

Condensed matter physics and string theory

Talk online: sachdev.physics.harvard.edu



Comfort zone of string theory

One, two, three... particles
(quarks, gluons, gravitons...)

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Works at critical points where particle
spectrum changes
(AdS/CFT correspondence)

Comfort zone of condensed matter

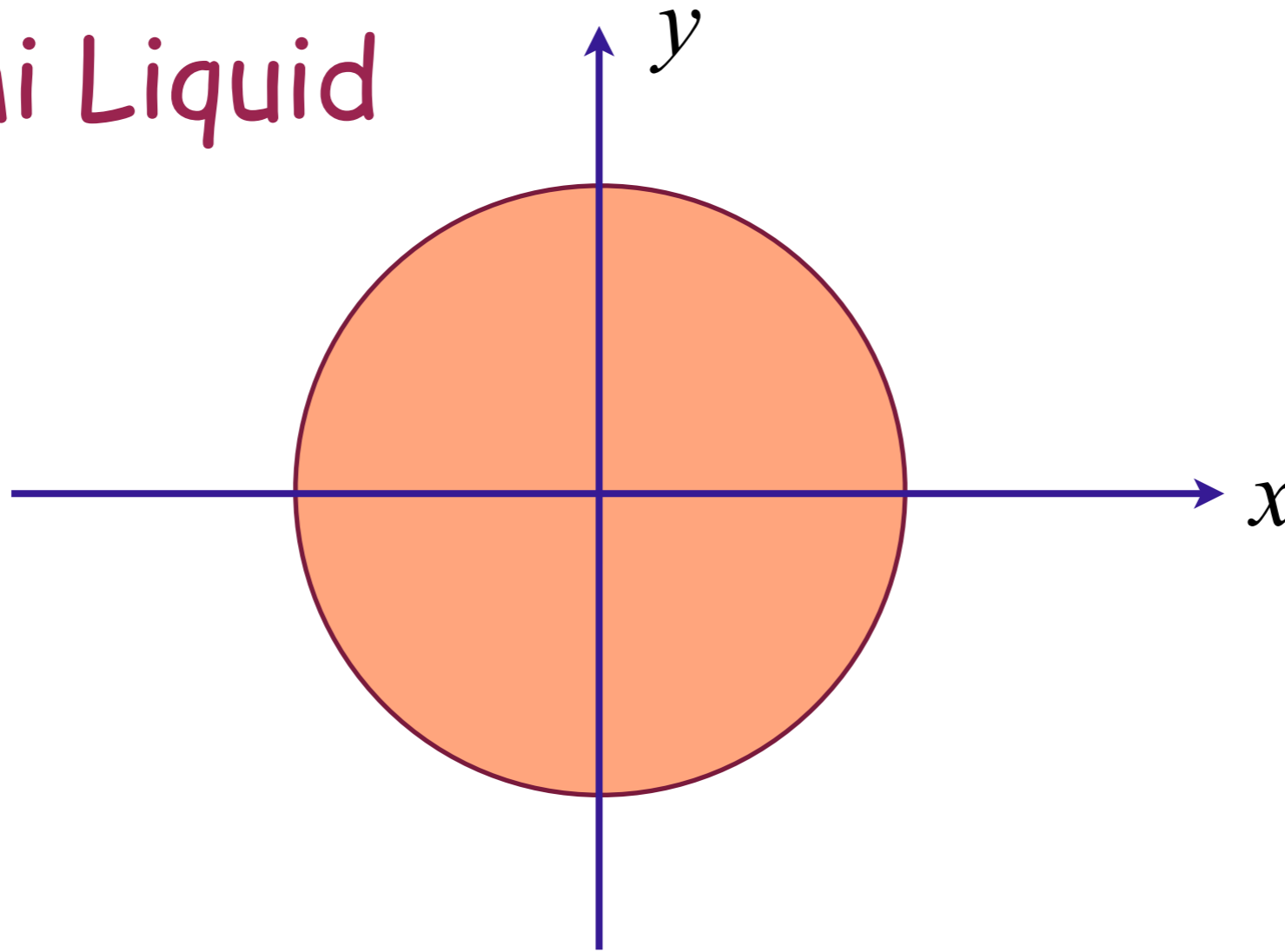
Infinite numbers of particles (non-zero density)
with weak interactions:
electrons, trapped cold atoms

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Quantum phases with universal
low energy properties, independent of
most interaction details:
Fermi liquids, solids, superfluids

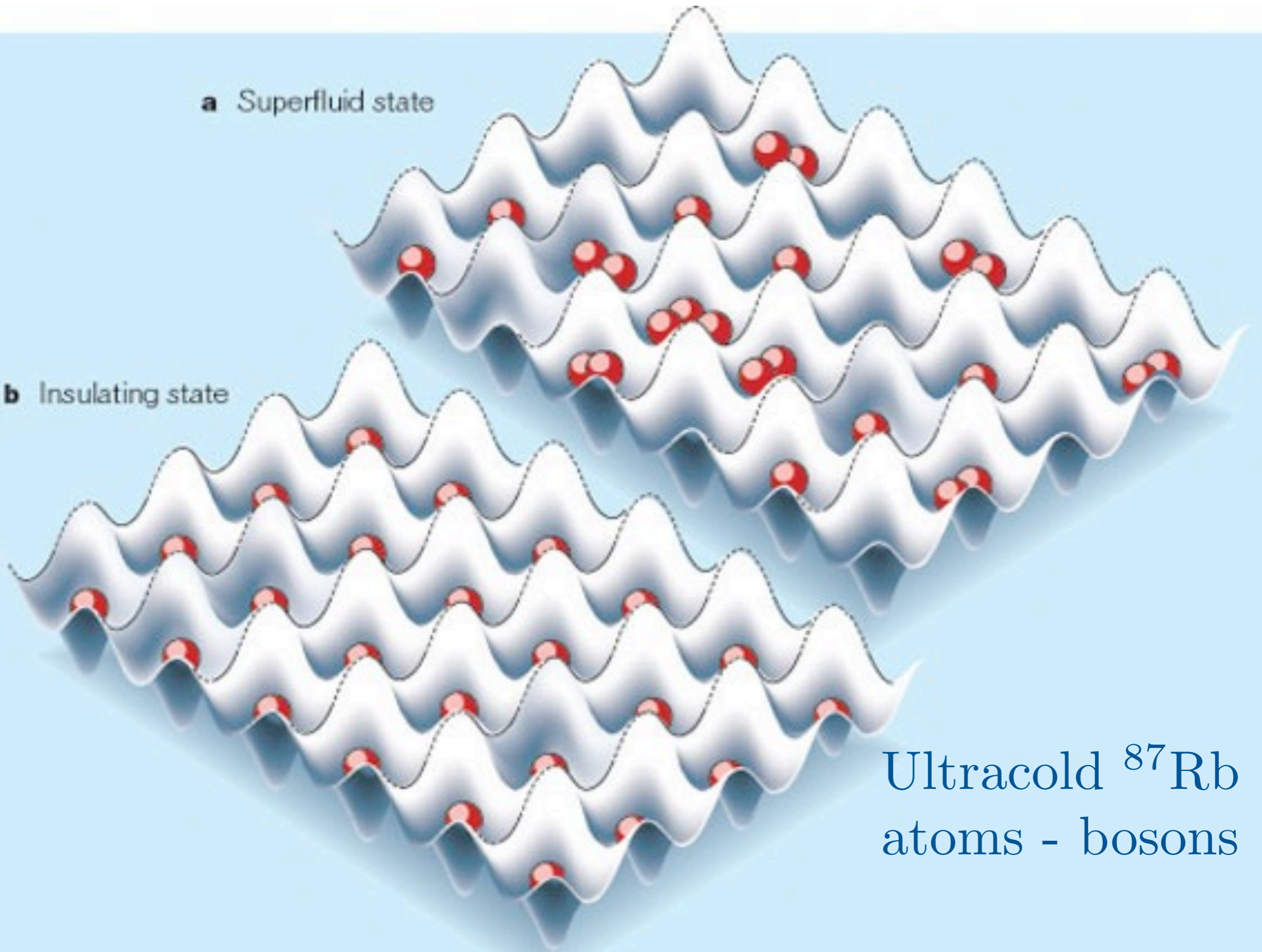
Fermi Liquid



In momentum space, Fermi surface separates occupied and empty electron states

a Superfluid state

b Insulating state



Ultracold ^{87}Rb
atoms - bosons

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Fermi liquid theory (and its cousins) allow
adiabatic continuation to
strong interaction regime

Move out of
comfort zones,
and use tools developed
in both fields

Strong coupling problems in condensed matter

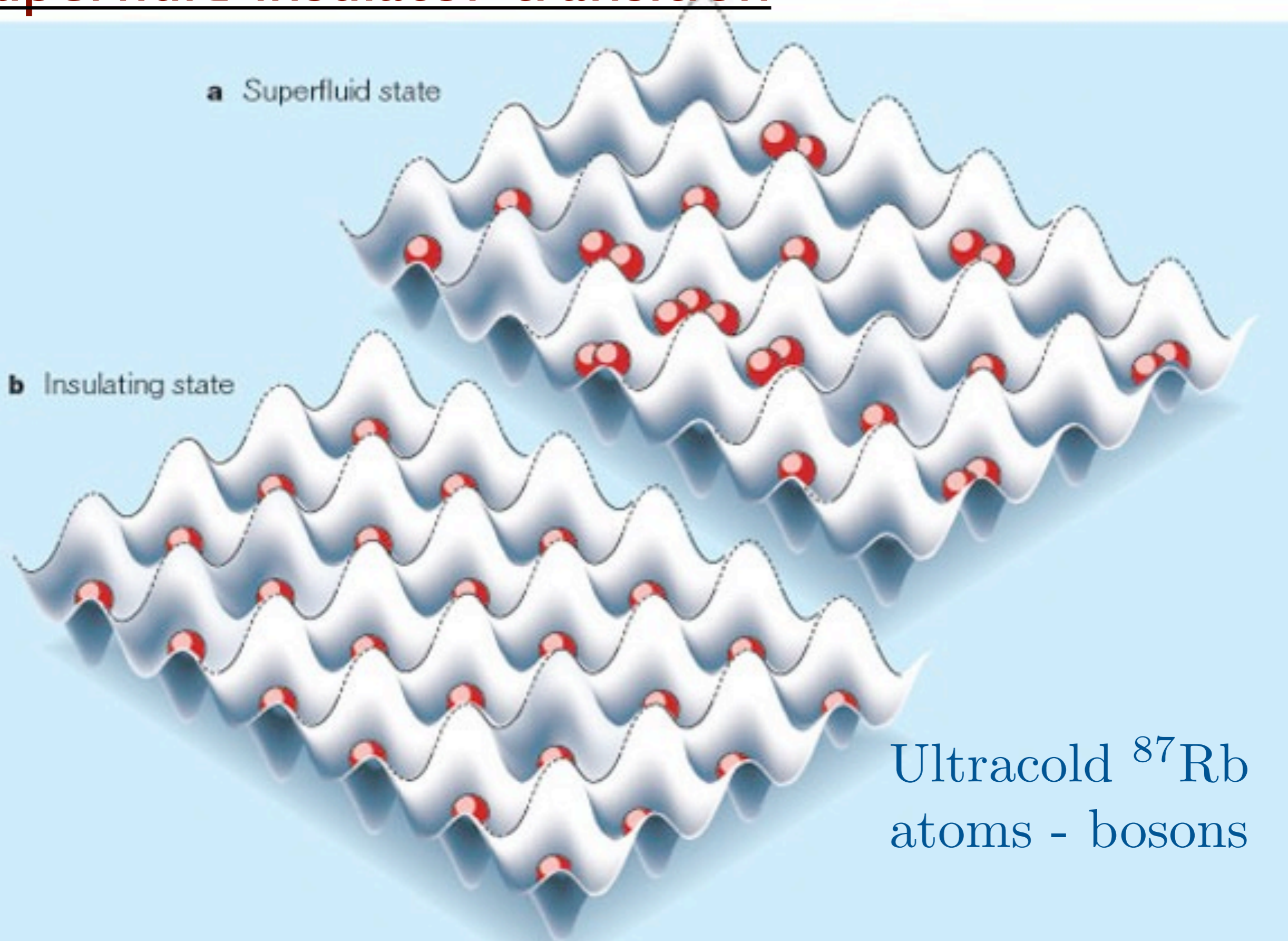
1. Electrical resistance at non-zero temperature near quantum critical points
2. Zero temperature quantum phase transitions of Fermi liquids

Strong coupling problems in condensed matter

1. Electrical resistance at non-zero temperature near quantum critical points

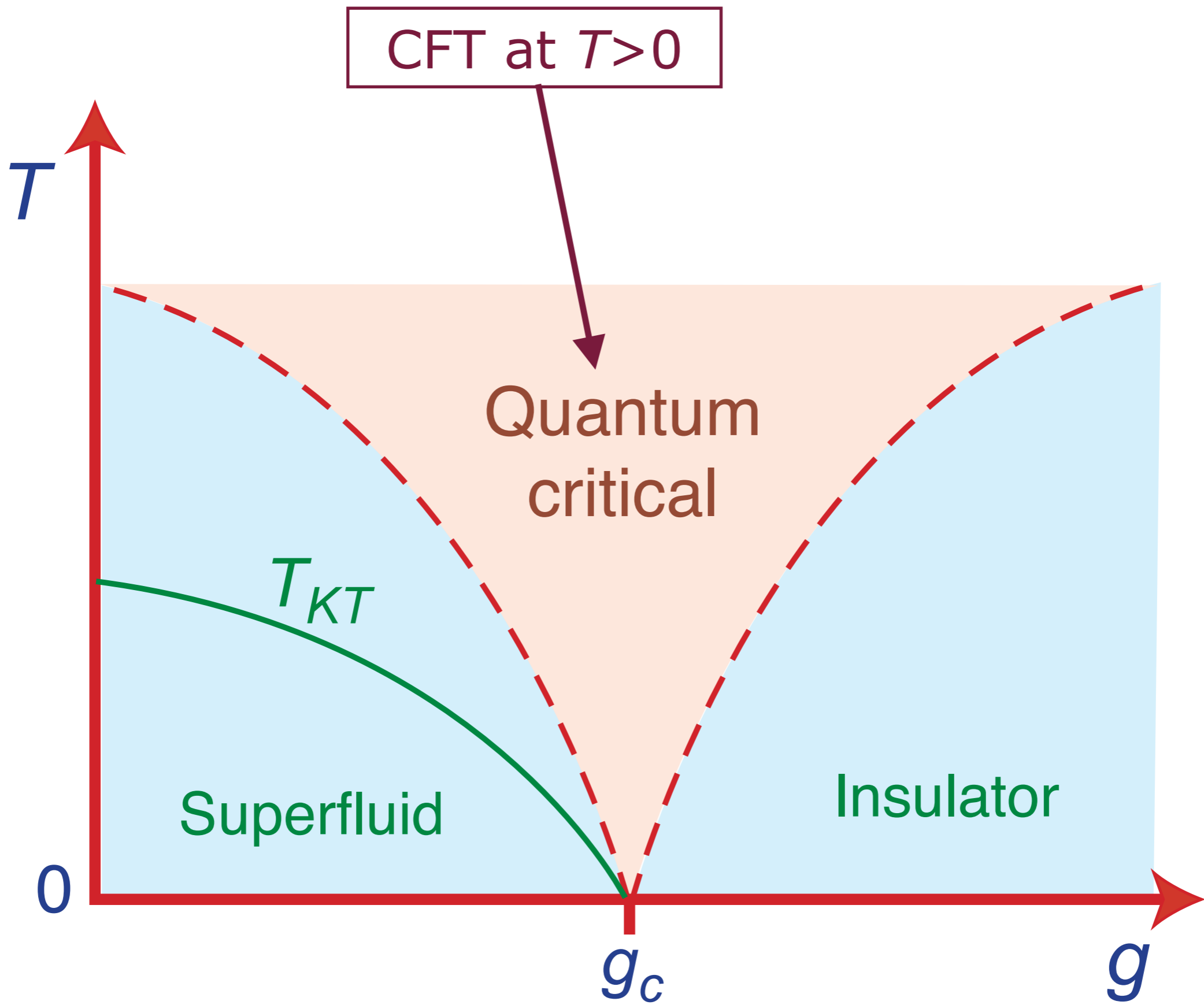
2. Zero temperature quantum phase transitions of Fermi liquids

Superfluid-insulator transition



Ultracold ^{87}Rb
atoms - bosons

M. Greiner, O. Mandel, T. Esslinger, T. W. Hänsch, and I. Bloch, *Nature* **415**, 39 (2002).



Quantum critical transport

Quantum “*perfect fluid*”
with shortest possible
relaxation time, τ_R

$$\tau_R \gtrsim \frac{\hbar}{k_B T}$$

K. Damle and S. Sachdev, *Phys. Rev. B* **56**, 8714 (1997).

Quantum critical transport

Transport co-efficients not determined
by collision rate, but by
universal constants of nature

Electrical conductivity

$$\sigma = \frac{4e^2}{h} \times [\text{Universal constant } \mathcal{O}(1)]$$

K. Damle and S. Sachdev, *Phys. Rev. B* **56**, 8714 (1997).

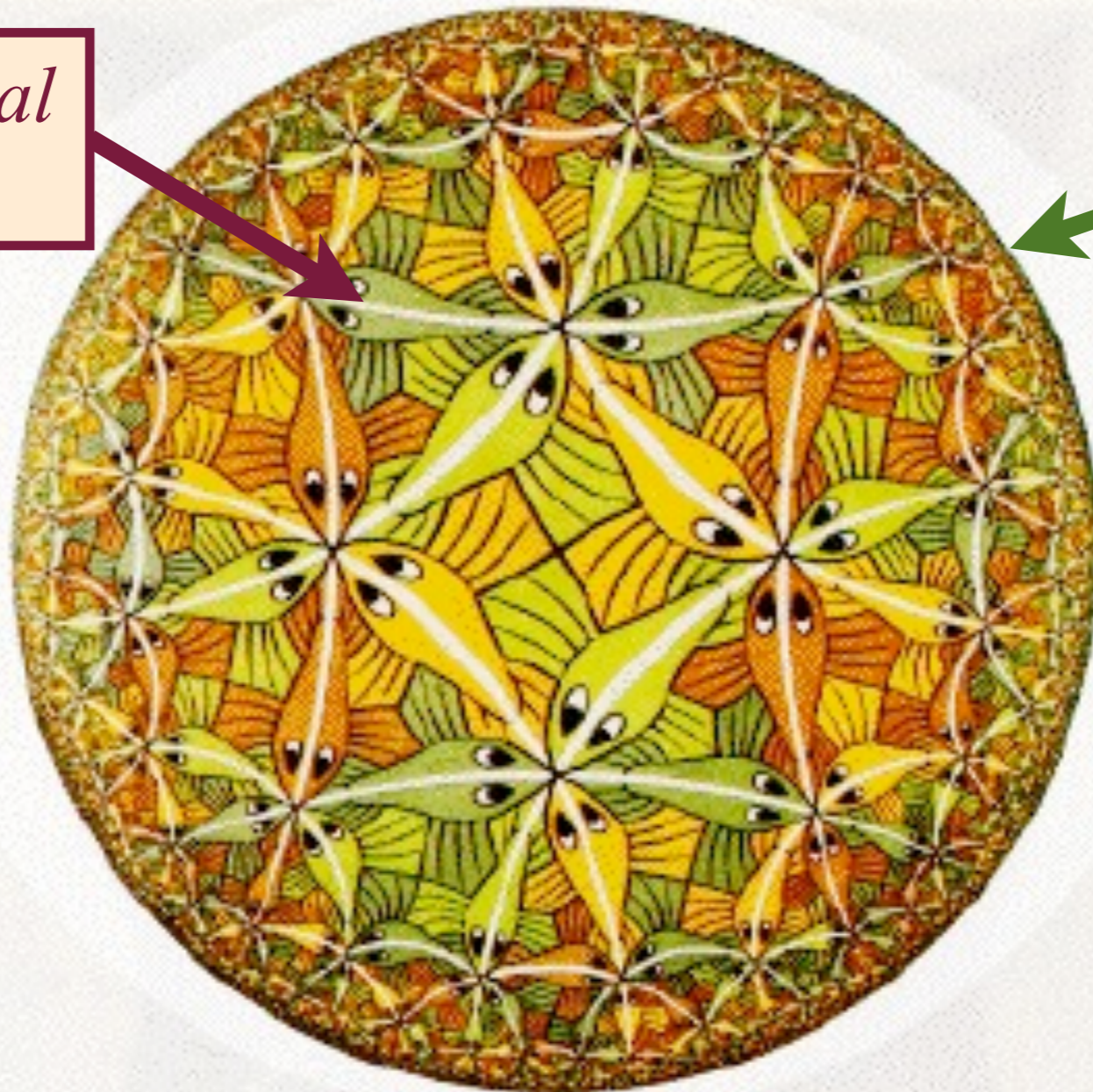
Requires long-time limit of correlations at non-zero temperature.

Cannot be obtained numerically by computers even for the simplest quantum problems. Is the sum of an exponentially large number of terms which oscillate in sign.

AdS/CFT correspondence

The quantum theory of a black hole in a 3+1-dimensional negatively curved AdS universe is holographically represented by a CFT (the theory of a quantum critical point) in 2+1 dimensions

*3+1 dimensional
AdS space*



A 2+1
dimensional
system at its
quantum
critical point

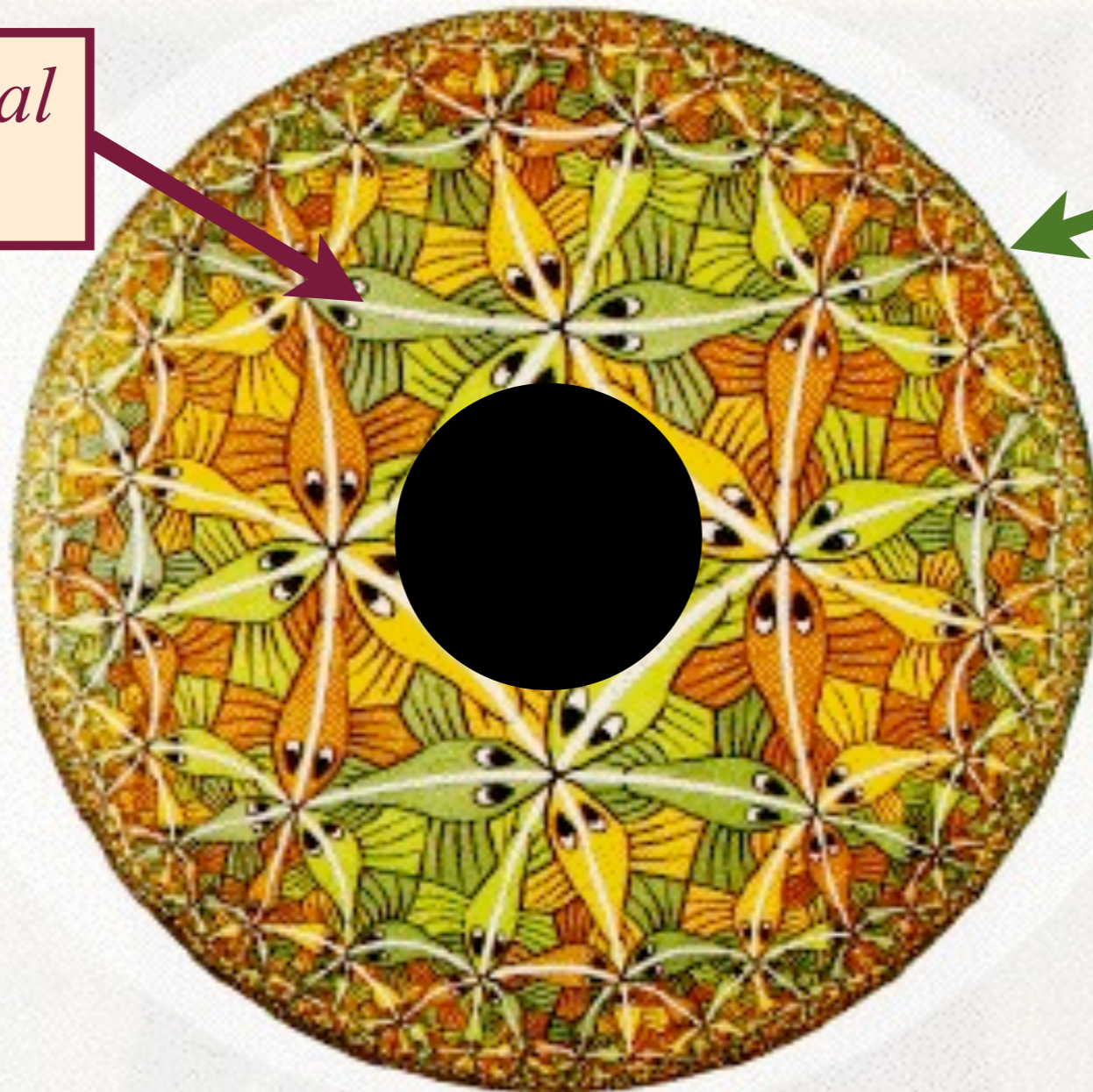
Maldacena, Gubser, Klebanov, Polyakov, Witten

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Quantum
criticality in
2+1
dimensions



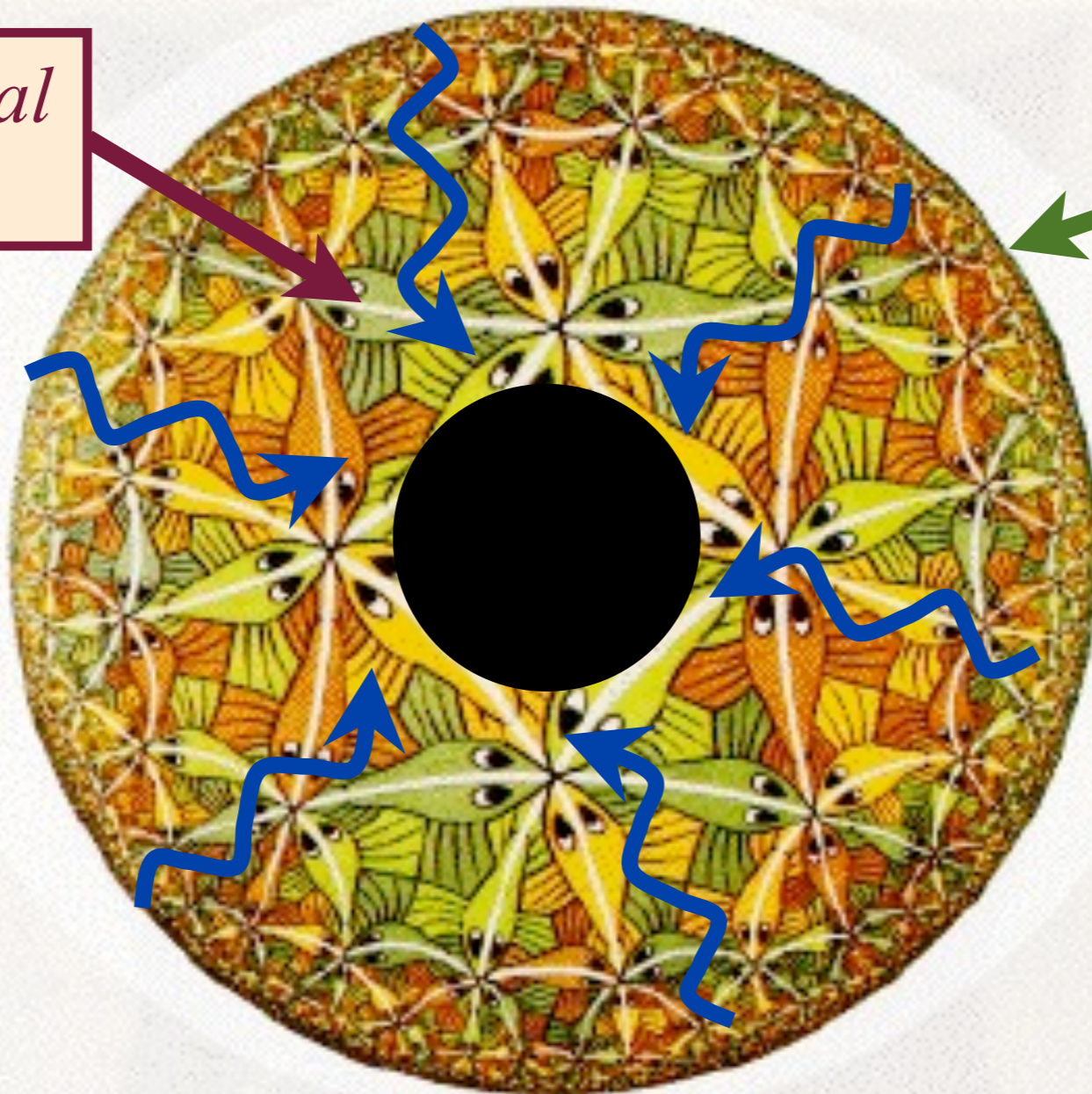
Black hole
temperature
=
temperature
of quantum
criticality

Maldacena, Gubser, Klebanov, Polyakov, Witten

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Quantum
criticality in
2+1
dimensions

Friction of
quantum
criticality =
waves
falling into
black hole

Kovtun, Policastro, Son

Strong coupling problems in condensed matter

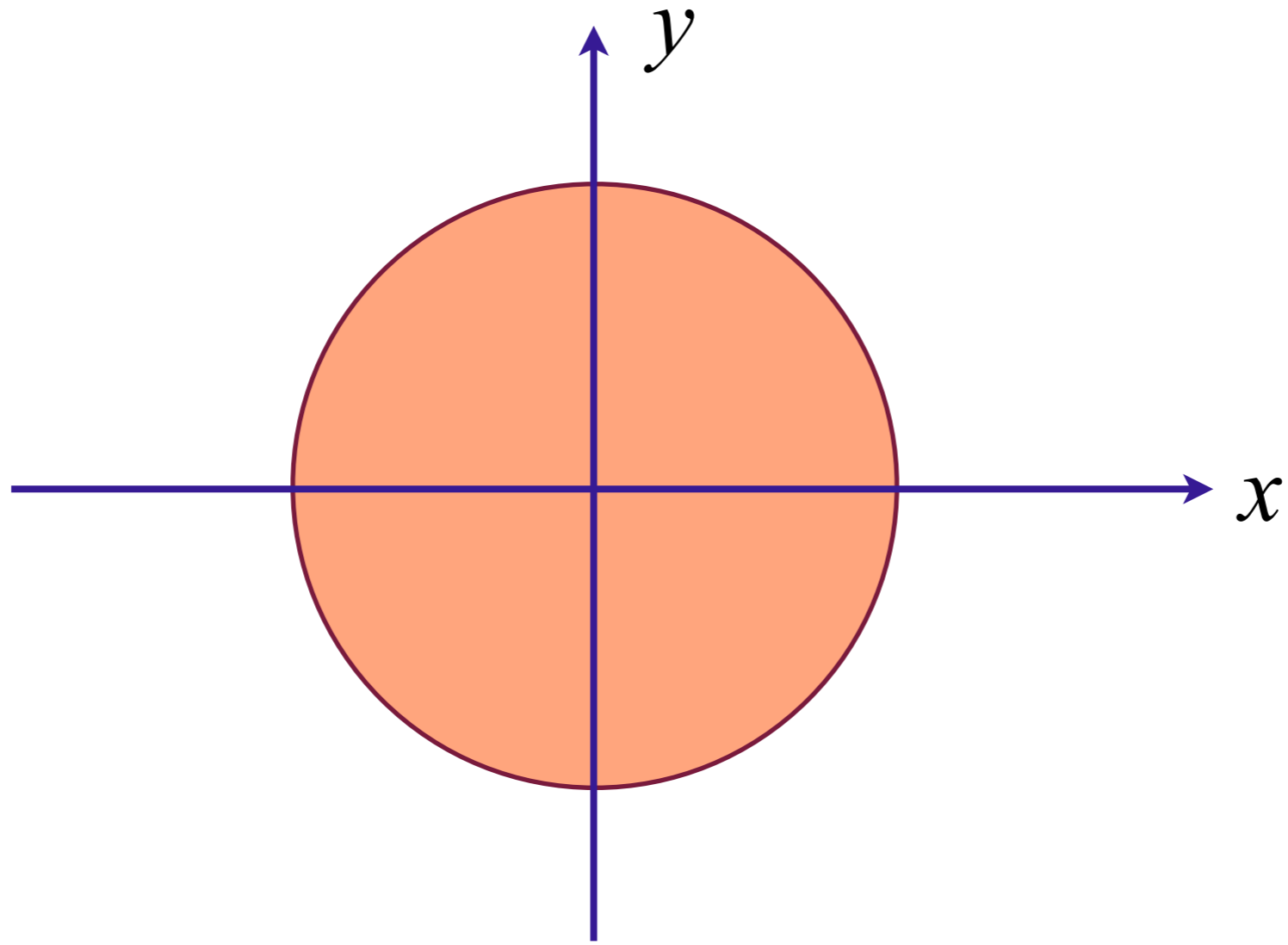
1. Electrical resistance at non-zero temperature near quantum critical points
2. Zero temperature quantum phase transitions of Fermi liquids

Strong coupling problems in condensed matter

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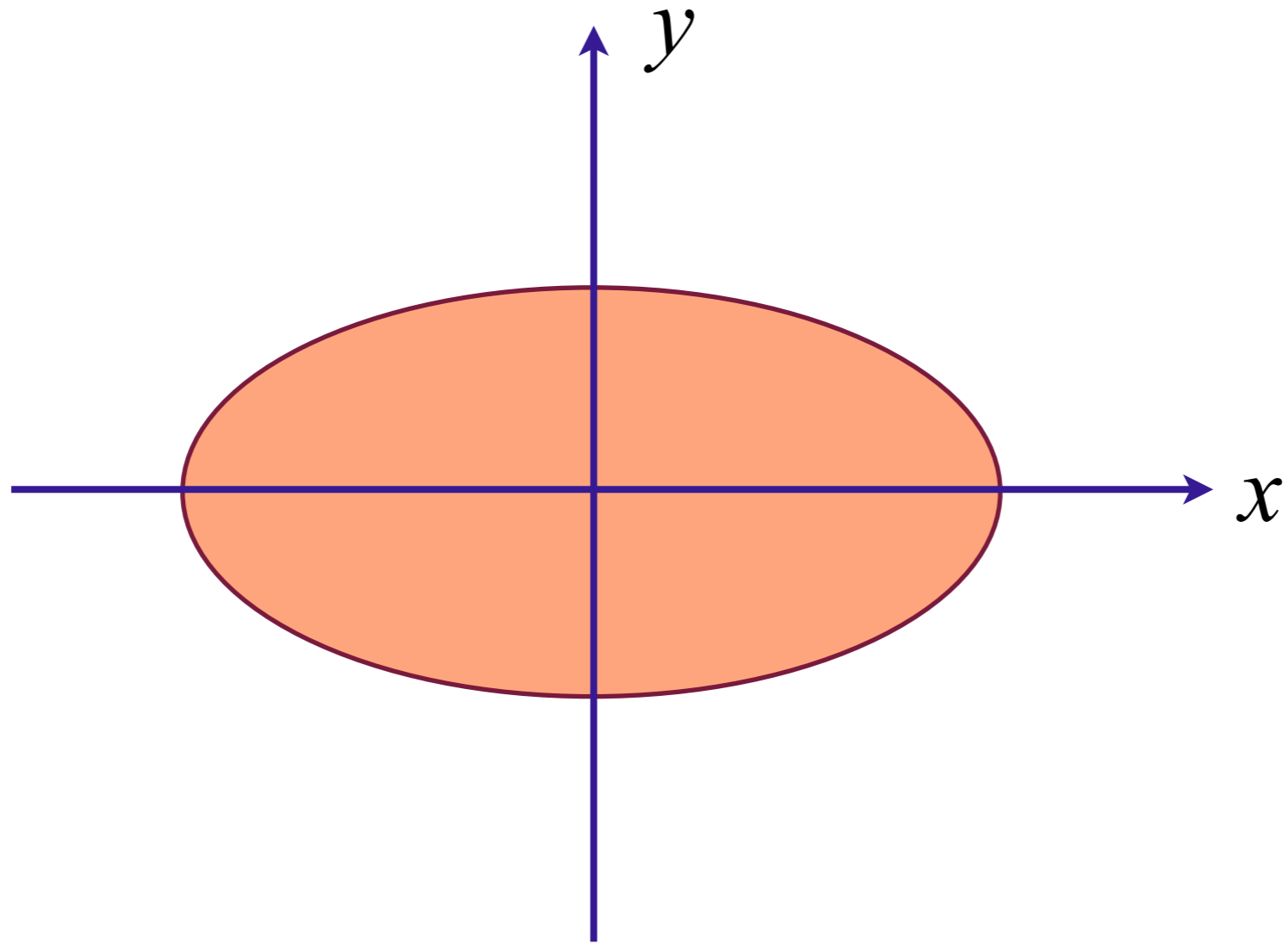
2. Zero temperature quantum phase transitions of Fermi liquids

Quantum criticality of Pomeranchuk instability



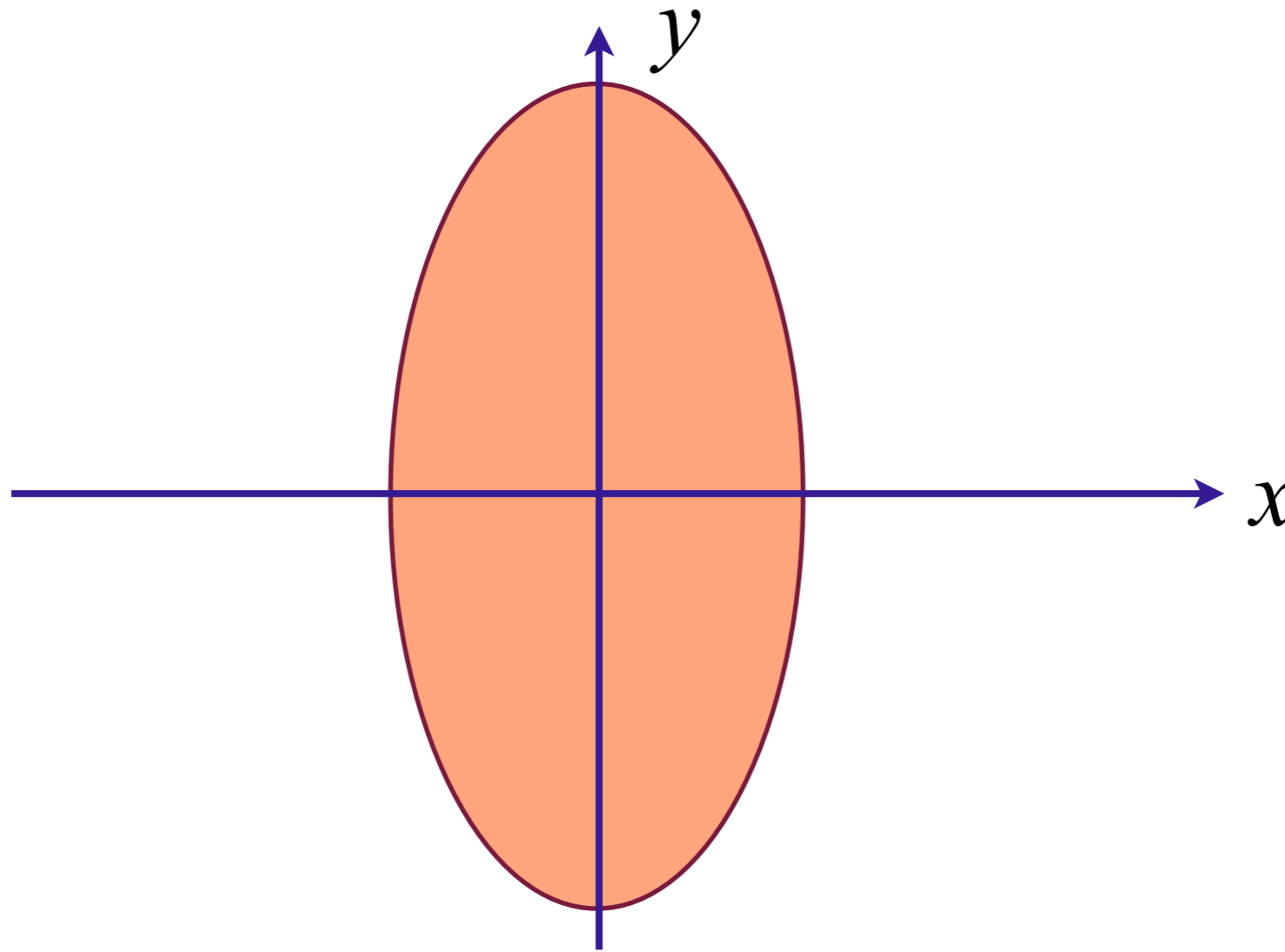
Fermi surface with full square lattice symmetry

Quantum criticality of Pomeranchuk instability



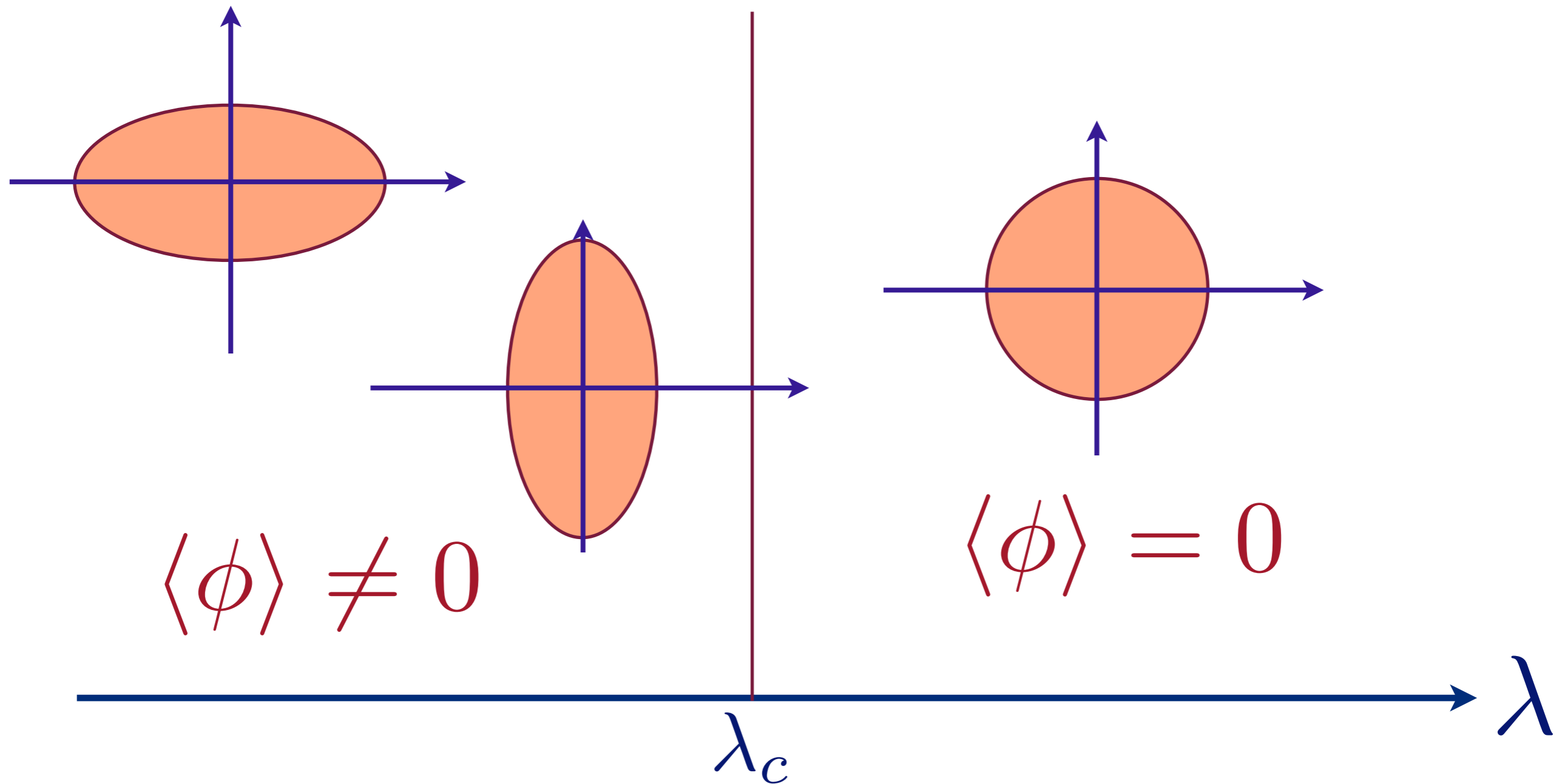
Spontaneous elongation along x direction:
Ising order parameter $\phi > 0$.

Quantum criticality of Pomeranchuk instability

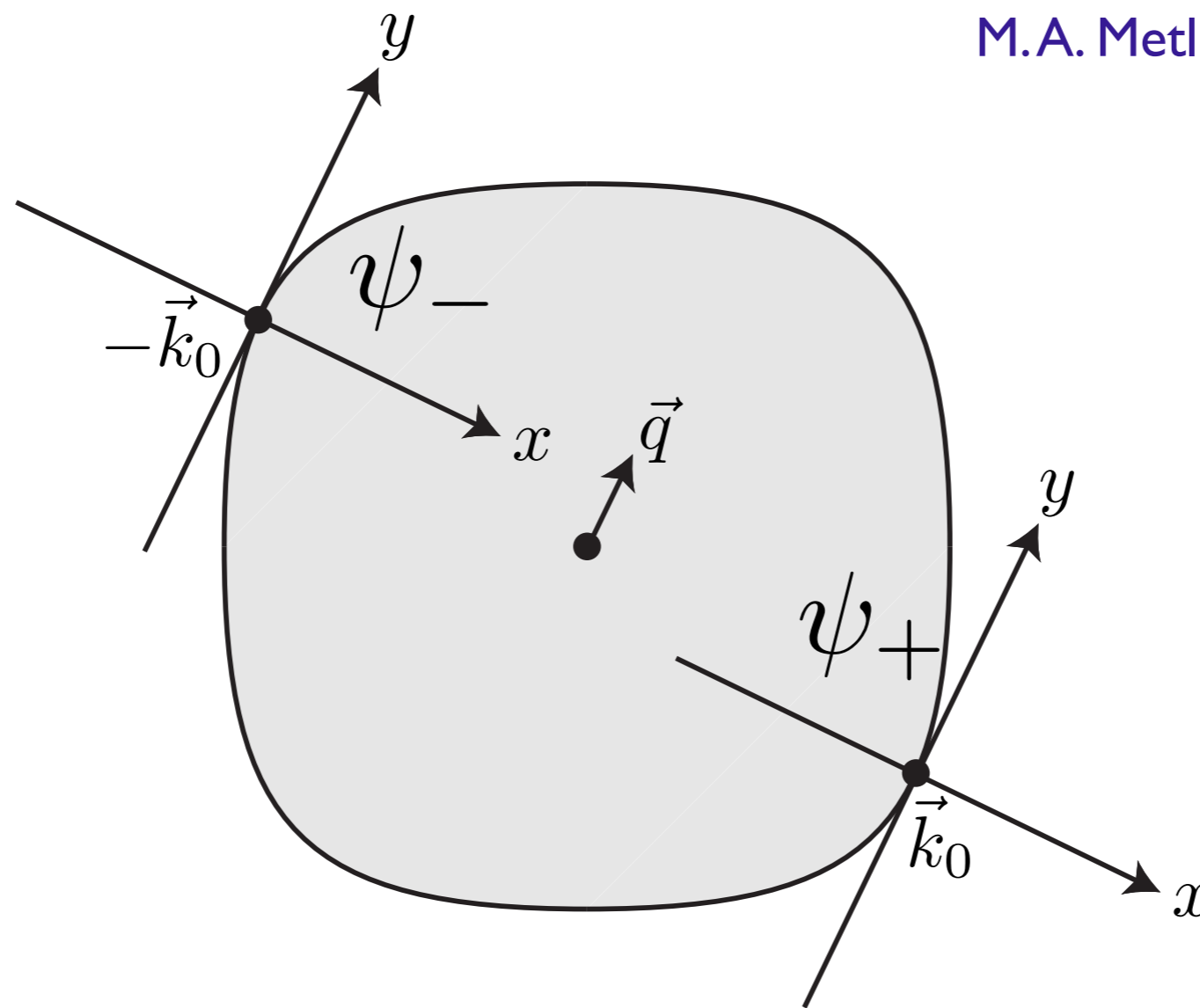


Spontaneous elongation along y direction:
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Quantum criticality of Pomeranchuk instability



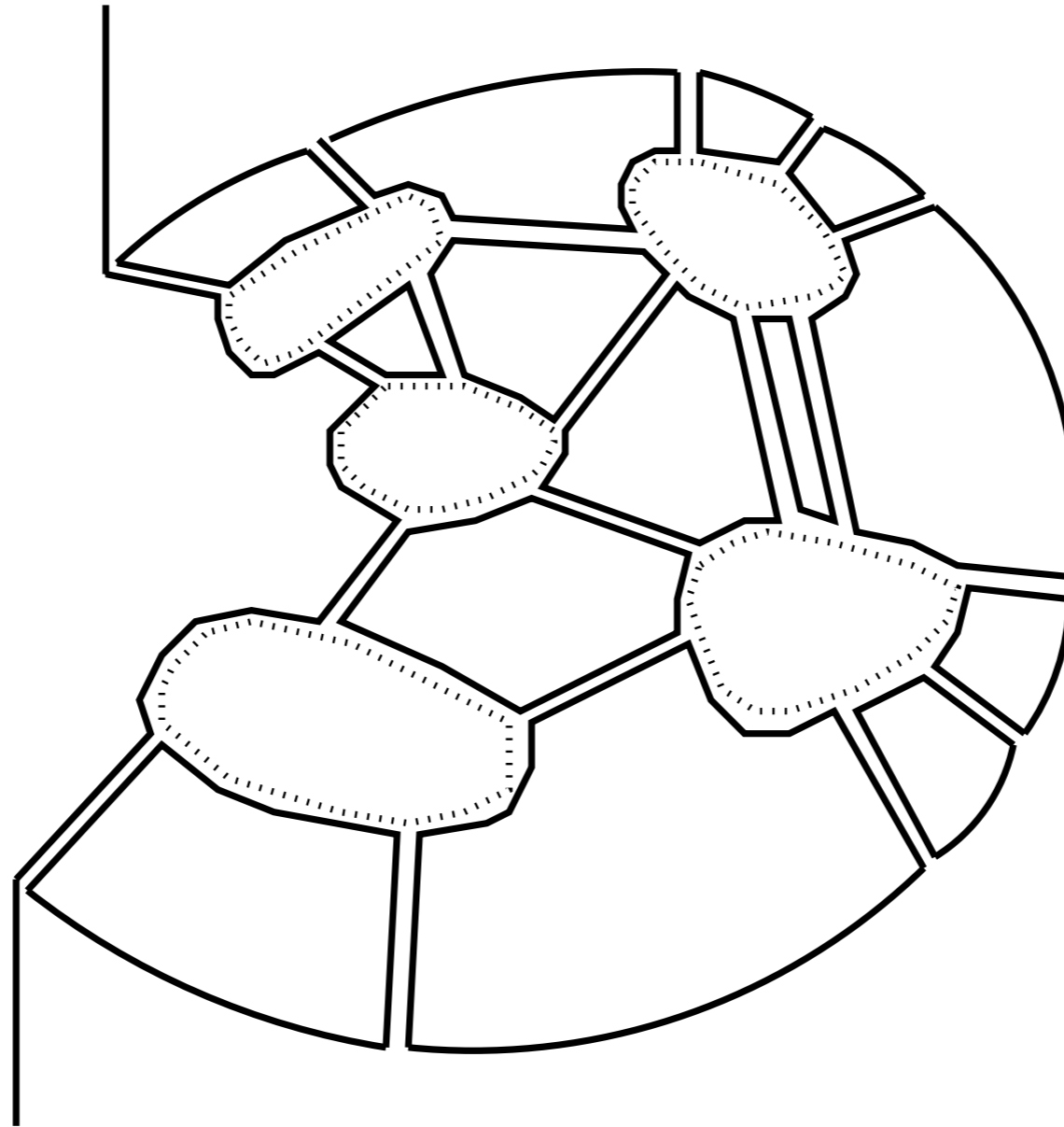
Pomeranchuk instability as a function of coupling λ



A ϕ fluctuation at wavevector \vec{q} couples most efficiently to fermions near $\pm\vec{k}_0$.

Infinite set of 2+1 dimensional quantum-critical field theories,
one for each pair of points on the Fermi surface.

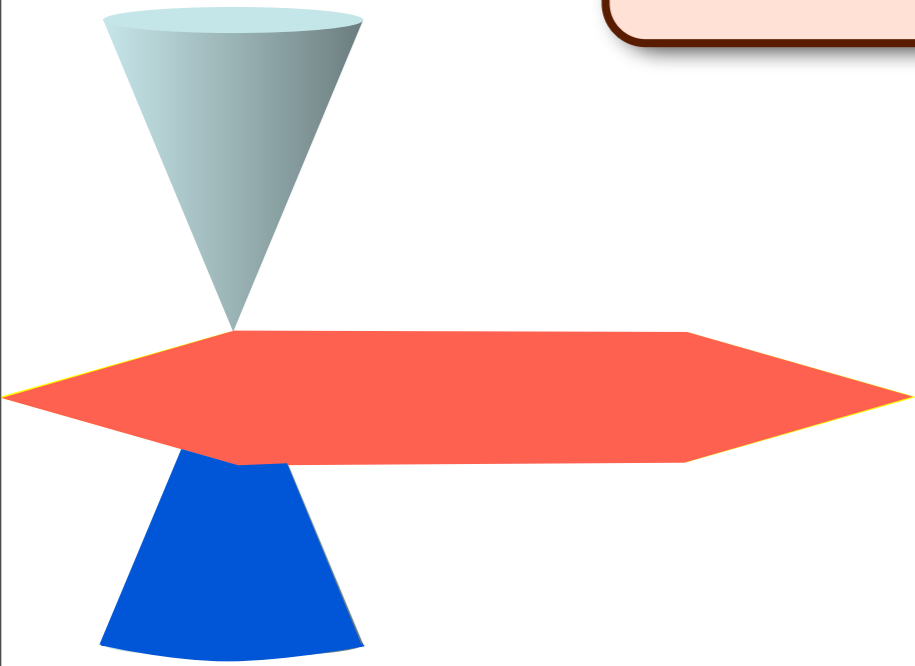
Ward identities ensure consistency
of redundant description



All planar graphs of ψ_+ alone
are as important as the leading
term

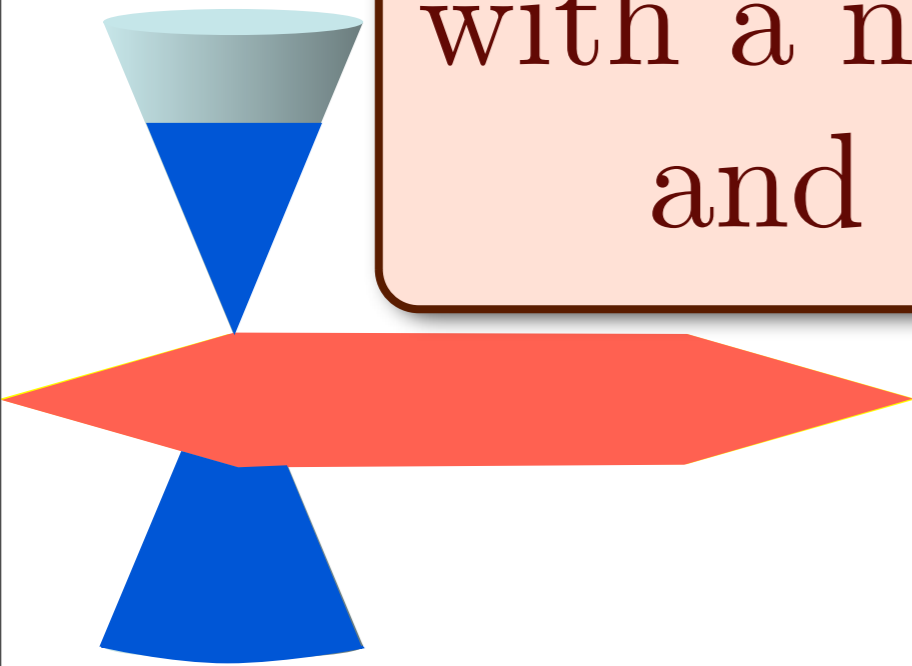
Sung-Sik Lee, *Physical Review B* **80**, 165102 (2009)

Conformal field theory
in $2+1$ dimensions at $T = 0$

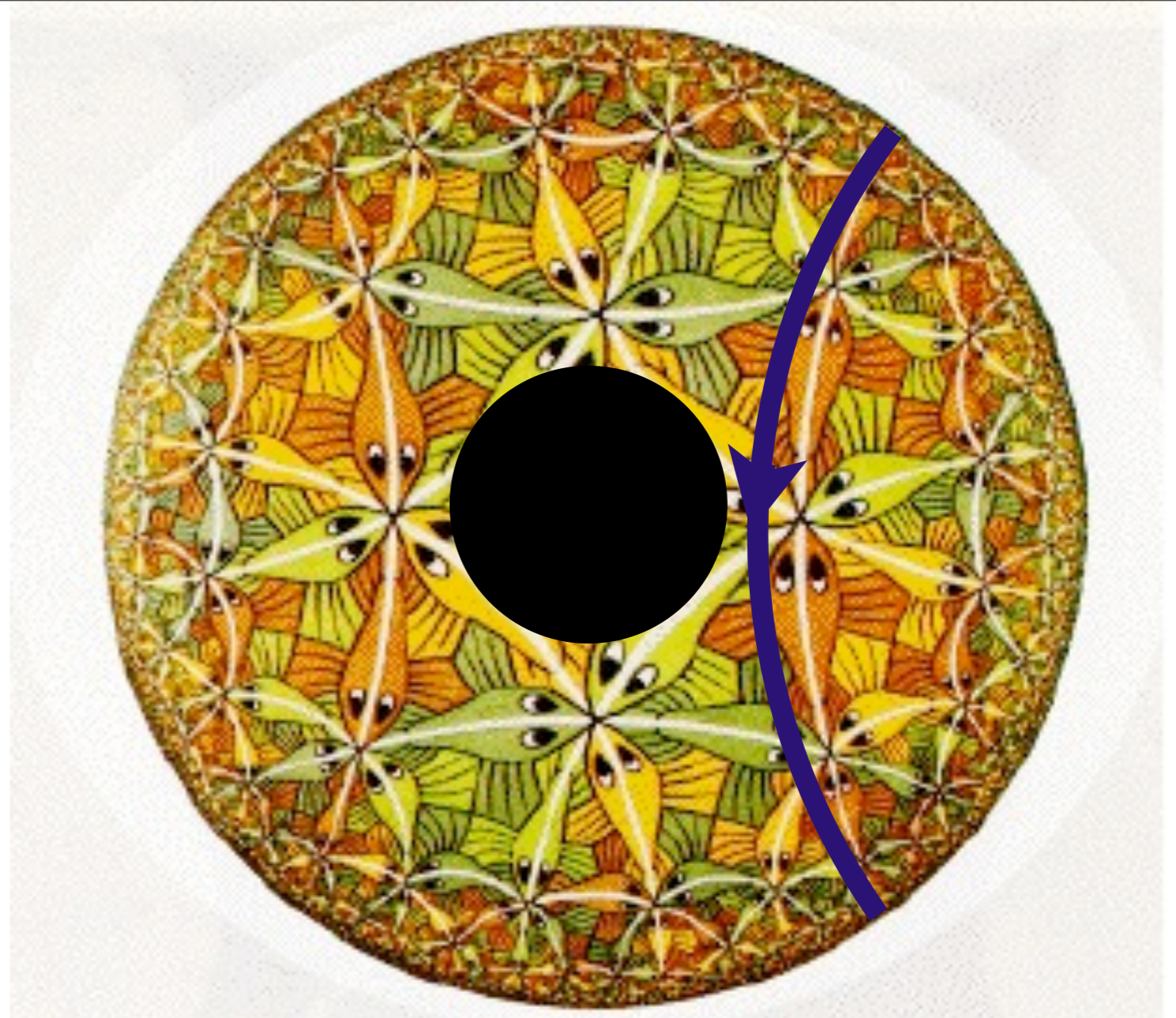
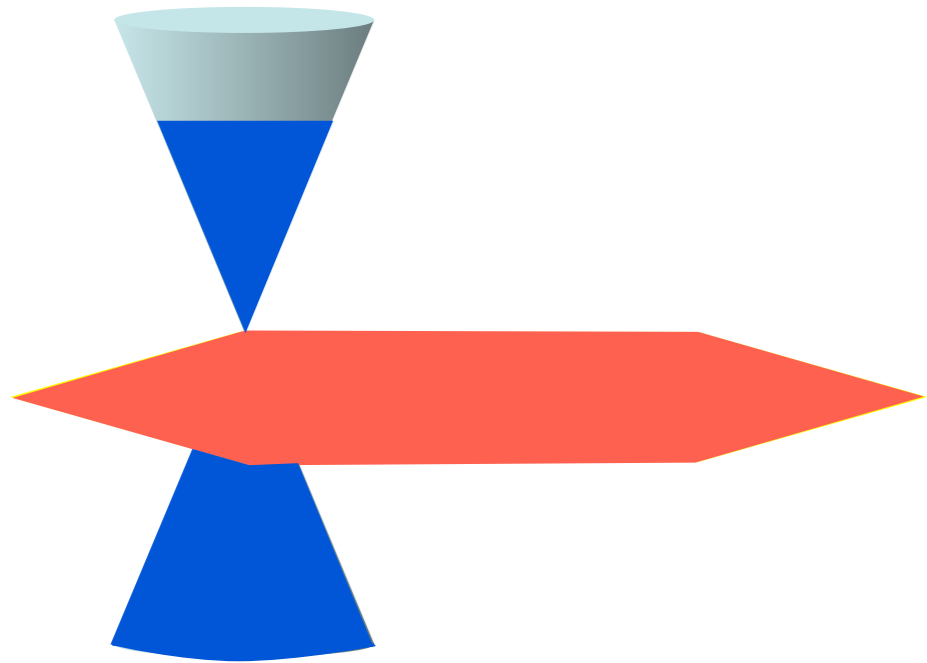


Einstein gravity
on AdS_4

Conformal field theory
in $2+1$ dimensions at $T > 0$,
with a non-zero chemical potential, μ
and applied magnetic field, B



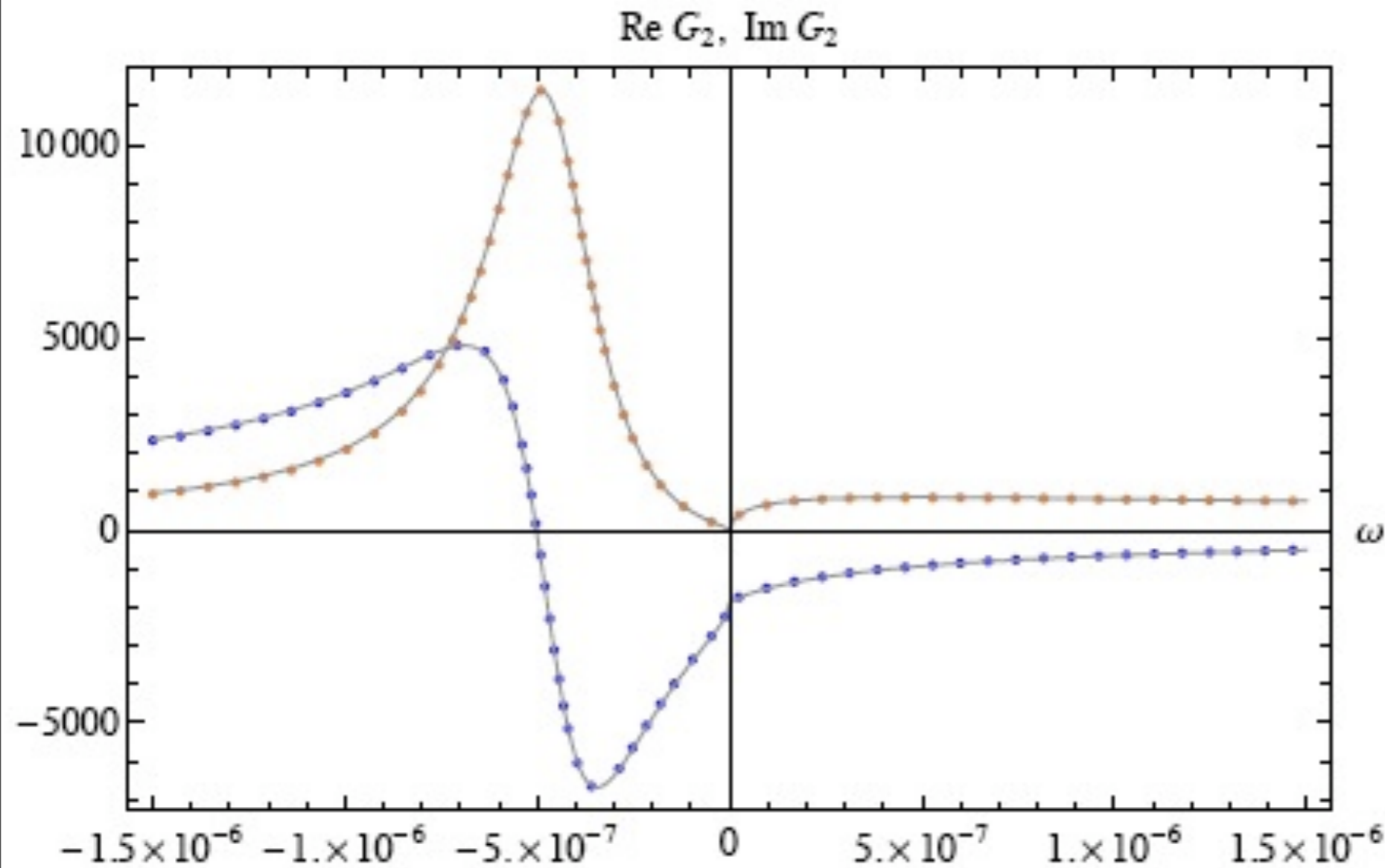
Einstein gravity on AdS_4
with a Reissner-Nordstrom
black hole carrying electric
and magnetic charges



Examine free energy and Green's function
of a probe particle

T. Faulkner, H. Liu, J. McGreevy, and D. Vegh, arXiv:0907.2694
F. Denef, S. Hartnoll, and S. Sachdev, arXiv:0908.1788

Green's function of a fermion



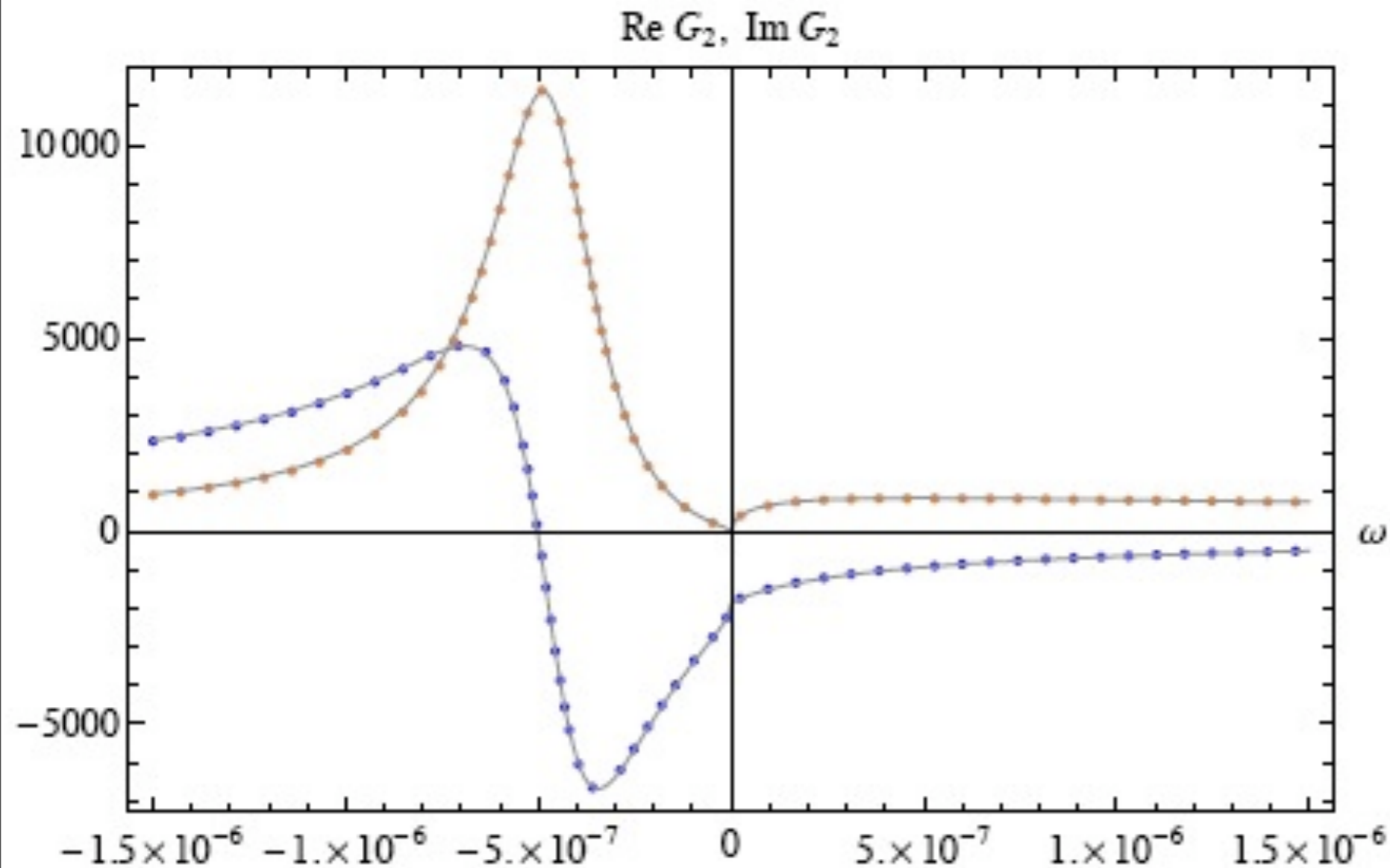
T. Faulkner, H. Liu,
J. McGreevy, and D. Vegh,
arXiv:0907.2694

$$G(k, \omega) \approx \frac{1}{\omega - v_F(k - k_F) - i\omega^\theta(k)}$$

See also M. Cubrovic, J. Zaanen, and K. Schalm, arXiv:0904.1993

F. Denef, S. Hartnoll, and S. Sachdev, arXiv:0908.1788

Green's function of a fermion



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Similar to non-Fermi liquid theories of Fermi surfaces coupled to gauge fields, and at quantum critical points