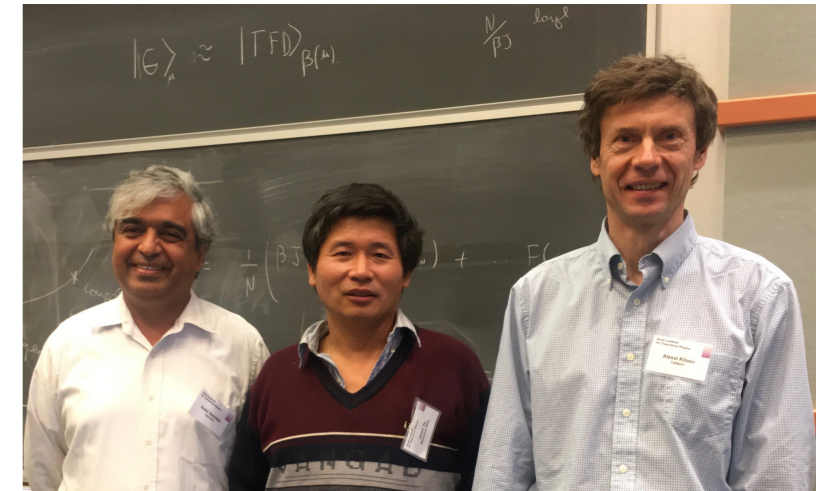


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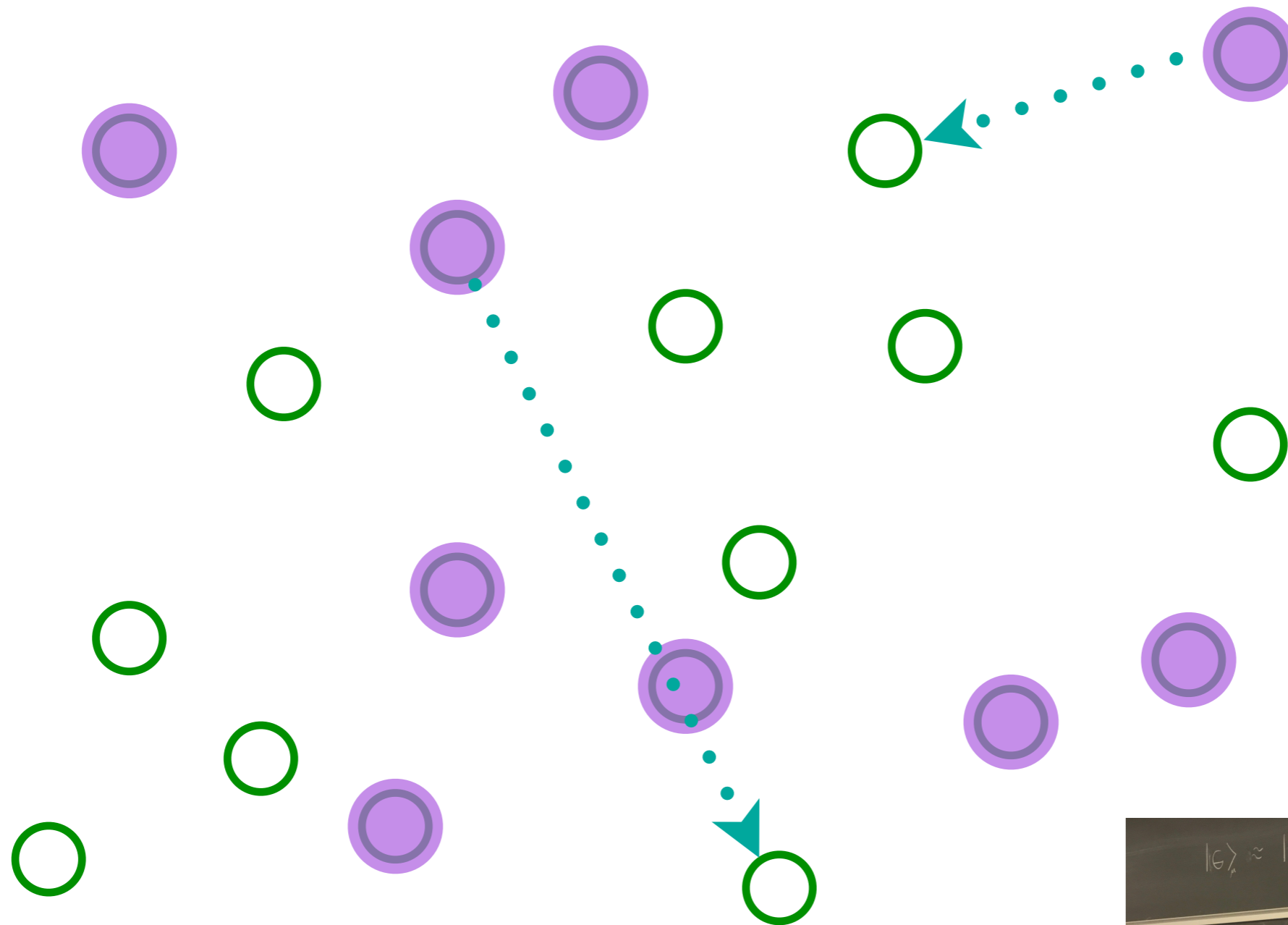


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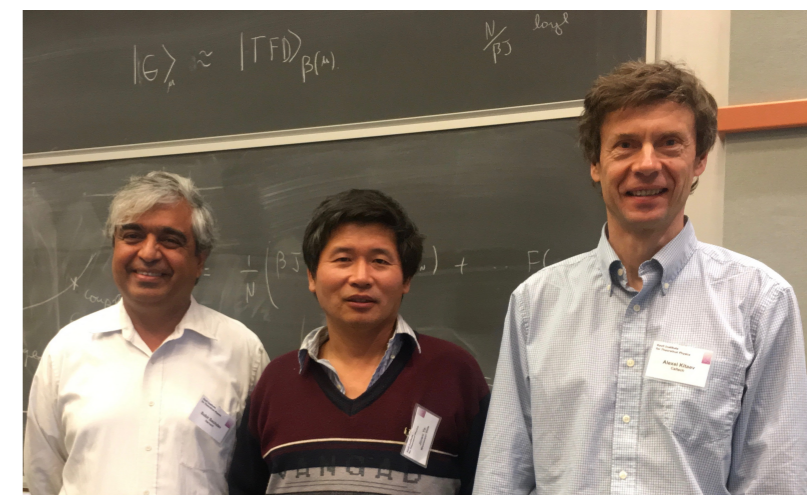


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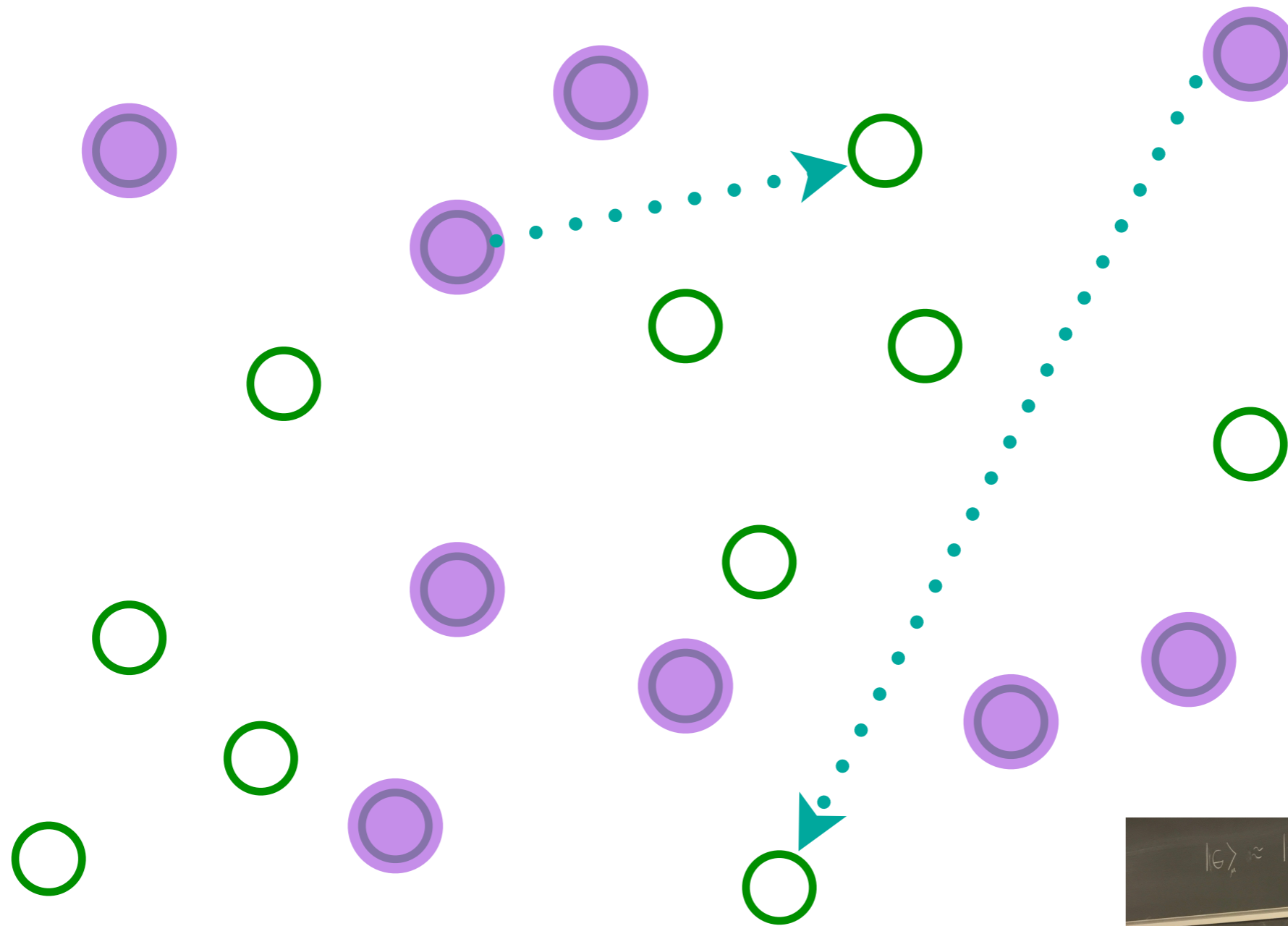


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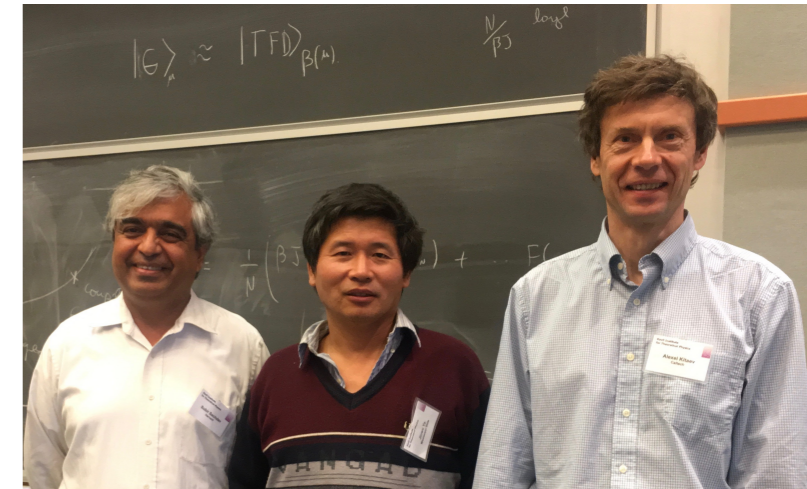


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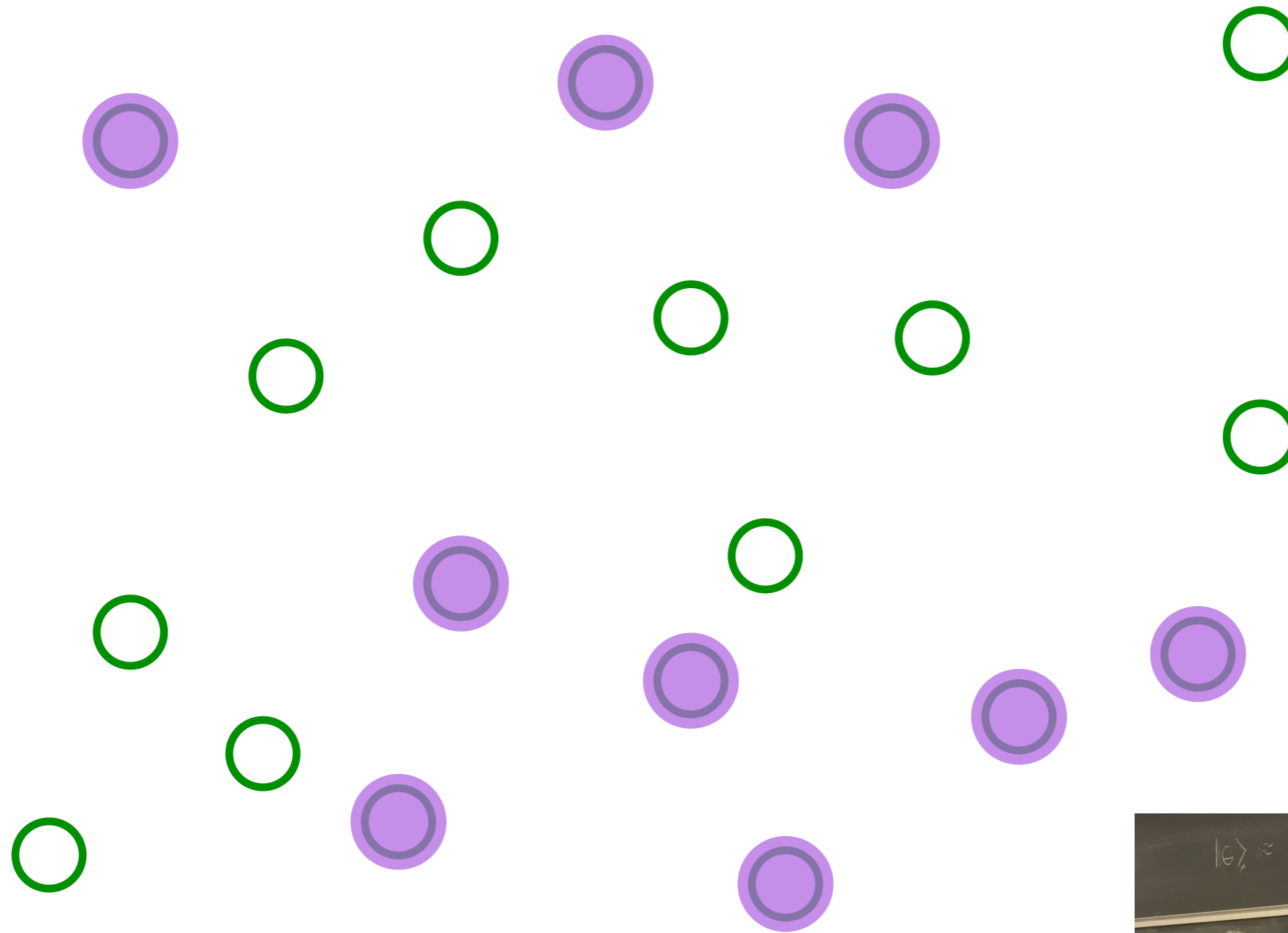


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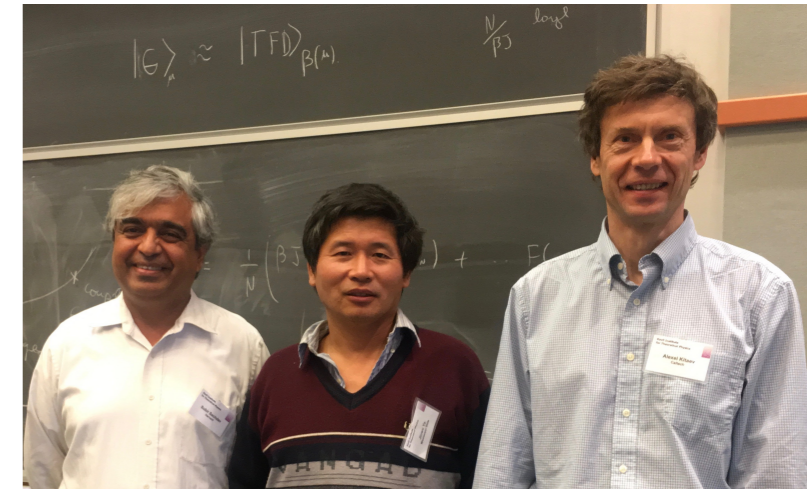


The SYK model

Sachdev, Ye (1993); Kitaev (2015)



Entangle electrons pairwise randomly



Complex multi-particle entanglement in the SYK model leads to a state without ‘quasiparticle’ excitations; *i.e.*

multiple excitations cannot be built by composing an elementary set of ‘quasiparticle’ excitations.

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Many-body chaos and thermal equilibration in the shortest possible Planckian time $\sim \frac{\hbar}{k_B T}$.

Quantum
entanglement

Black
holes

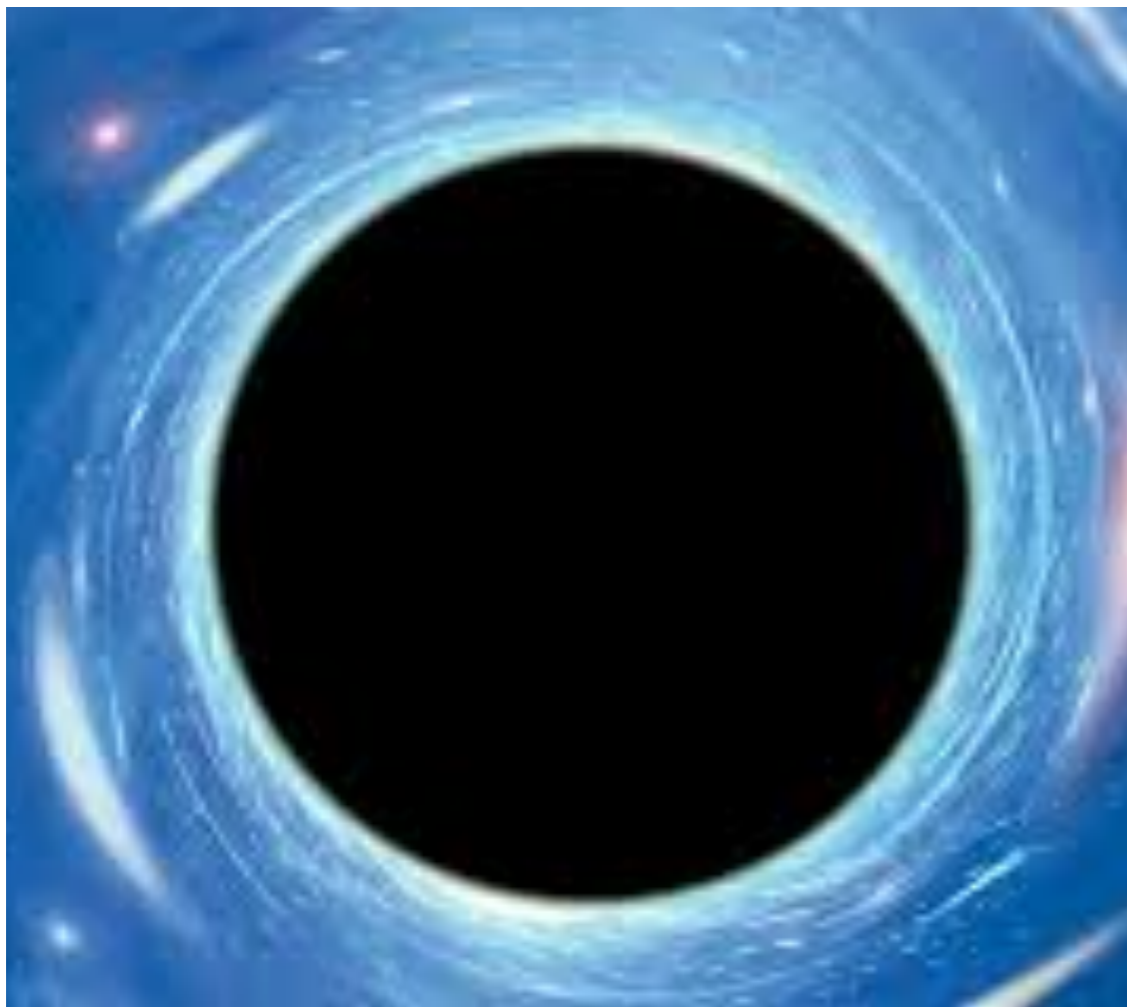
Hologram ?

A simple
many-particle
(SYK) model



Maxwell's electromagnetism and Einstein's general relativity allow black hole solutions with a net charge

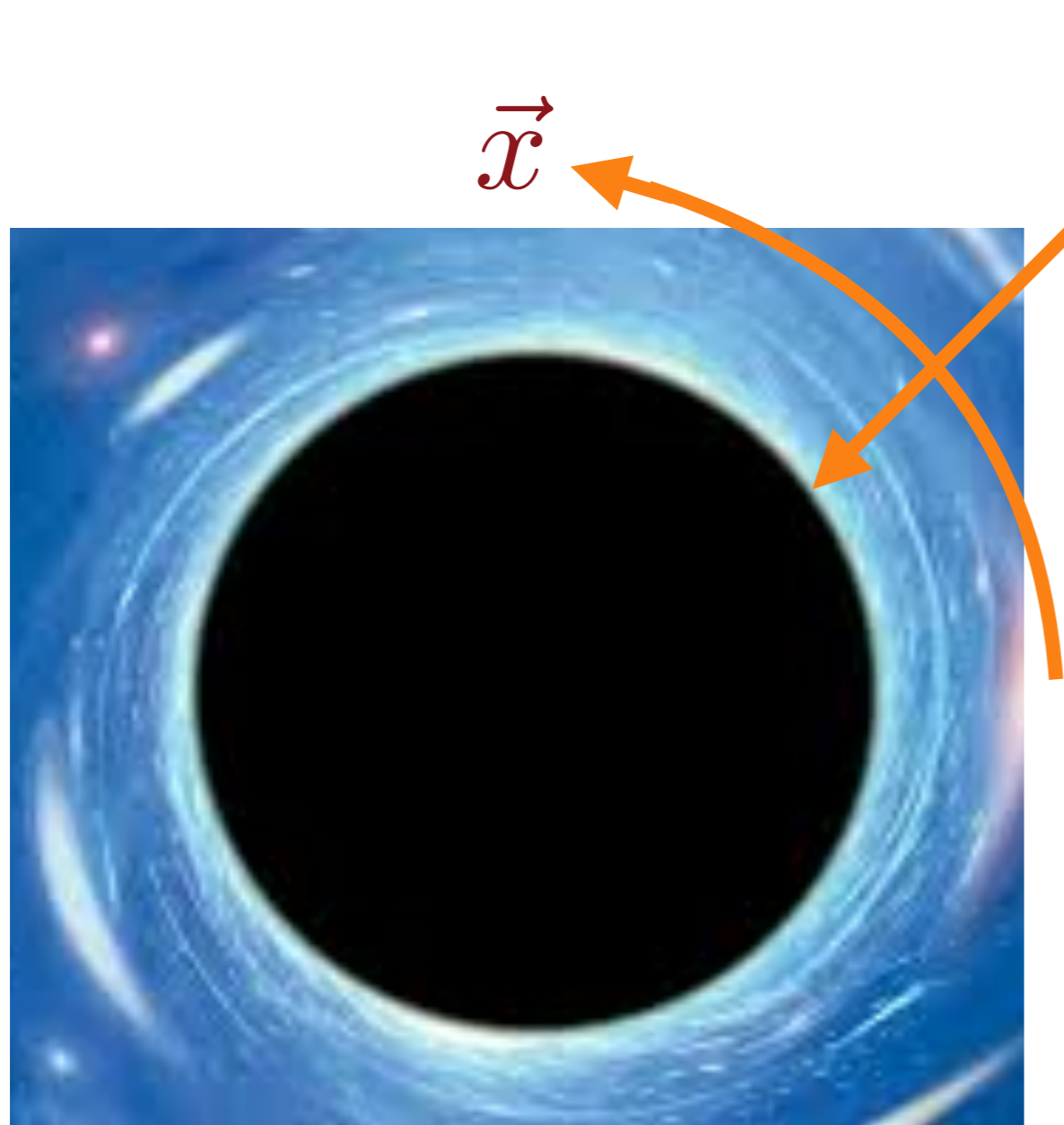
(Similar considerations also apply
to rapidly rotating black holes,
Moitra, Sake, Trivedi, Vishal (2019))





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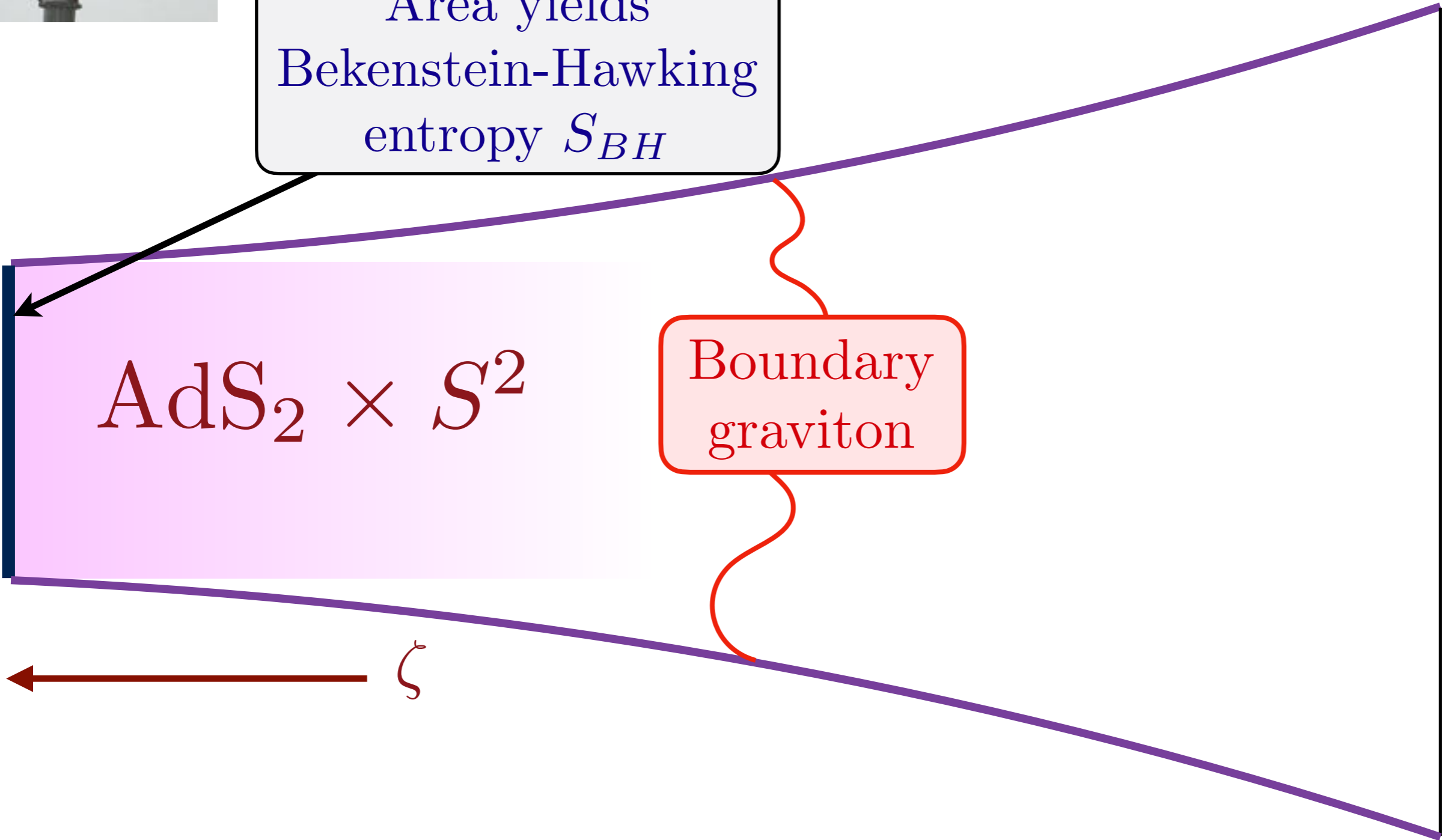


Zooming into the near-
horizon region of a
charged black hole at
low temperature, yields
a gravitational theory
in one space (ζ) and
one time dimension

SYK model and charged black holes



Horizon
Area yields
Bekenstein-Hawking
entropy S_{BH}



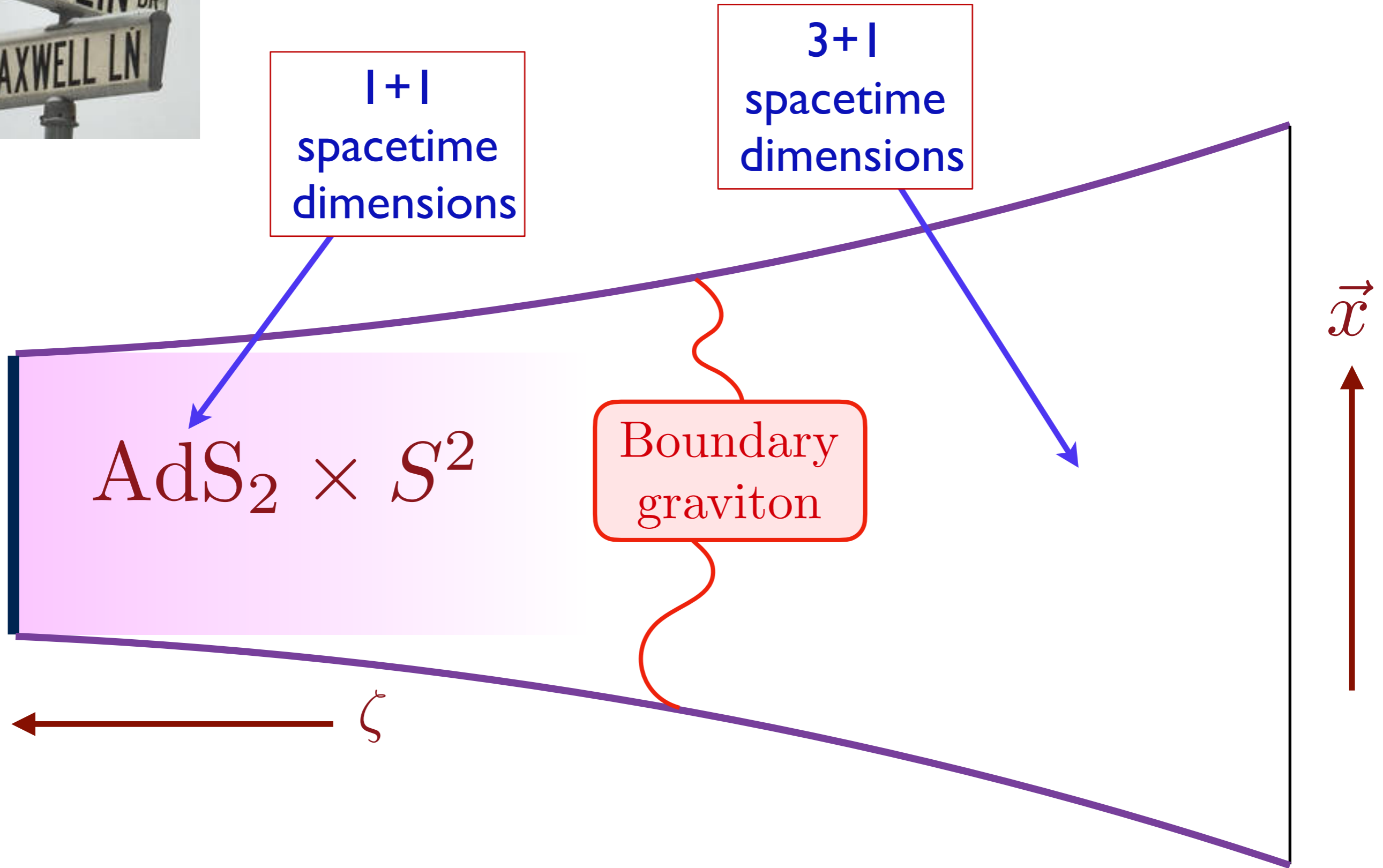
$AdS_2 \times S^2$

Boundary
graviton

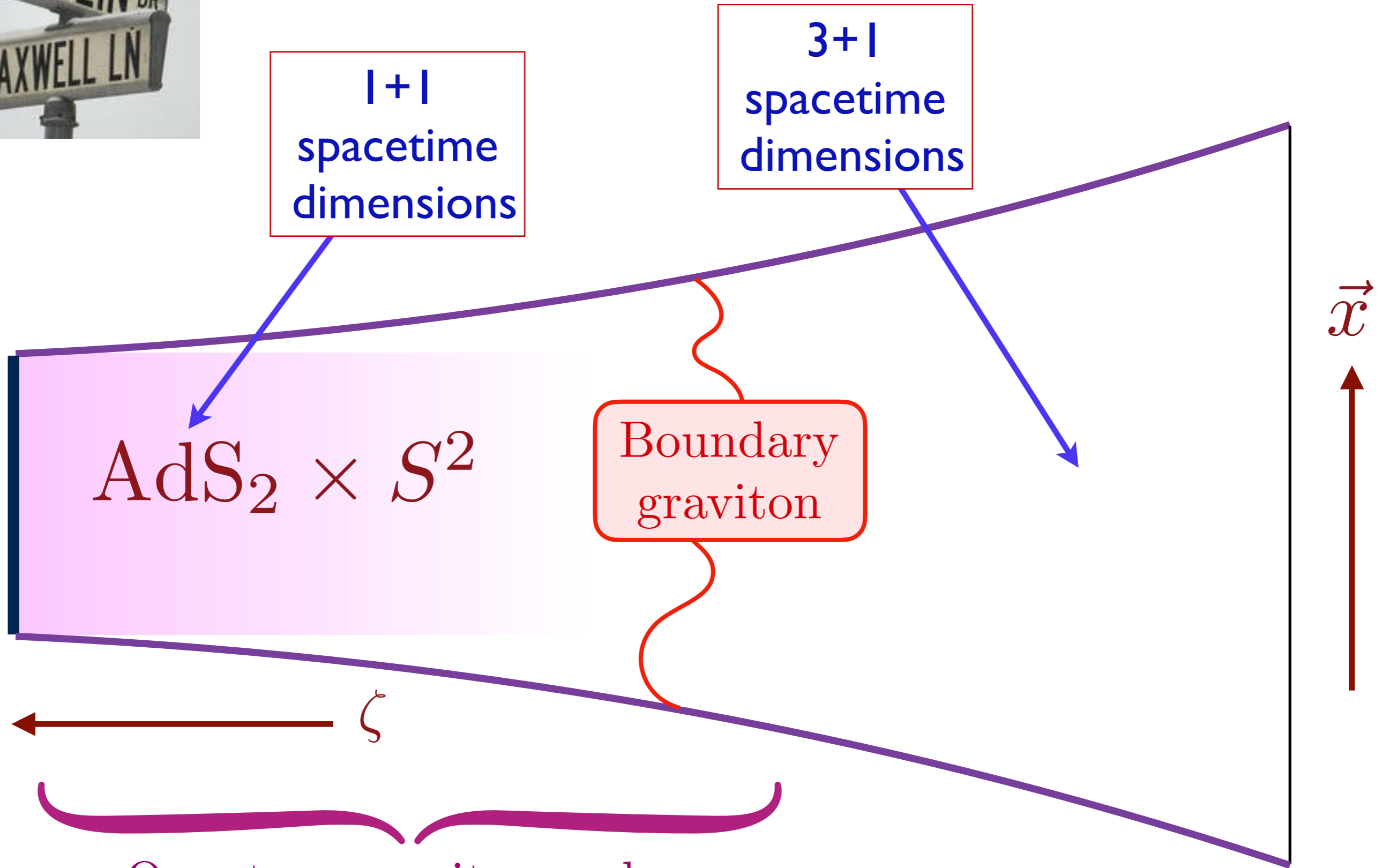
\vec{x}

ζ

SYK model and charged black holes



SYK model and charged black holes



Quantum gravity can be exactly solved in this region!

SYK model and charged black holes

Thermodynamics of charged quantum black holes

$$\int \mathcal{D}g_{\mu\nu} \exp \left(-\frac{1}{\hbar} \mathcal{S}_{\text{Einstein-Maxwell theory}}^{(3+1)}[g_{\mu\nu}] \right) T \rightarrow 0,$$
$$\approx \int \mathcal{D}g_{\mu\nu} \exp \left(-\frac{1}{\hbar} \mathcal{S}_{\text{Gravity of AdS}_2 \text{ and boundary}}^{(1+1)}[g_{\mu\nu}] \right)$$

SYK model and charged black holes

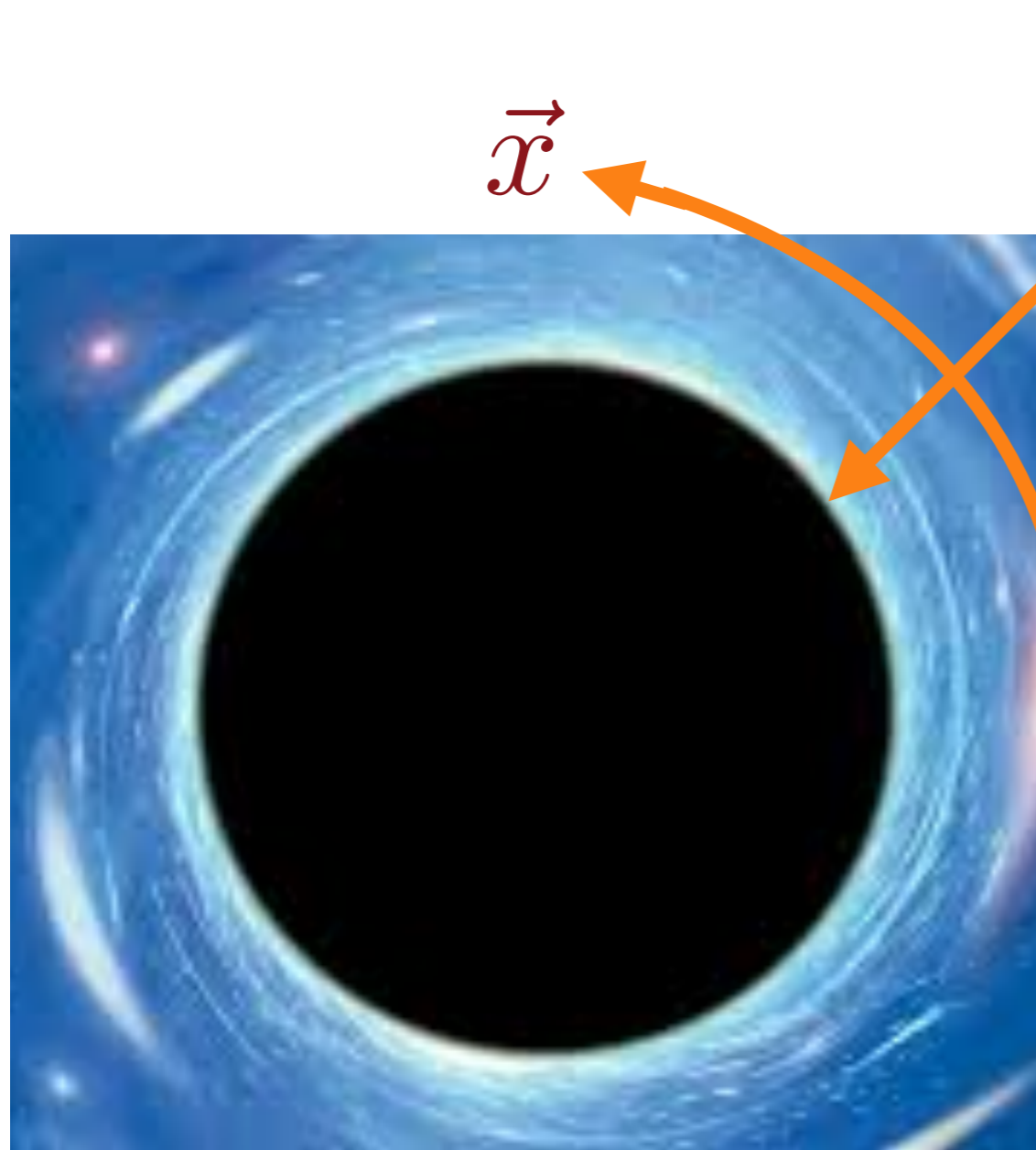
Thermodynamics of charged quantum black holes

$$\begin{aligned} & \int \mathcal{D}g_{\mu\nu} \exp \left(-\frac{1}{\hbar} \mathcal{S}_{\text{Einstein-Maxwell theory}}^{(3+1)}[g_{\mu\nu}] \right) \quad T \rightarrow 0, \\ & \approx \int \mathcal{D}g_{\mu\nu} \exp \left(-\frac{1}{\hbar} \mathcal{S}_{\text{Gravity of AdS}_2 \text{ and boundary}}^{(1+1)}[g_{\mu\nu}] \right) \\ & = \exp(S_{BH}) \times \exp \left(-\frac{1}{T} \times \text{Free energy of SYK model} \right) \end{aligned}$$

The hologram of the 1+1 dimensional gravity near the horizon of a charged black hole is the 0+1 dimensional SYK model



Maxwell's electromagnetism
and Einstein's general relativity
allow black hole solutions with a net charge



The near-horizon
1+1D-gravity theory is
precisely that of the
low T limit of the
SYK models

Quantum entanglement

A simple many-particle (SYK) model

Charged black holes

Low temperatures

Quantum gravity in 1+1 dimensions

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Many-body chaos and
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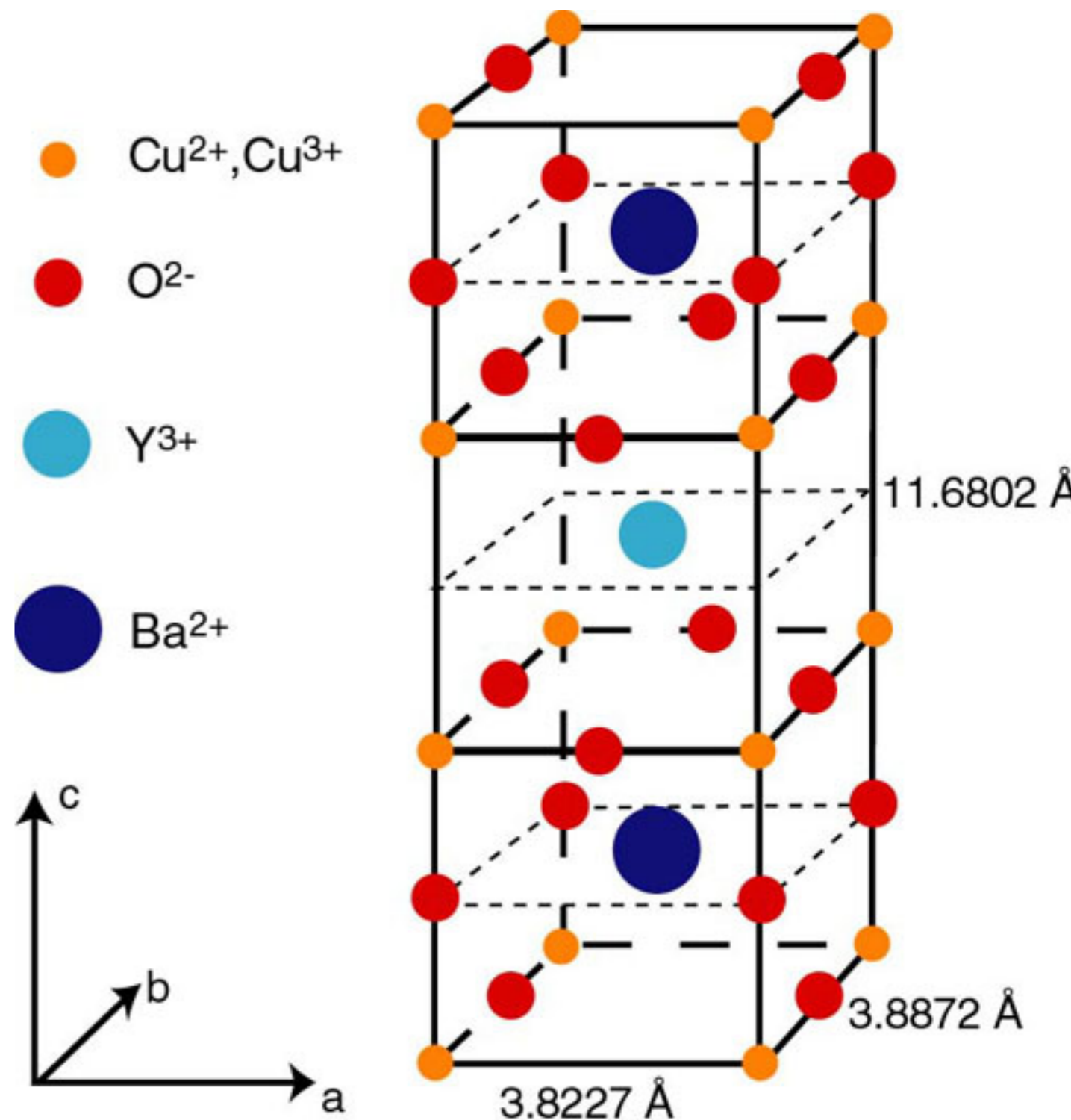
**Quantum
entanglement**

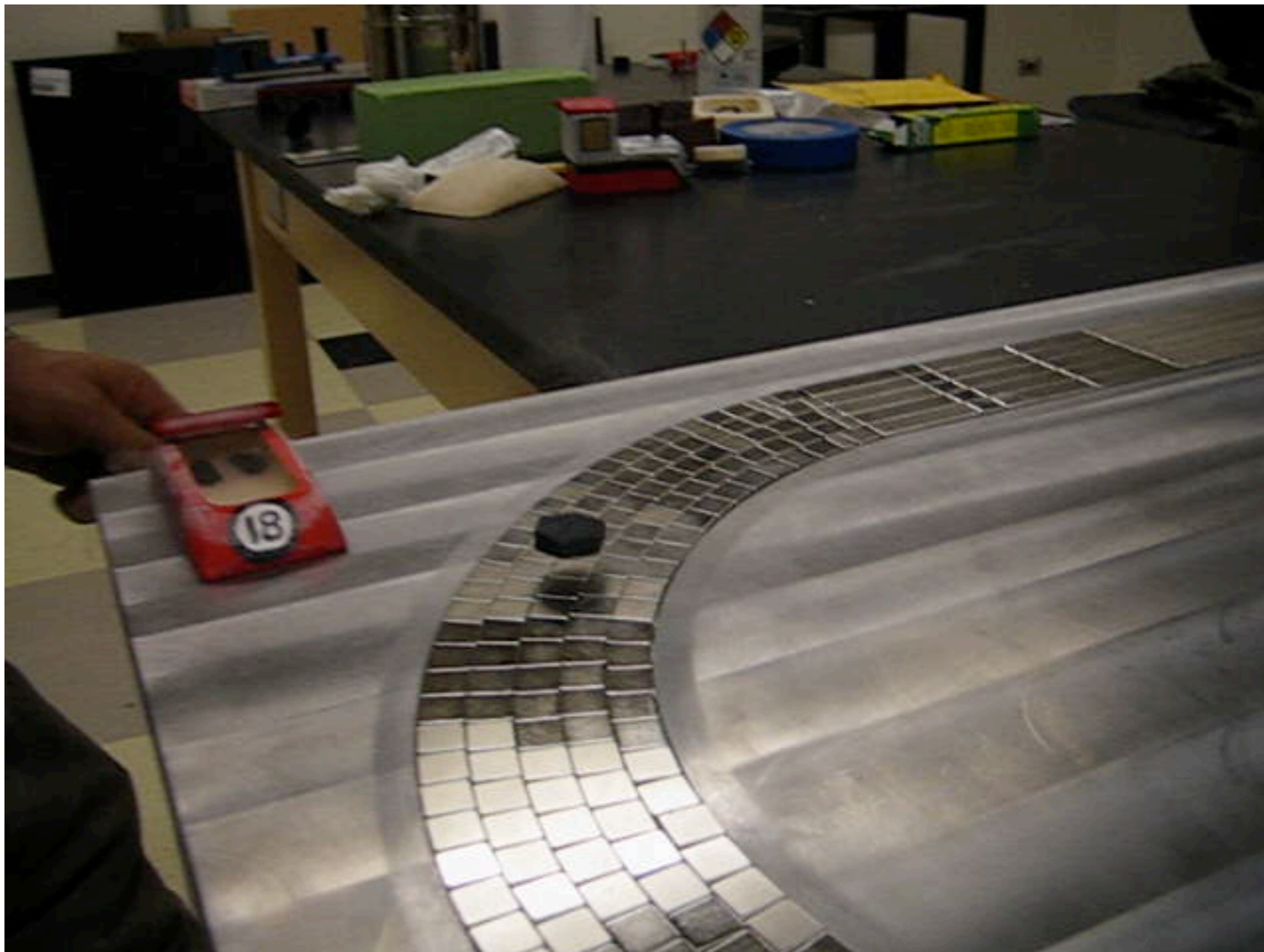
**Charged
black holes**

**A simple
many-particle
(SYK) model**

**Copper-based
superconductors**

High temperature superconductors

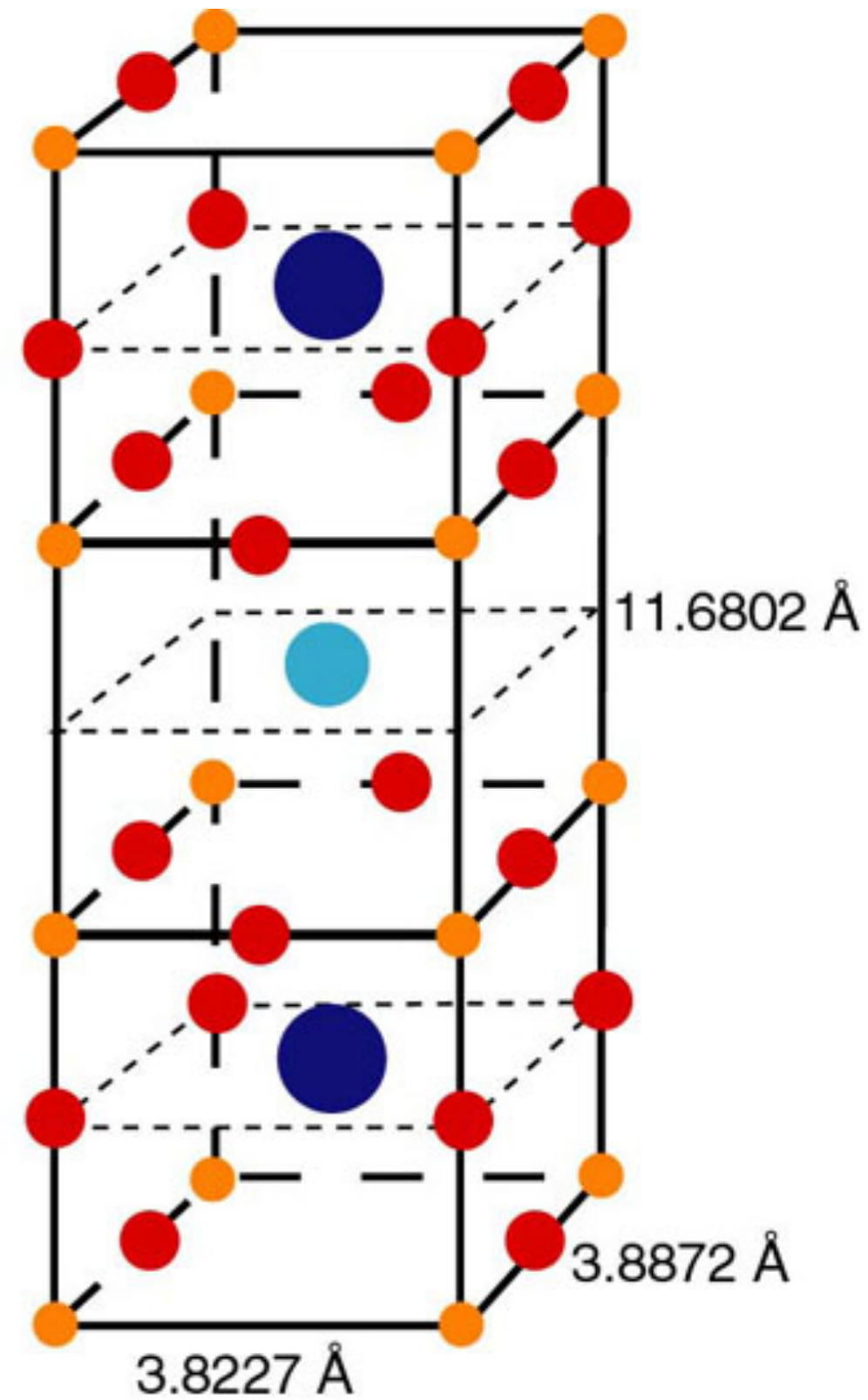
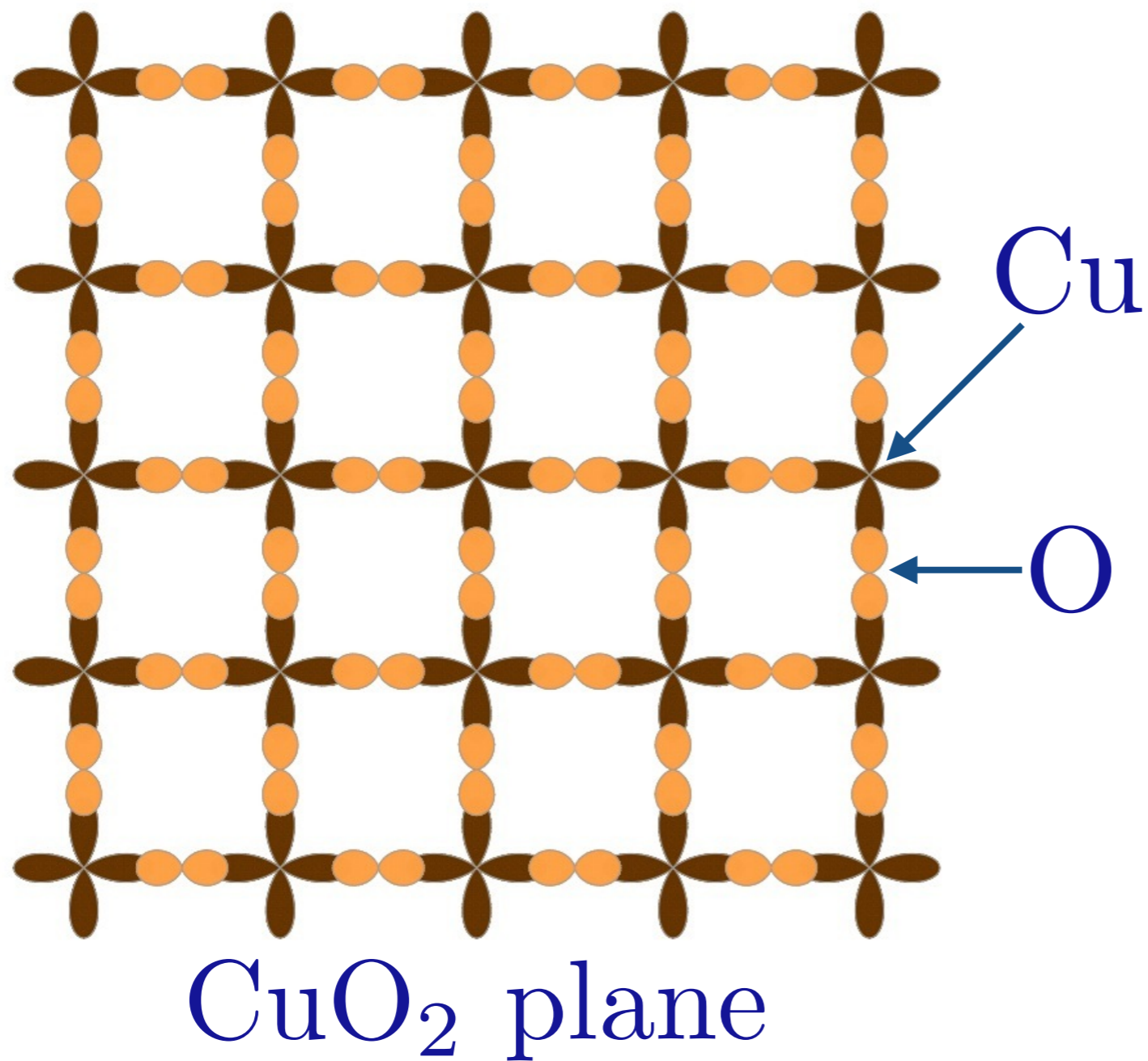




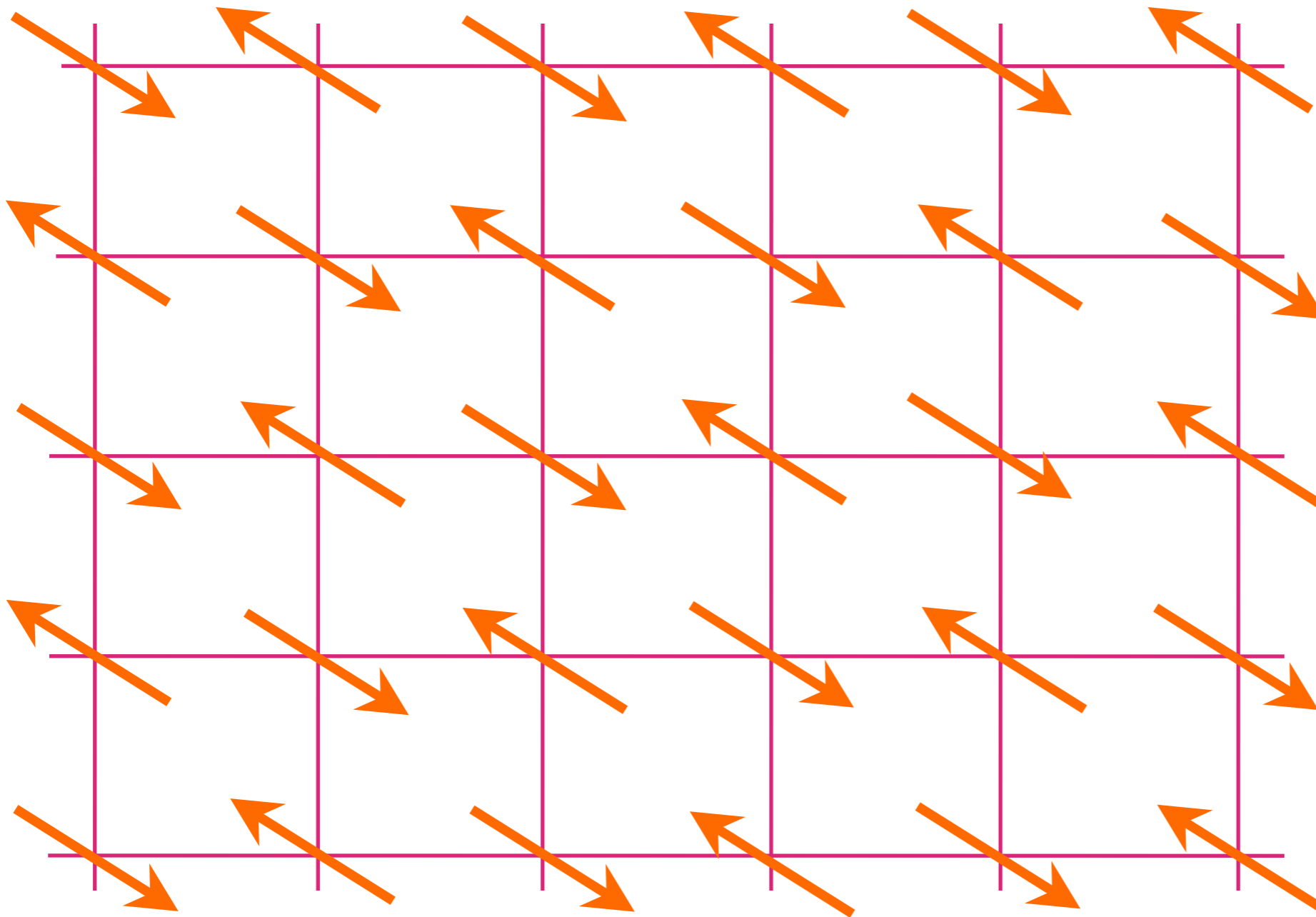
Nd-Fe-B magnets, YBaCuO superconductor

Julian Hetel and Nandini Trivedi, Ohio State University

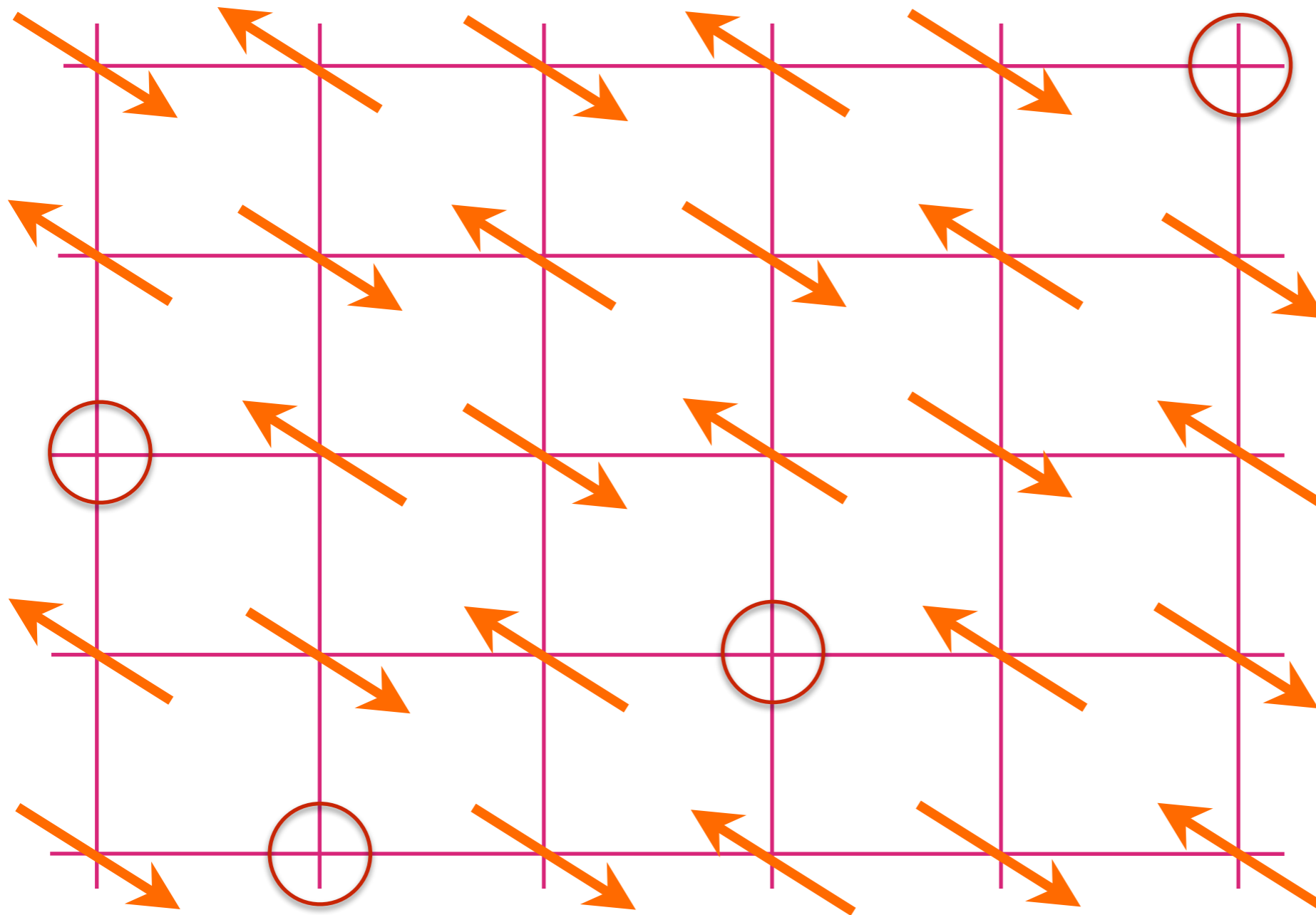
High temperature superconductors



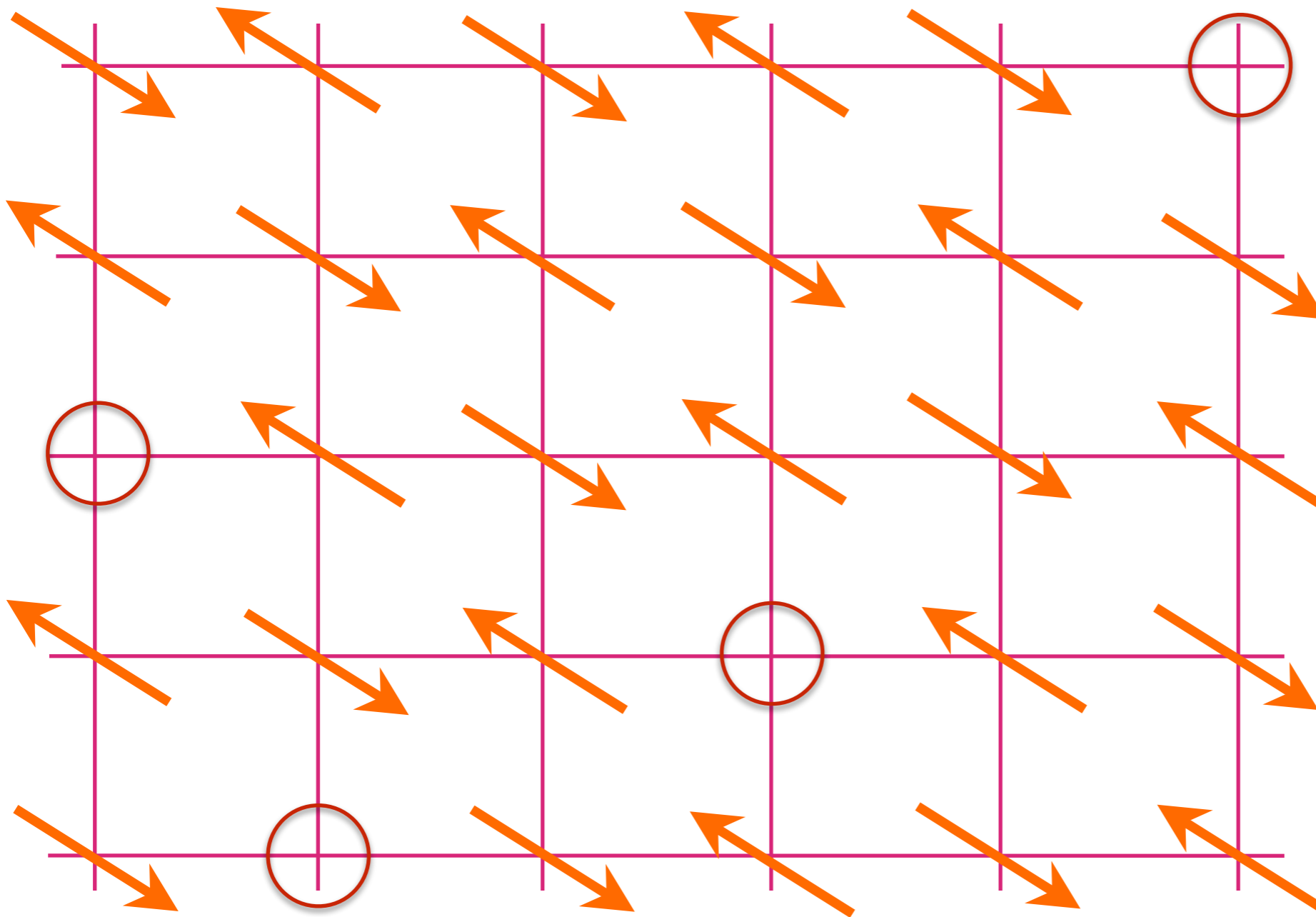
Insulating antiferromagnet



Antiferromagnet doped with hole density p

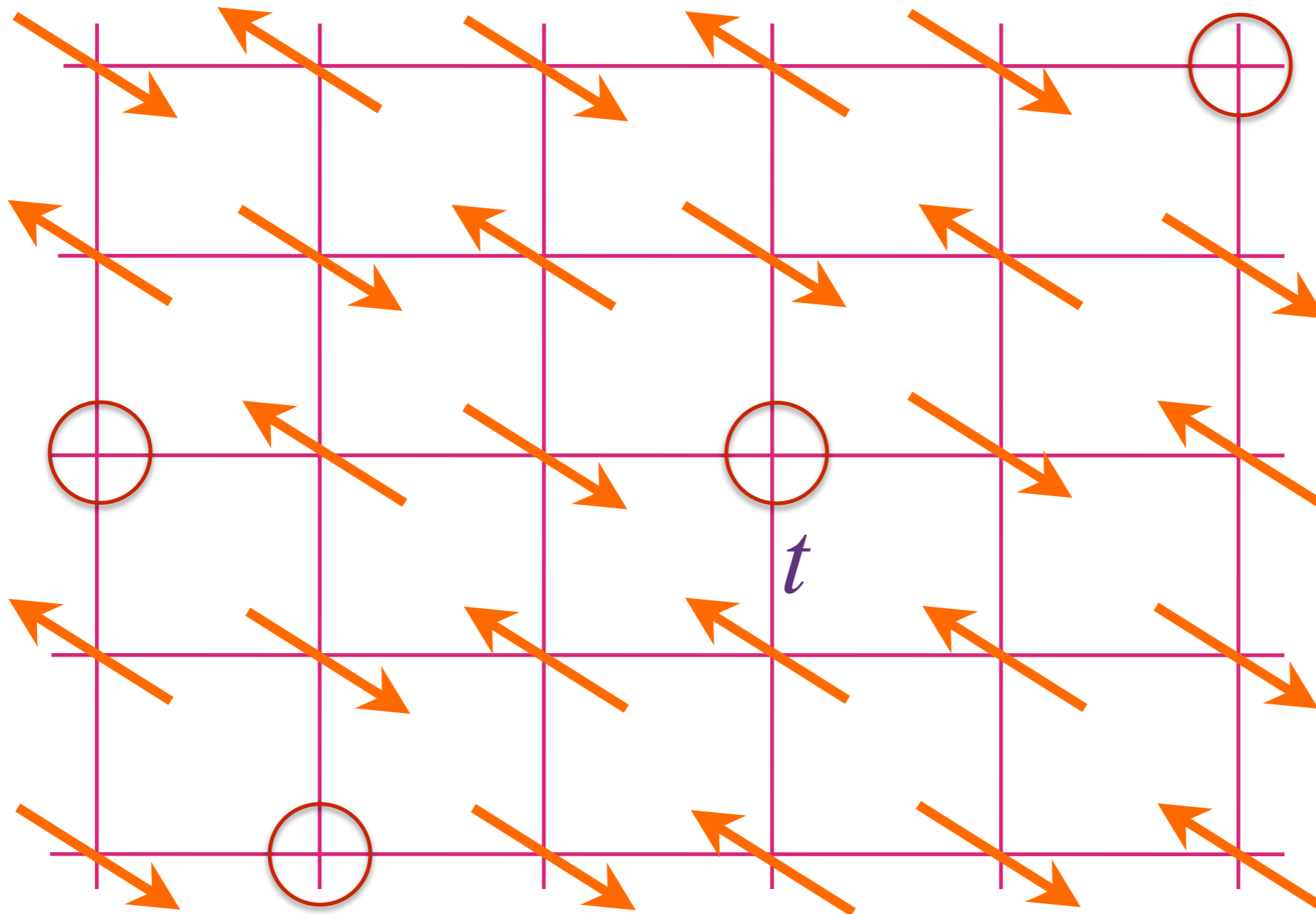


Real-space view



p mobile holes in a background of
fluctuating spins

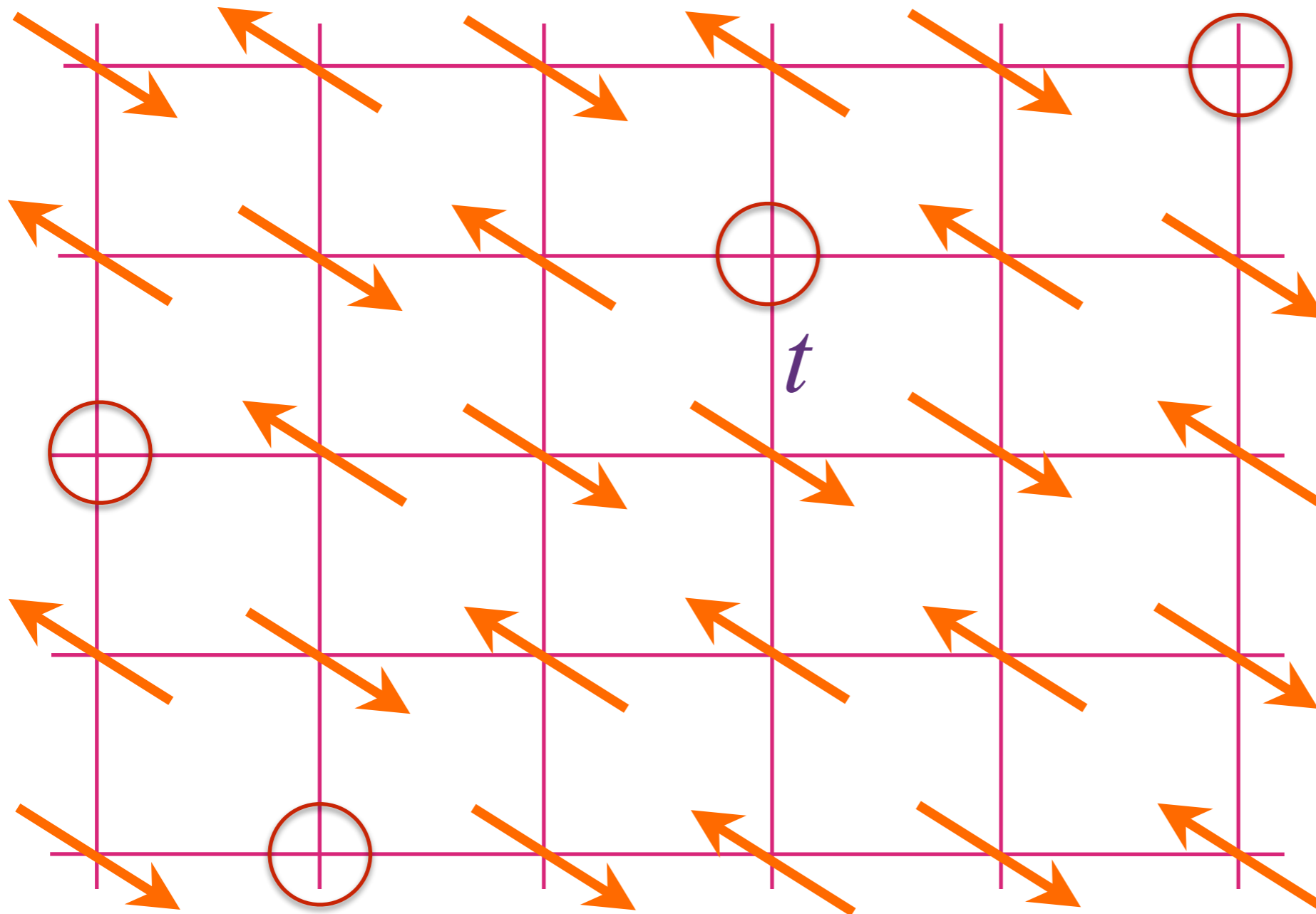
Real-space view



Baskaran,
Anderson (1988)

p mobile holes in a background of
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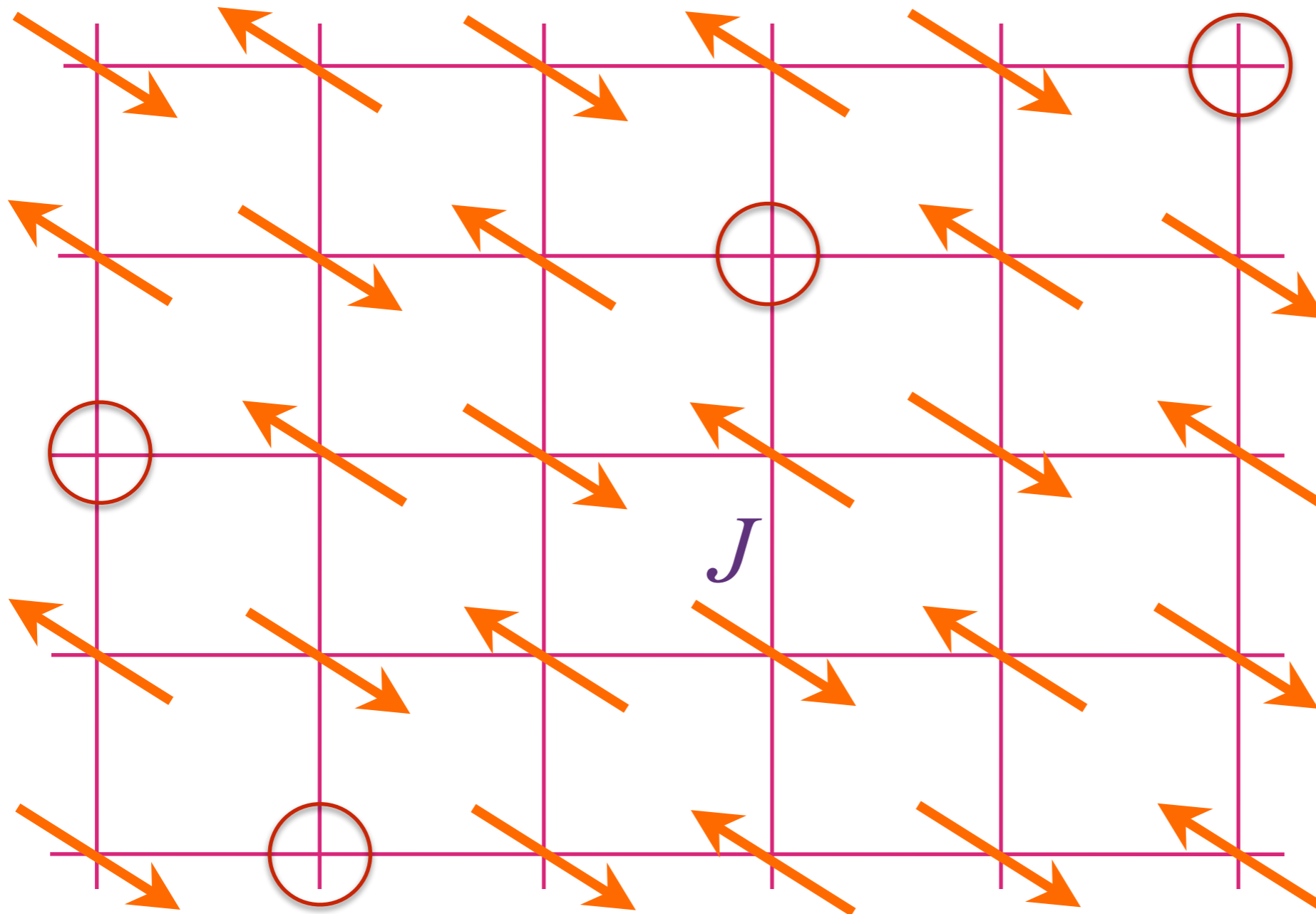
Real-space view



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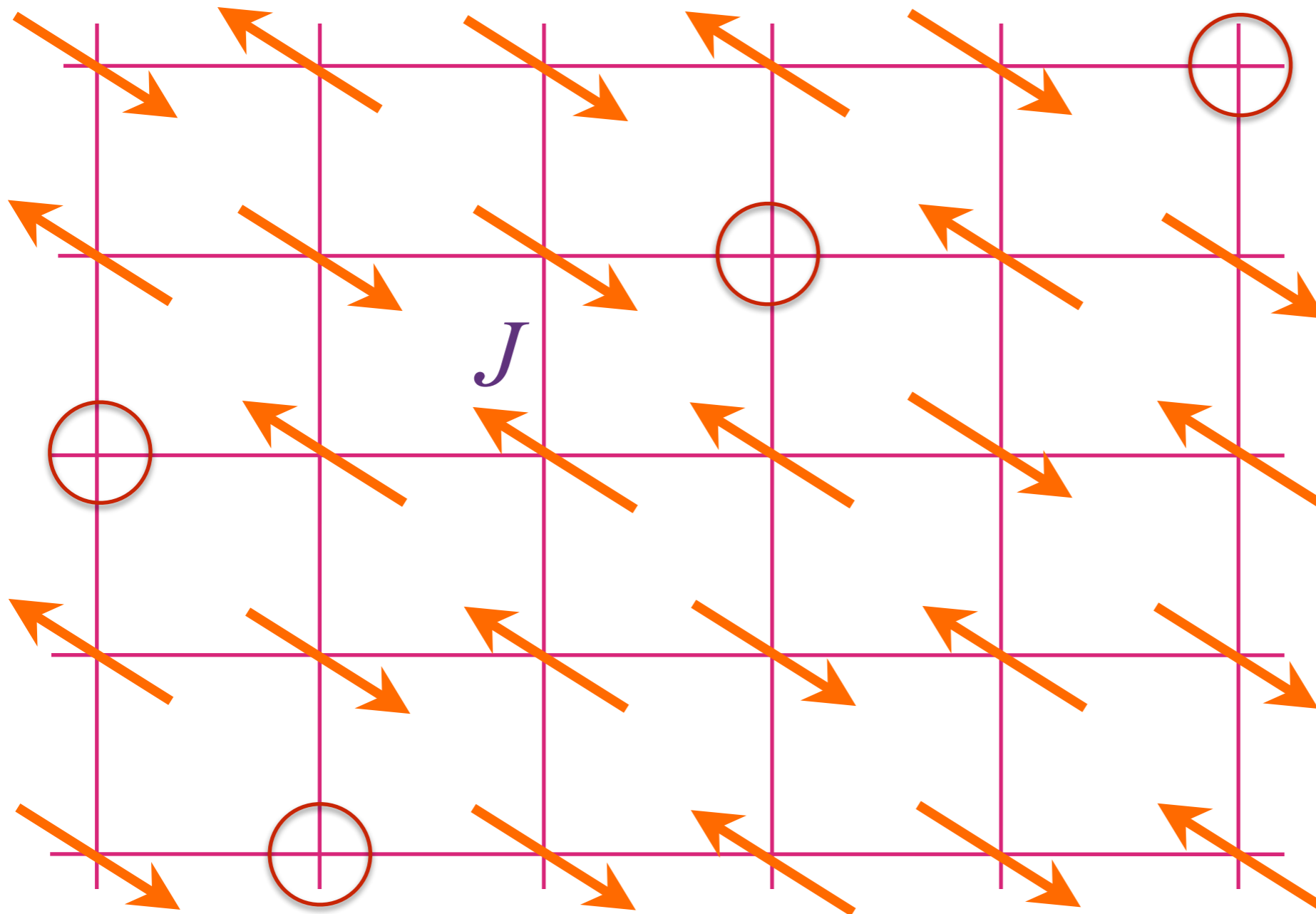
Real-space view



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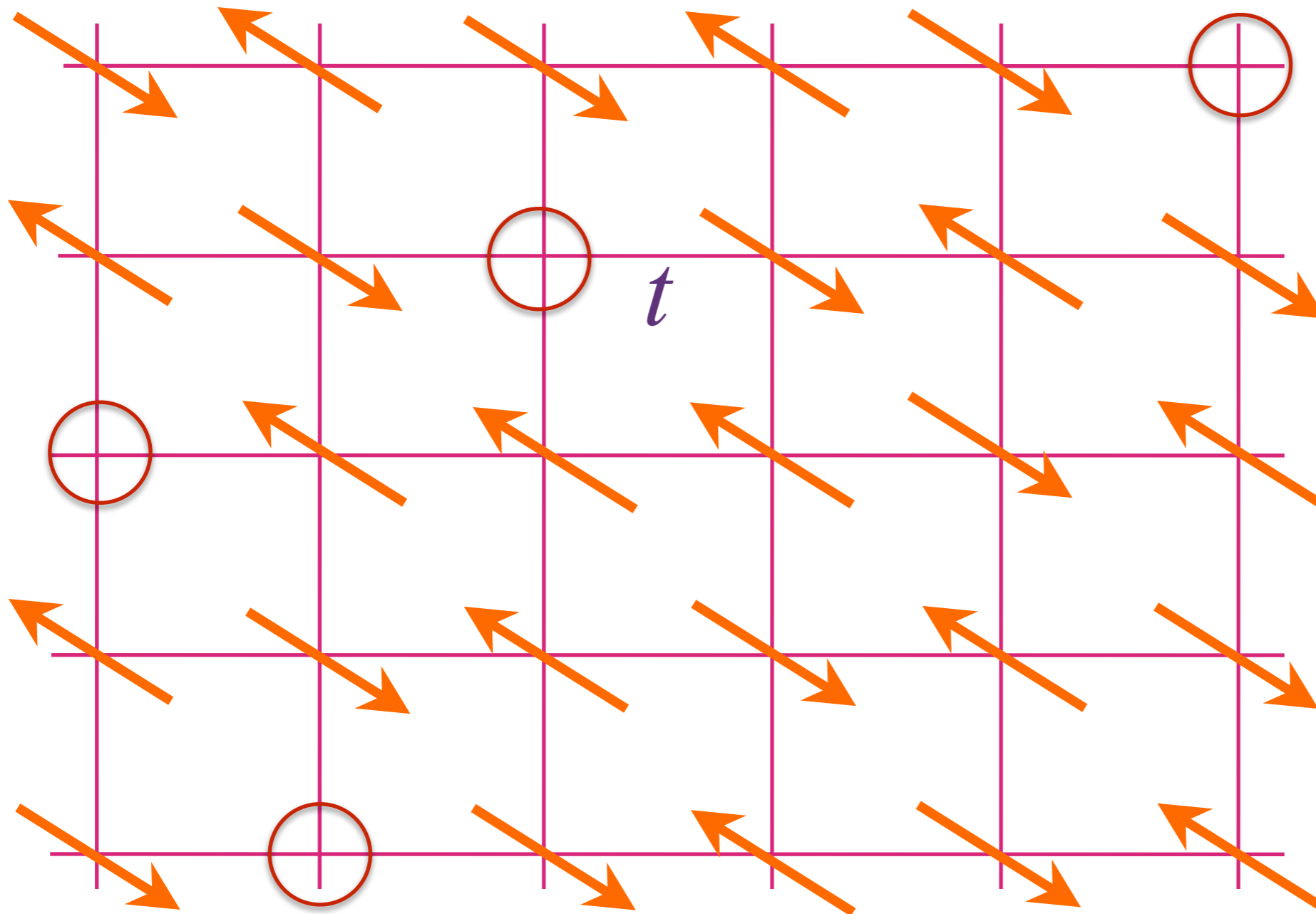
Real-space view



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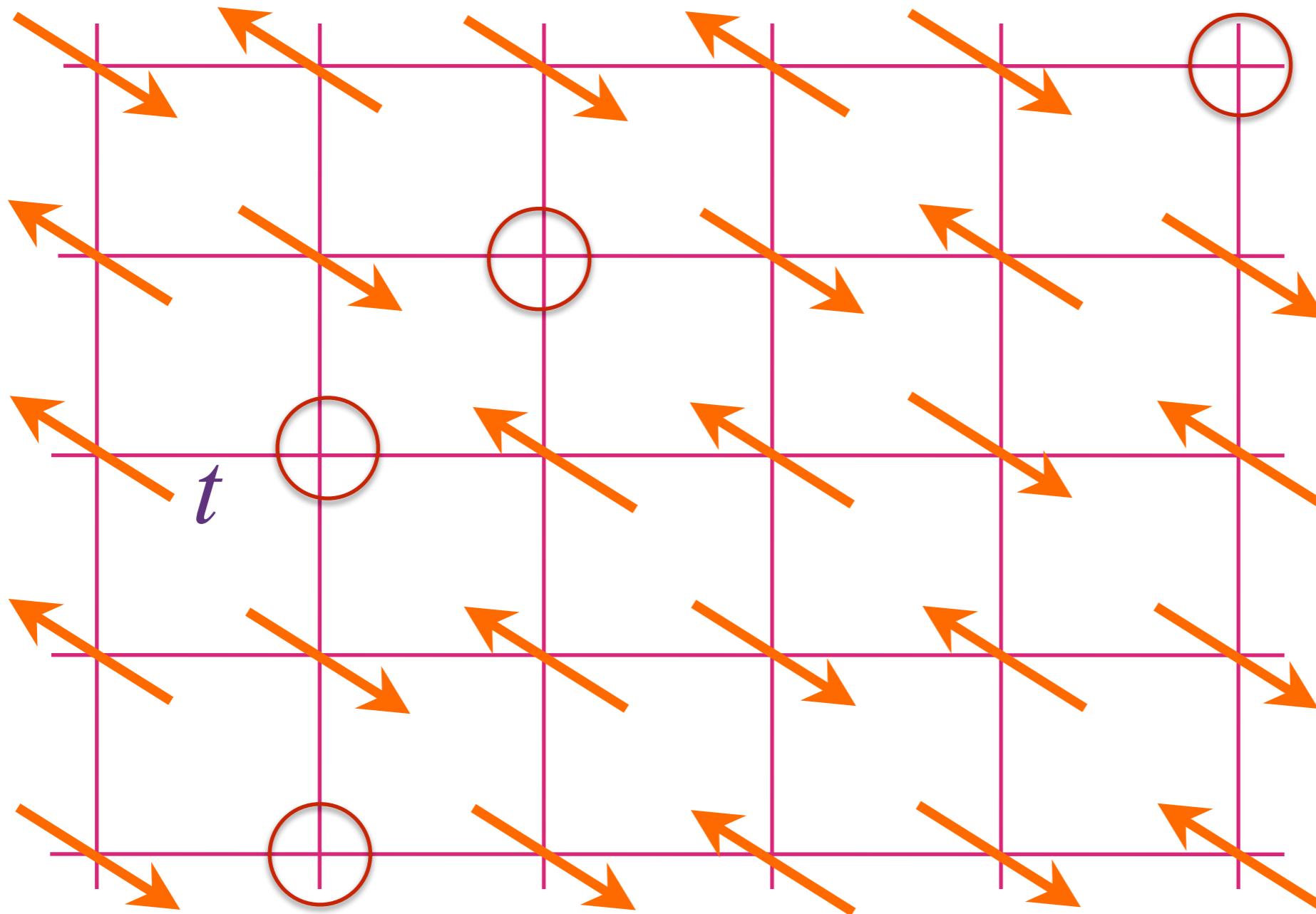
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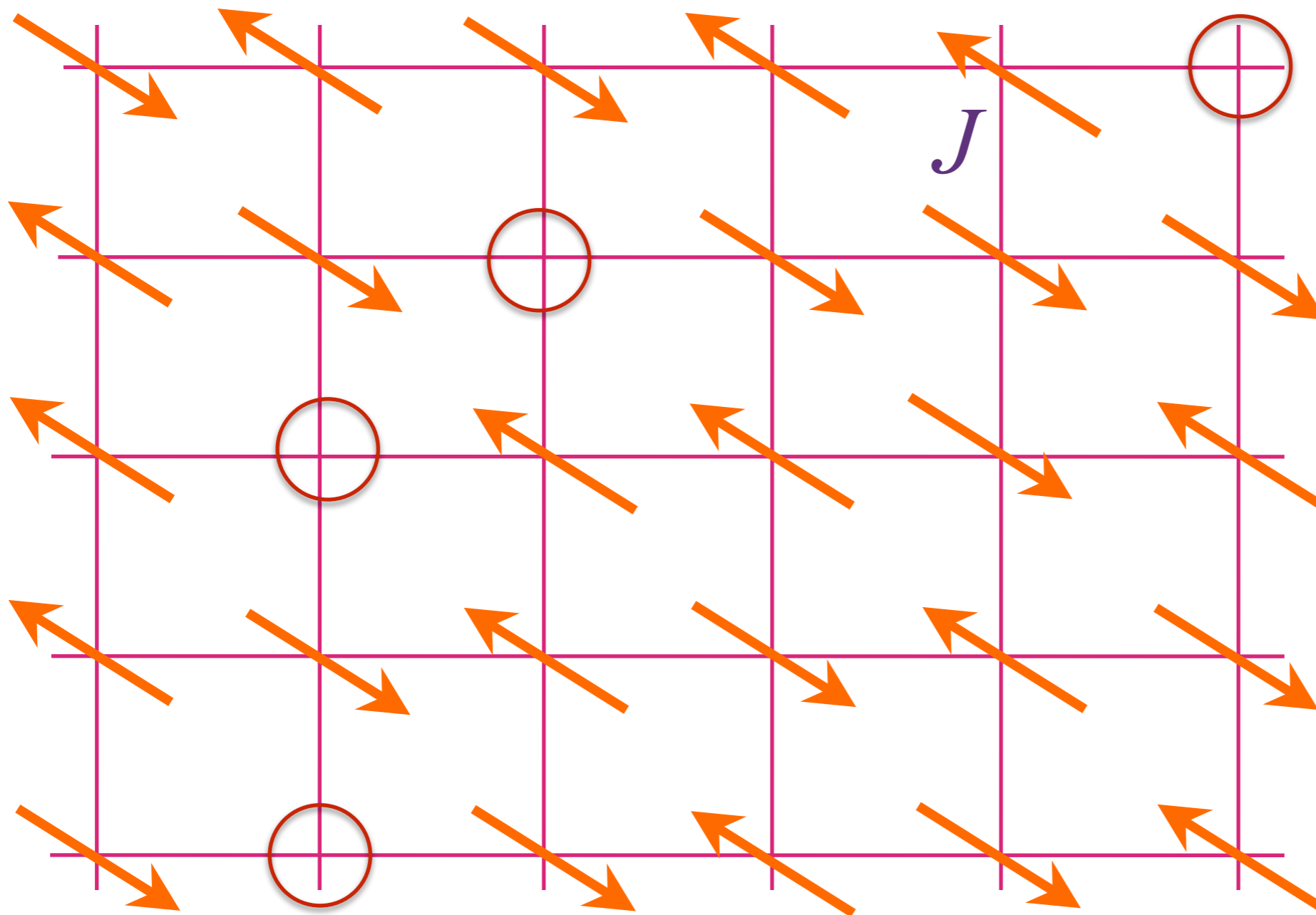
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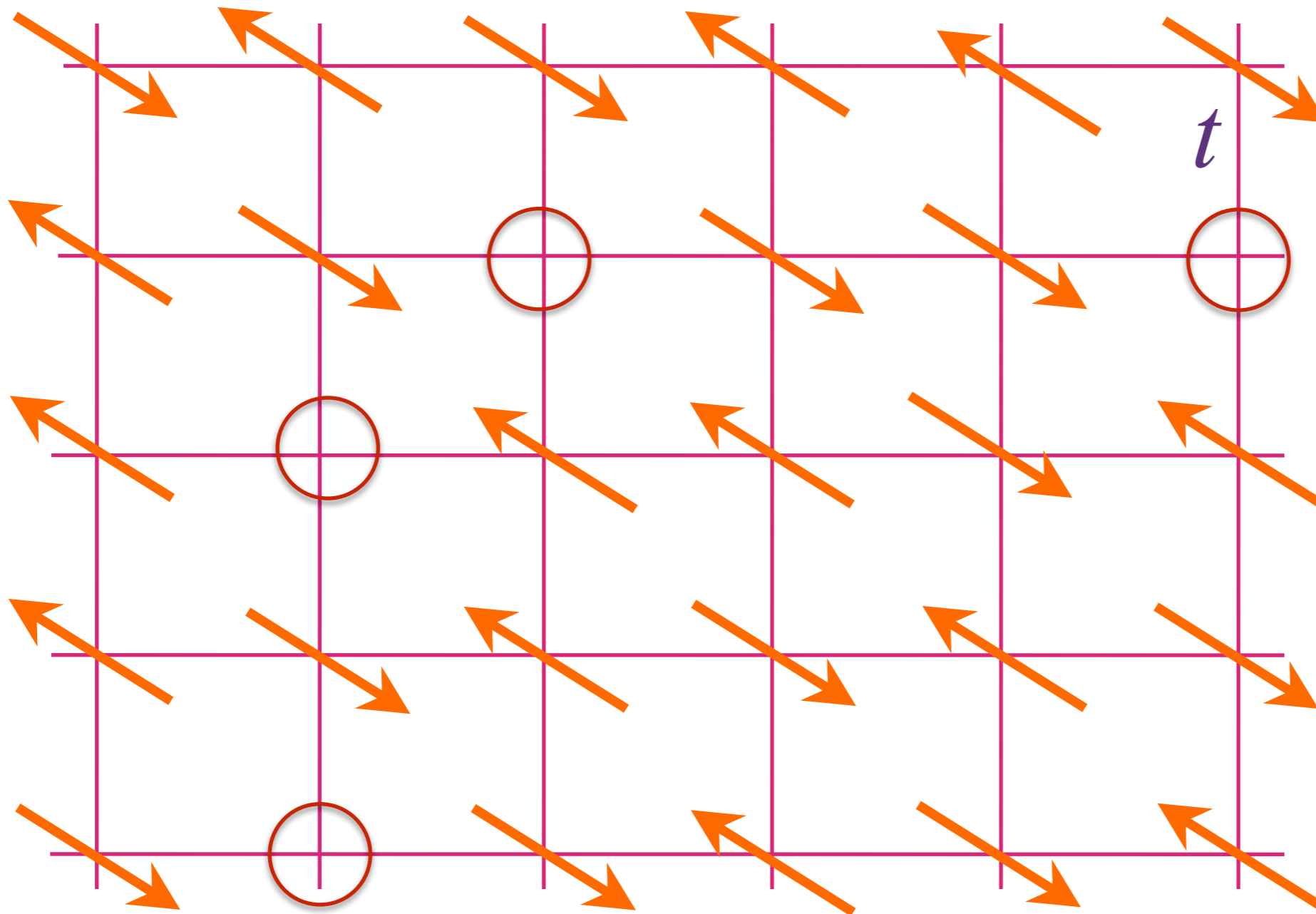
Real-space view



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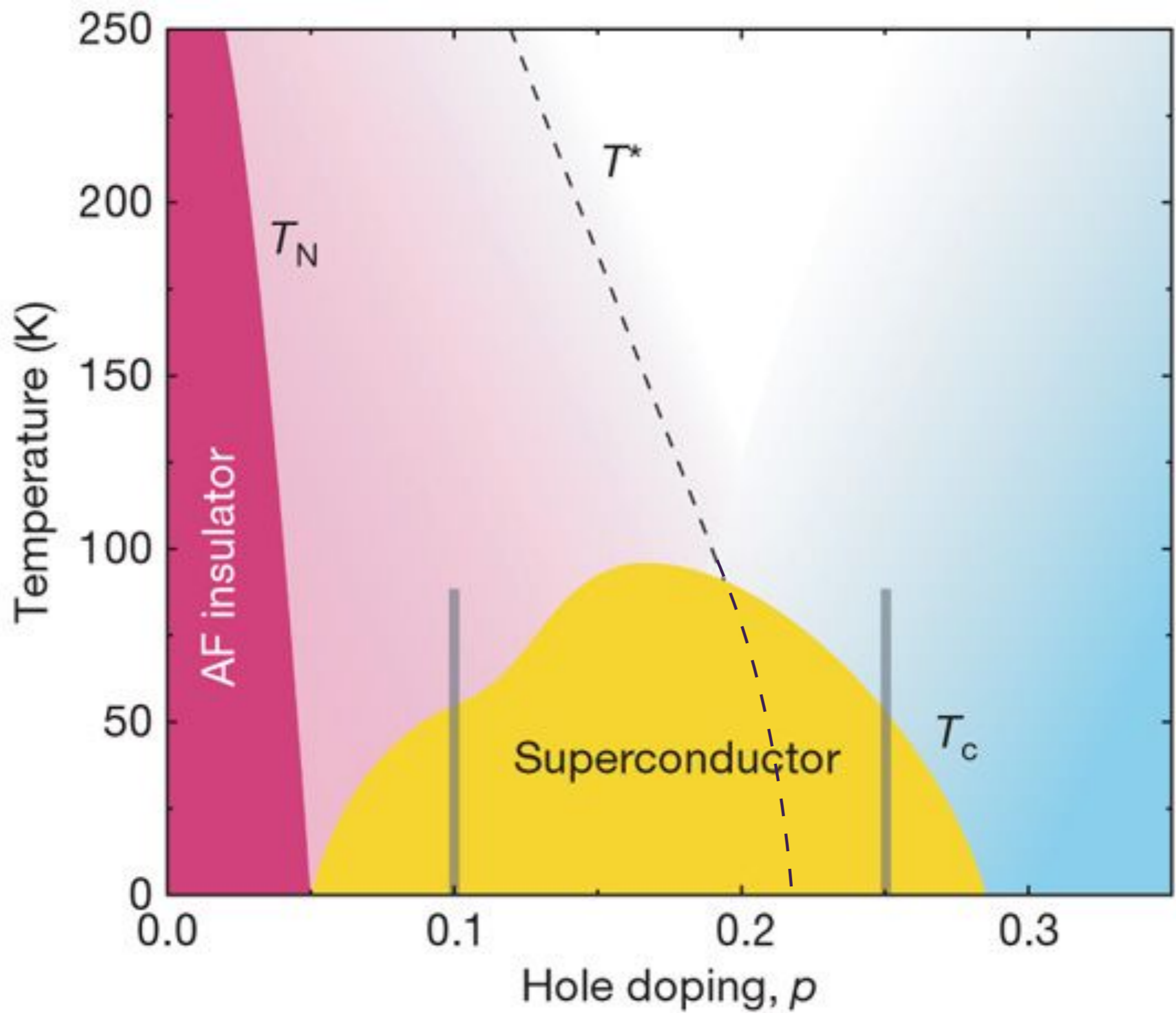
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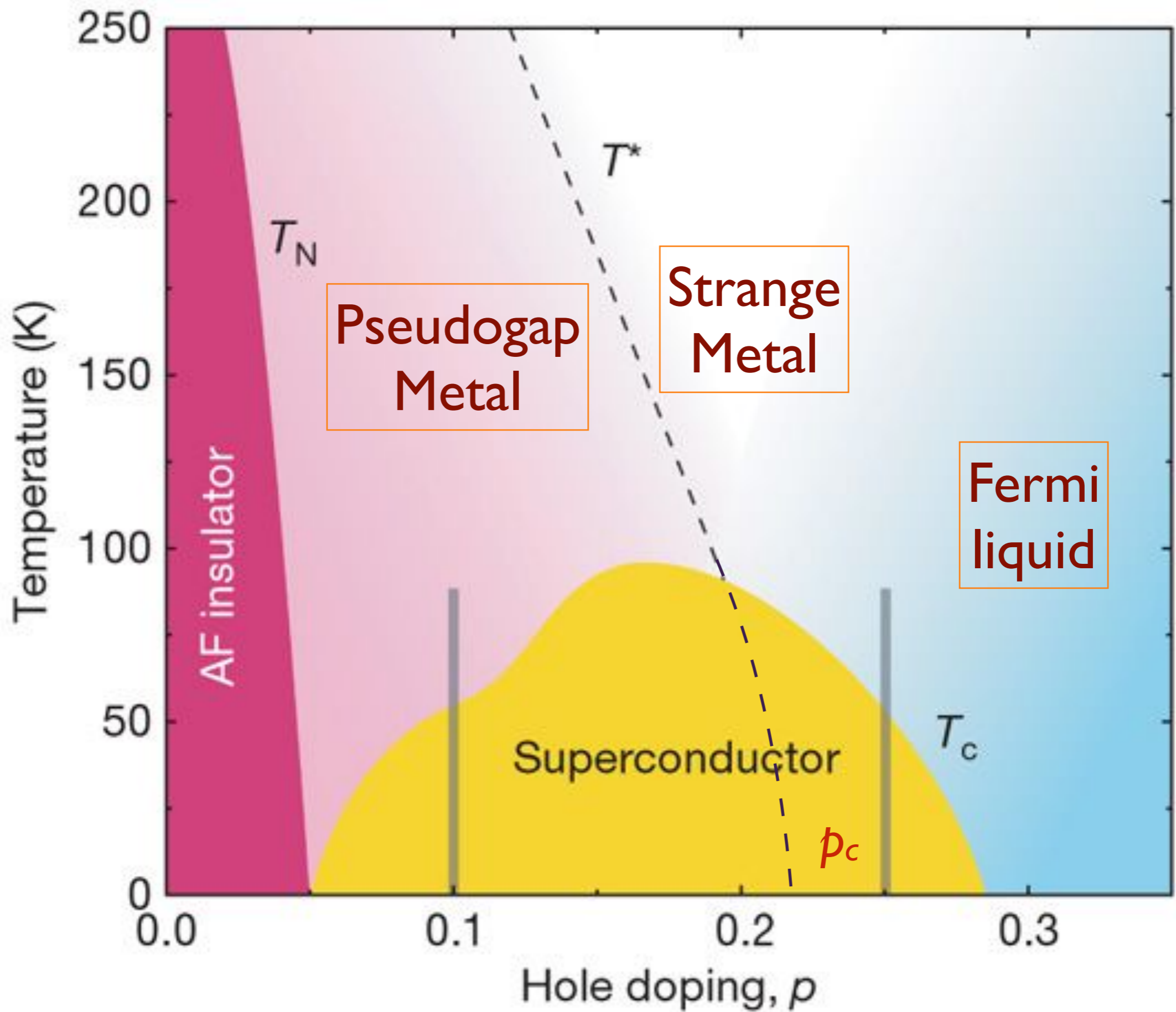
Real-space view

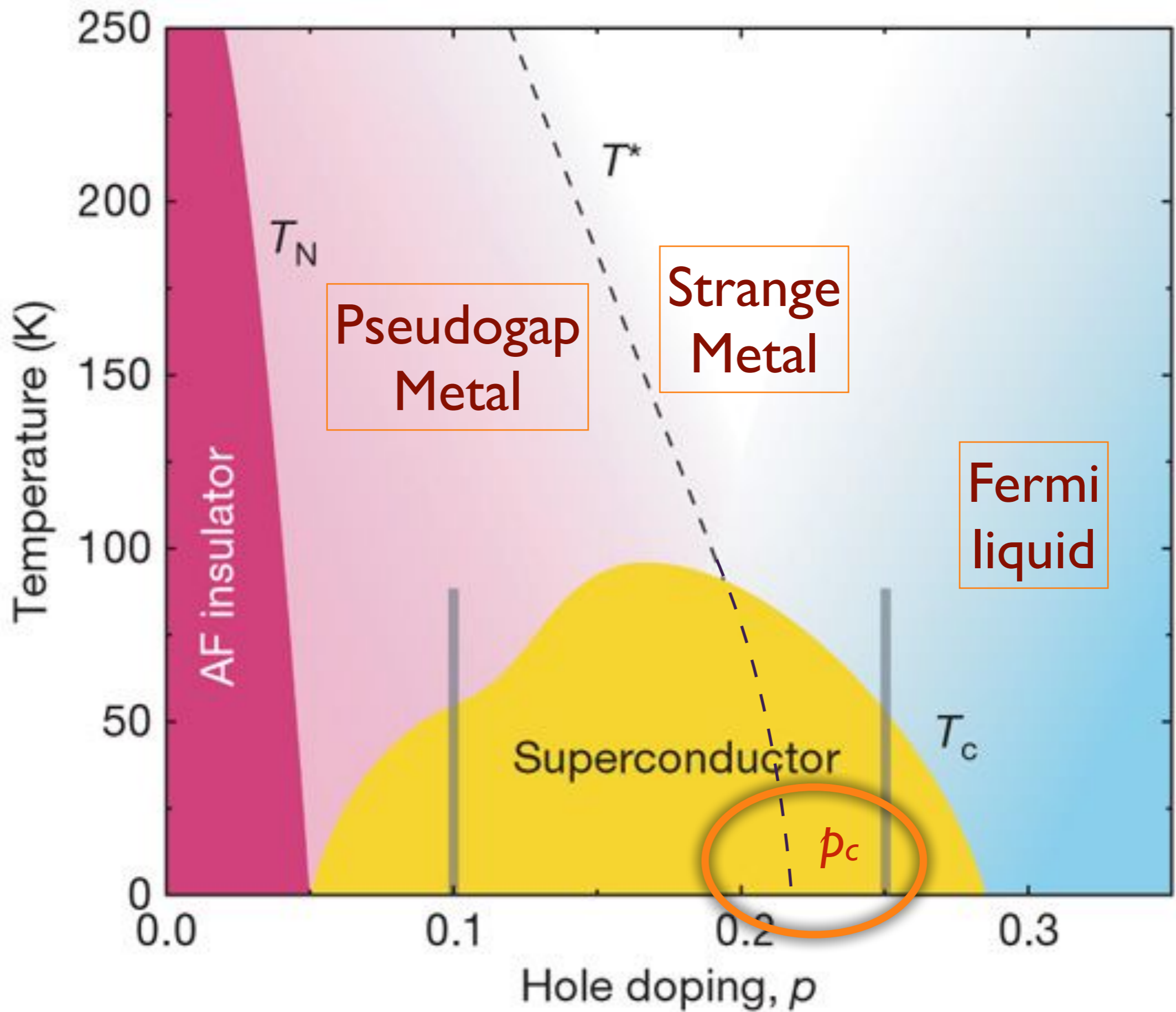


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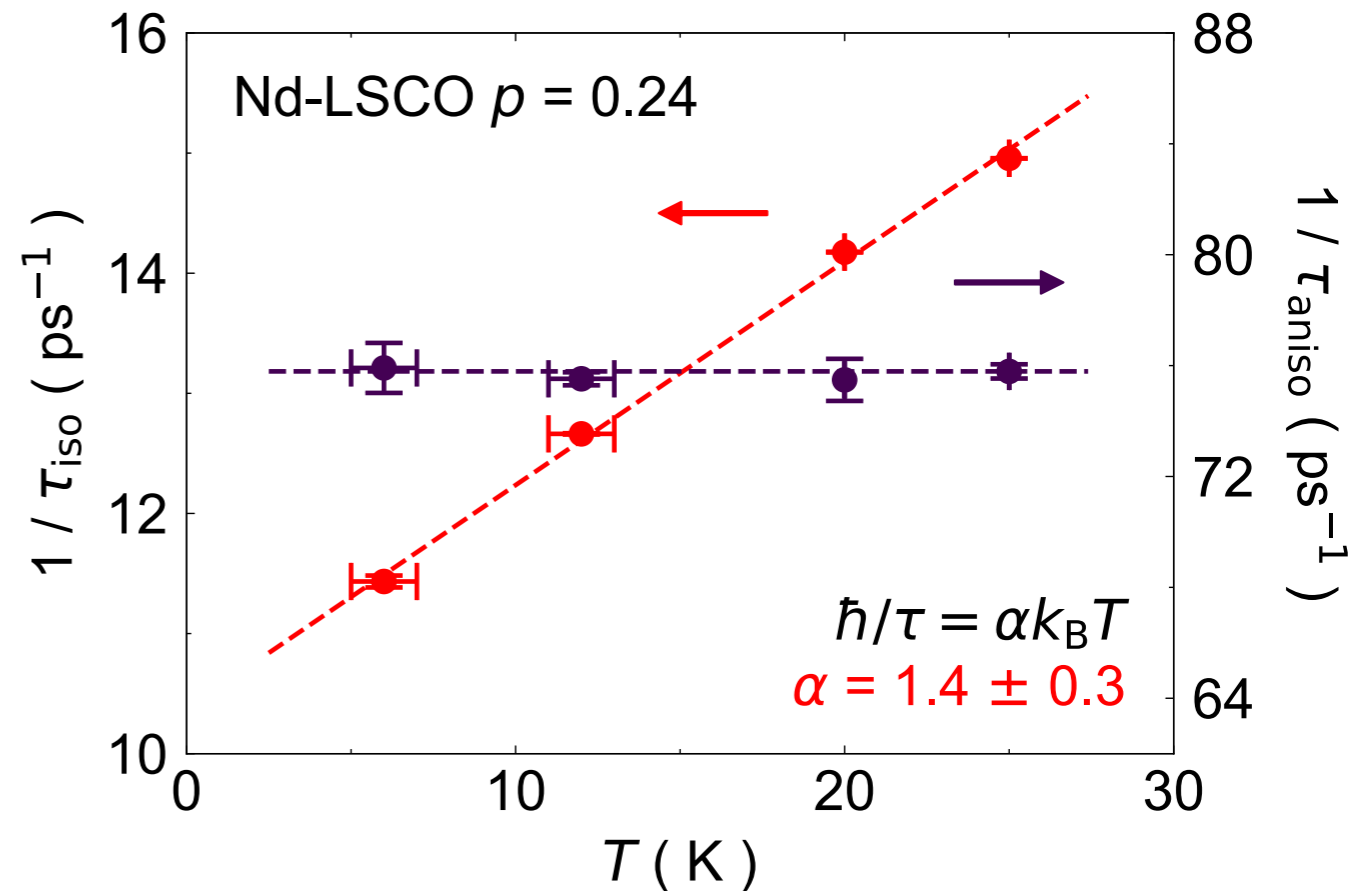
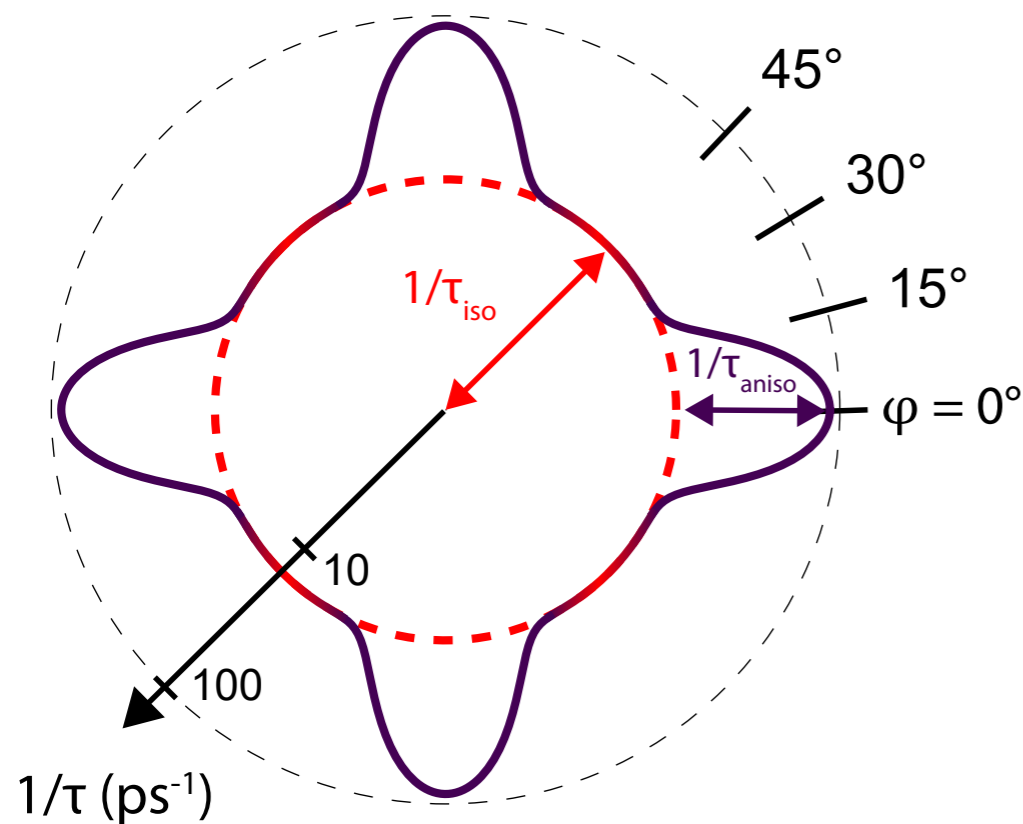




Measurement of the Planckian Scattering Rate

G. Grissonnanche, Y. Fang, A. Legros, S. Verret, F. Laliberté, C. Collignon, J. Zhou, D. Graf, P. Goddard, L. Taillefer, B. J. Ramshaw, arXiv:2011.13054

Angle-dependent magnetoresistance in Nd-LSCO near $p = p_c \approx 0.23$.

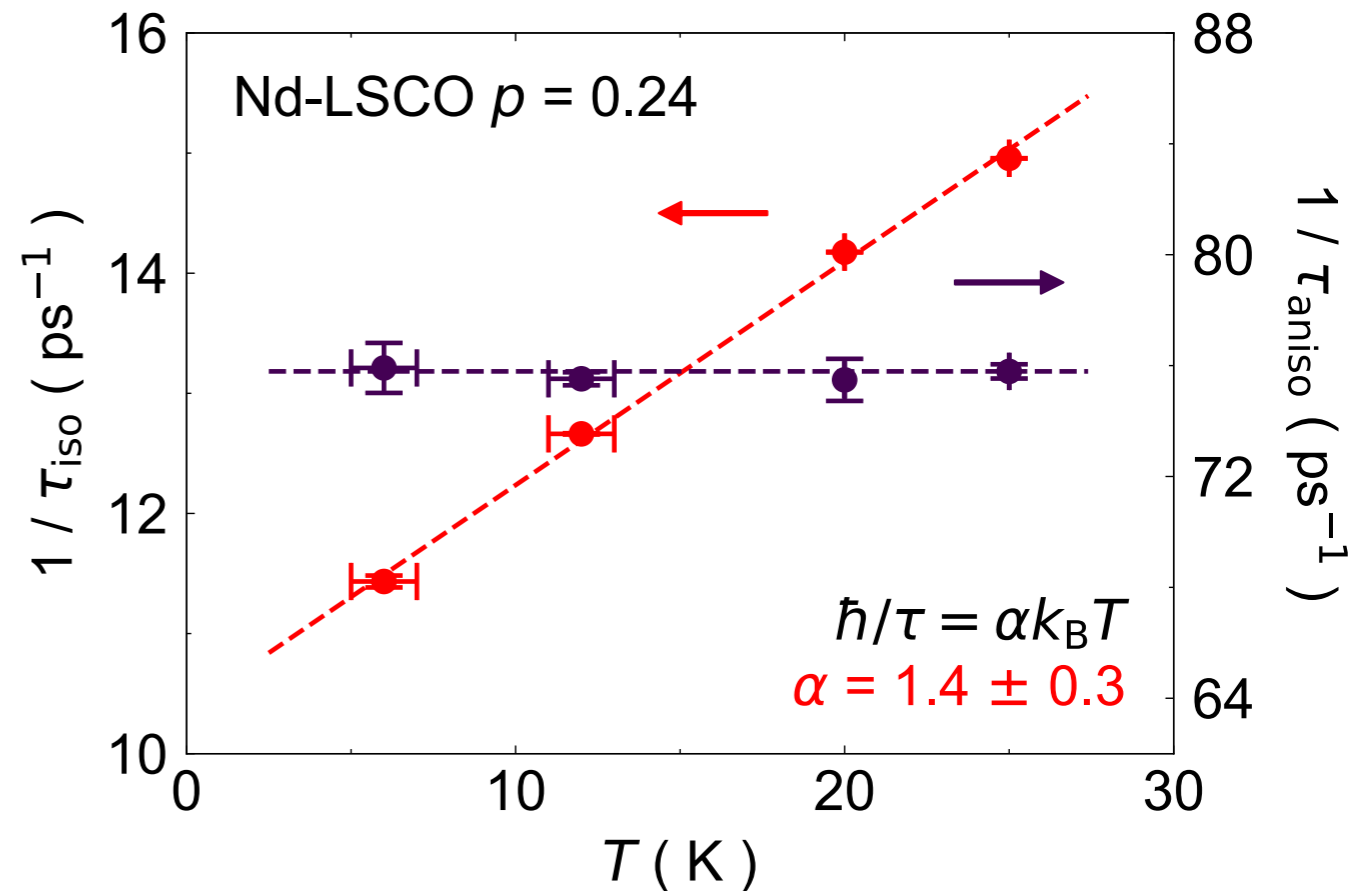
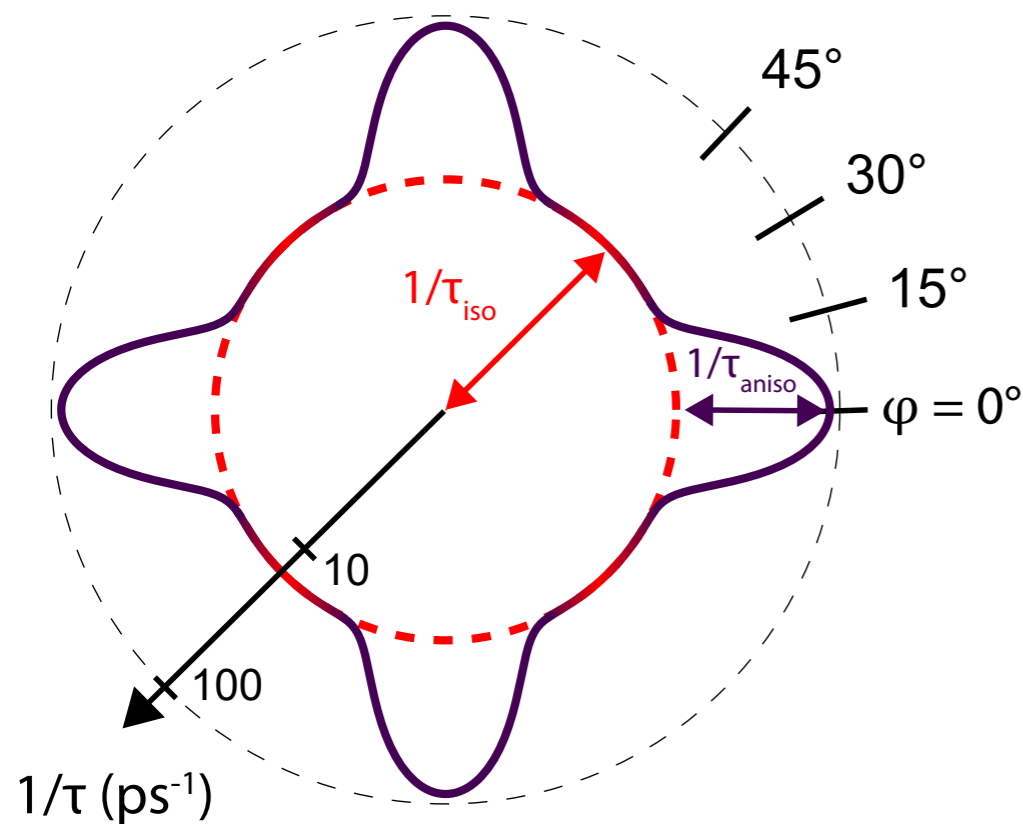


$$\frac{1}{\tau} = \frac{1}{\tau_{aniso}(\vec{k})} + \frac{\alpha}{\hbar} k_B T$$

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Henry Shackleton



Alexander Wietek



Antoine Georges

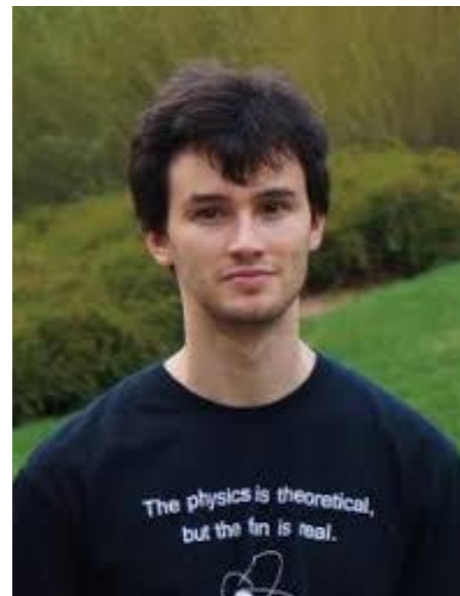
[arXiv:2012.06589](https://arxiv.org/abs/2012.06589)



Maria Tikhanovskaya



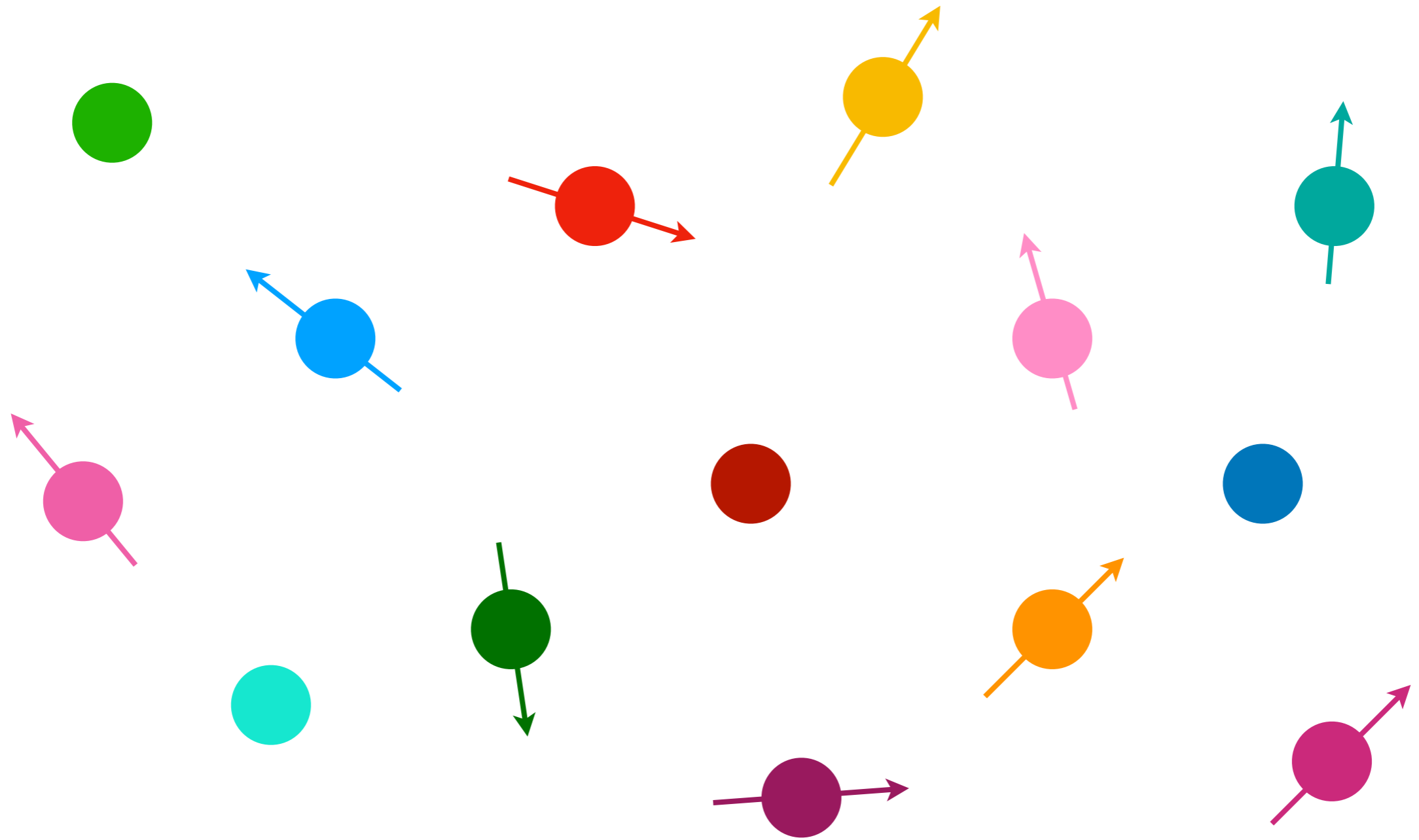
Haoyu Guo



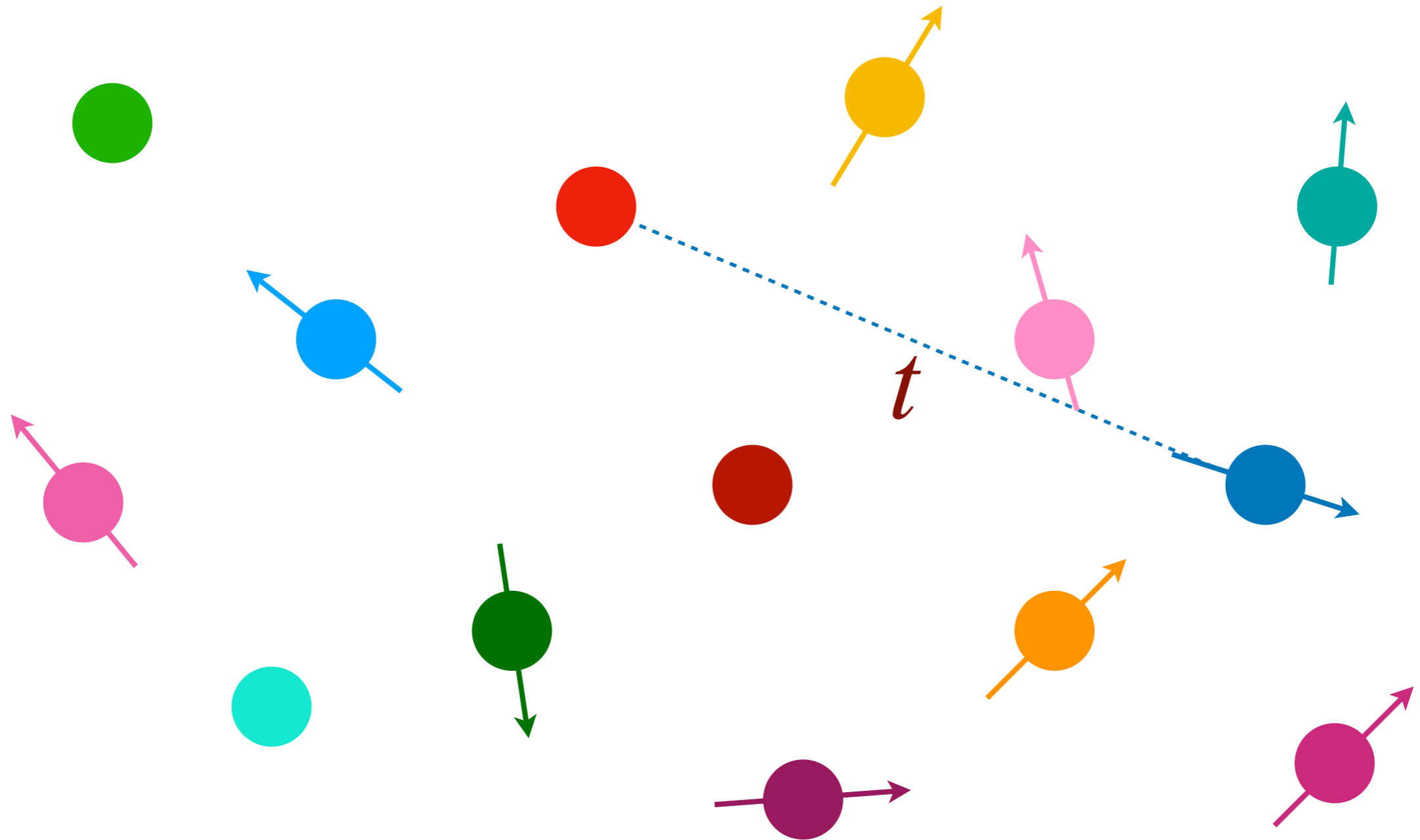
Grigory Tarnopolsky

arXiv:2010.09742
arXiv:2012.14449

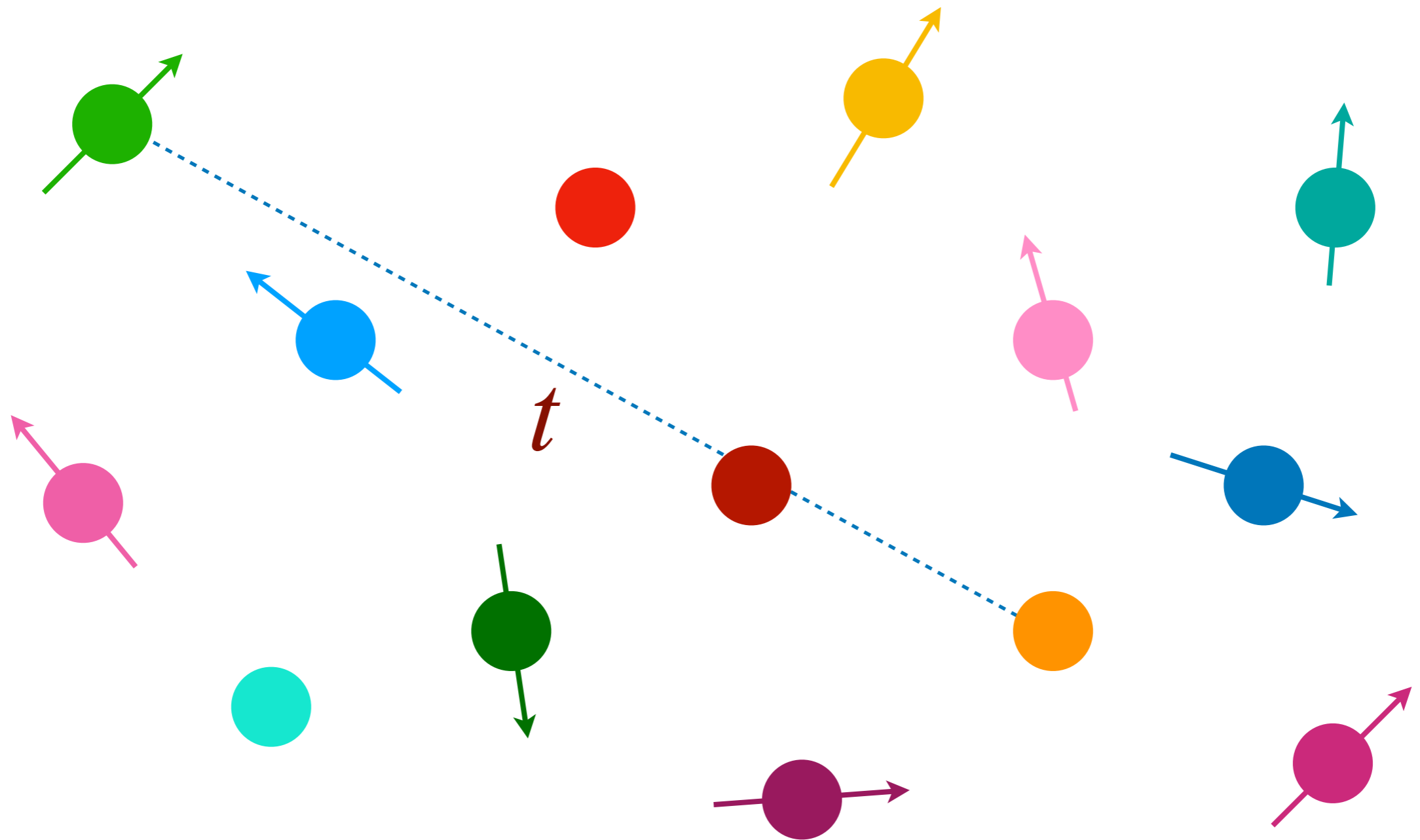
Random t - J model



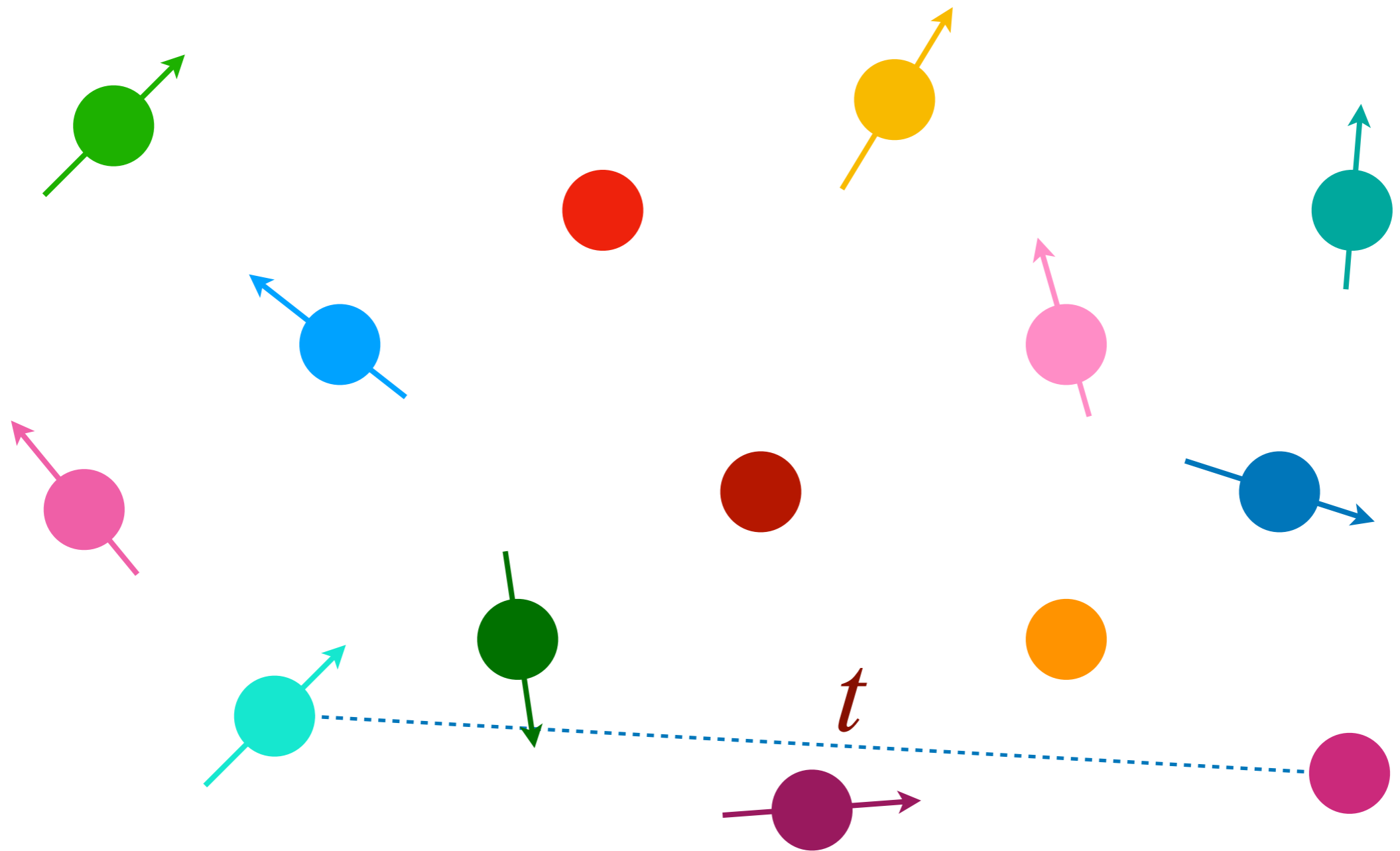
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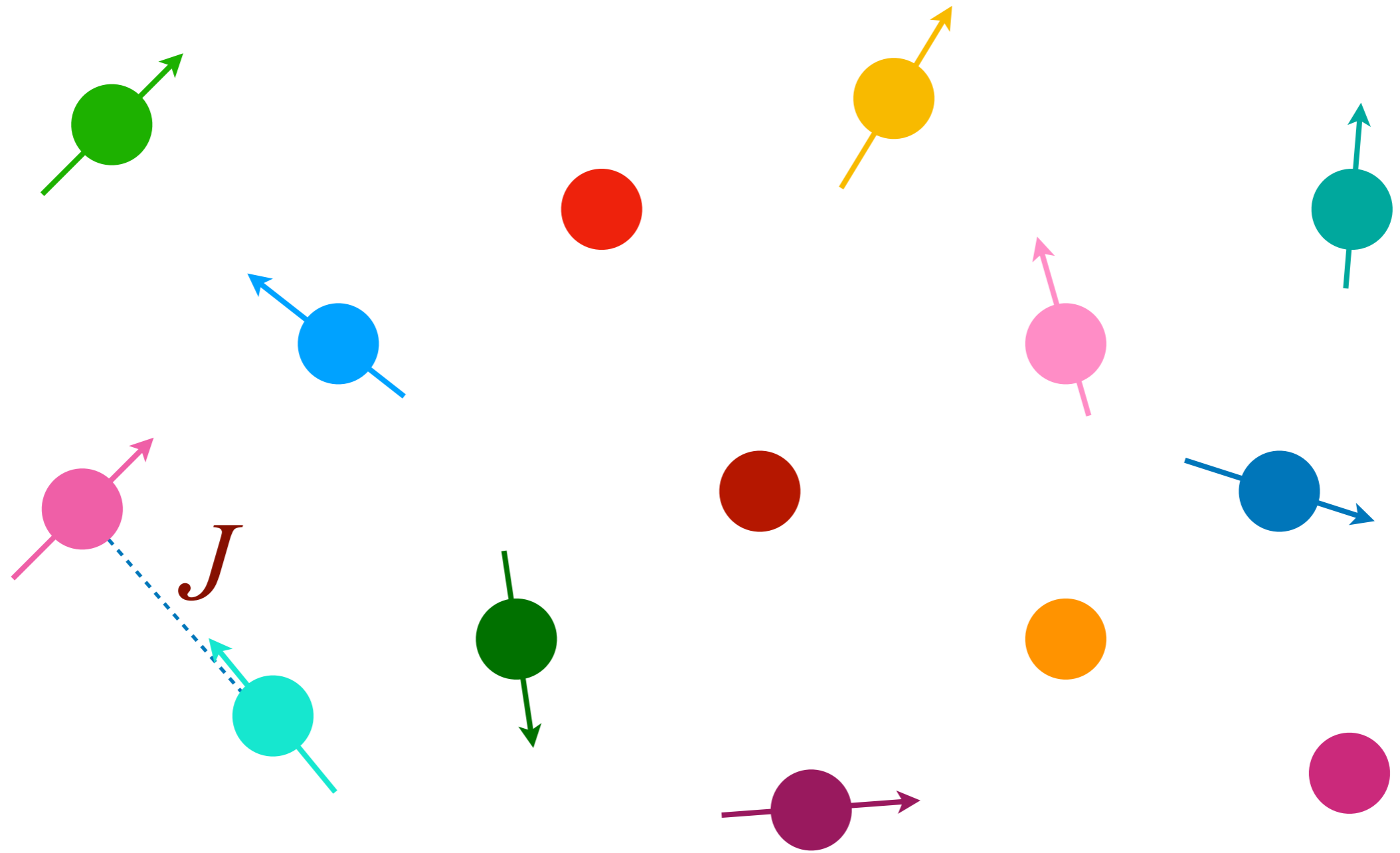
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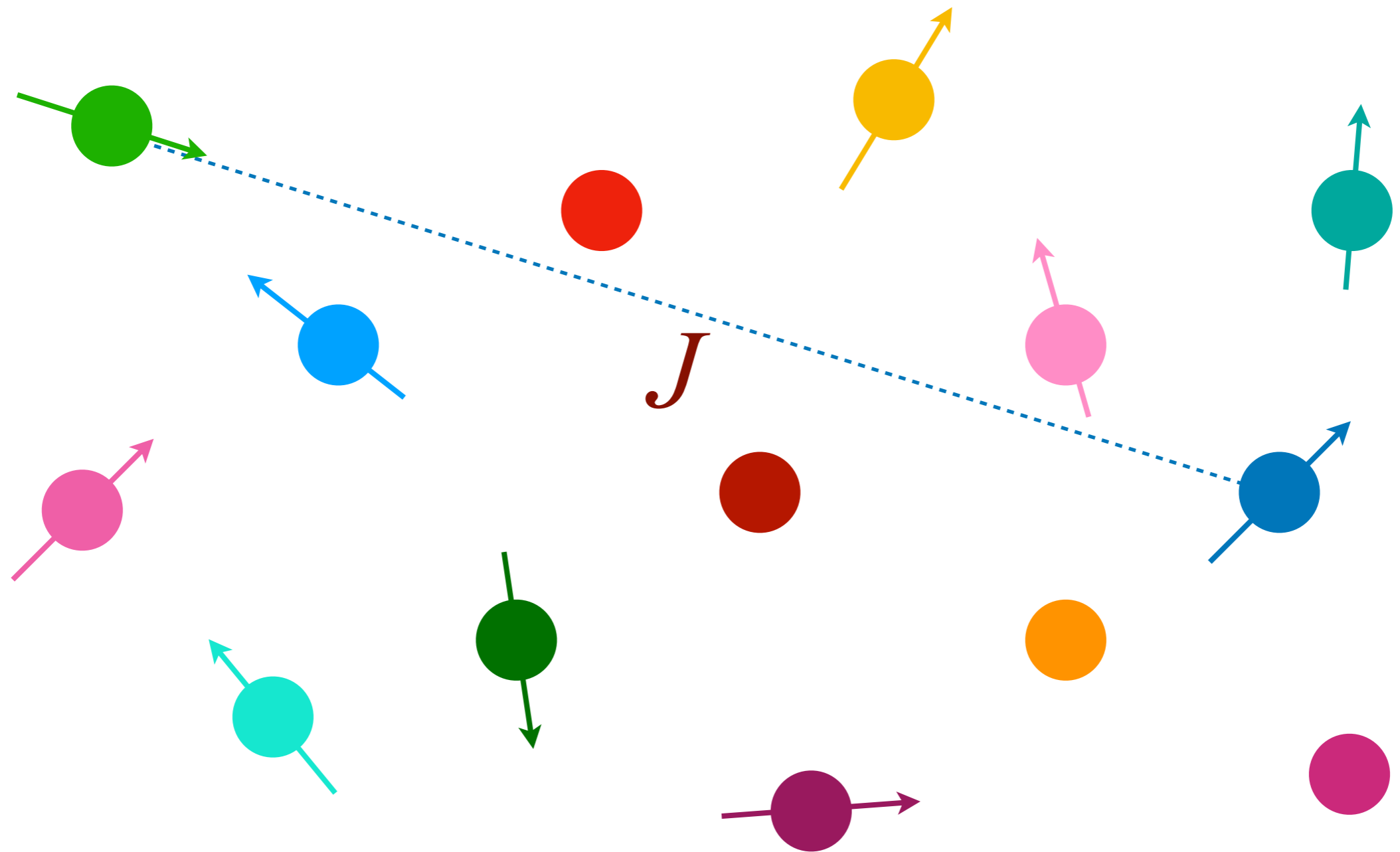
Random t - J model



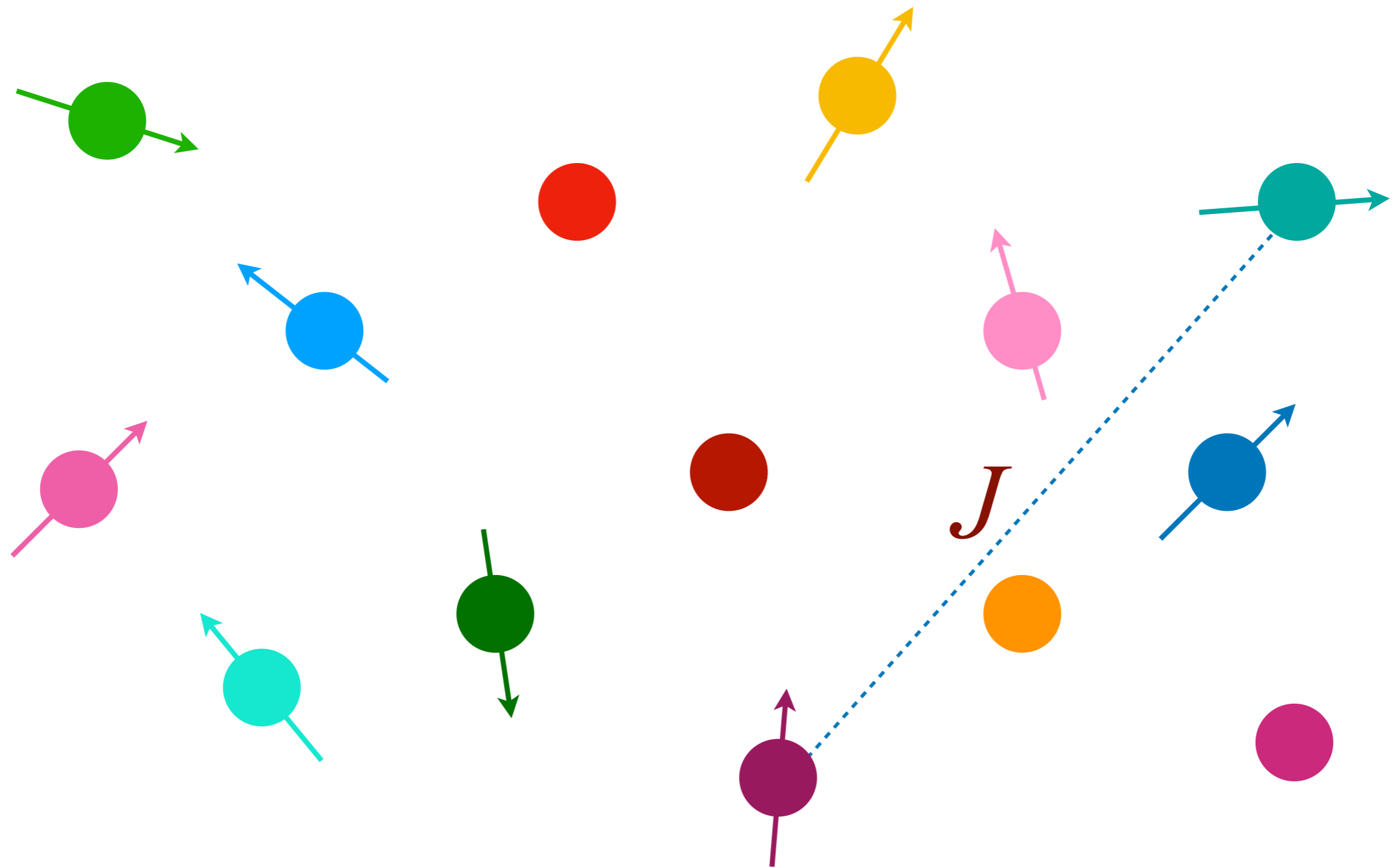
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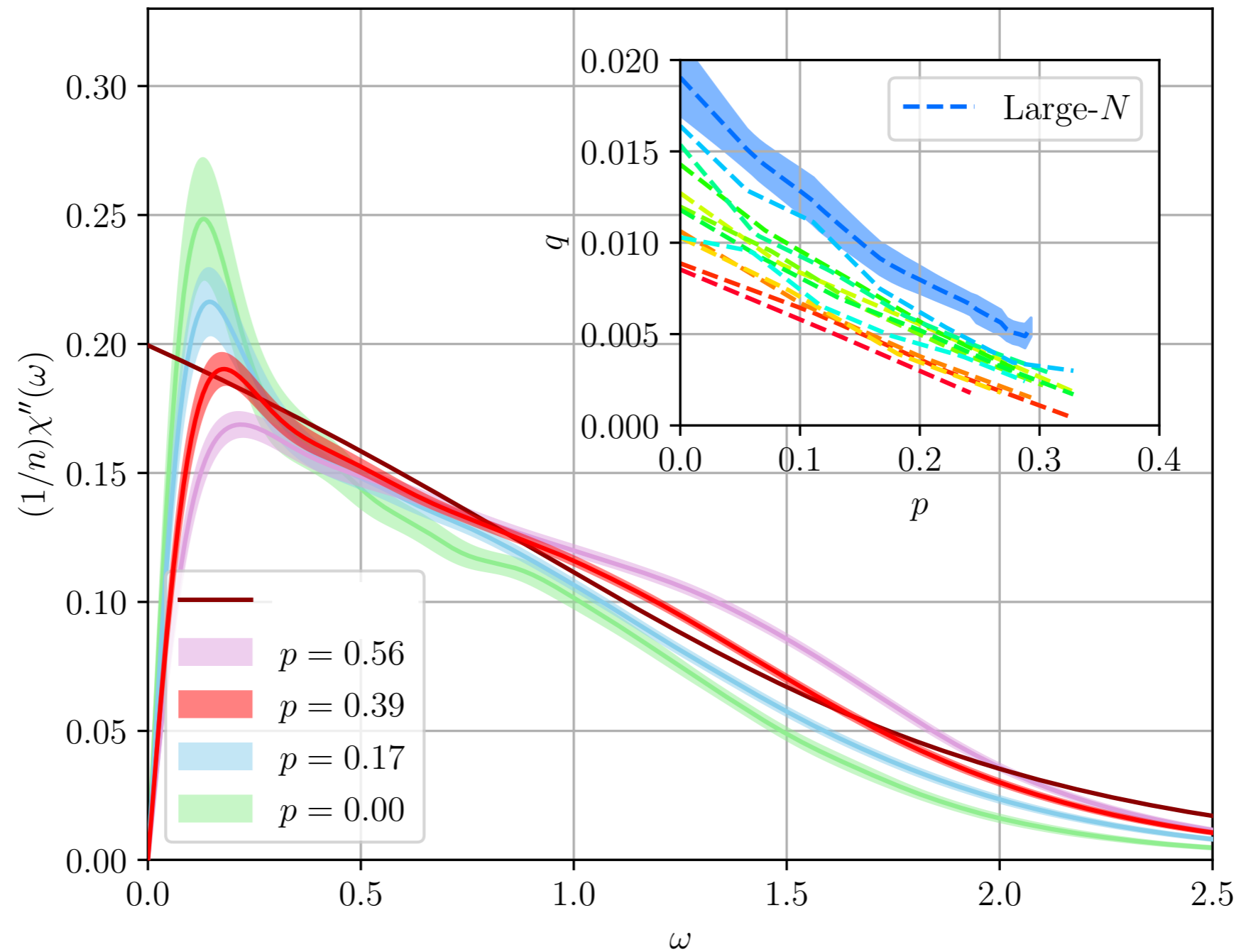


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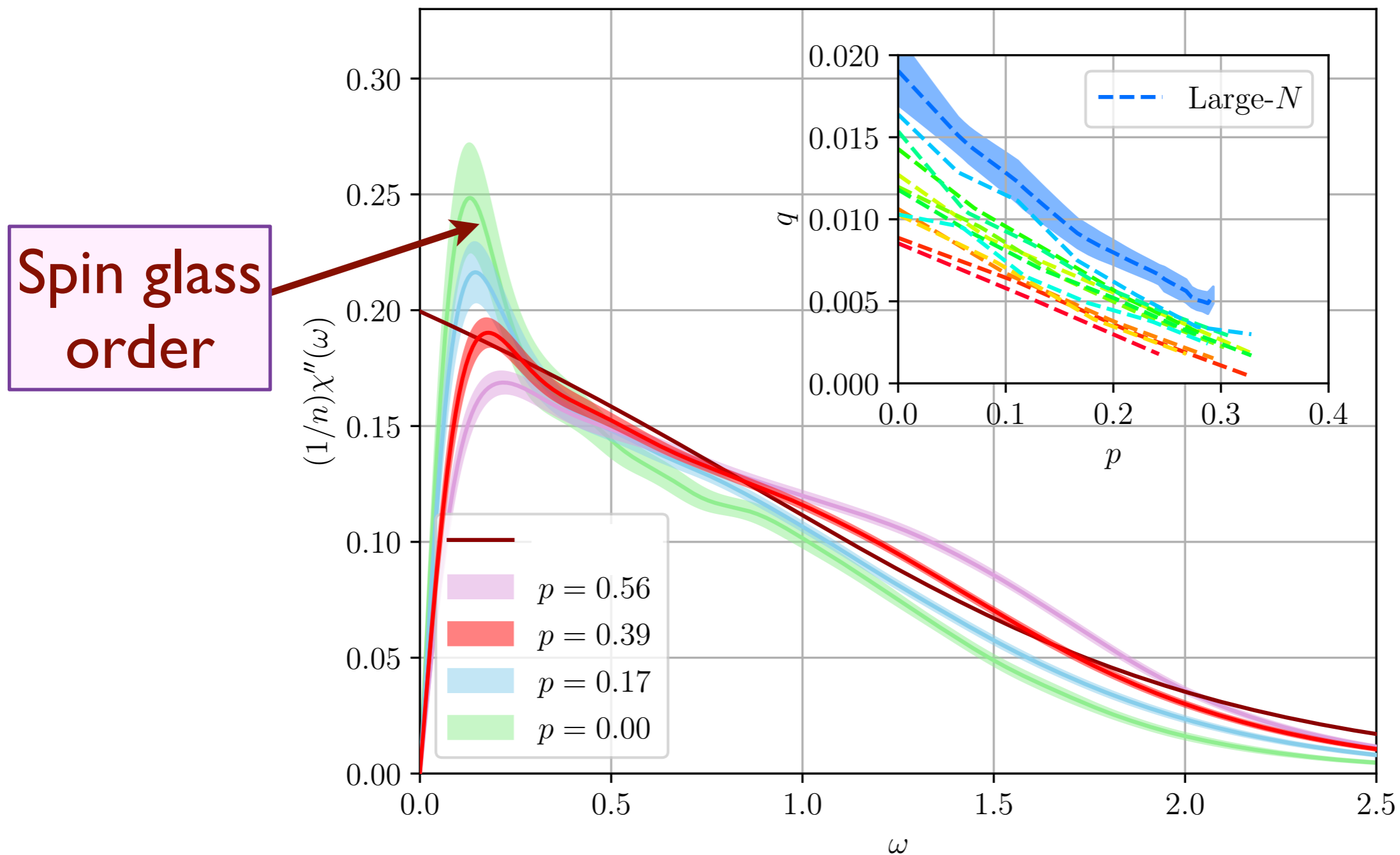
Dynamic spin susceptibility

Probability to flip an electron spin while absorbing energy $\hbar\omega$



Dynamic spin susceptibility

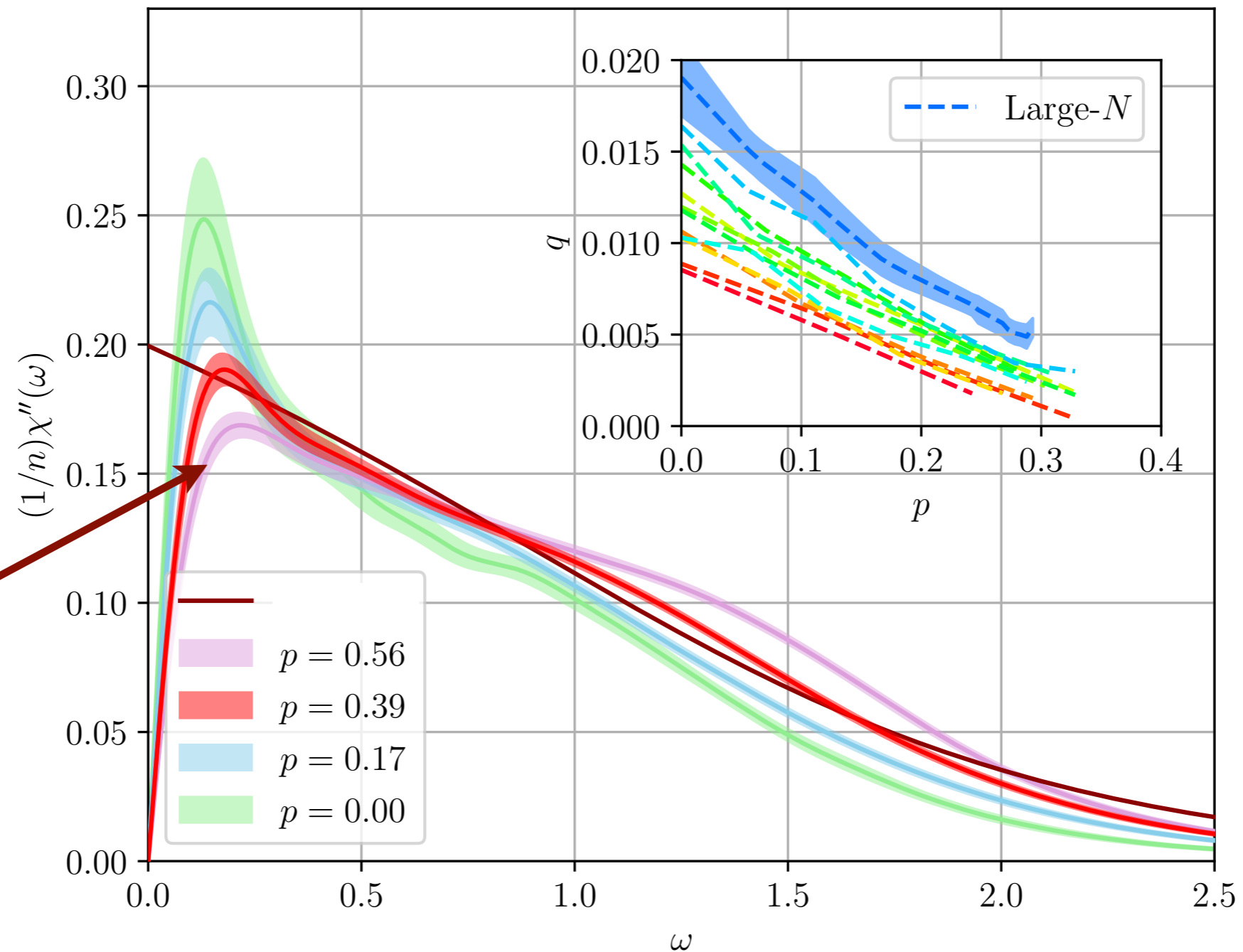
Probability to flip an electron spin while absorbing energy $\hbar\omega$



Spin glass order q non-zero for $p < p_c \approx 0.4$

Dynamic spin susceptibility

Probability to flip an electron spin while absorbing energy $\hbar\omega$

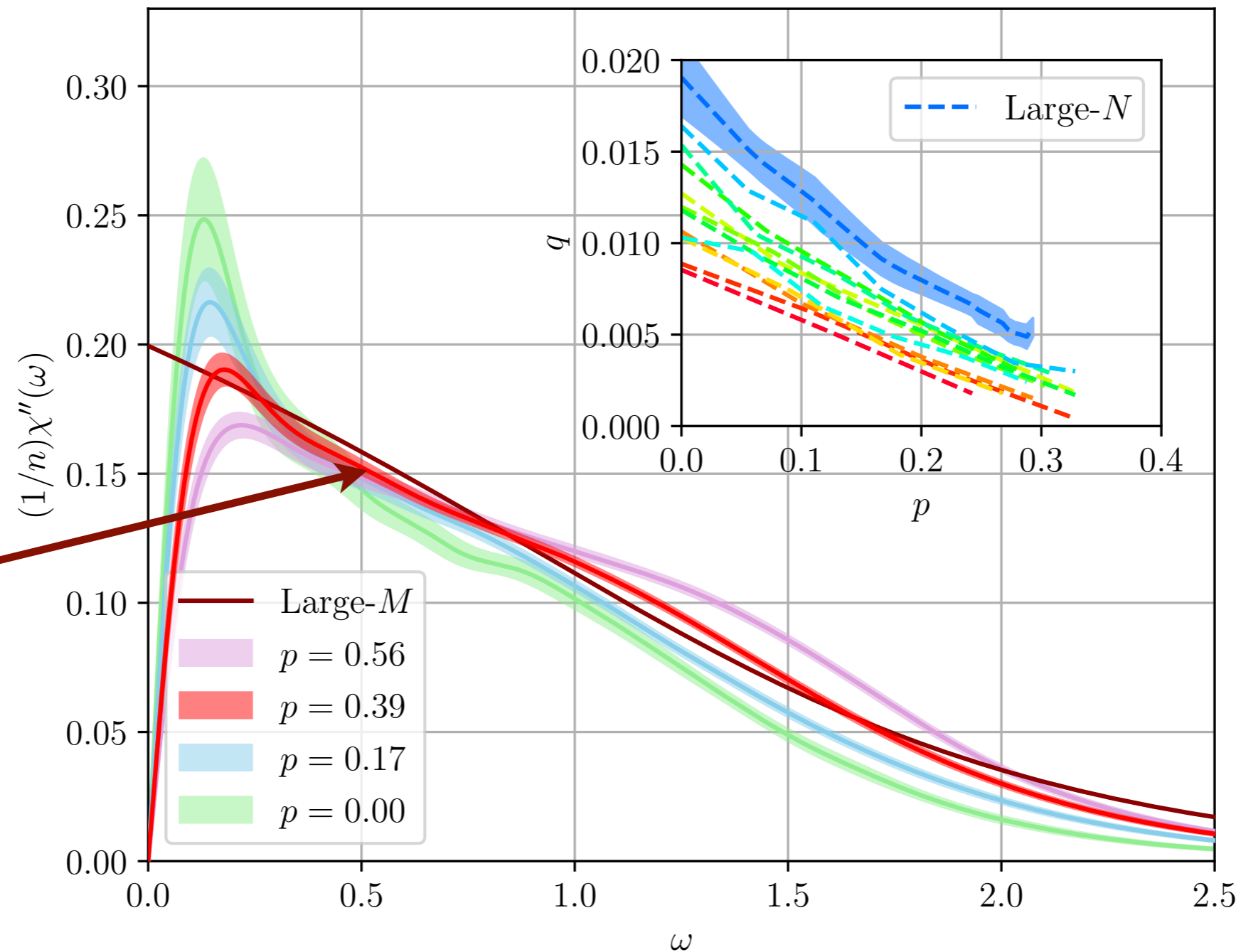


Ordinary
metal

Spin susceptibility and other properties
match those of an ordinary metal $p > p_c$

Dynamic spin susceptibility

Probability to flip an electron spin while absorbing energy $\hbar\omega$

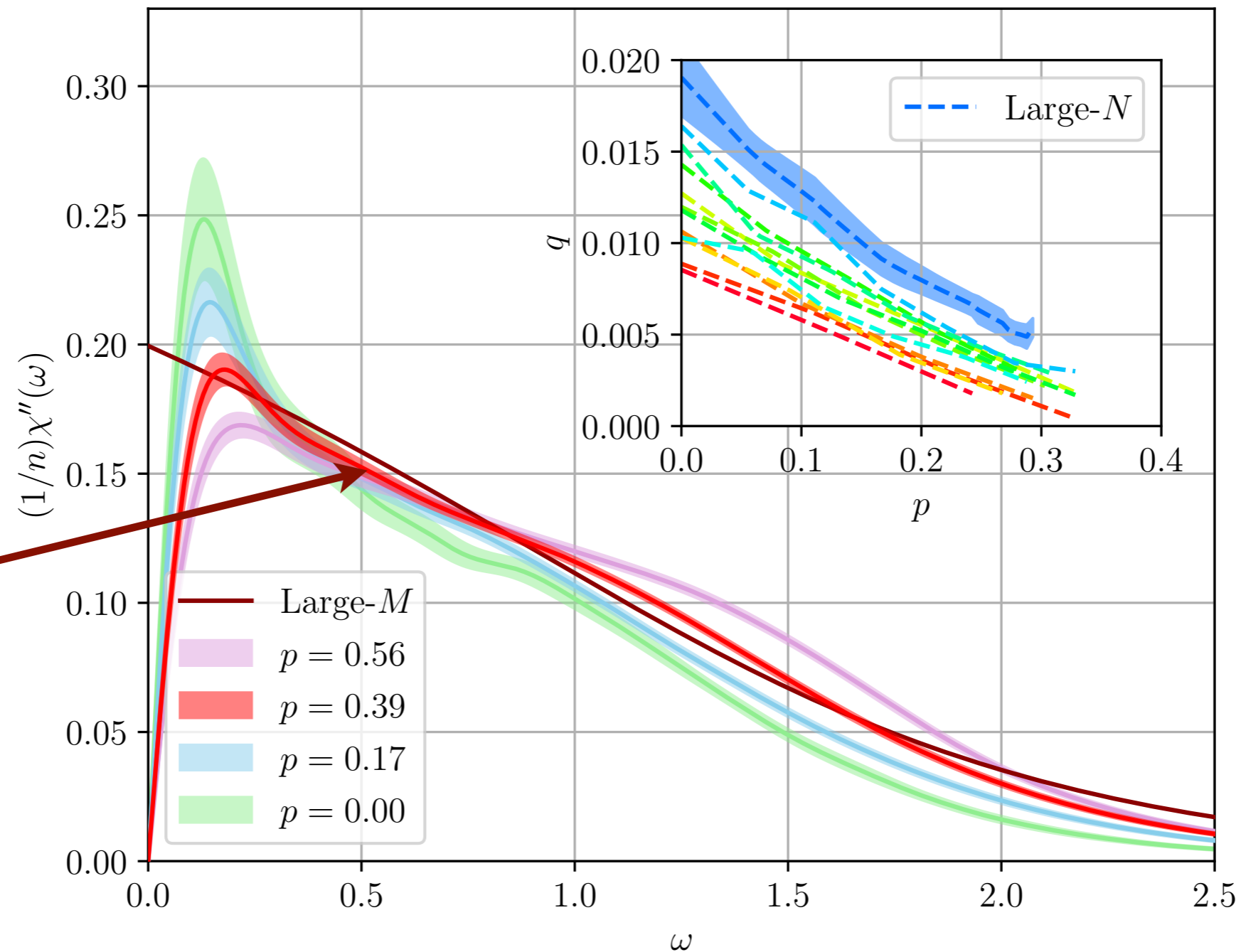


Critical
point

Critical spin susceptibility matches the SYK model!
Planckian dissipation in time $\sim \hbar/(k_B T)$,
and frequency dependence $\sim \text{sgn}(\omega) [1 - \mathcal{C}\gamma|\omega| + \dots]$
matches contribution of boundary graviton.

Dynamic spin susceptibility

Probability to flip an electron spin while absorbing energy $\hbar\omega$



Critical
point

D. G. Joshi,
Chenyuan Li,
G. Tarnopolsky,
A. Georges, and
S. Sachdev,
PRX **10**, 021033
(2020)

SYK criticality can be understood by the fractionalization of the electron into 'partons' carrying its spin and charge.
These partons obey an SYK-like model

**Quantum
entanglement**

**Charged
black holes**

**A simple
many-particle
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**Copper-based
superconductors**

**Quantum
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2D
quantum
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