



Quantum Phase Transitions

Subir Sachdev

Talks online at <http://sachdev.physics.harvard.edu>

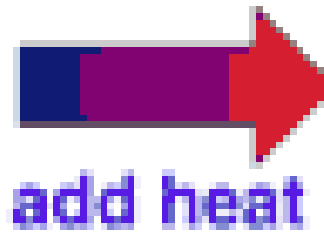
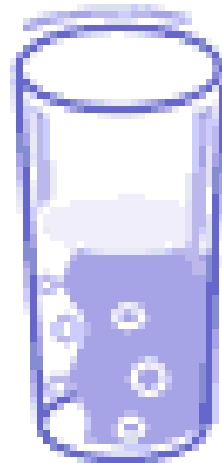
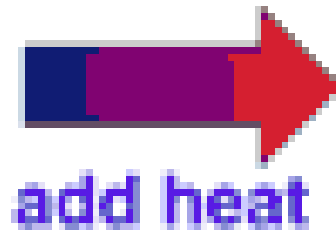


What is a “phase transition” ?

ice

water

steam



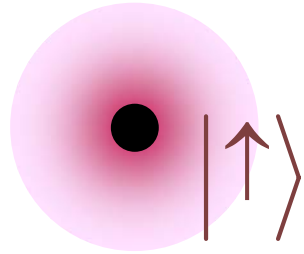
A change in the collective properties of a macroscopic number of atoms

What is a “*quantum phase transition*” ?

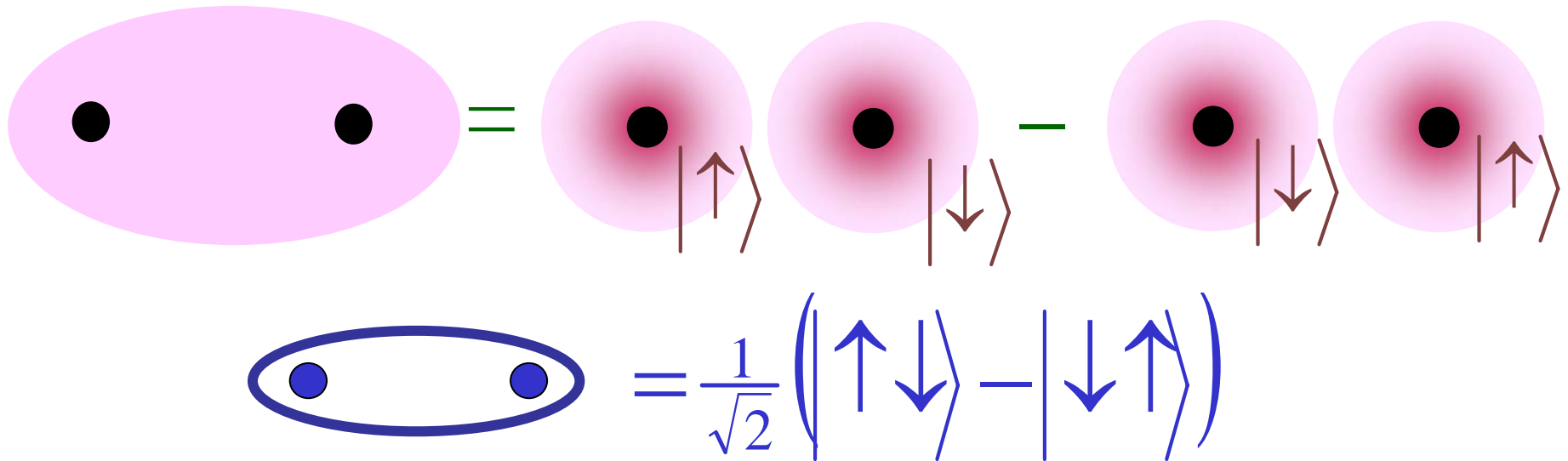
Change in the nature of
entanglement in a macroscopic
quantum system.

Entanglement

Hydrogen atom:



Hydrogen molecule:



Superposition of two electron states leads to non-local correlations between spins

Outline

Quantum phase transitions

1. Spin ordering in “Han purple”
2. Entanglement at the critical point: physical consequences at non-zero temperatures
 - (a) Double-layer antiferromagnet
 - (b) Superfluid-insulator transition
 - (c) Hydrodynamics via mapping to quantum theory of black holes.
3. Entanglement of valence bonds
4. Conclusions

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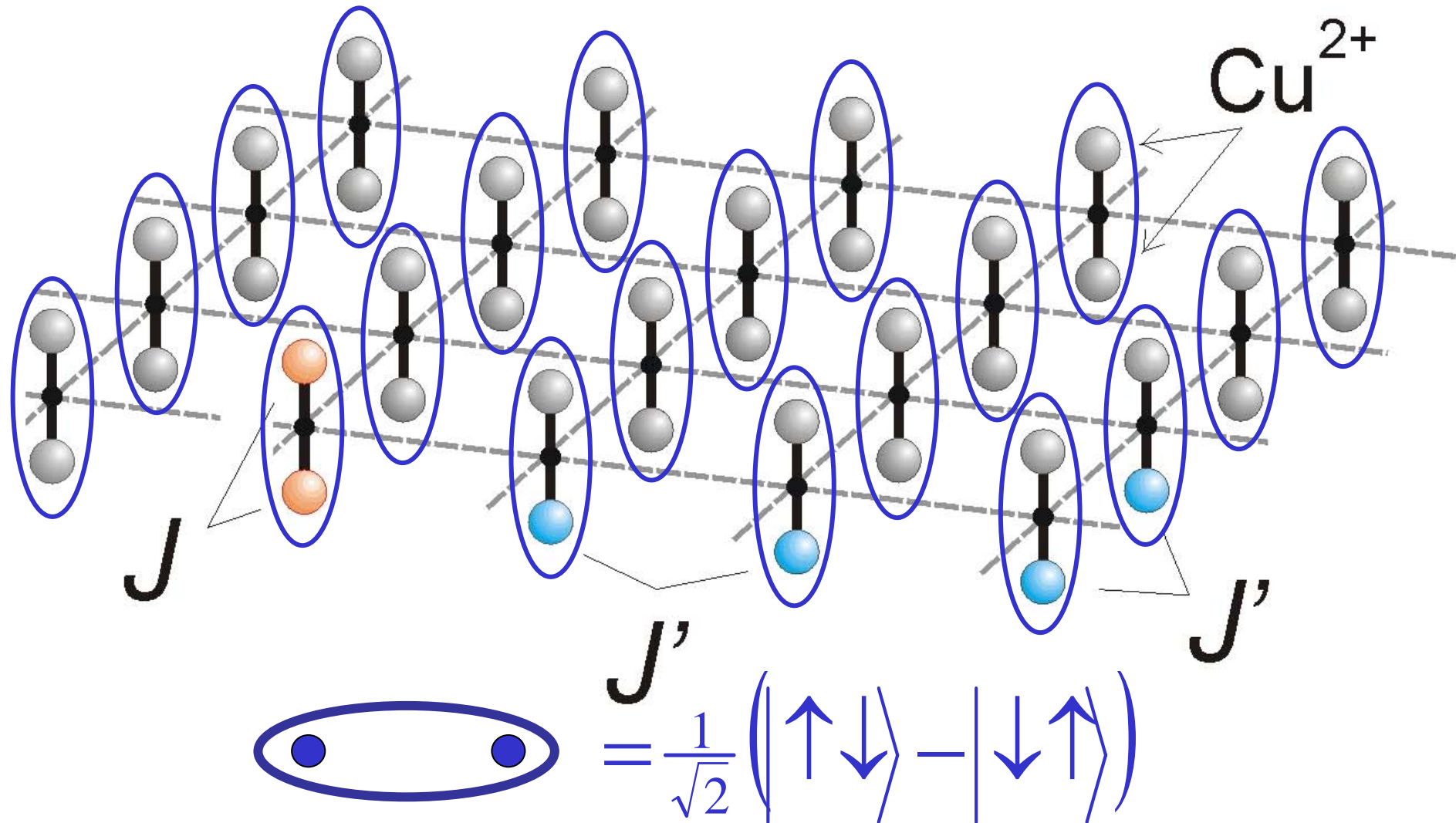
Chinese Terracotta warriors (479-221 BC)



Han Purple – $\text{BaCuSi}_2\text{O}_6$

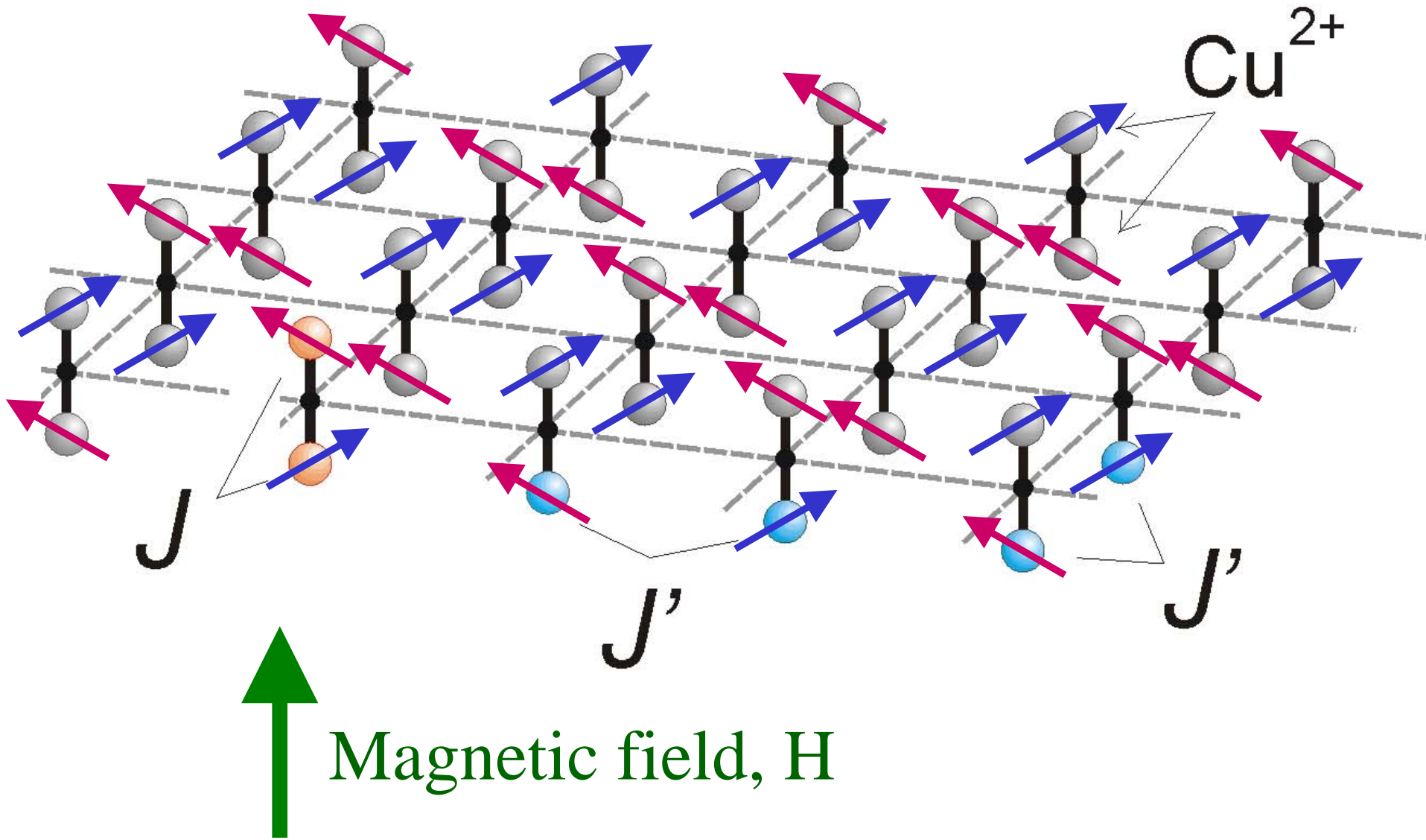
Weak magnetic field

Han Purple – $\text{BaCuSi}_2\text{O}_6$



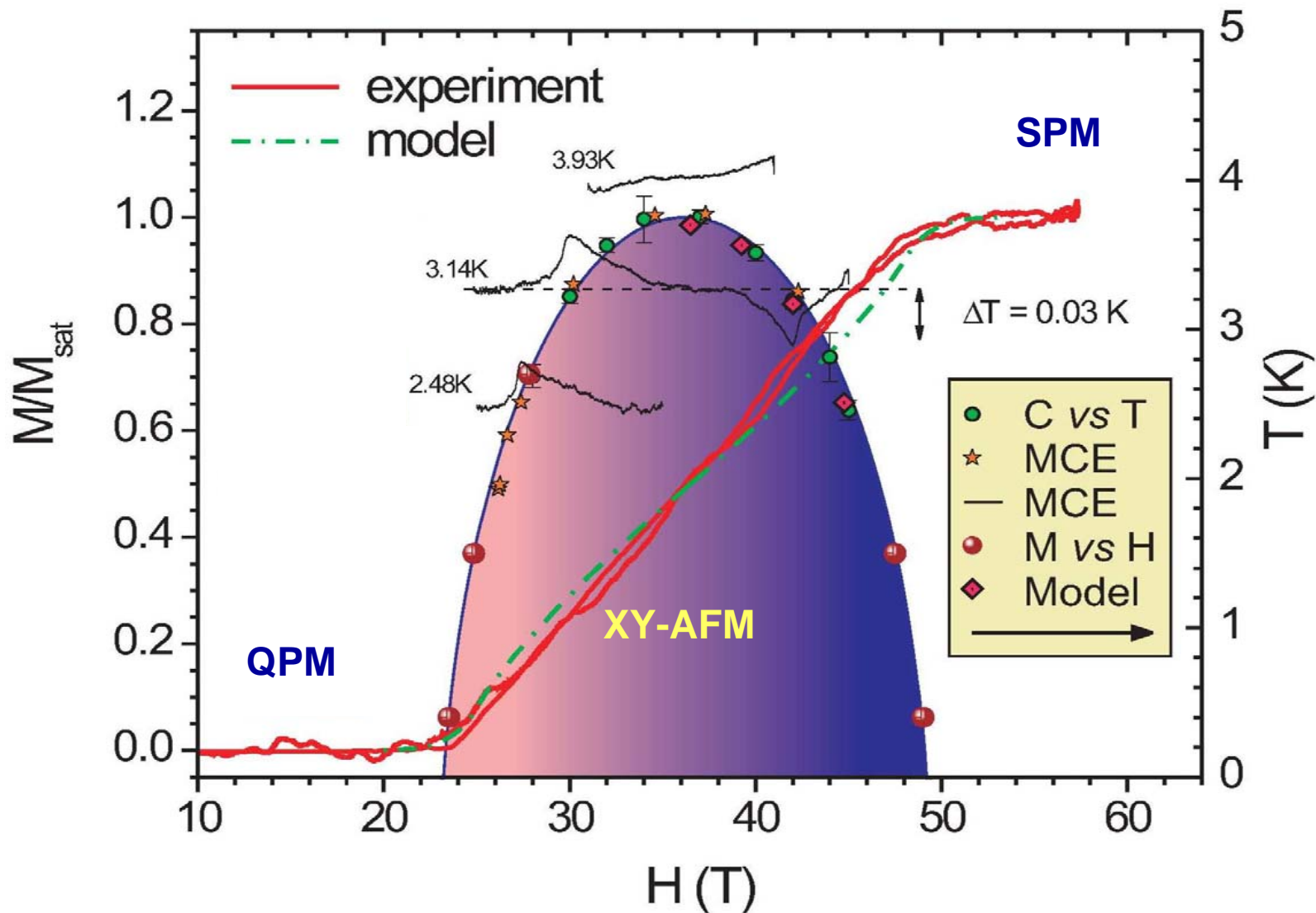
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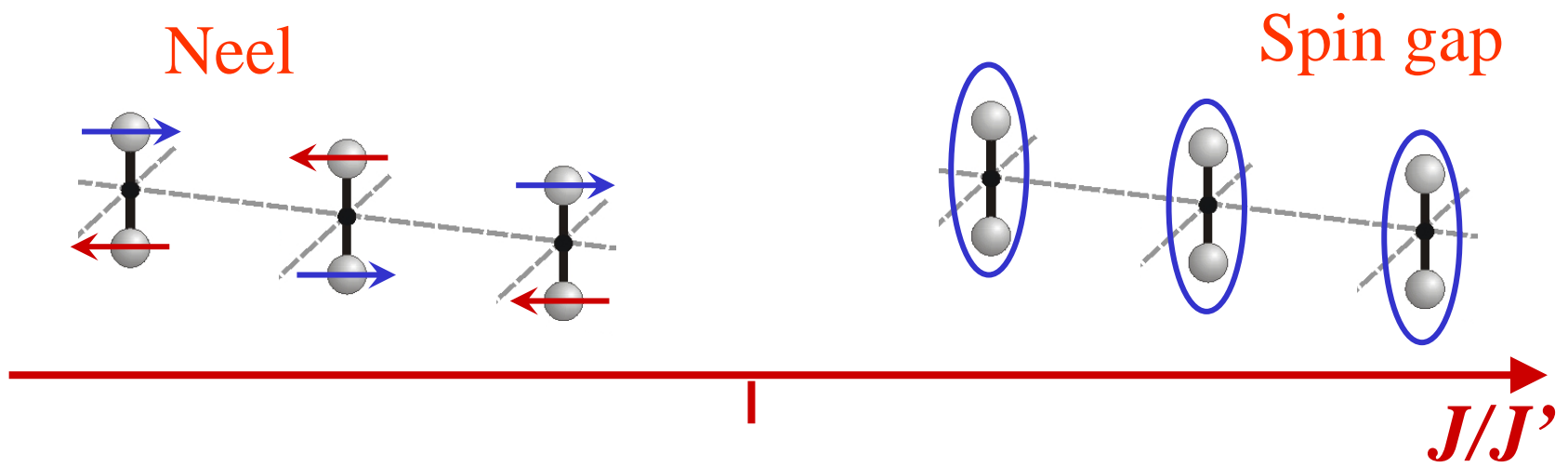
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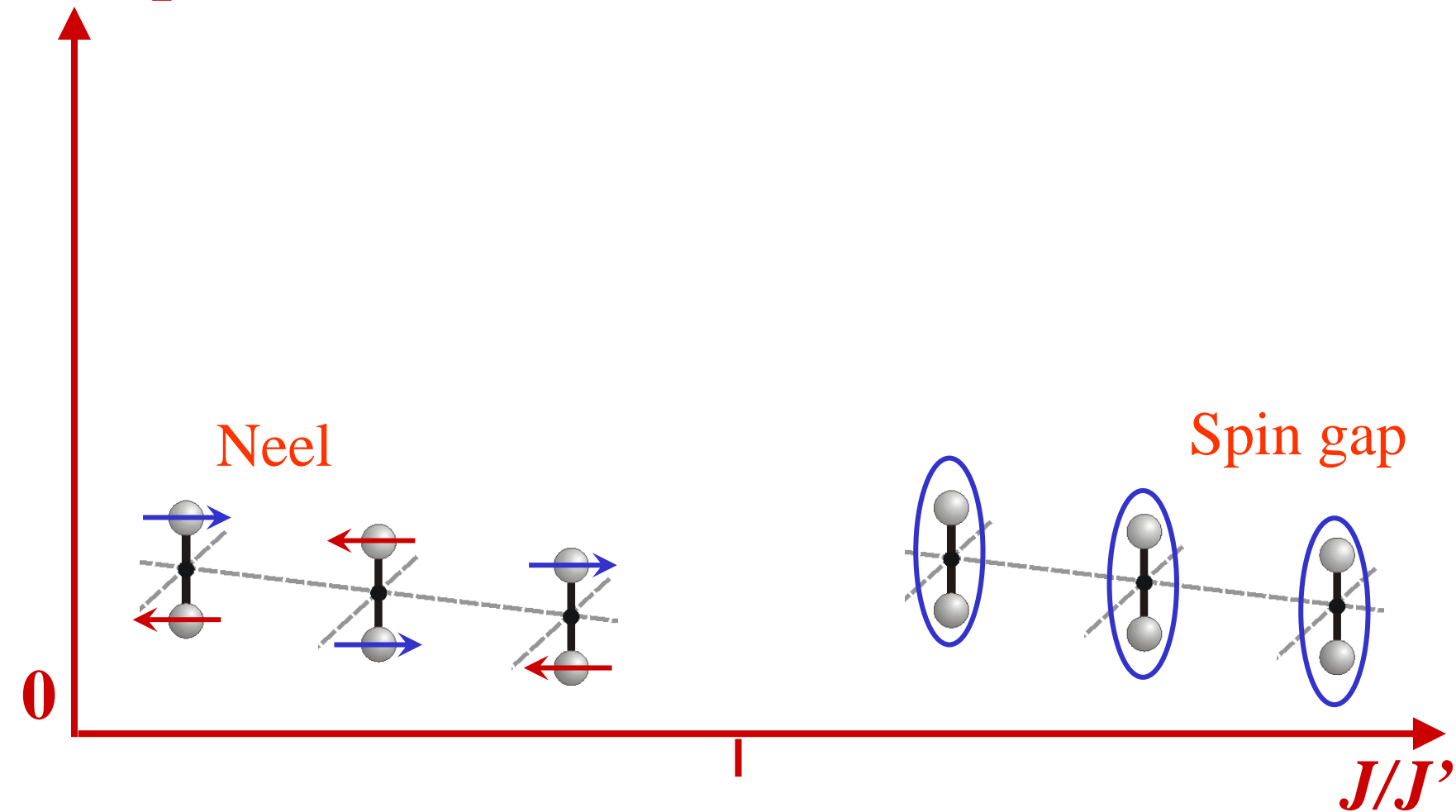
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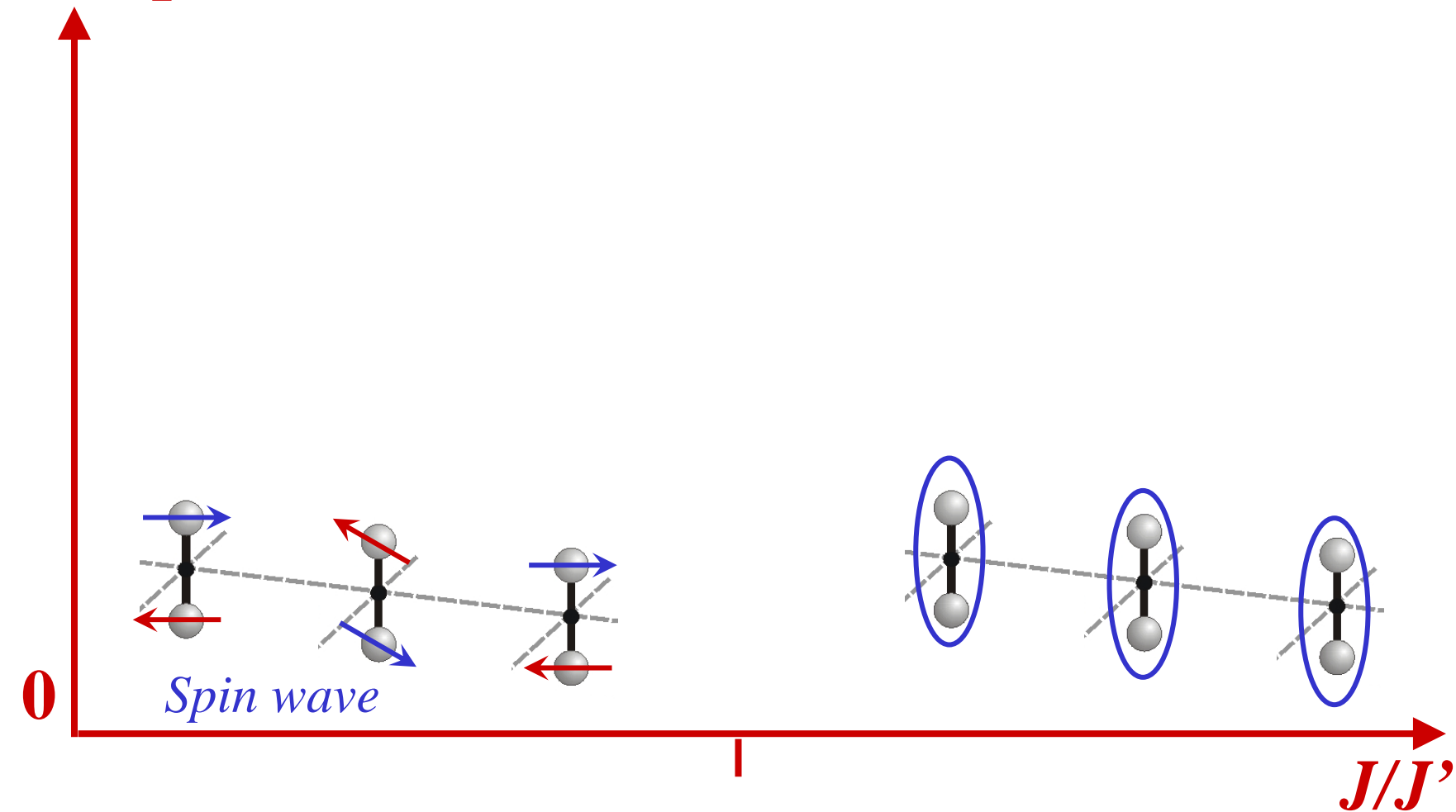
Vary the ratio J/J'

Temperature, T



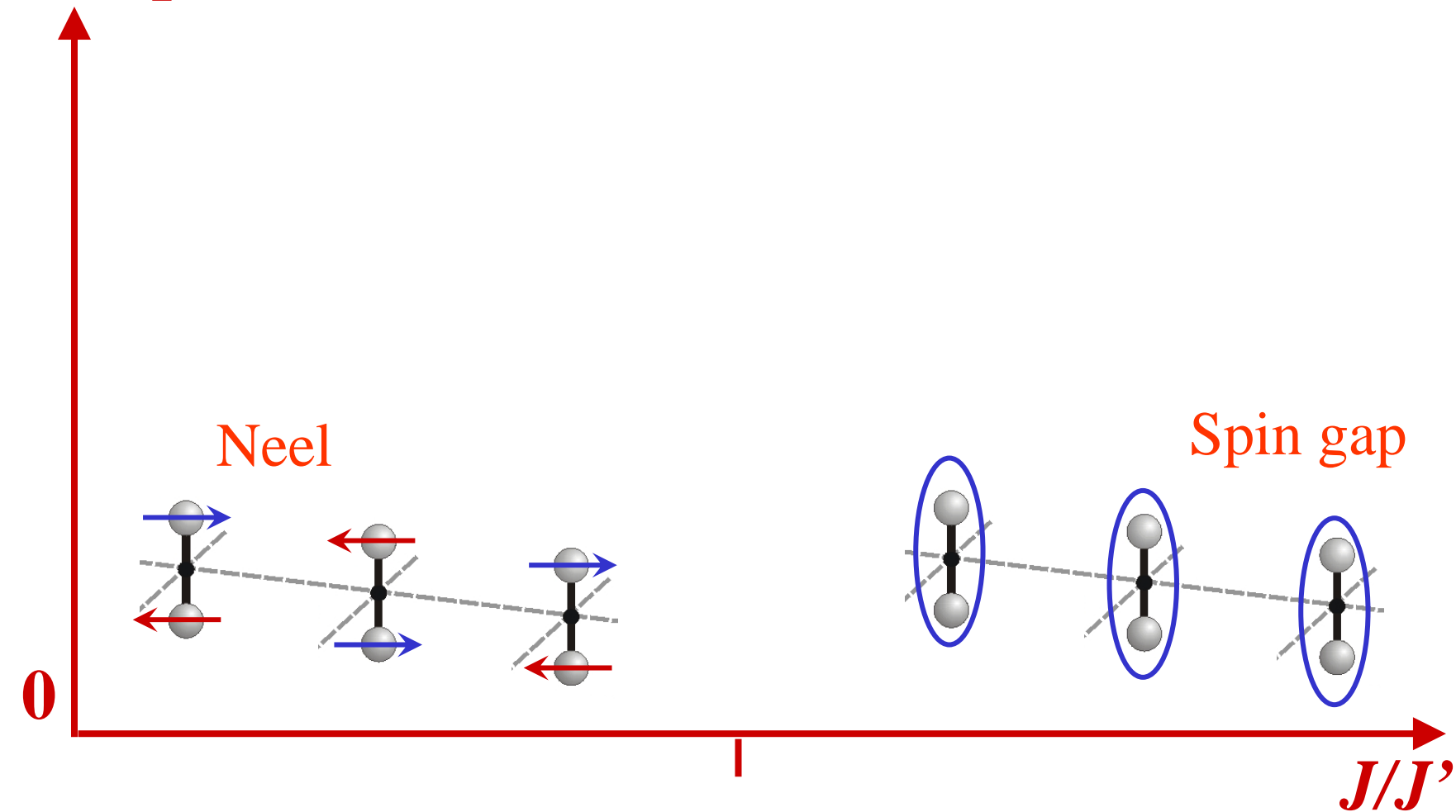
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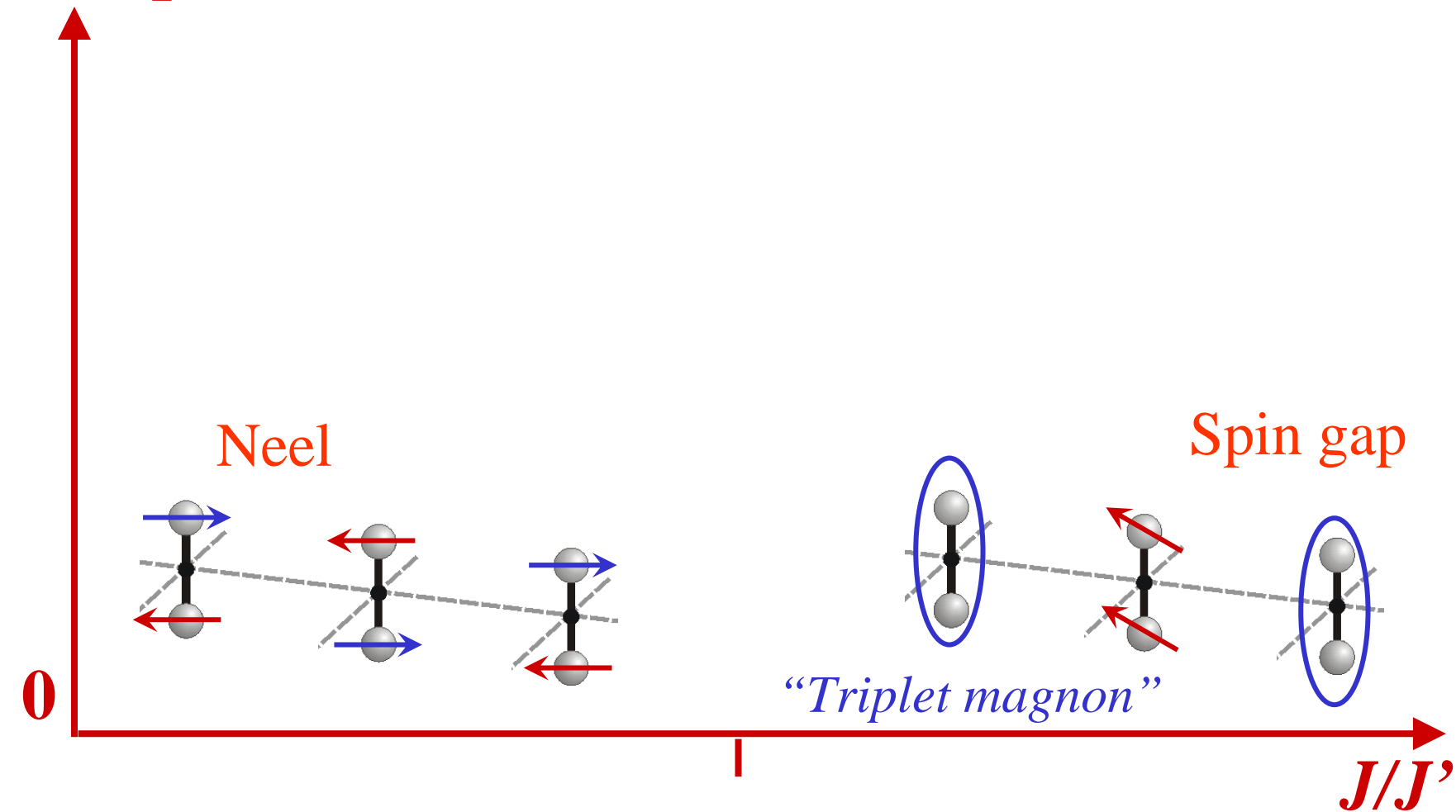
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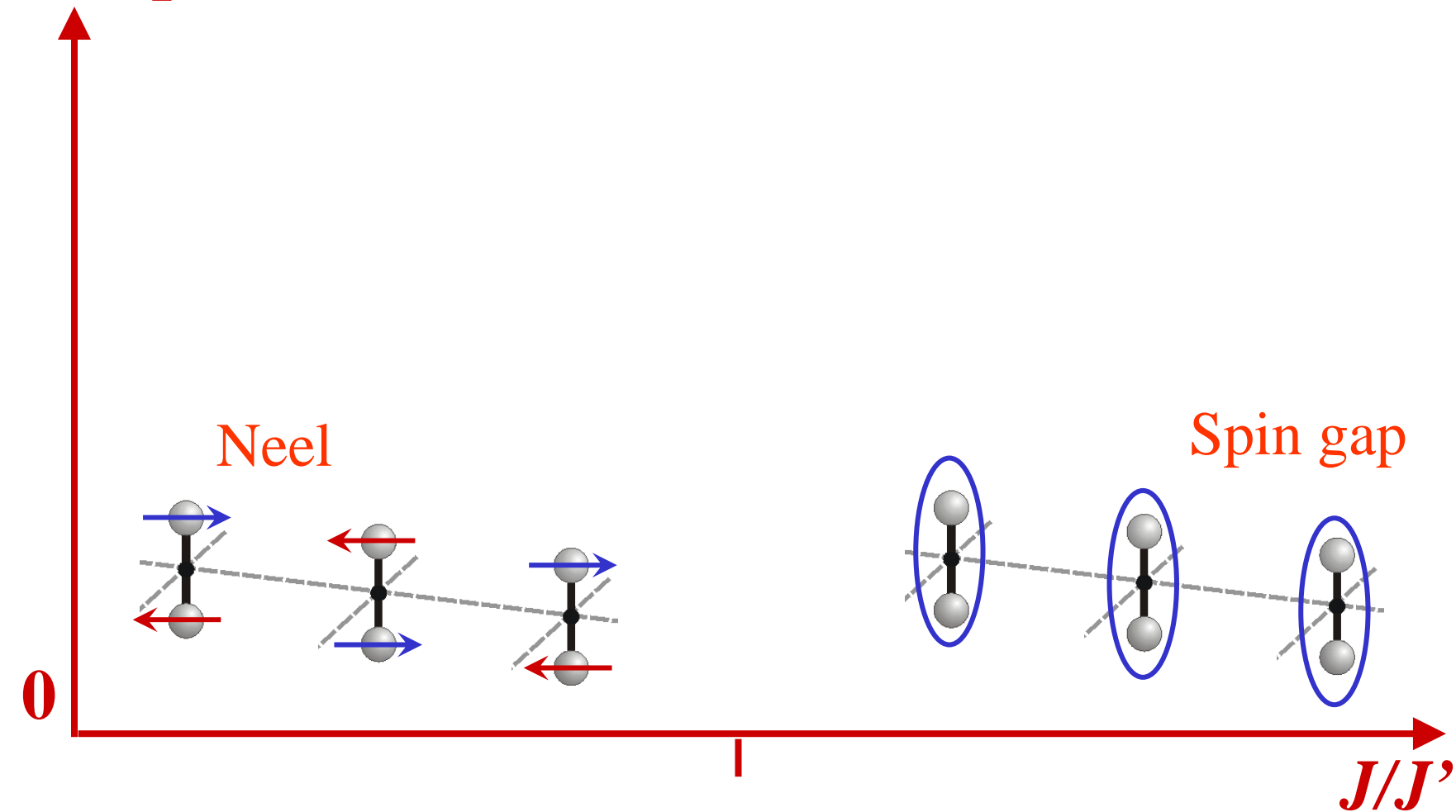
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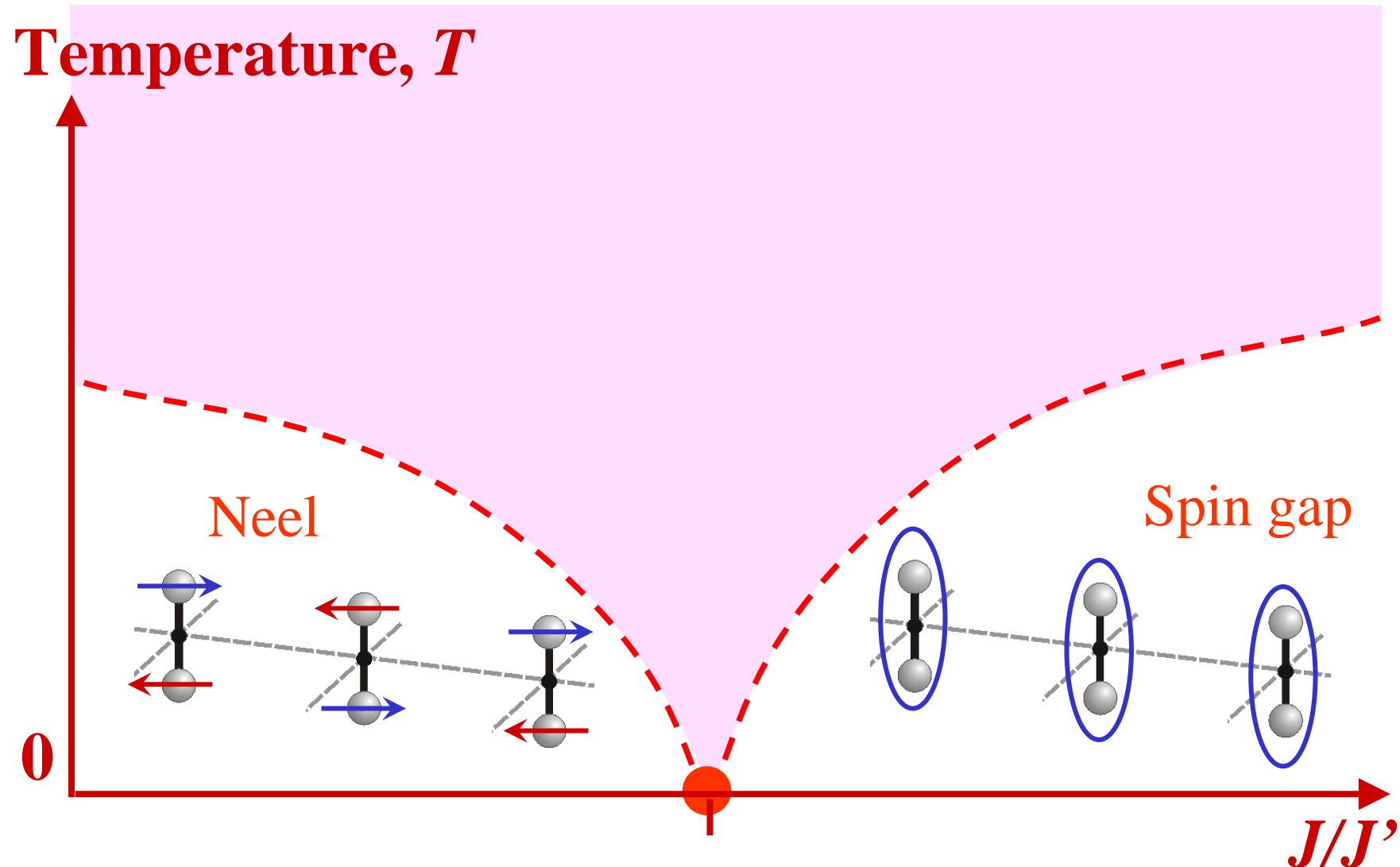


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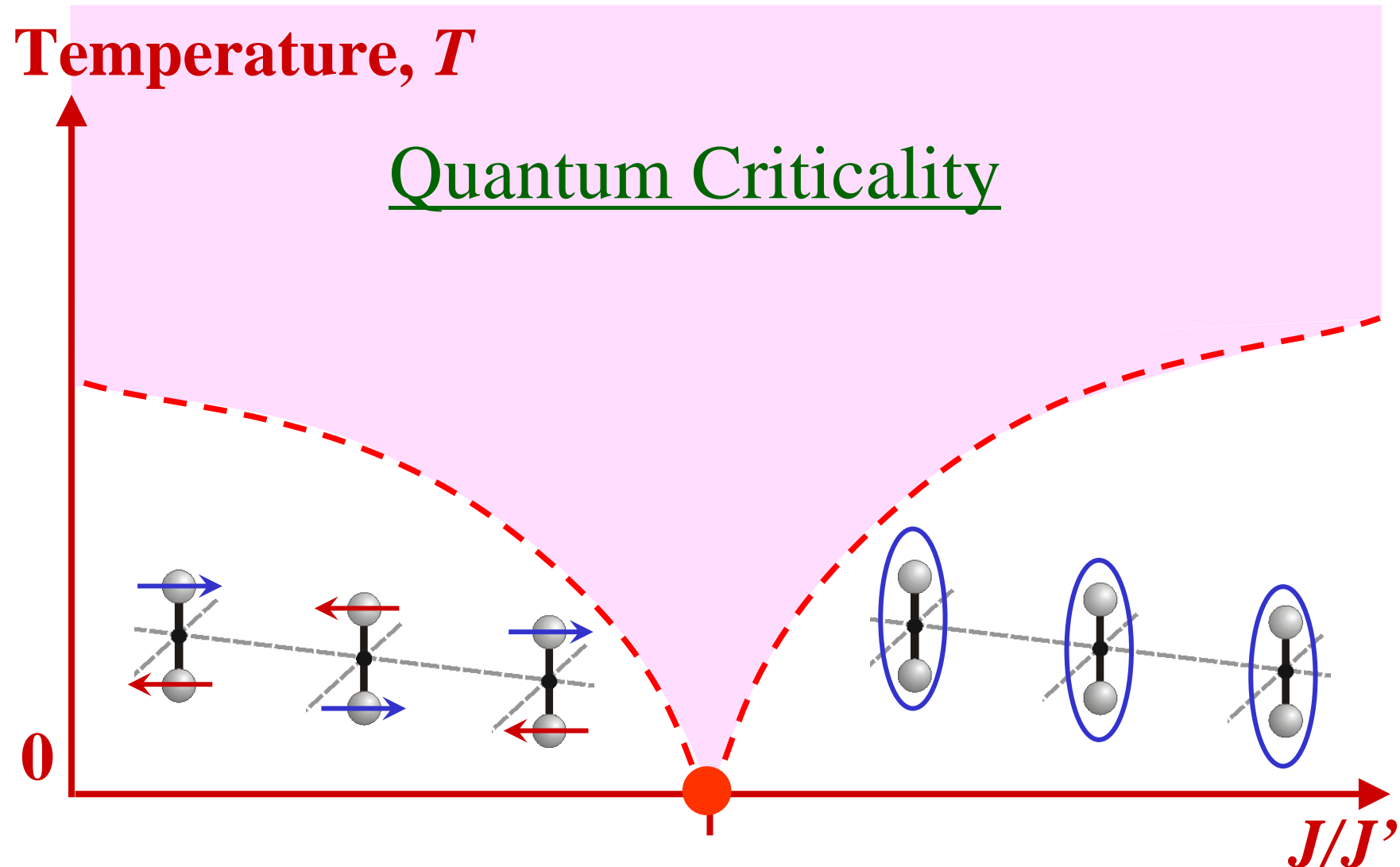
Temperature, T



Vary the ratio J/J'



Vary the ratio J/J'

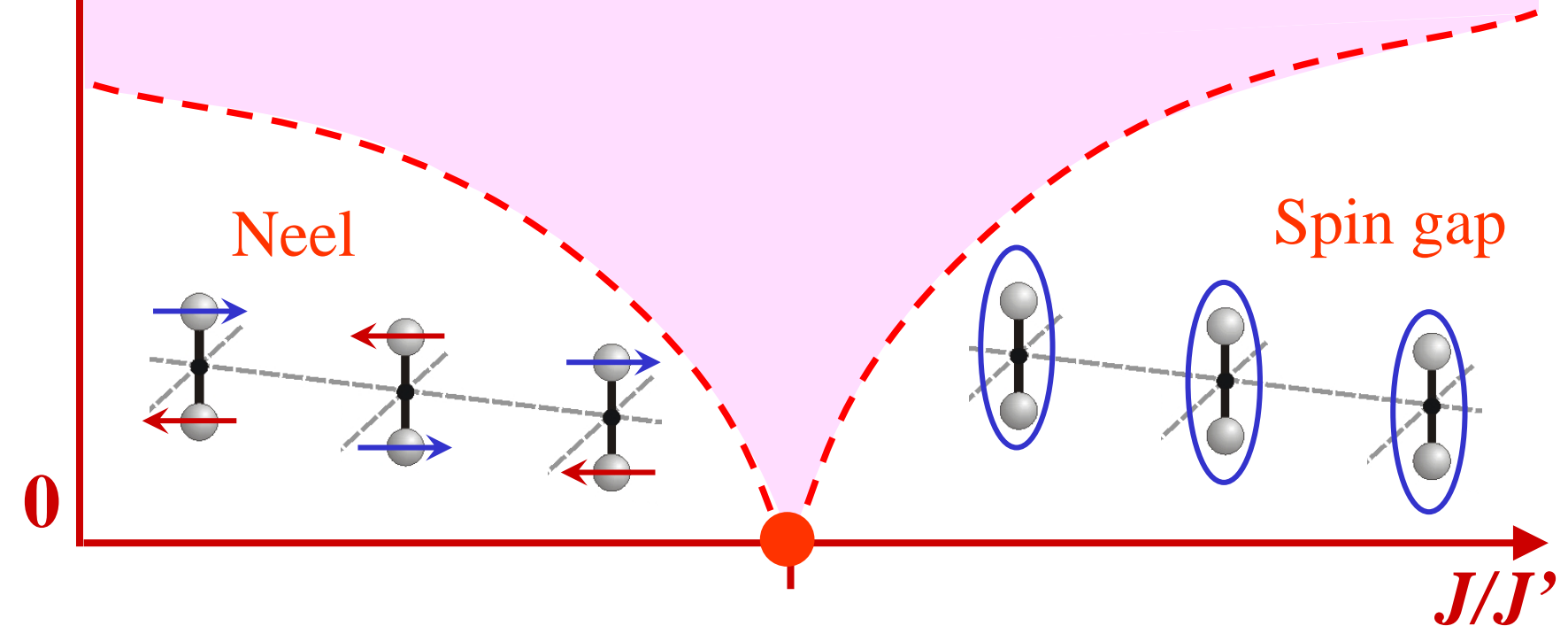


Vary the ratio J/J'

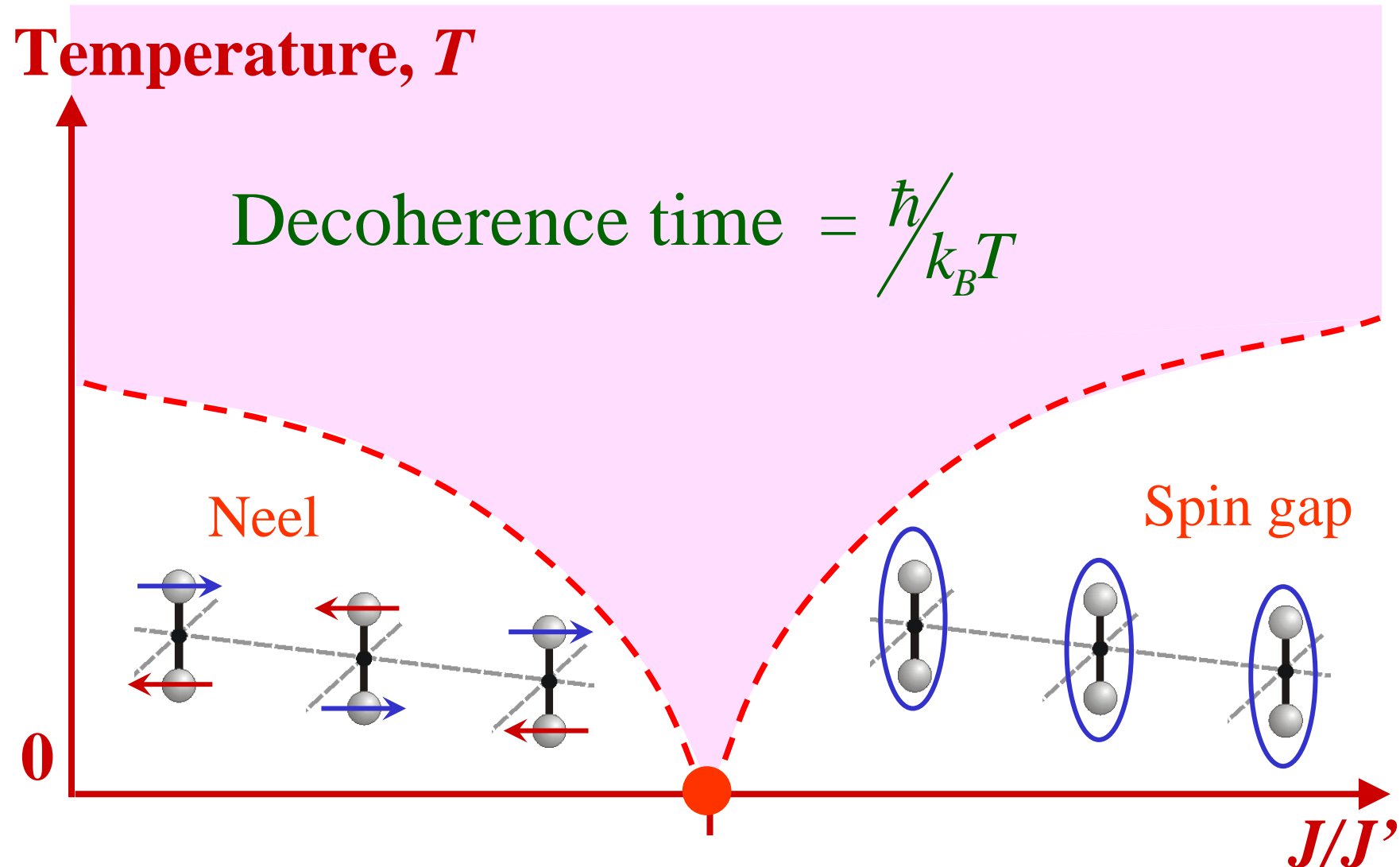
Temperature, T

Quantum Criticality

Thermal excitations interact via a universal S matrix.



Vary the ratio J/J'



Decoherence time = $\frac{\hbar}{k_B T}$

Vary the ratio J/J'

Temperature, T

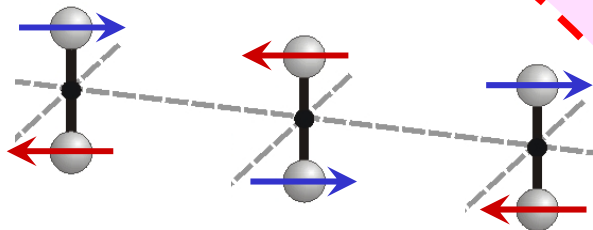
Quantum critical transport

Spin diffusion constant

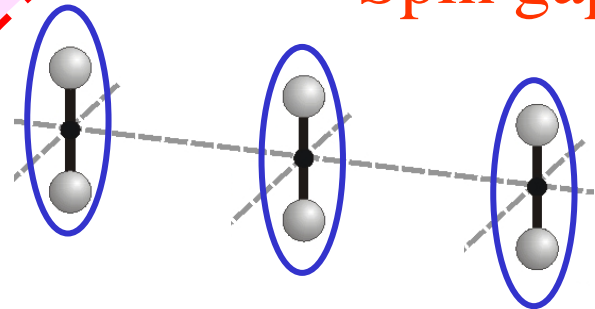
$$D_s = \Theta \frac{c^2}{k_B T}$$

where Θ is a universal number

Neel



Spin gap



0

J/J'

Vary the ratio J/J'

Outline

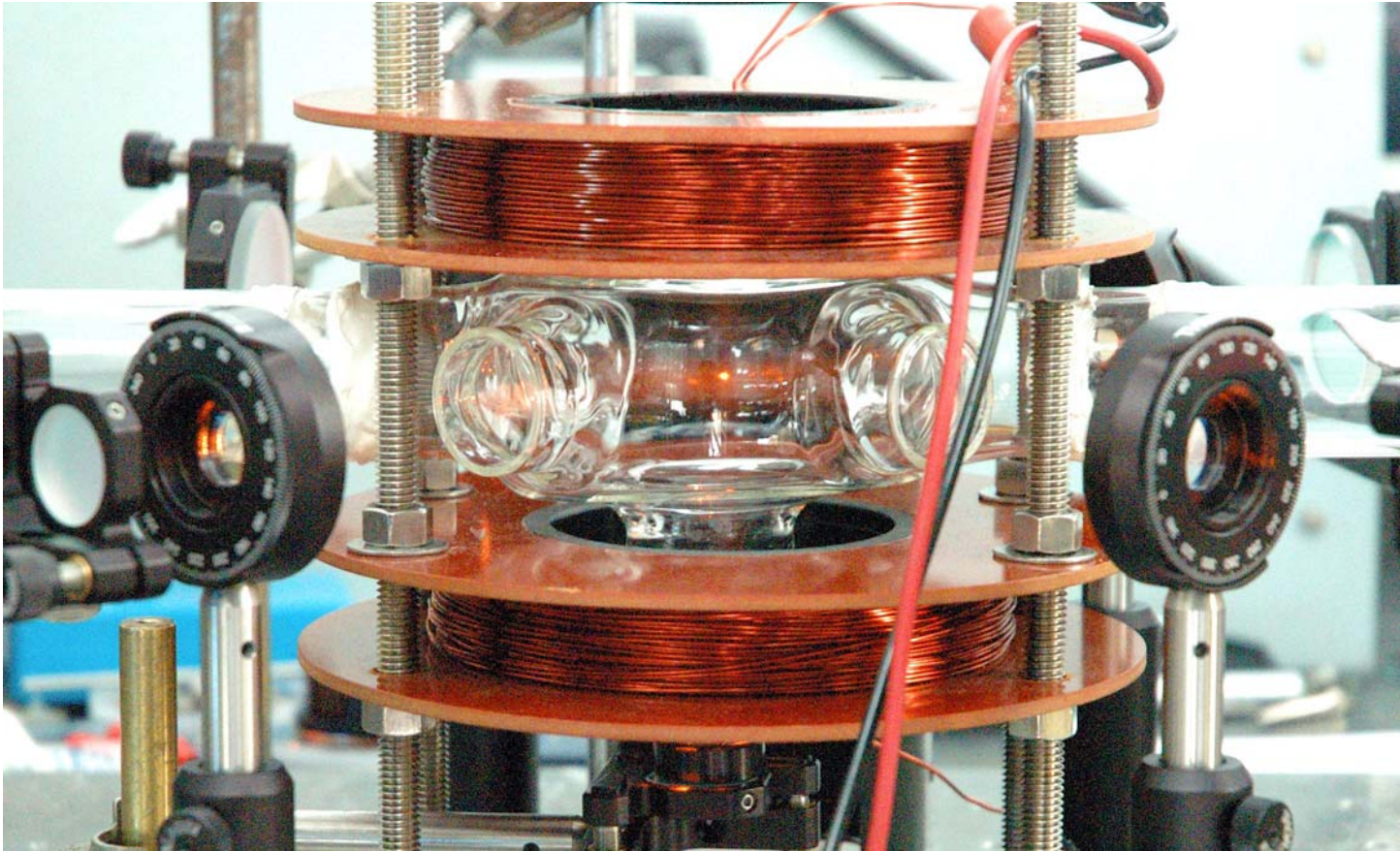
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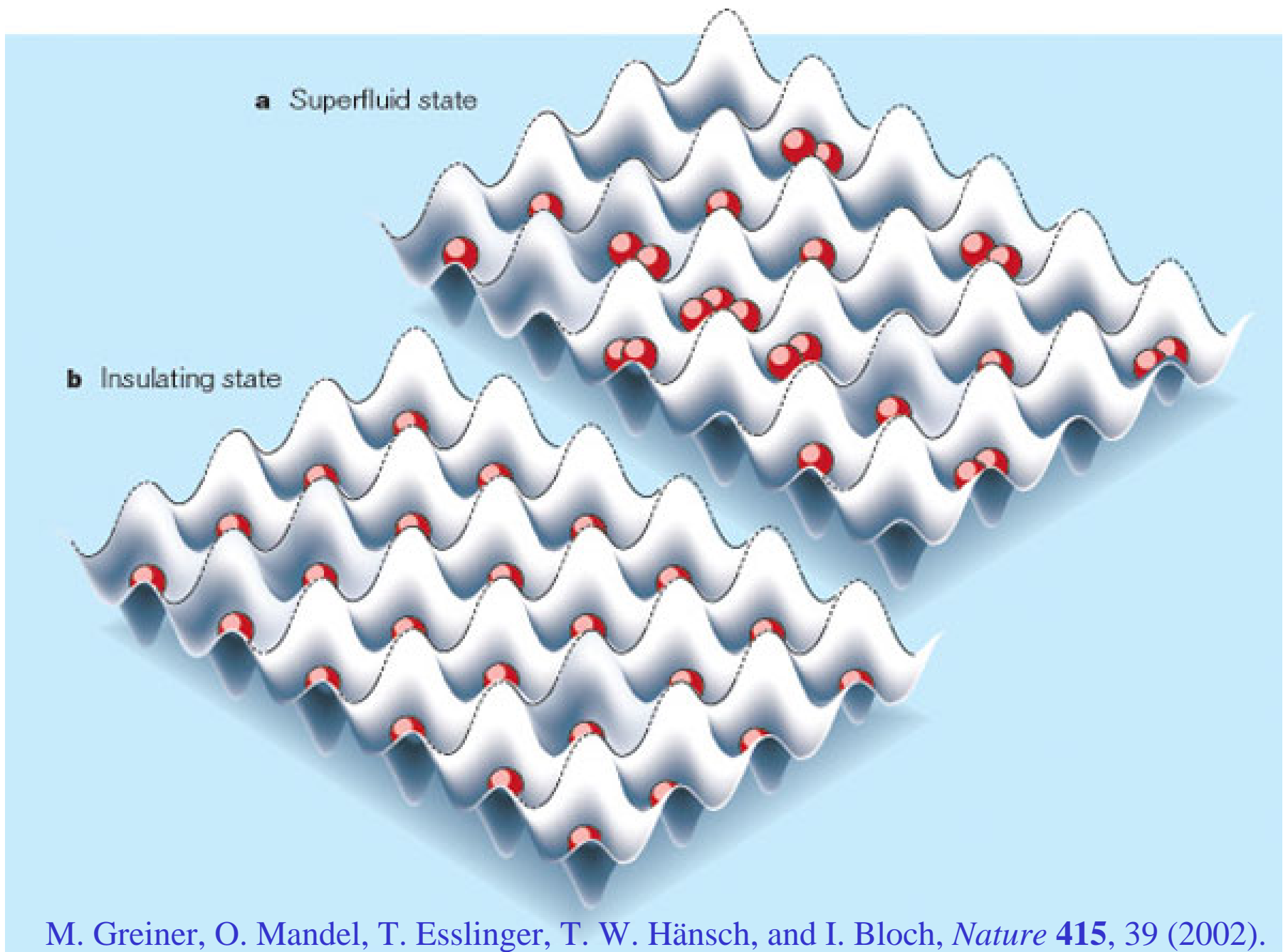
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Trap for ultracold ^{87}Rb atoms



M. Greiner, O. Mandel, T. Esslinger, T. W. Hänsch, and I. Bloch, *Nature* **415**, 39 (2002).

The Bose-Einstein condensate in a periodic potential

$$|G\rangle = \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right|$$

Lowest energy state for many atoms

$$\begin{aligned}
 |BEC\rangle &= |G\rangle|G\rangle|G\rangle \\
 &= \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \\
 &\quad + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| + \dots 27 \text{ terms}
 \end{aligned}$$

Large fluctuations in number of atoms in each potential well
 – *superfluidity* (atoms can “flow” without dissipation)

Breaking up the Bose-Einstein condensate

$$|G\rangle = \left| \left| \begin{array}{c} \circ \\ | \\ | \end{array} \right\rangle + \left| \left| \begin{array}{c} | \\ \circ \\ | \end{array} \right\rangle + \left| \left| \begin{array}{c} | \\ | \\ \circ \end{array} \right\rangle \right.$$

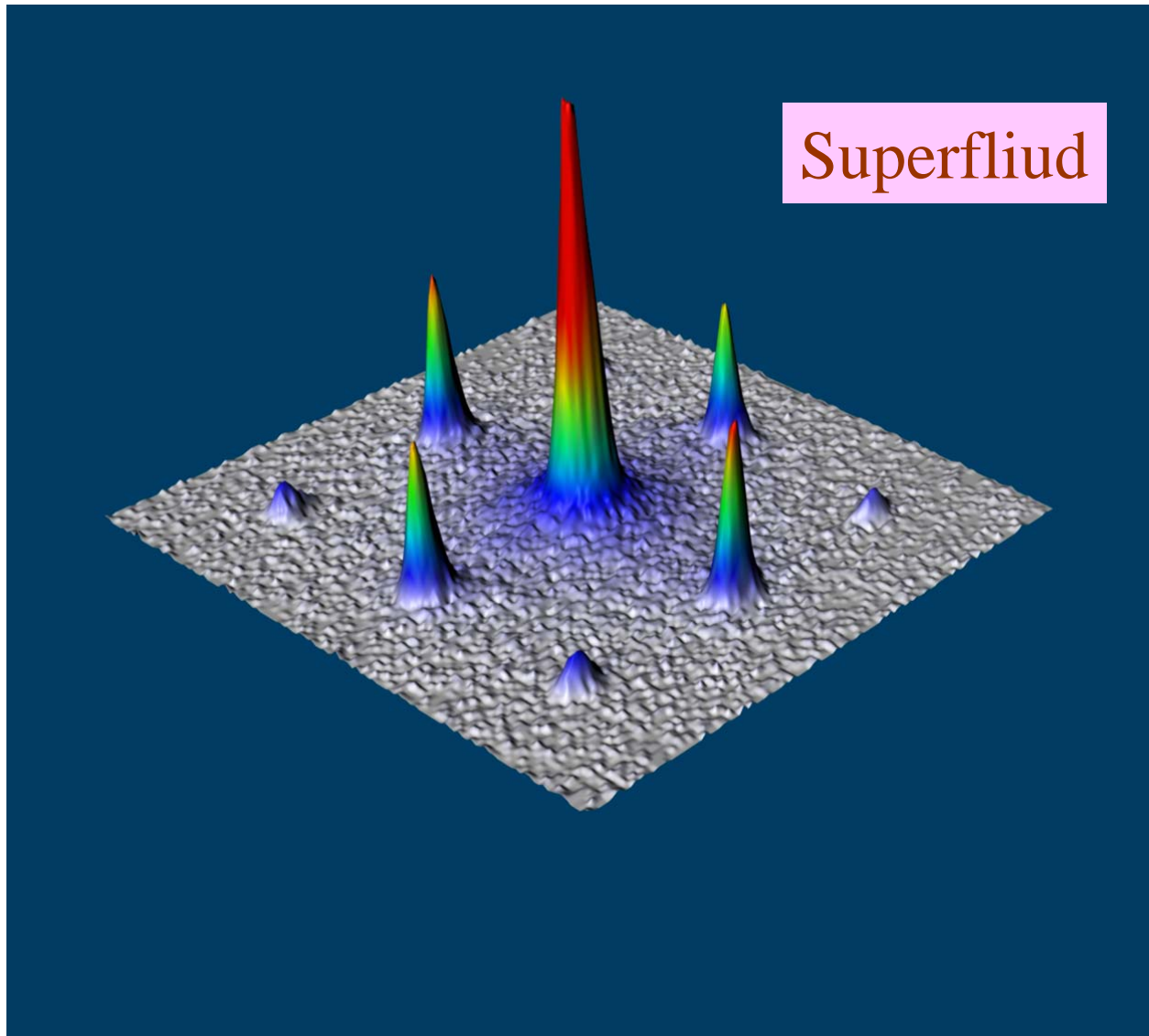
Lowest energy state for many atoms

$$\begin{aligned} |MI\rangle = & \left| \left| \begin{array}{c} \circ \\ \circ \\ \circ \end{array} \right\rangle + \left| \left| \begin{array}{c} \circ \\ \circ \\ \circ \end{array} \right\rangle + \left| \left| \begin{array}{c} \circ \\ \circ \\ \circ \end{array} \right\rangle \right. \\ & + \left| \left| \begin{array}{c} \circ \\ \circ \\ \circ \end{array} \right\rangle + \left| \left| \begin{array}{c} \circ \\ \circ \\ \circ \end{array} \right\rangle + \left| \left| \begin{array}{c} \circ \\ \circ \\ \circ \end{array} \right\rangle \right. \end{aligned}$$

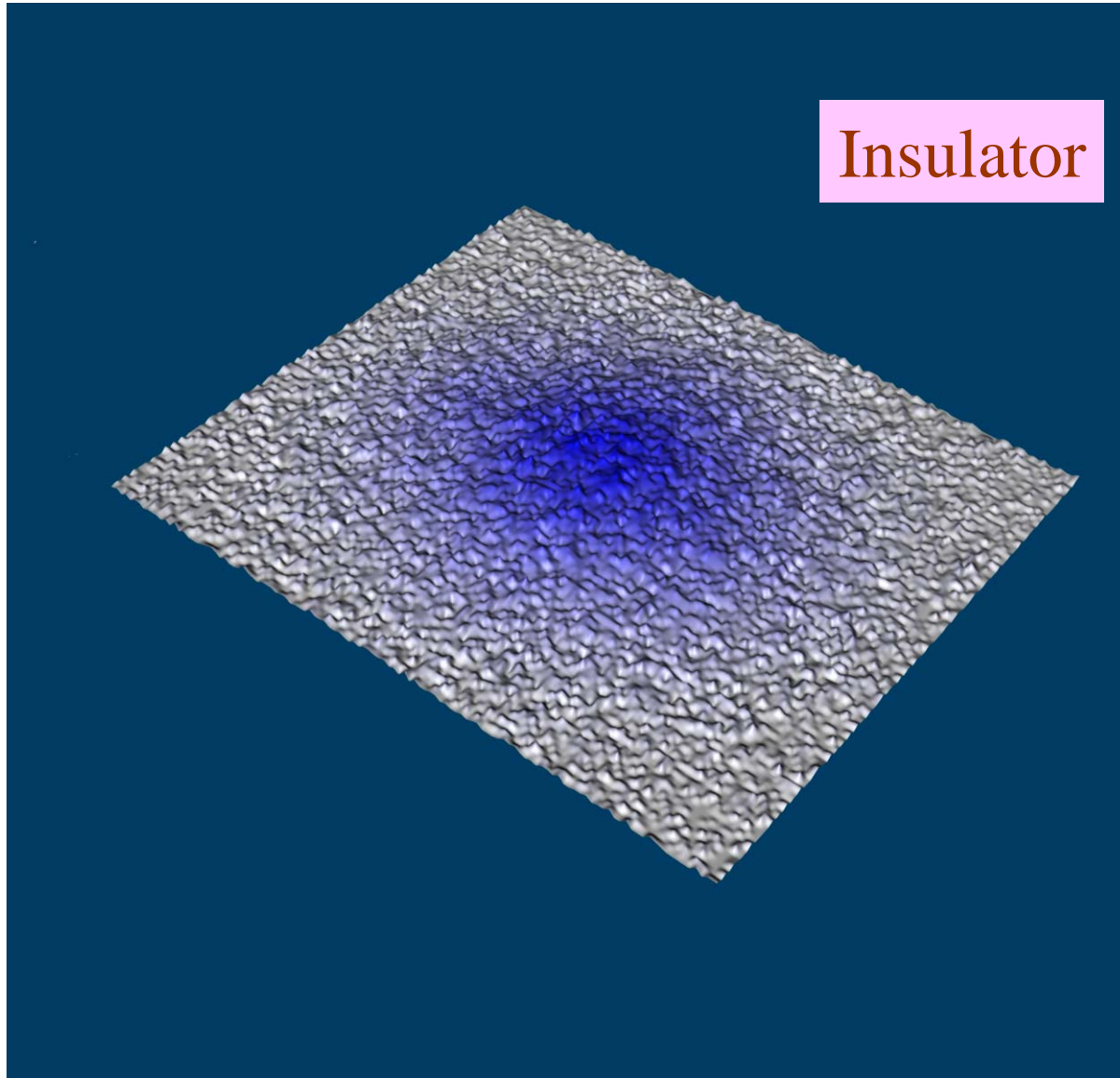
By tuning repulsive interactions between the atoms, states with multiple atoms in a potential well can be suppressed.

The lowest energy state is then a *Mott insulator* – it has negligible number fluctuations, and atoms cannot “flow”

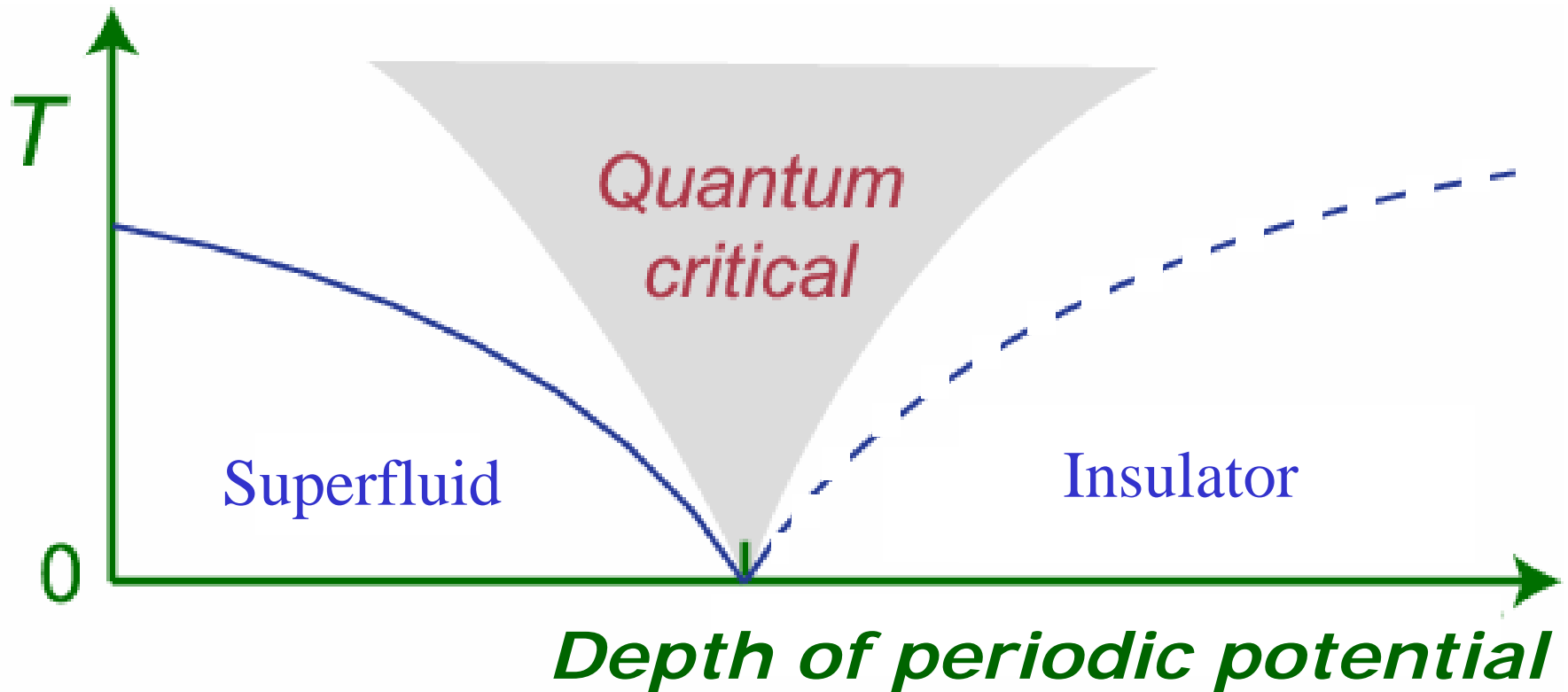
Velocity distribution of ^{87}Rb atoms



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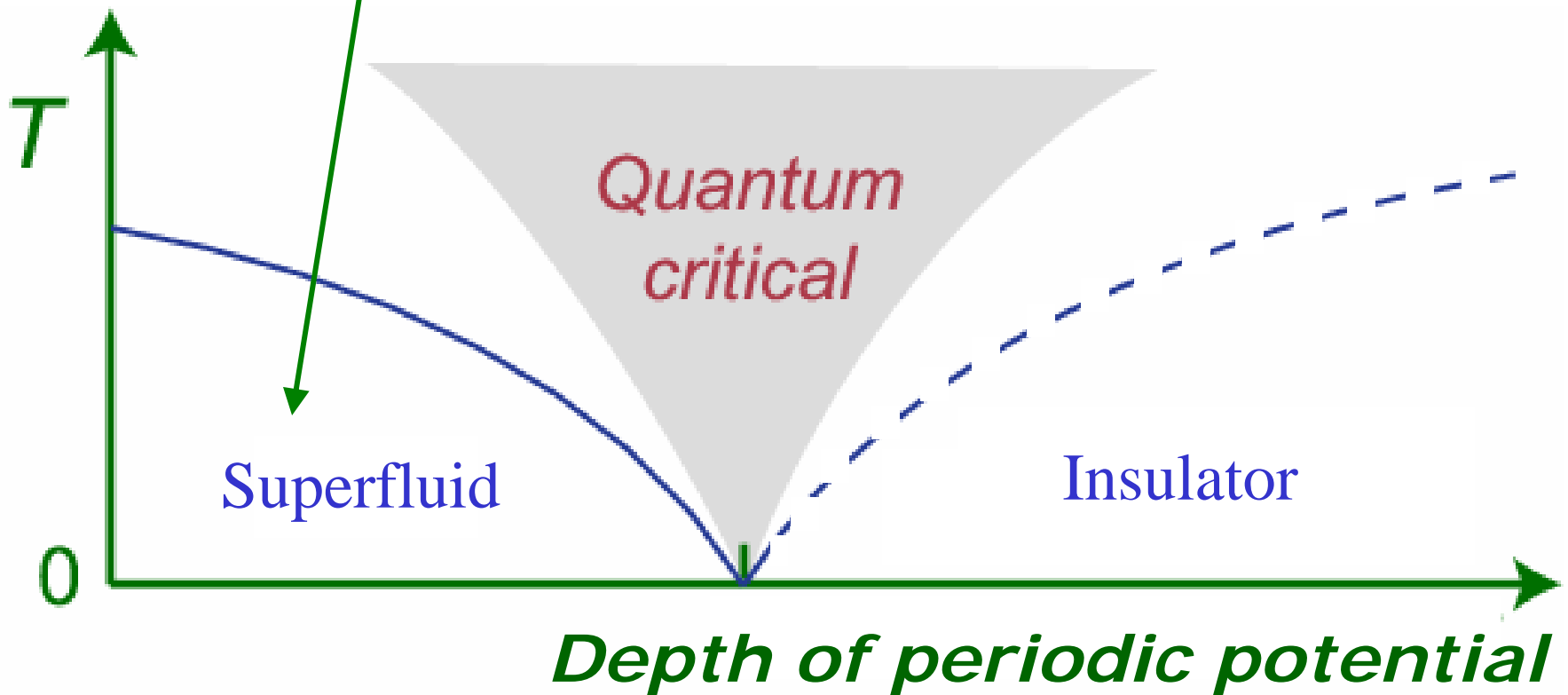


Non-zero temperature phase diagram



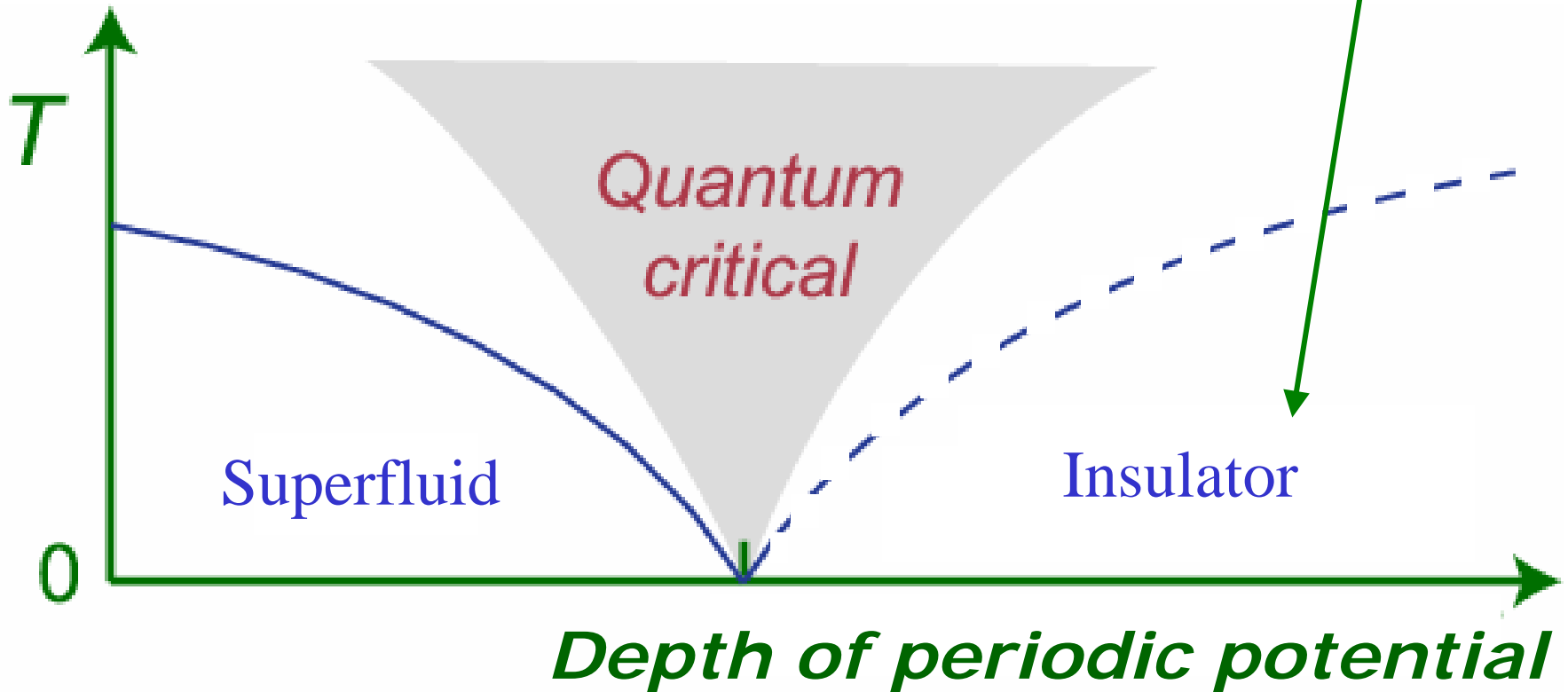
Non-zero temperature phase diagram

Dynamics of the classical
Gross-Pitaevski equation



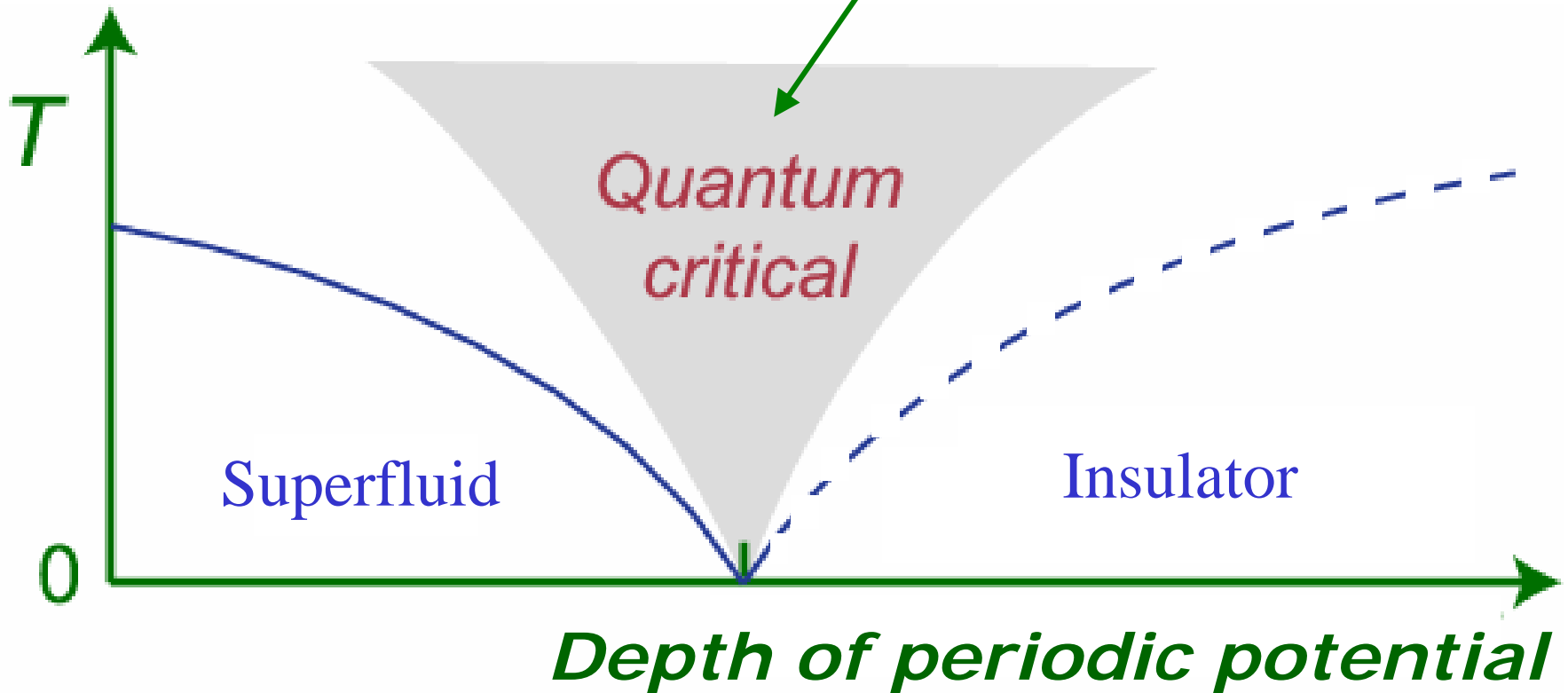
Non-zero temperature phase diagram

Dilute Boltzmann gas of
particle and holes



Non-zero temperature phase diagram

No wave or quasiparticle description



Resistivity of Bi films

Conductivity σ

$$\sigma_{\text{Superconductor}} (T \rightarrow 0) = \infty$$

$$\sigma_{\text{Insulator}} (T \rightarrow 0) = 0$$

$$\sigma_{\text{Quantum critical point}} (T \rightarrow 0) \approx \frac{4e^2}{h}$$

D. B. Haviland, Y. Liu, and A. M. Goldman,
Phys. Rev. Lett. **62**, 2180 (1989)

M. P. A. Fisher, *Phys. Rev. Lett.* **65**, 923 (1990)

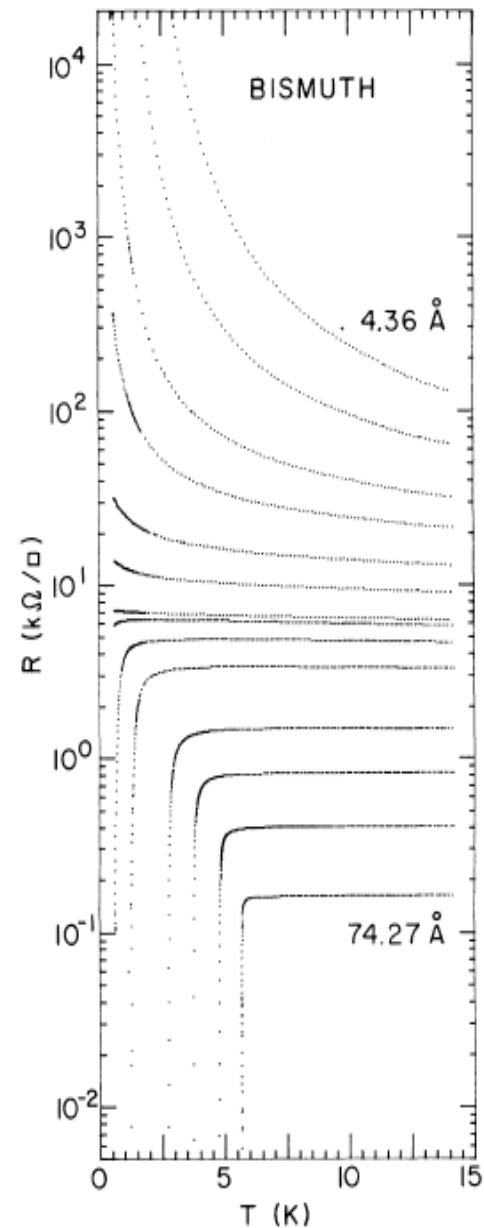
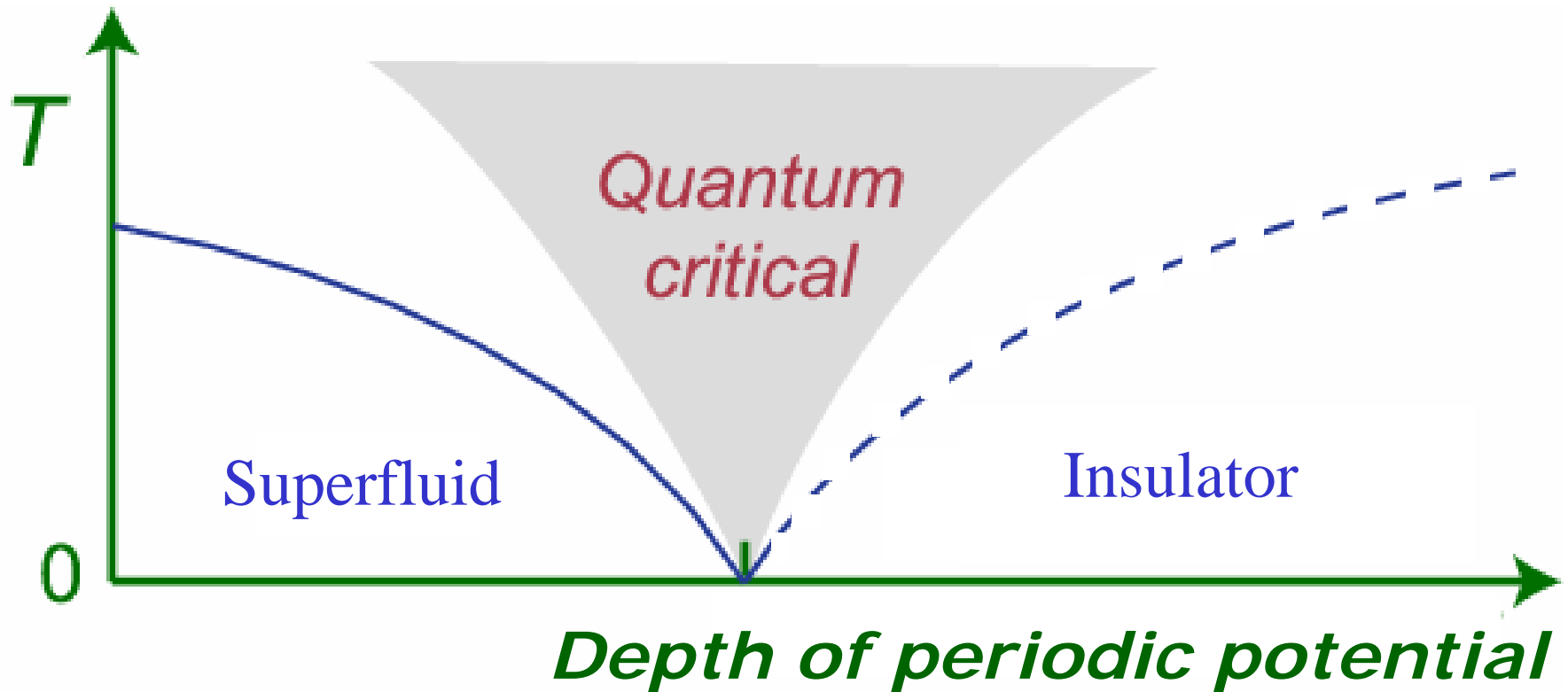
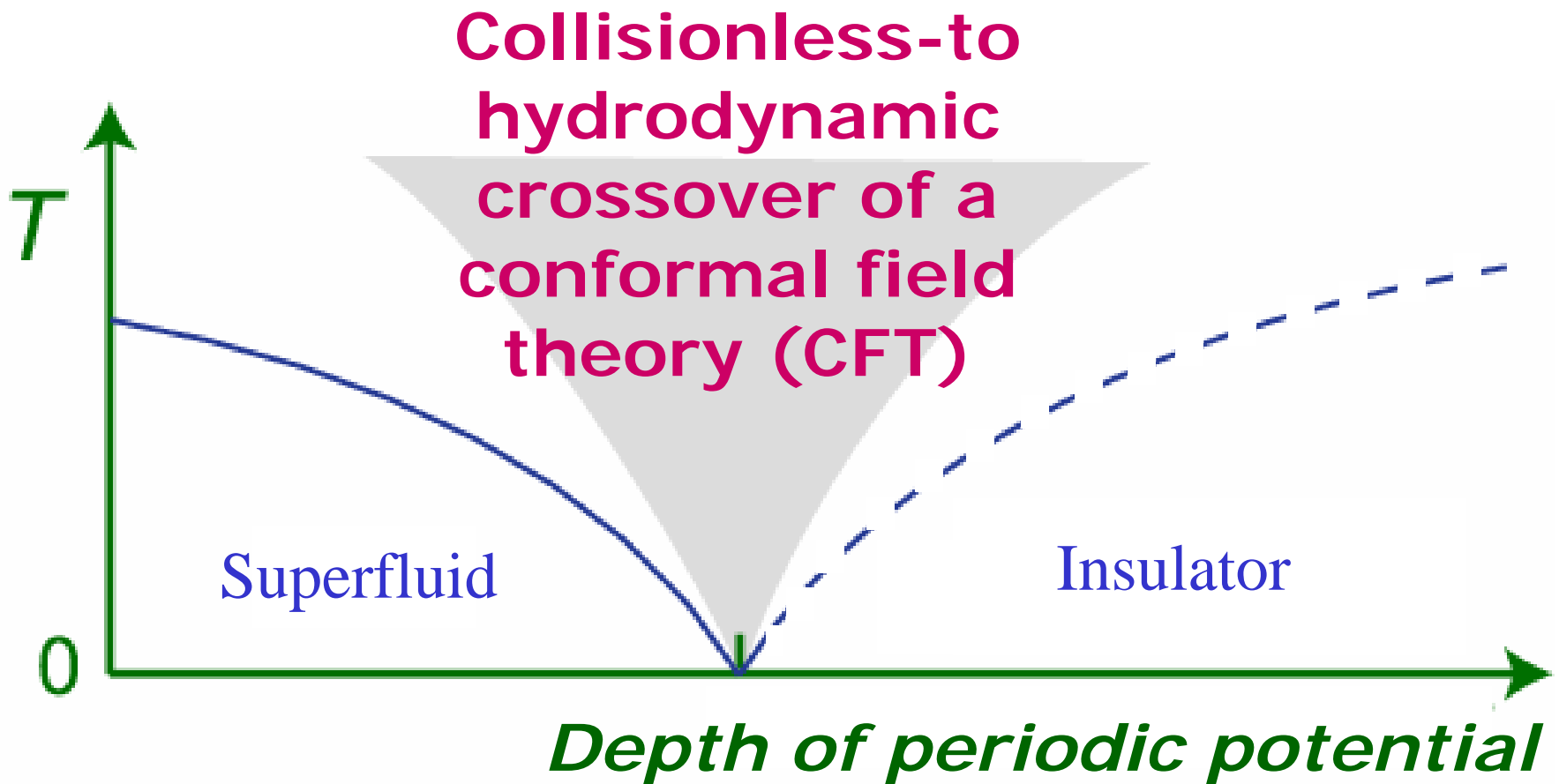


FIG. 1. Evolution of the temperature dependence of the sheet resistance $R(T)$ with thickness for a Bi film deposited onto Ge. Fewer than half of the traces actually acquired are shown. Film thicknesses shown range from 4.36 to 74.27 Å.

Non-zero temperature phase diagram

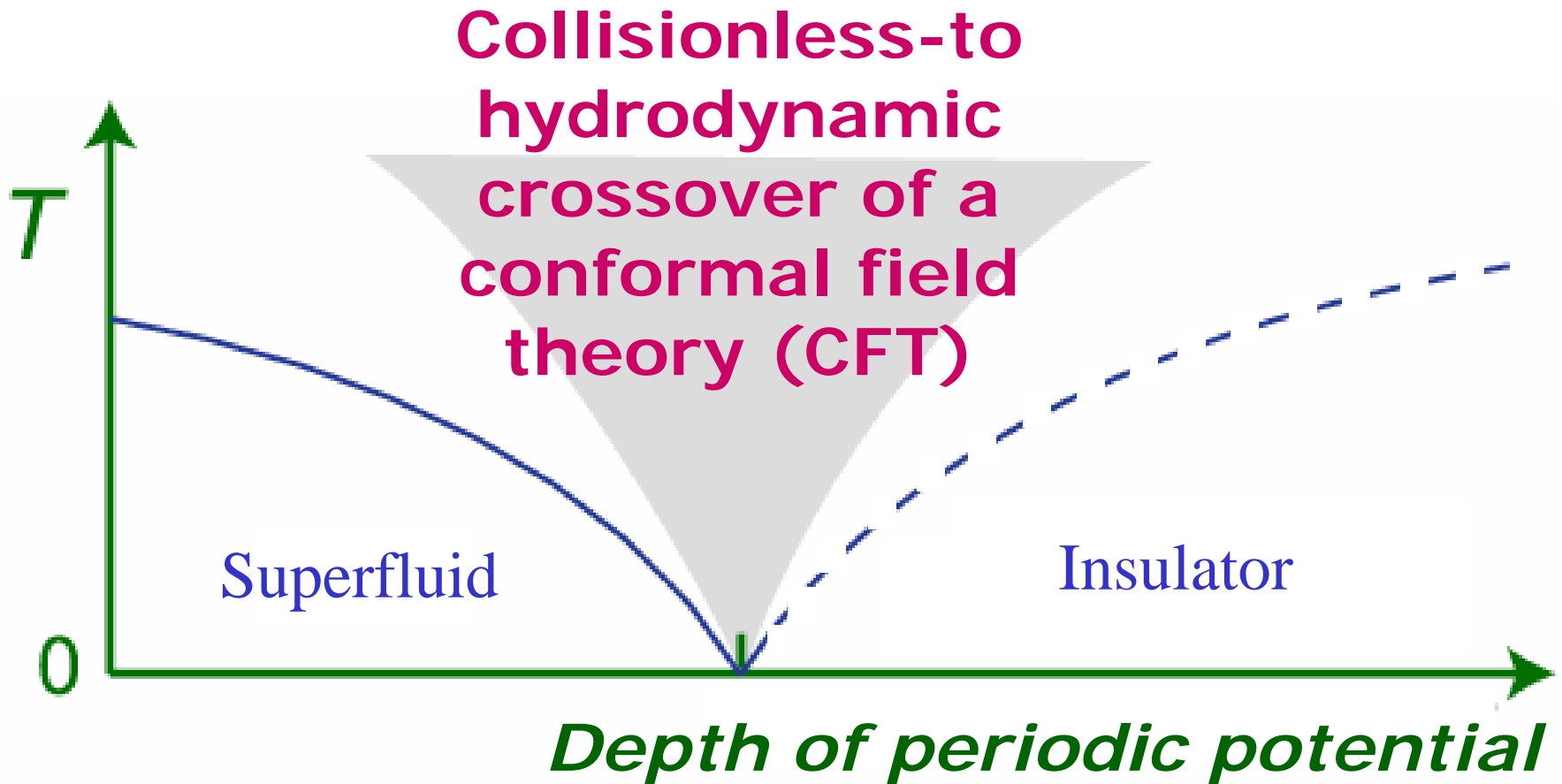


Non-zero temperature phase diagram



Non-zero temperature phase diagram

Needed: Cold atom experiments in this regime

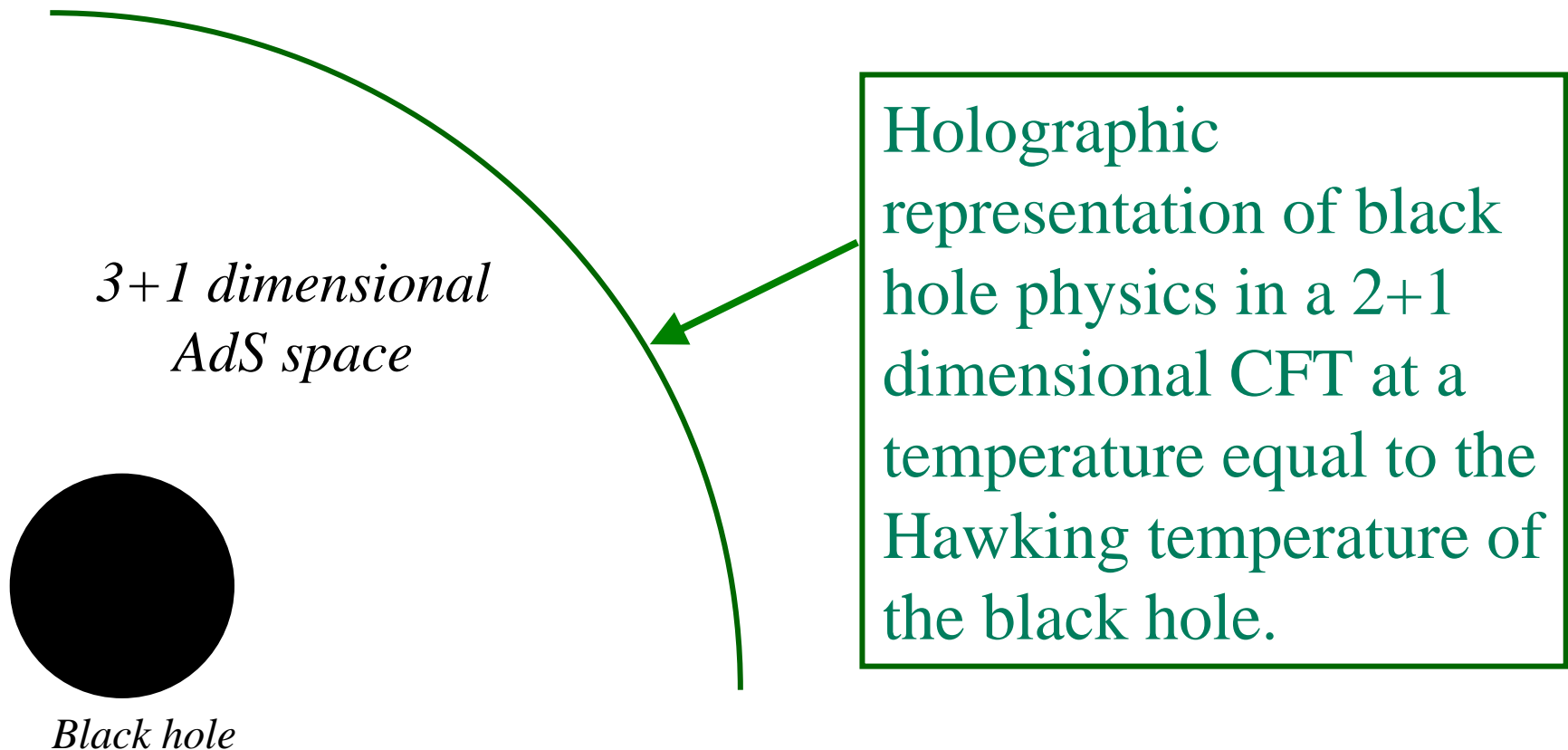


Hydrodynamics of a conformal field theory (CFT)

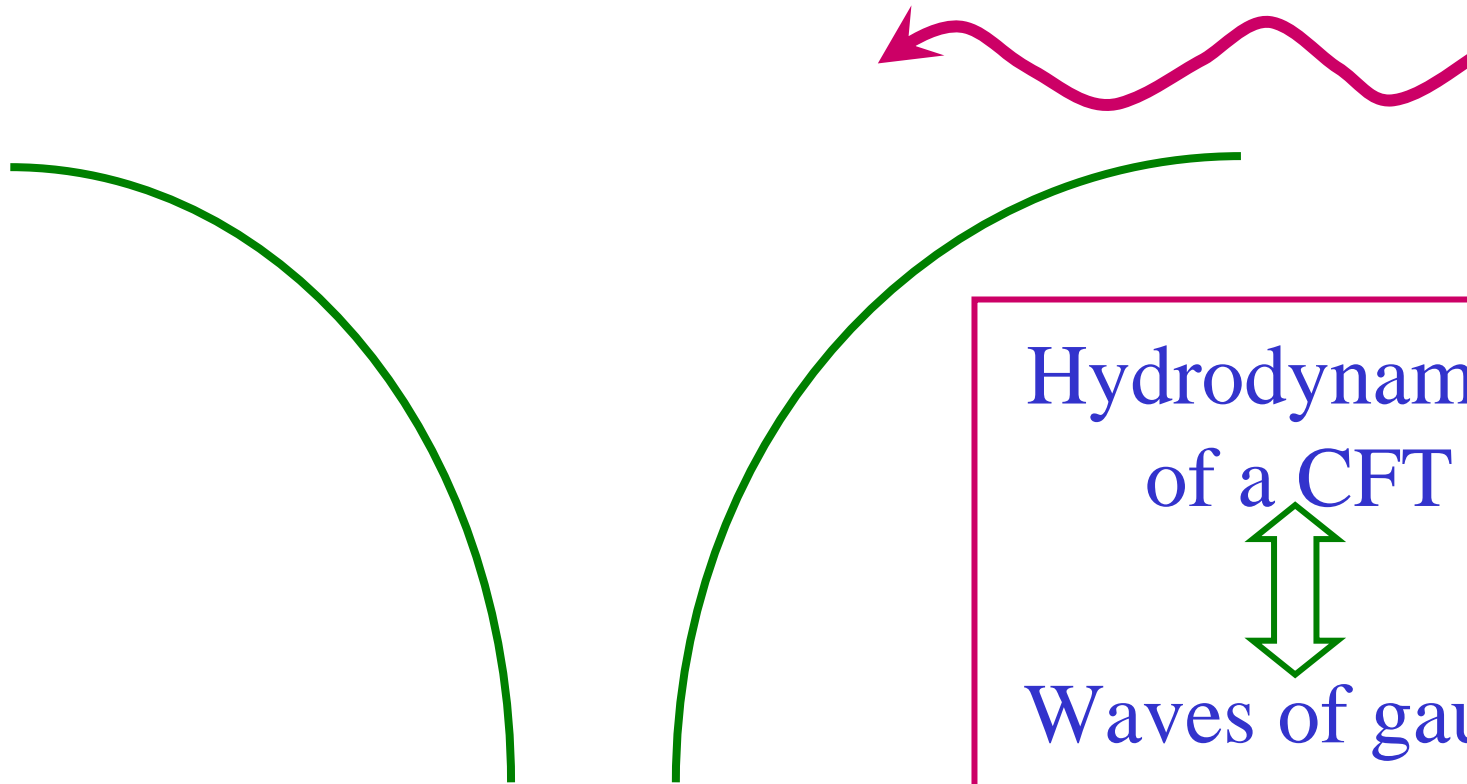
Maldacena's AdS/CFT correspondence relates the hydrodynamics of CFTs to the quantum gravity theory of the horizon of a black hole in Anti-de Sitter space.

Hydrodynamics of a conformal field theory (CFT)

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Hydrodynamics of a conformal field theory (CFT)



Hydrodynamics
of a CFT



Waves of gauge
fields in a curved
background

Hydrodynamics of a conformal field theory (CFT)

The scattering cross-section of the thermal excitations is universal and so transport coefficients are universally determined by $k_B T$

Charge diffusion constant

$$D_c = \Theta \frac{c^2}{k_B T}$$

Conductivity

$$\sigma = \Theta \frac{4e^2}{h}$$

Hydrodynamics of a conformal field theory (CFT)

For the (unique) CFT with a $SU(N)$ gauge field and 16 supercharges, we know the exact diffusion constant associated with a global $SO(8)$ symmetry:

Spin diffusion constant

$$D_s = \frac{3}{4\pi} \frac{c^2}{k_B T}$$

Spin conductivity

$$\sigma = \frac{N^{3/2}}{3\sqrt{2}\pi}$$

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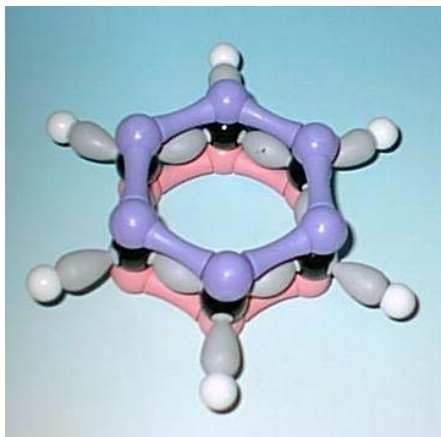
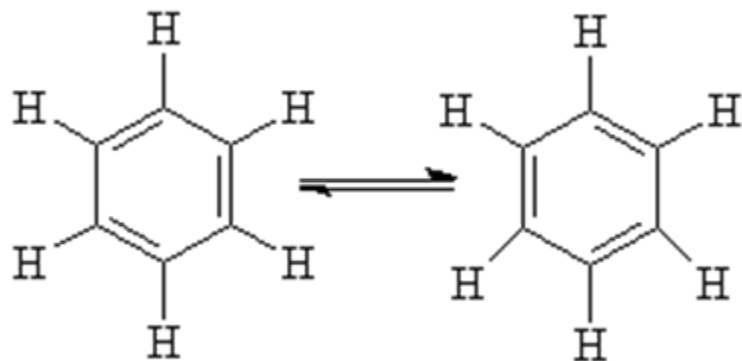
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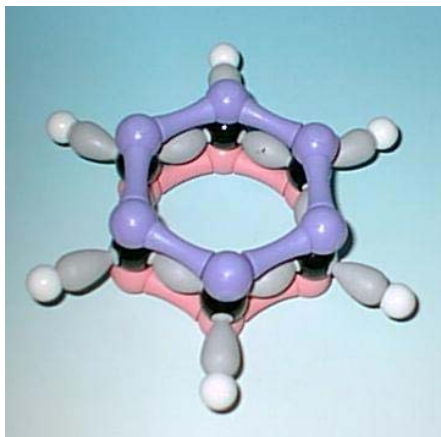
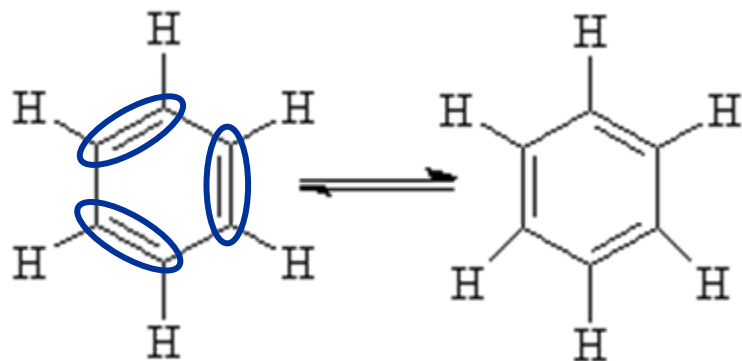
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Valence bonds in benzene



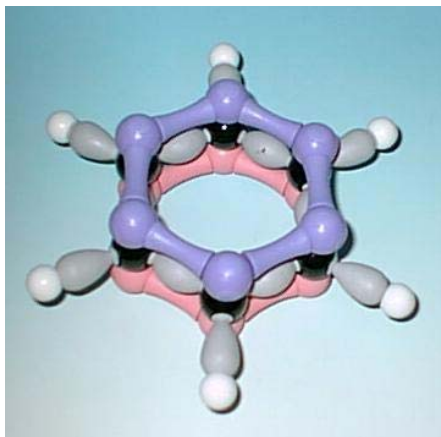
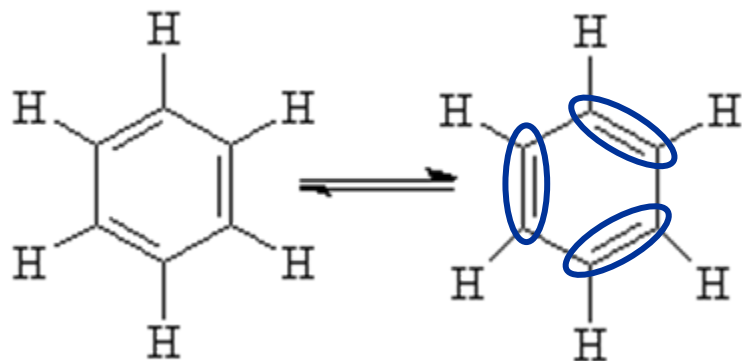
Resonance in benzene leads to a symmetric configuration of valence bonds
(*F. Kekulé, L. Pauling*)

Valence bonds in benzene



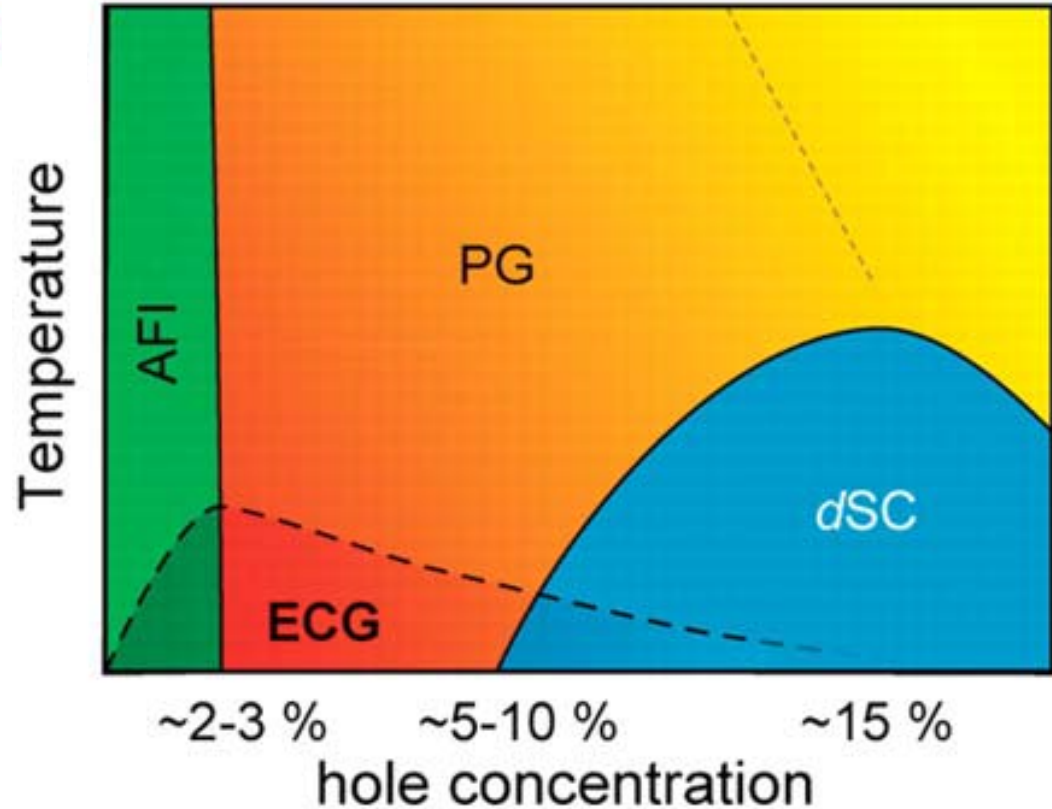
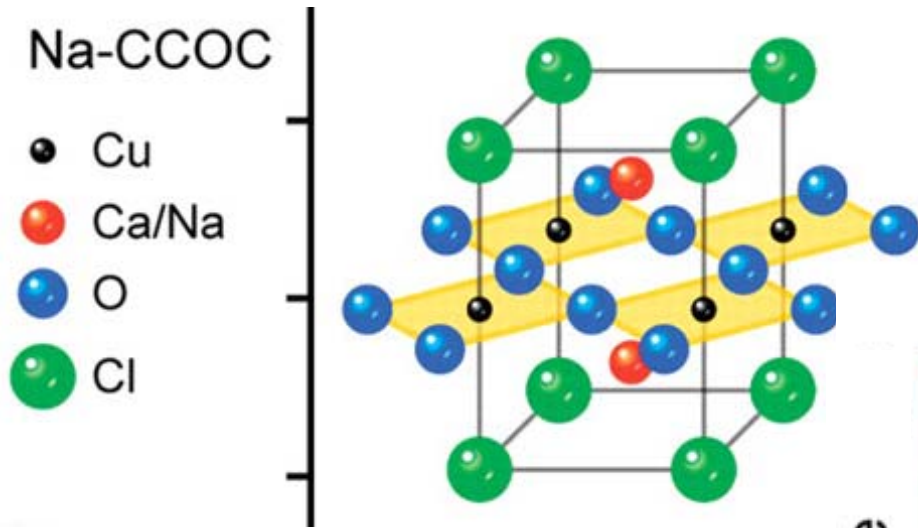
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Valence bonds in benzene

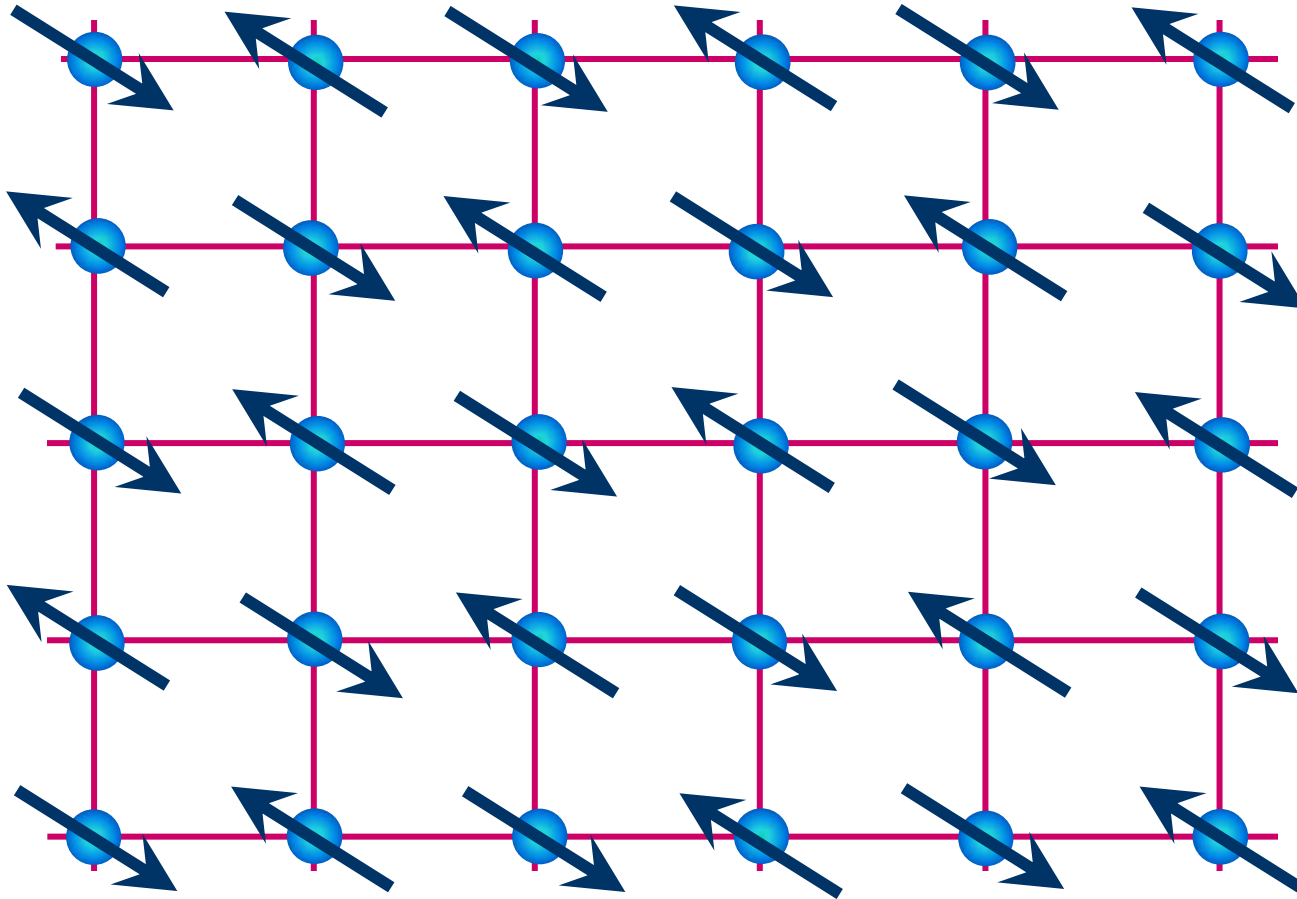


Resonance in benzene leads to a symmetric configuration of valence bonds
(*F. Kekulé, L. Pauling*)

Temperature-doping phase diagram of the cuprate superconductors

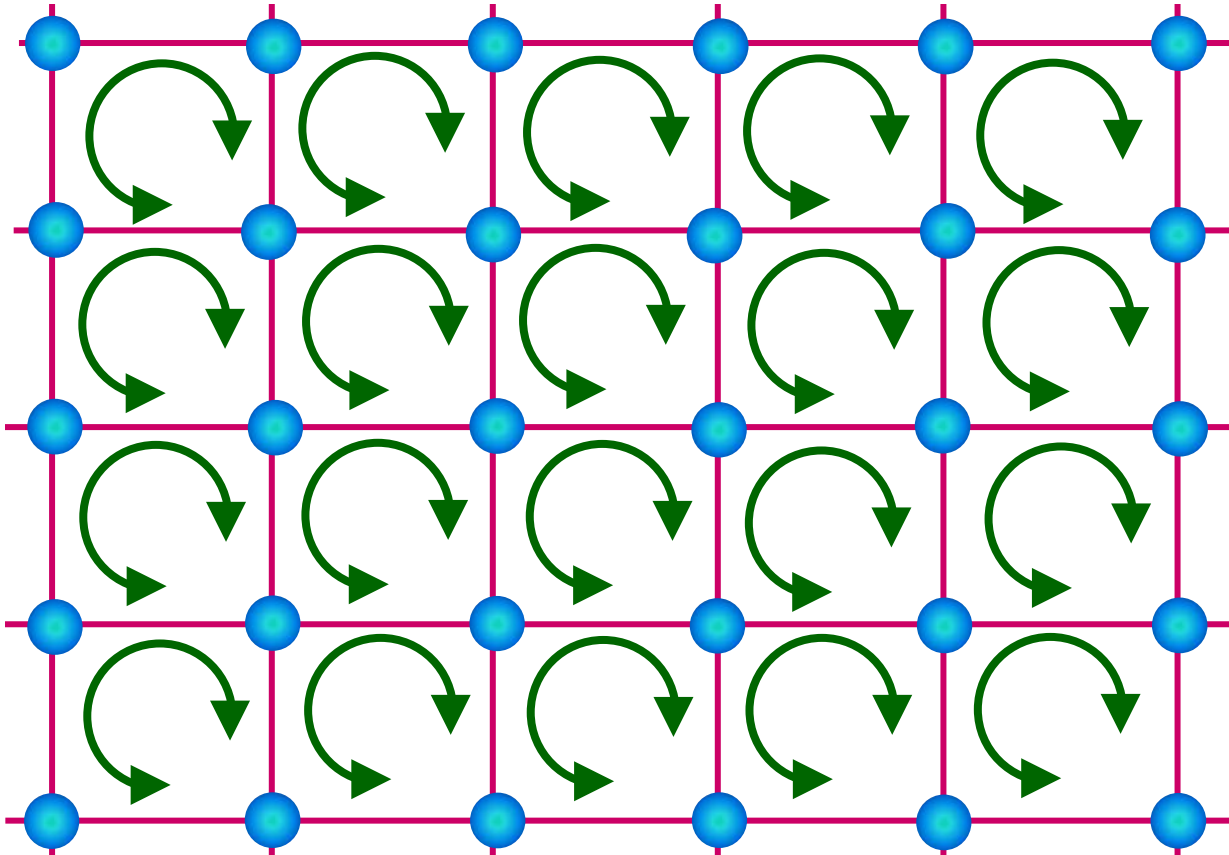


Antiferromagnetic (Neel) order in the insulator



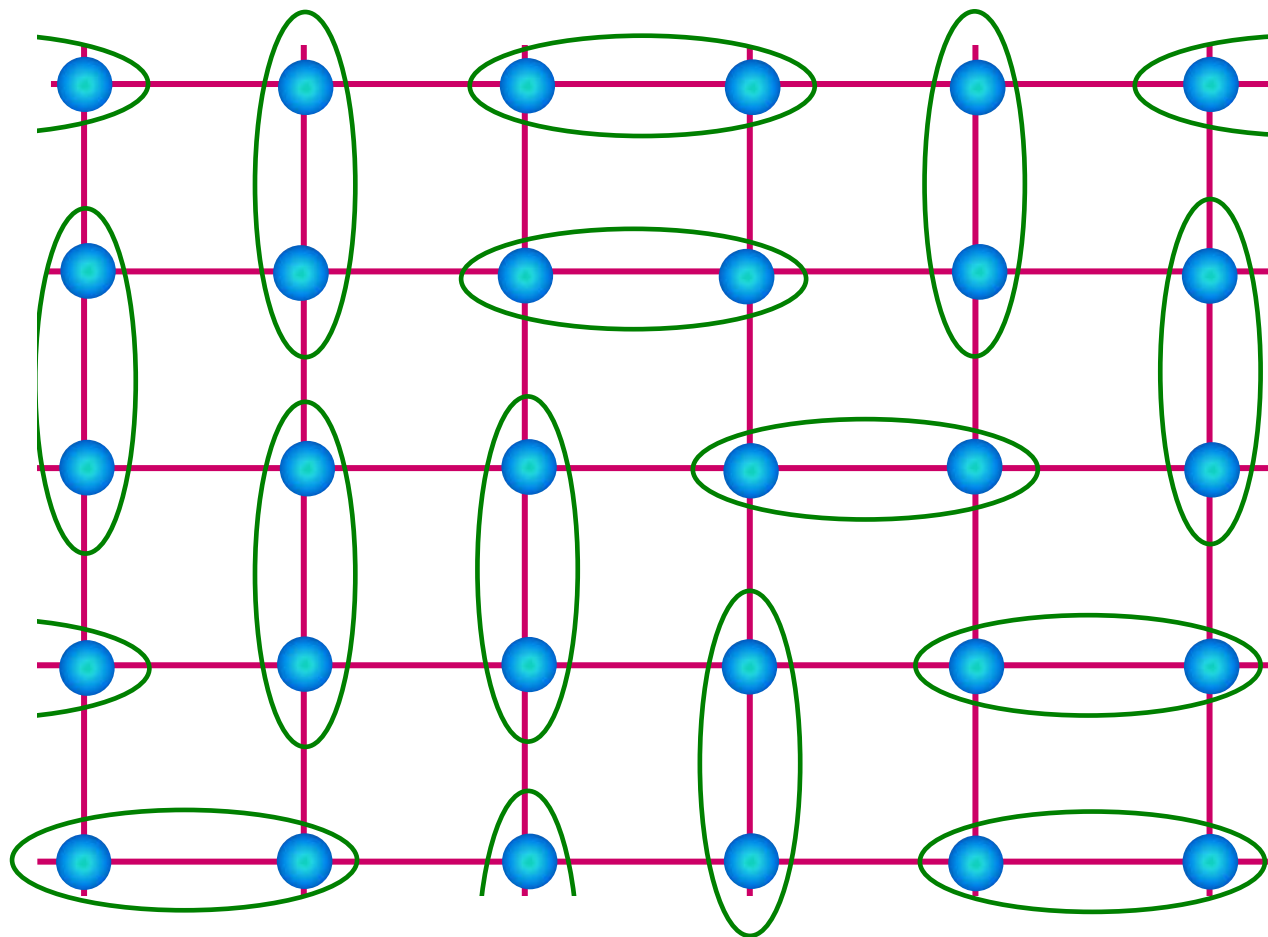
$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j \quad ; \quad \vec{S}_i \Rightarrow \text{spin operator with } S=1/2$$

Induce formation of valence bonds by
e.g. ring-exchange interactions

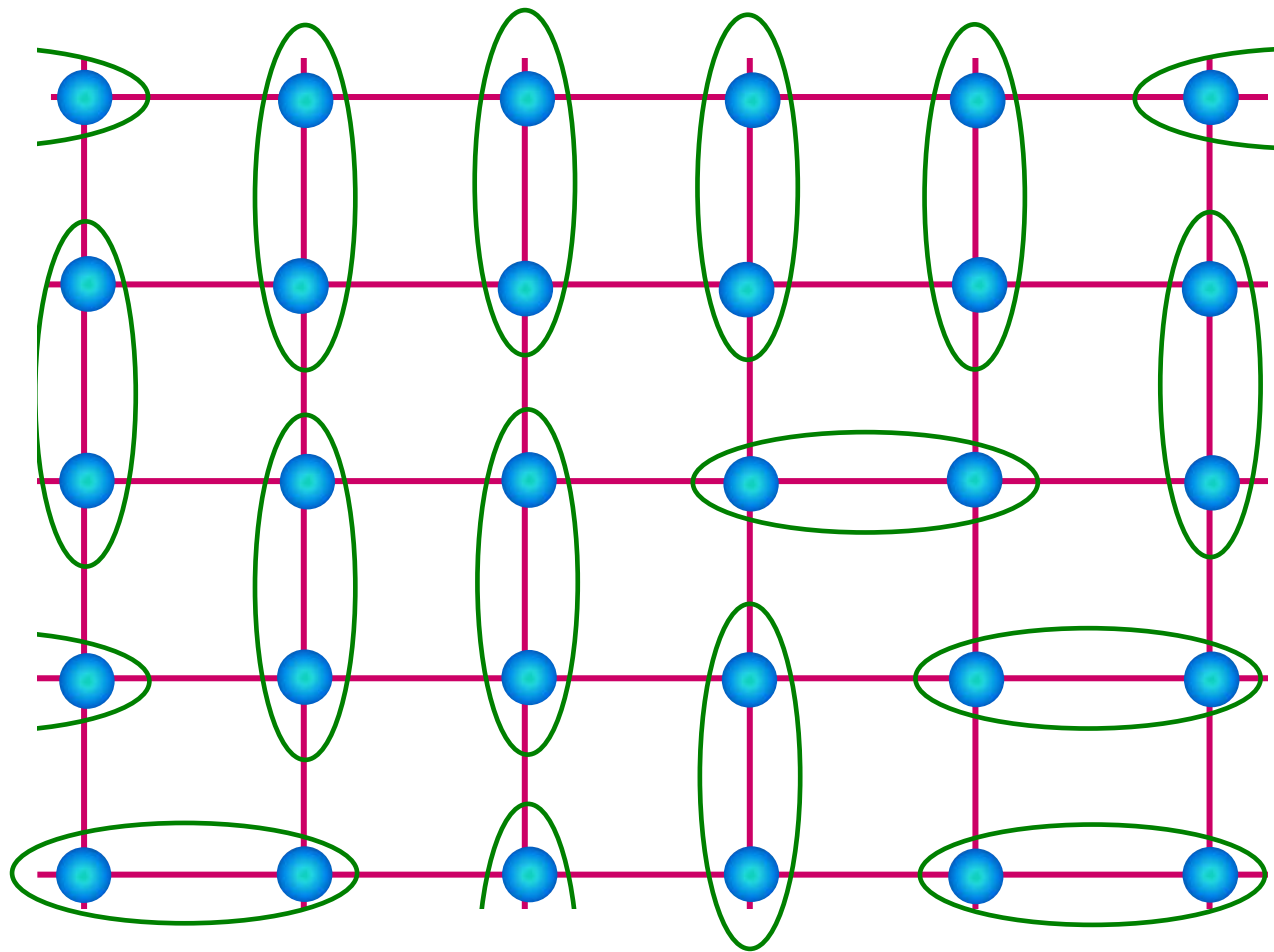


$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} 4\text{-spin exchange}$$

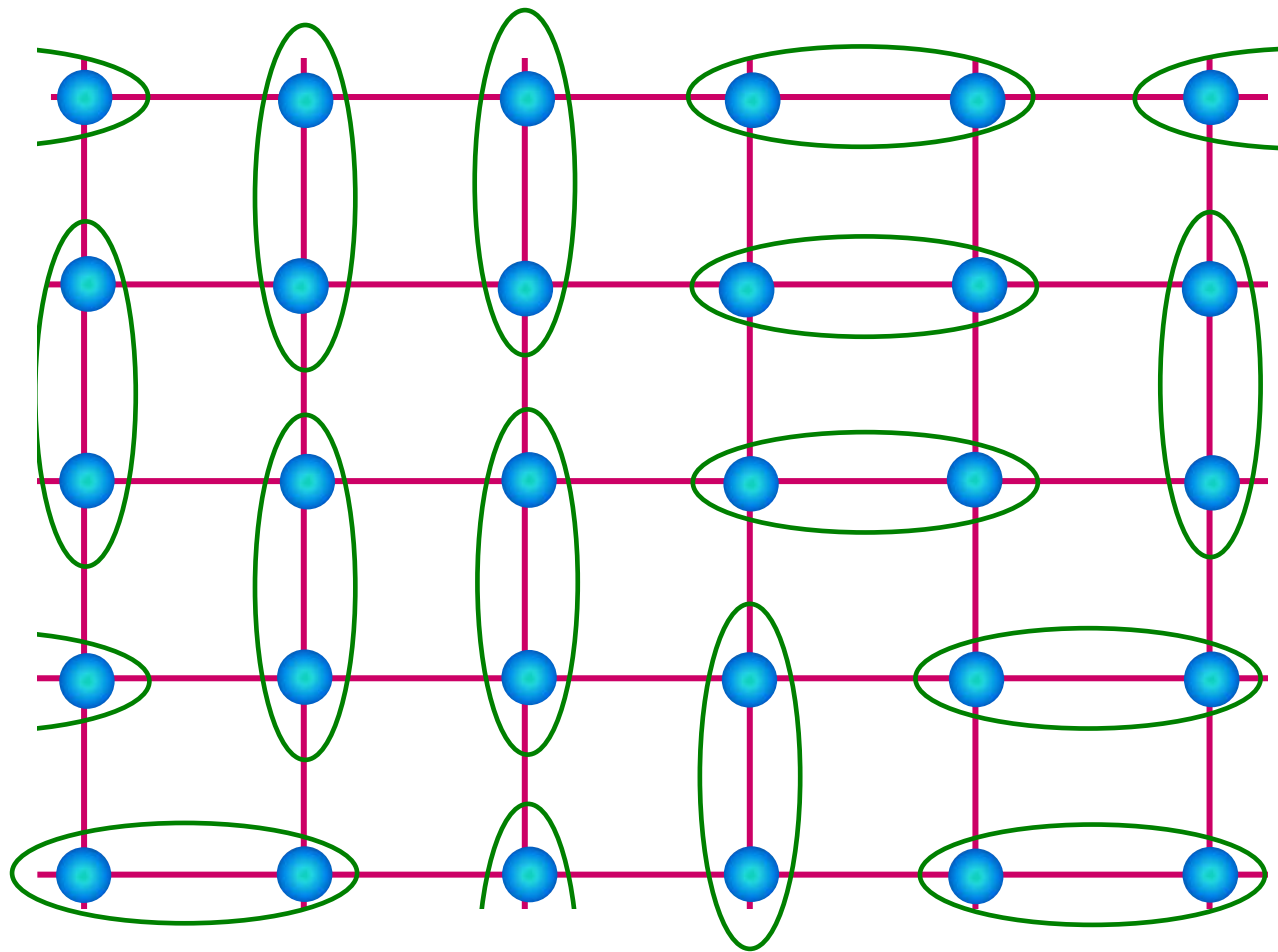
As in H_2 and benzene, each electron wants to pair up with another electron and form a valence bond



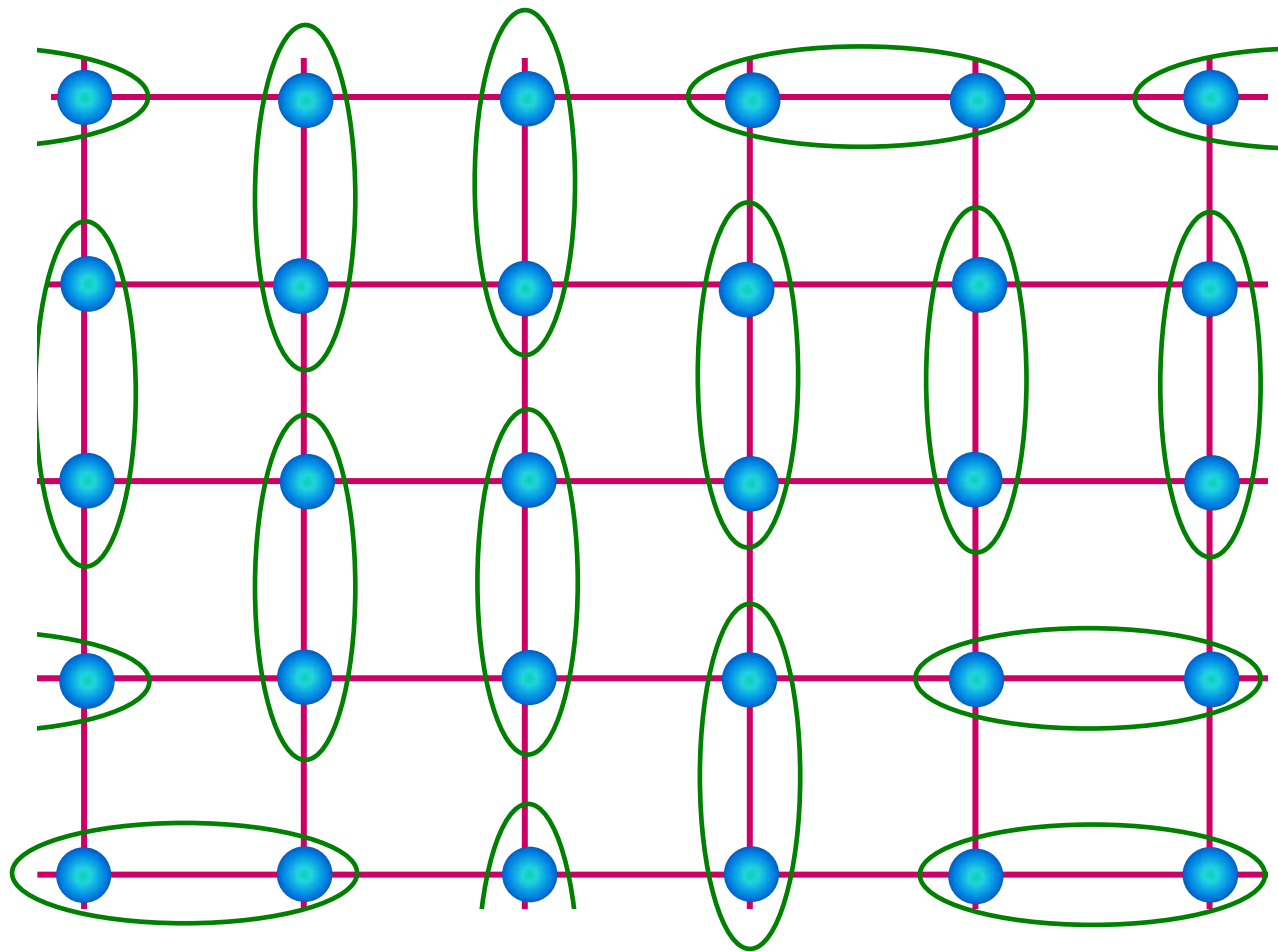
$$\begin{array}{c} \text{Oval with two blue spheres} \\ = \\ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \end{array}$$



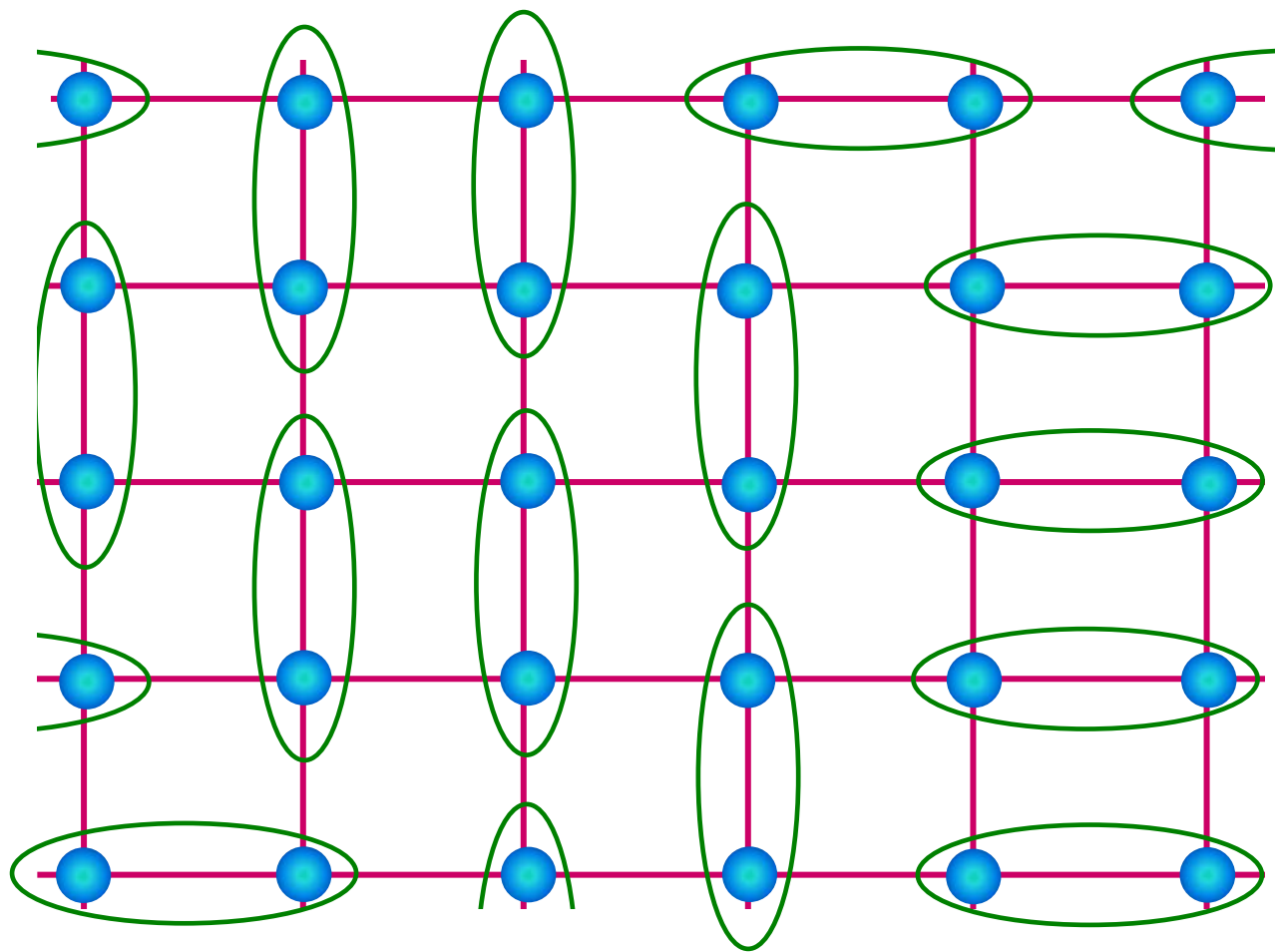
$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$



$$\begin{array}{c}
 \text{Diagram of two blue spheres in a green oval} \\
 = \\
 \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)
 \end{array}$$

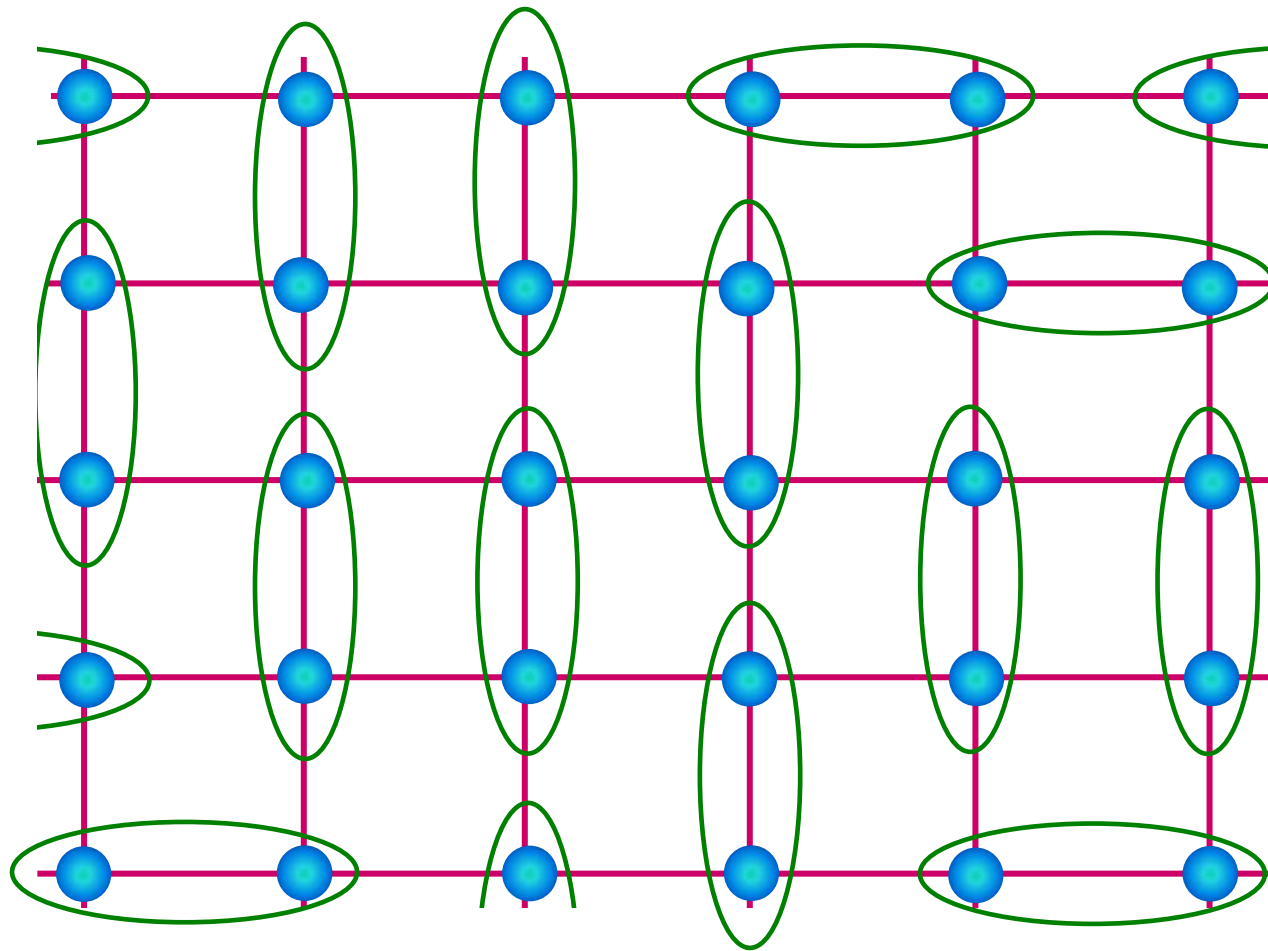


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 \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)
 \end{array}$$



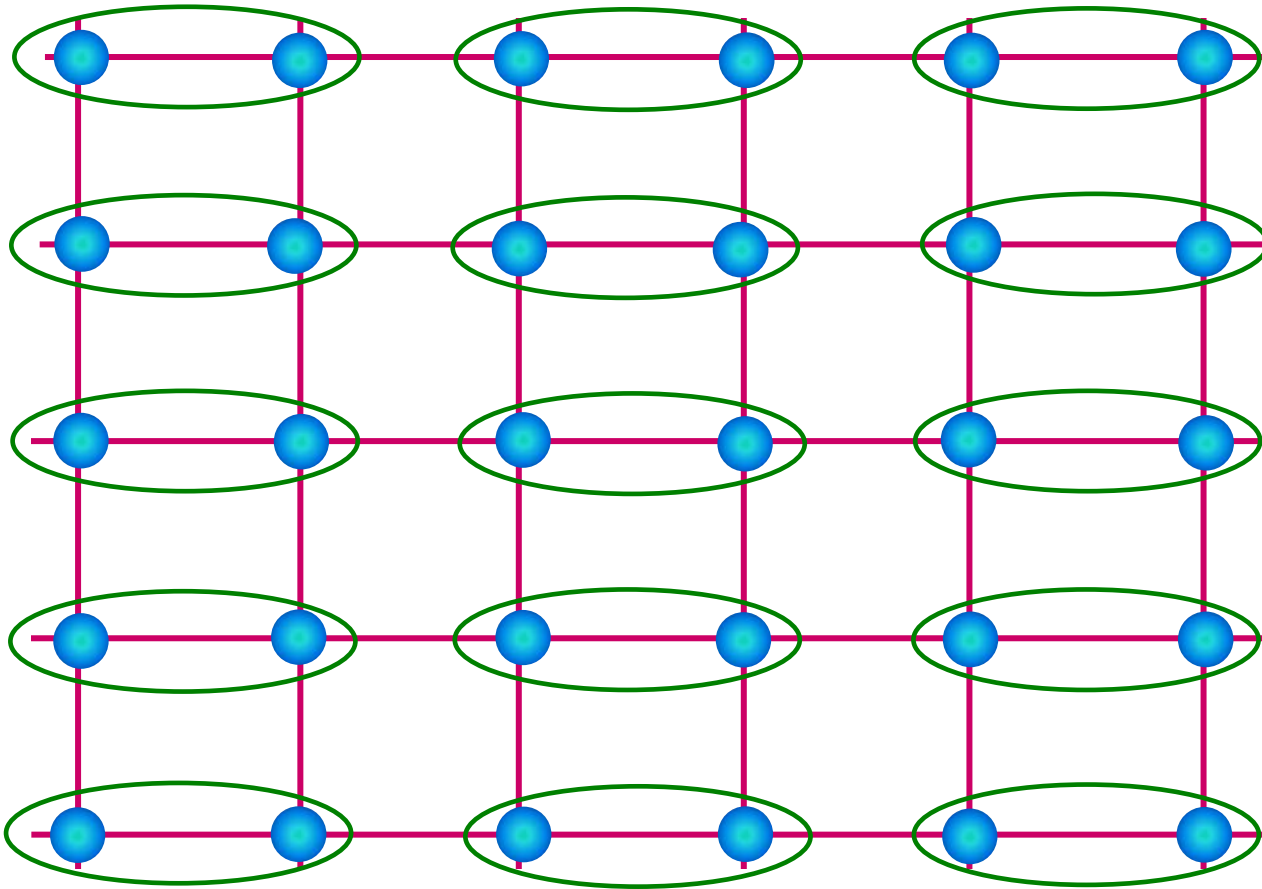
$$\begin{array}{c}
 \text{Oval with two blue spheres} \\
 = \\
 \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)
 \end{array}$$

Entangled liquid of valence bonds (Resonating valence bonds – RVB)



$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Valence bond solid (VBS)

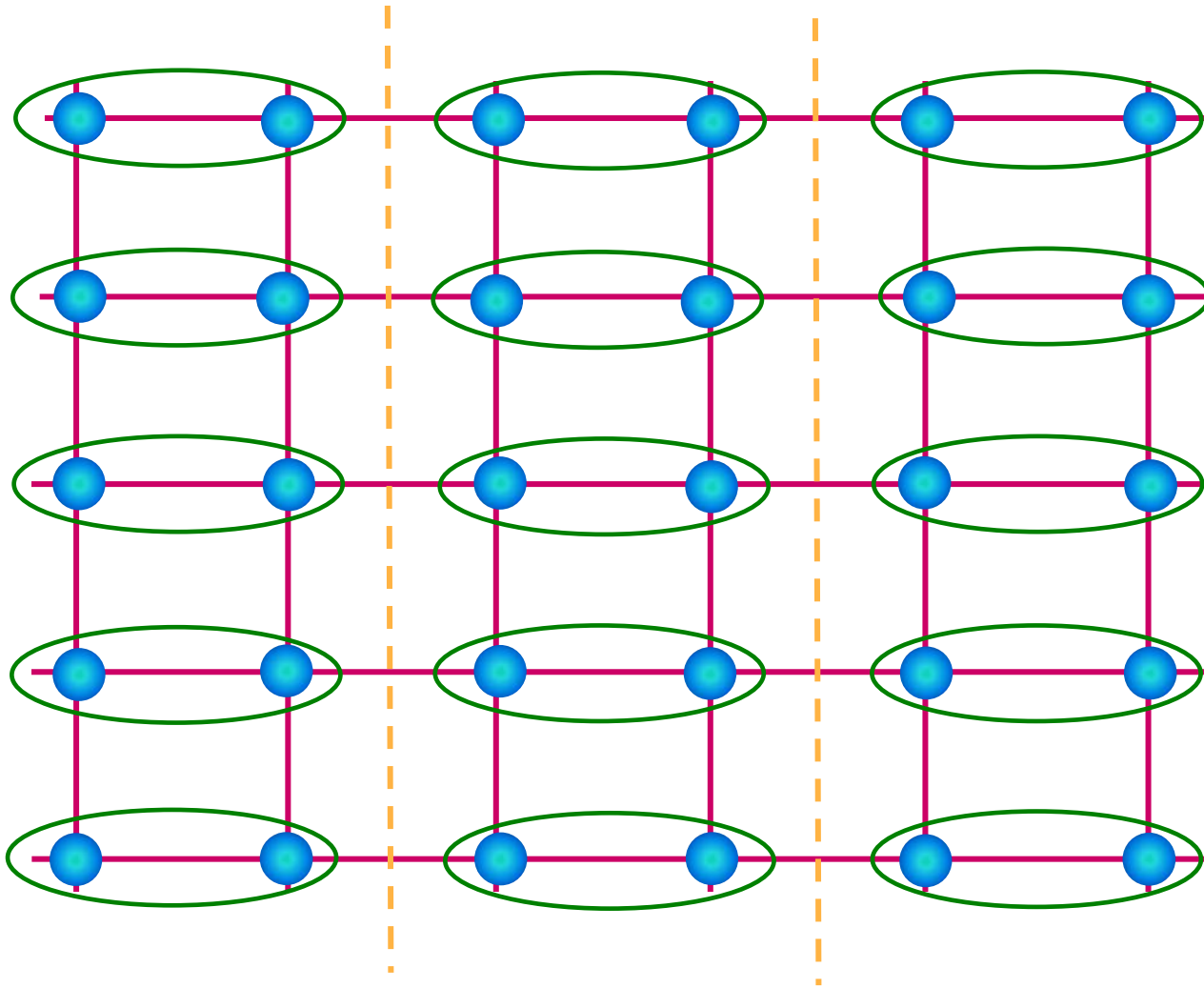


$$\text{green oval with two blue spheres} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

N. Read and S. Sachdev,
Phys. Rev. Lett. **62**, 1694
(1989).

R. Moessner and S. L.
Sondhi, *Phys. Rev. B* **63**,
224401 (2001).

Valence bond solid (VBS)



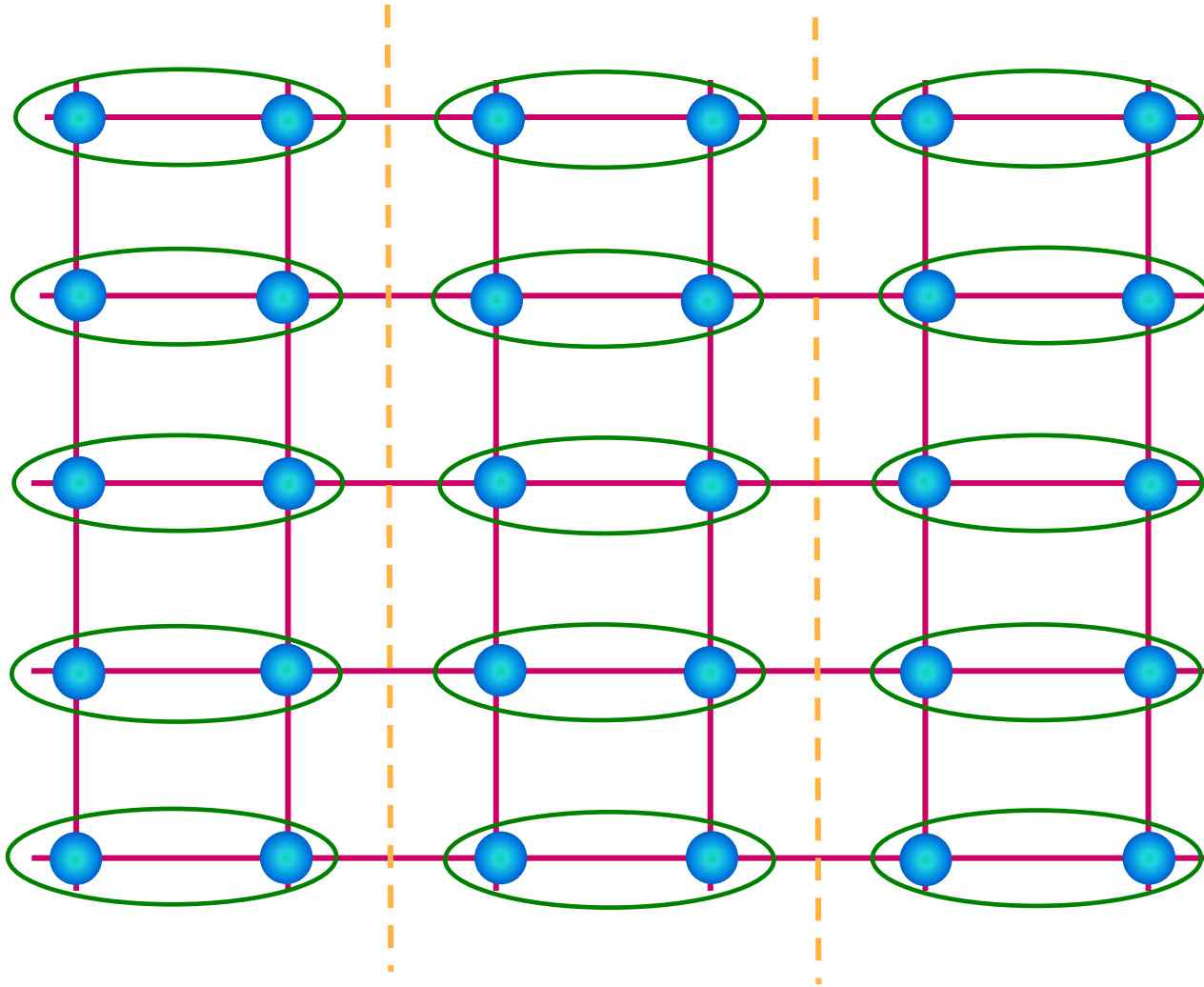
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Valence bond solid (VBS)

More possibilities for entanglement with nearby states



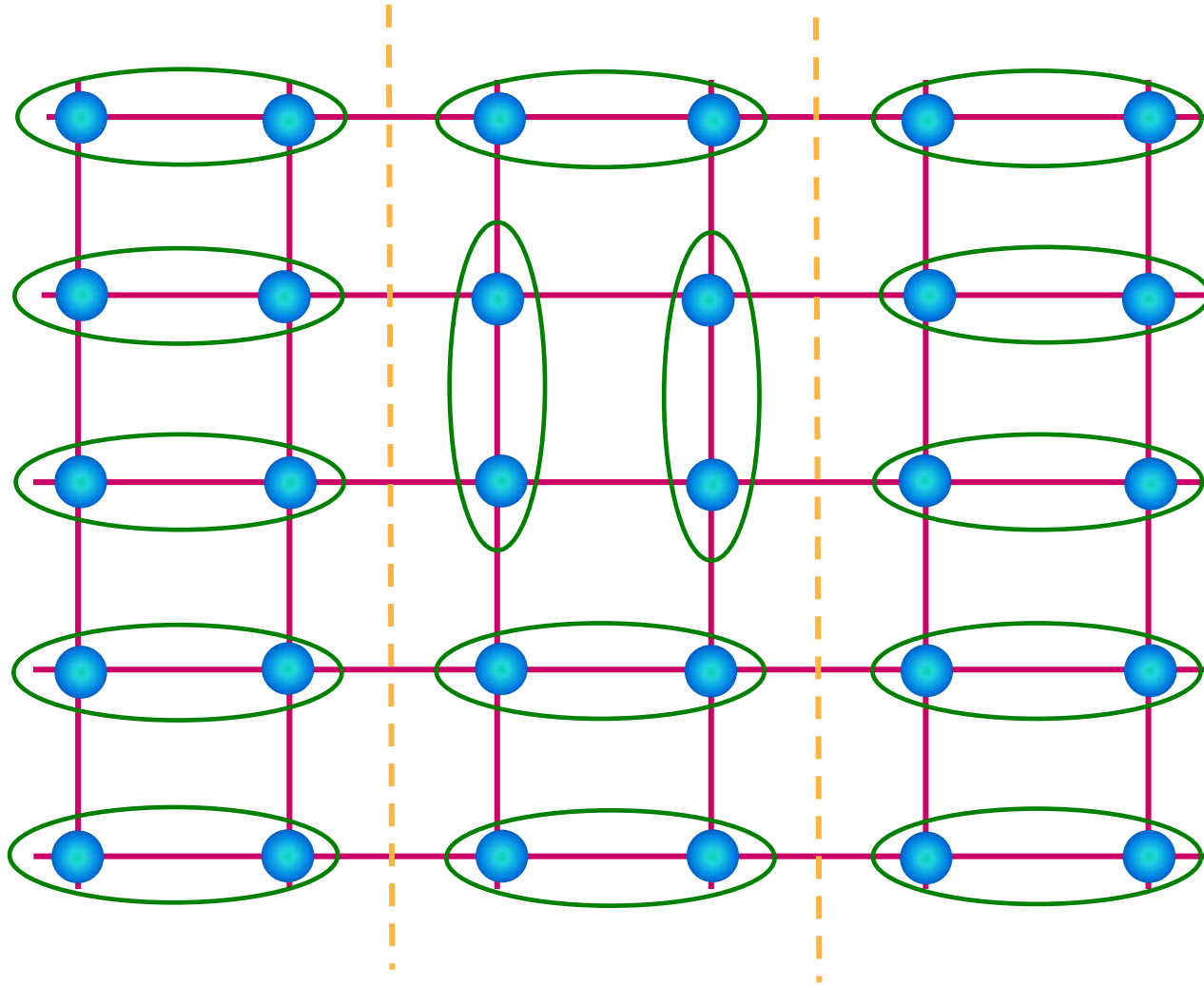
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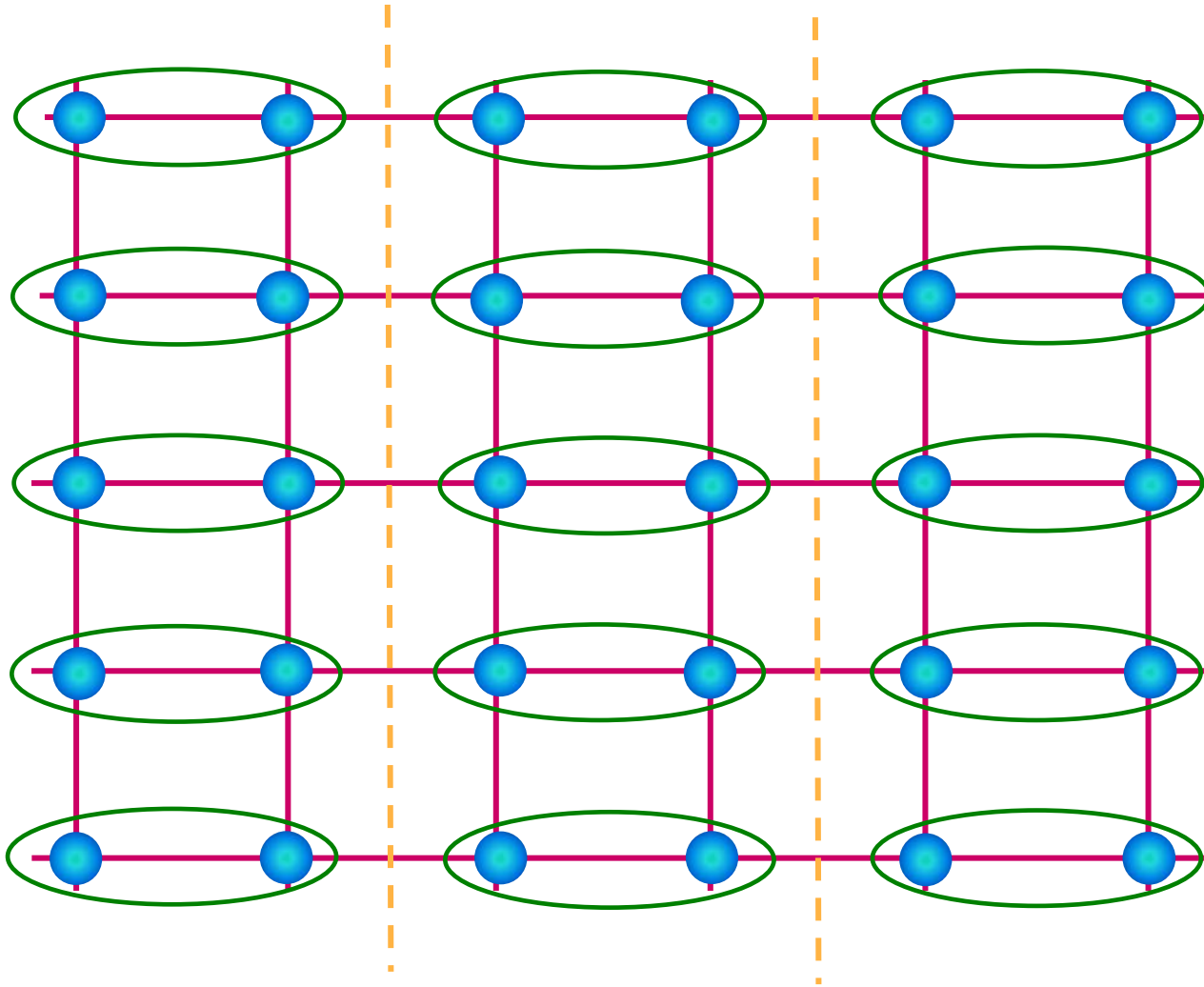
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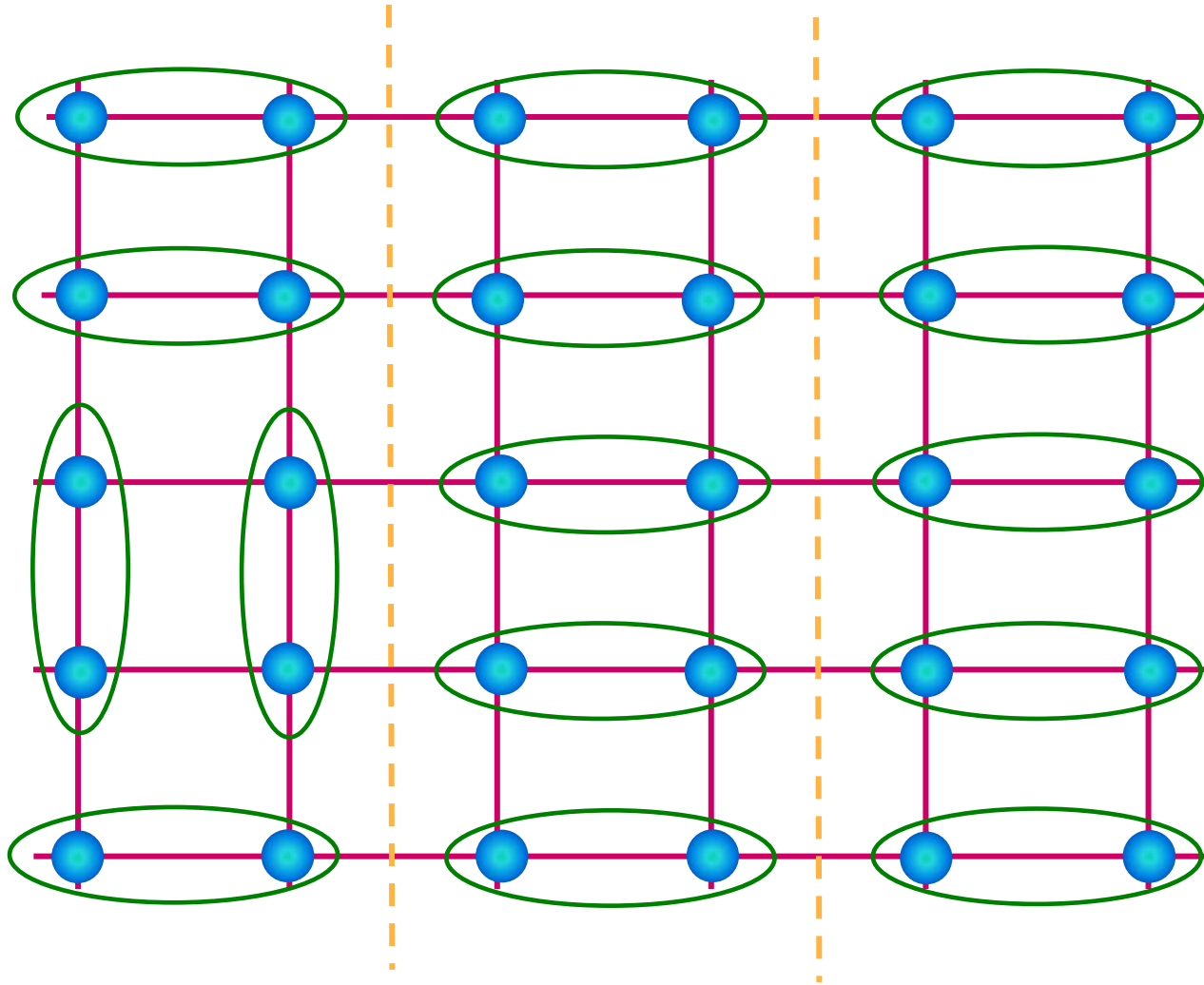
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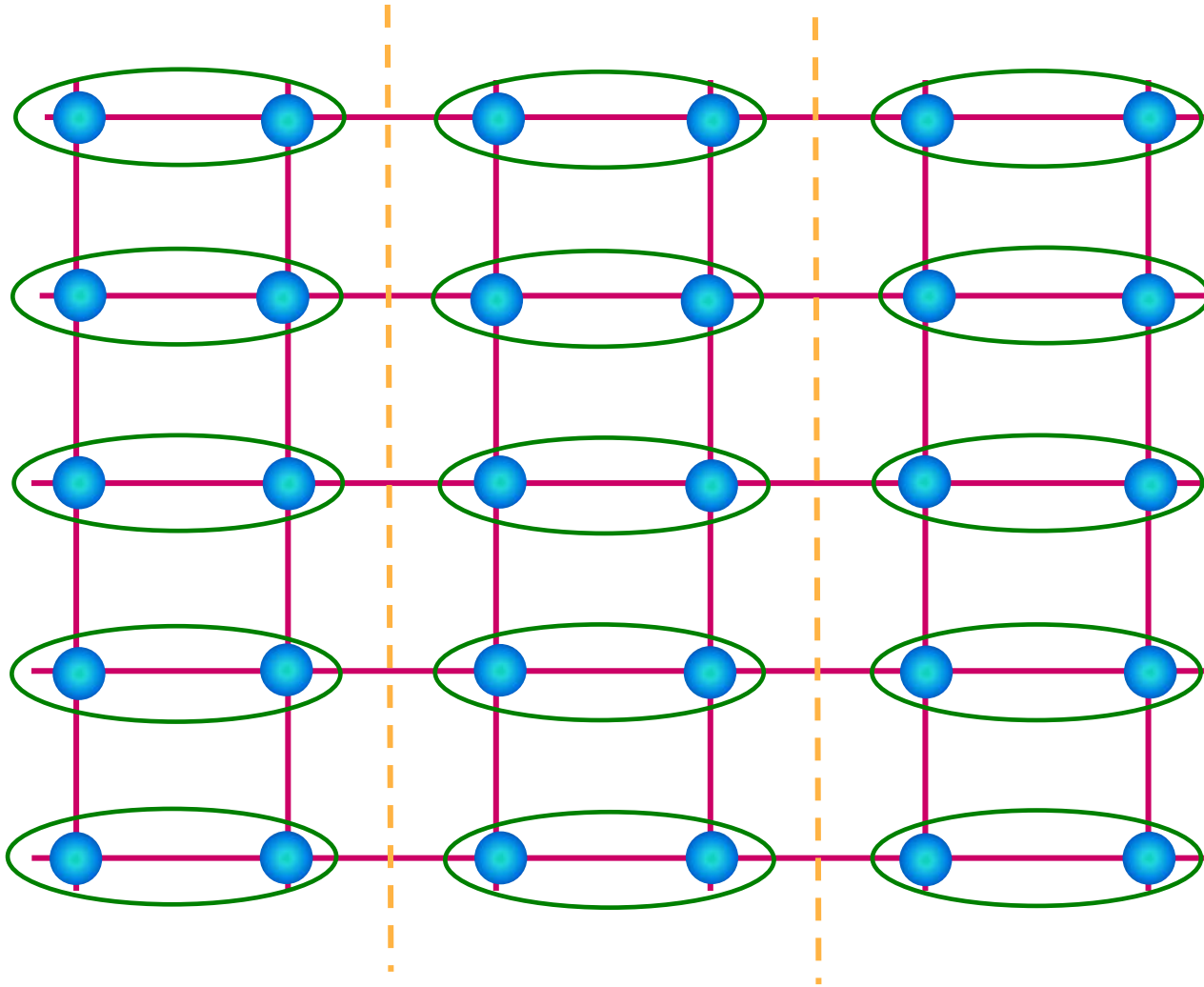
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More possibilities for entanglement with nearby states



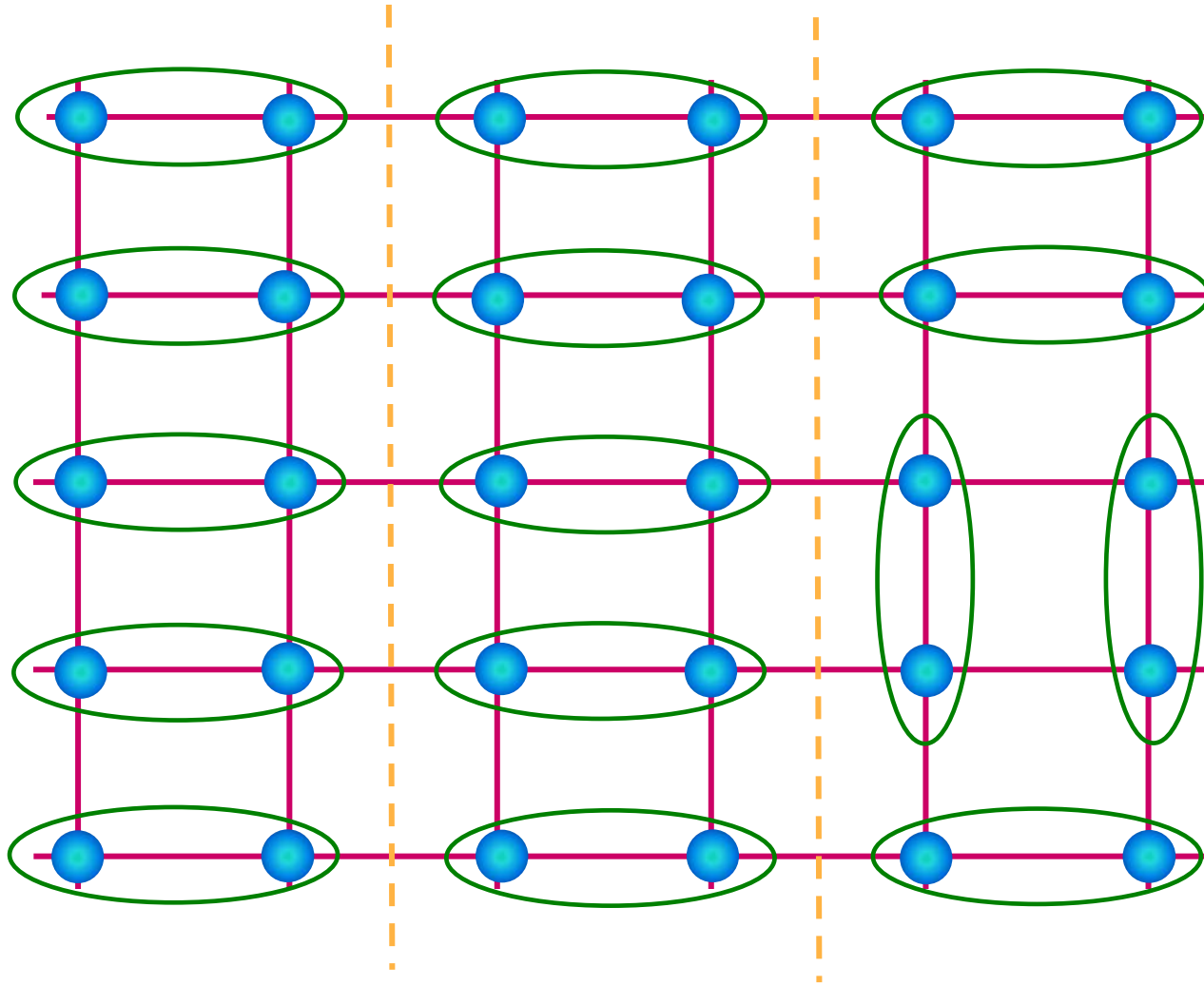
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More possibilities for entanglement with nearby states



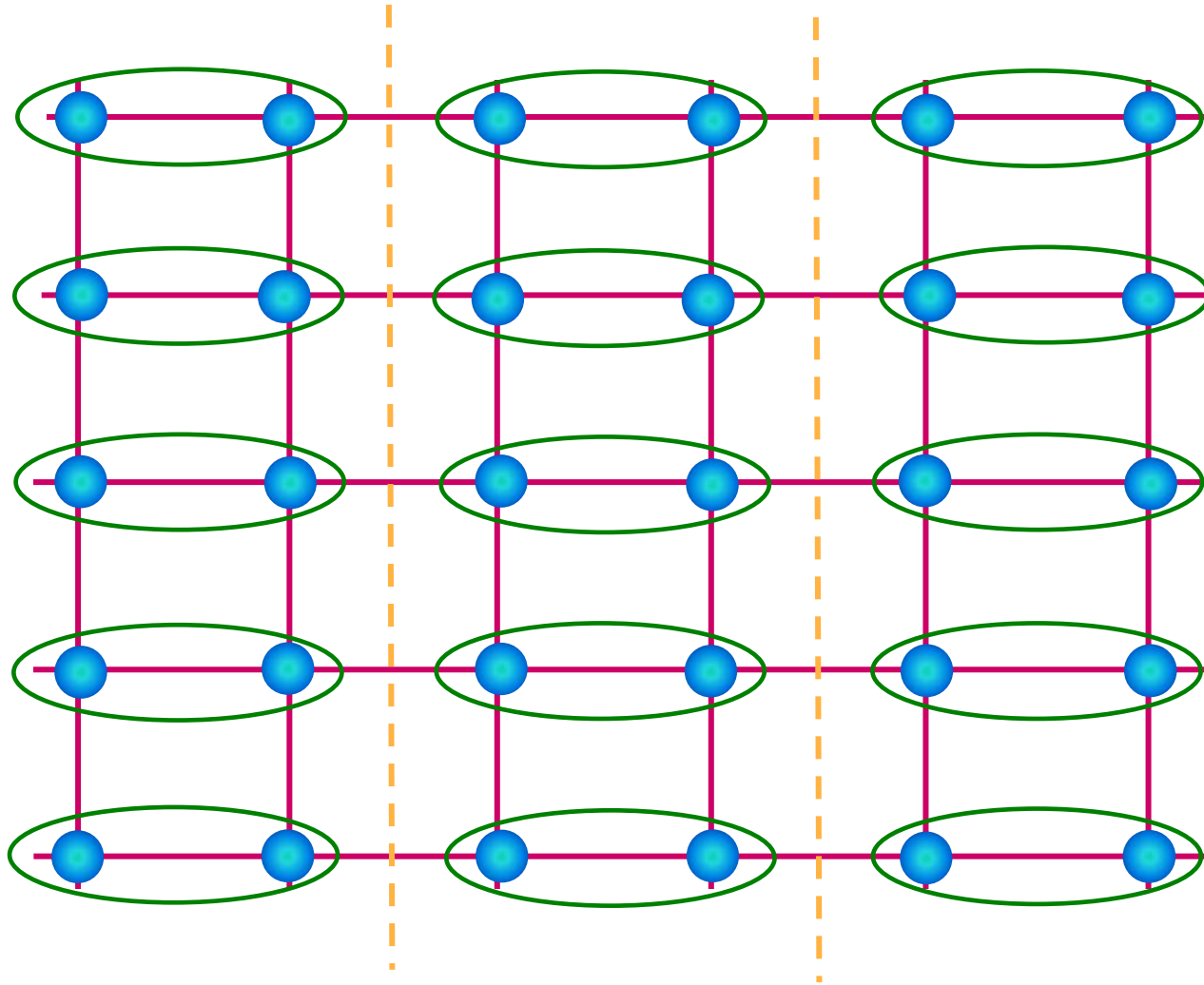
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More possibilities for entanglement with nearby states

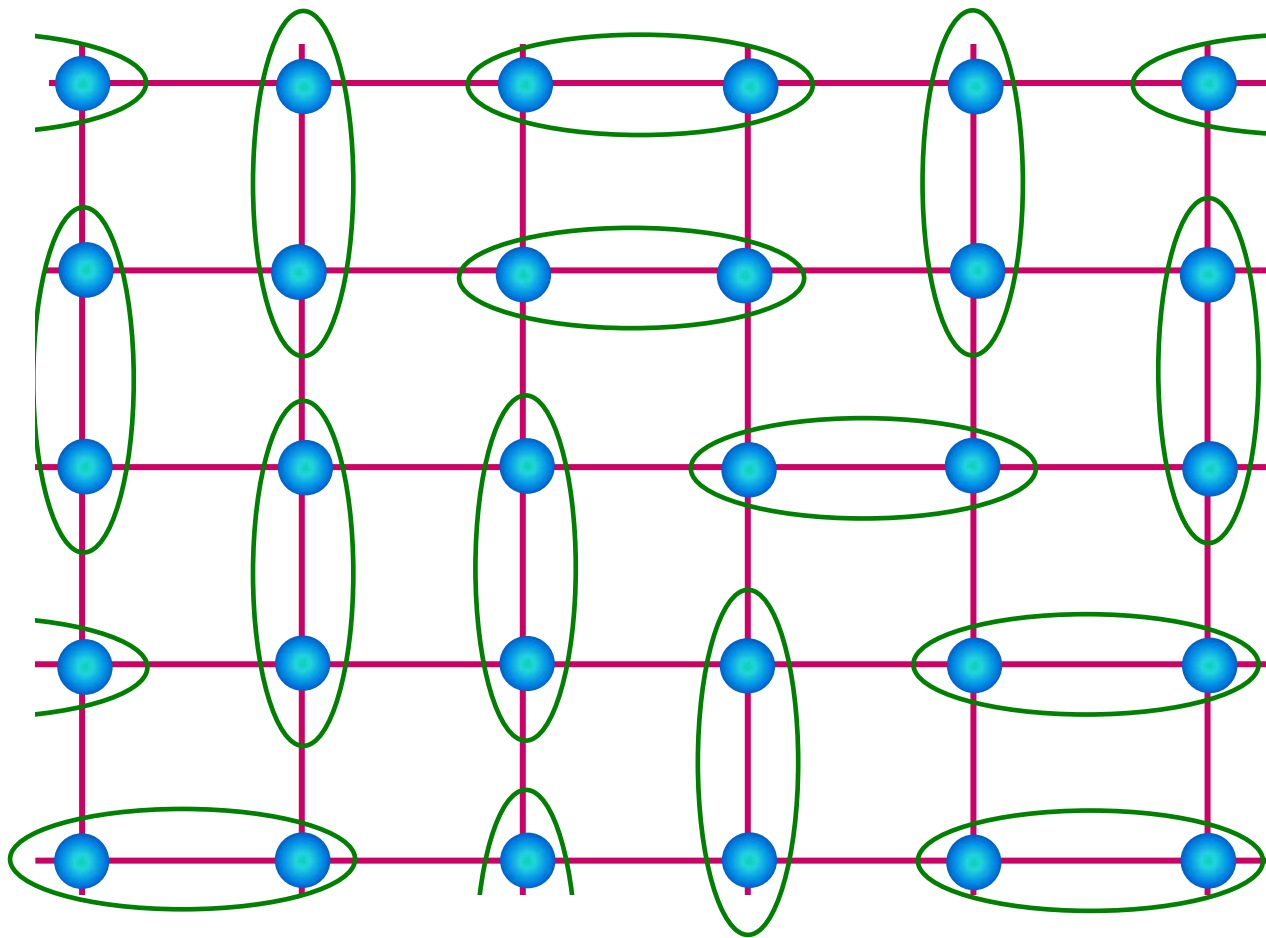


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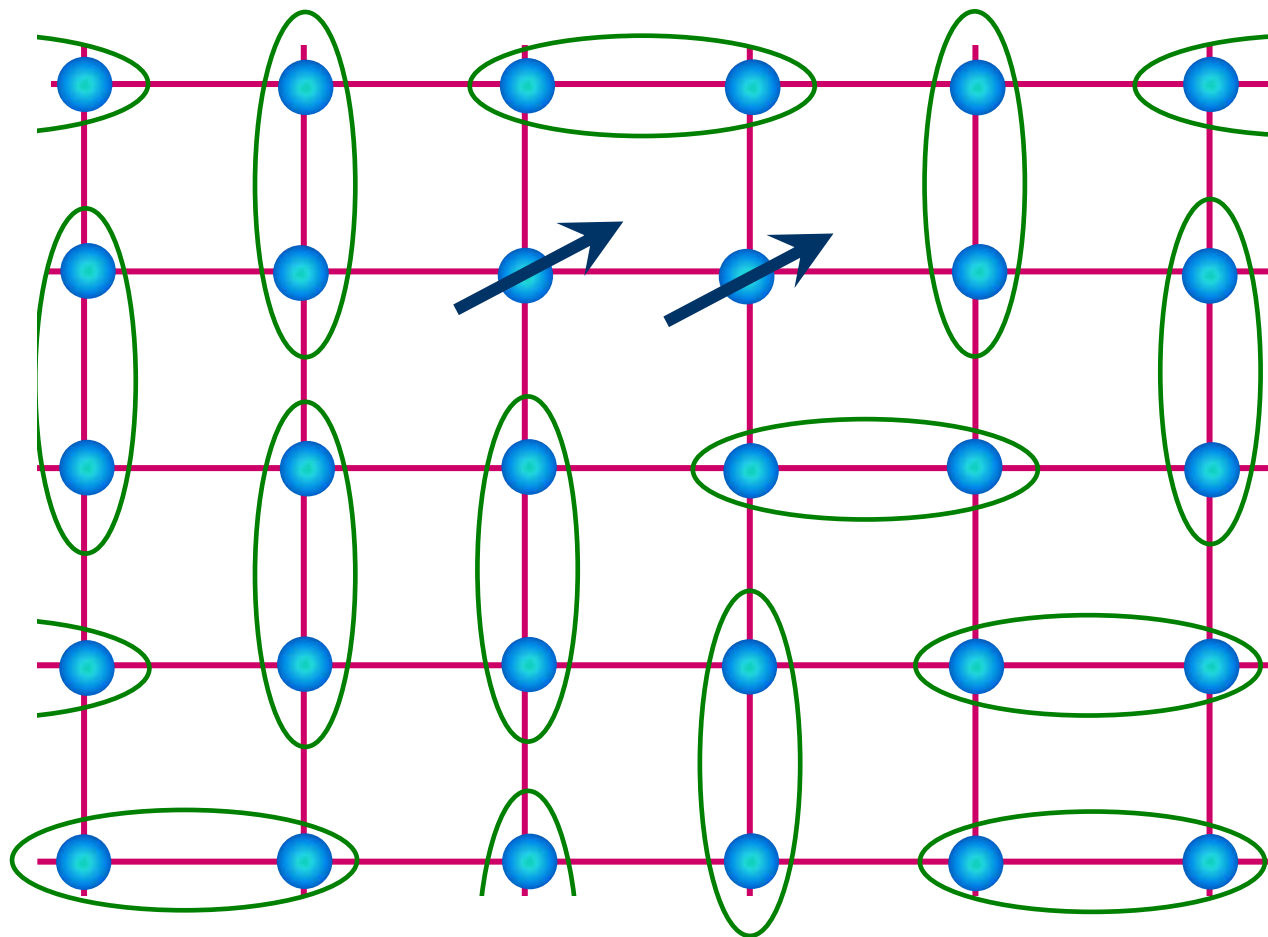
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Excitations of the RVB liquid



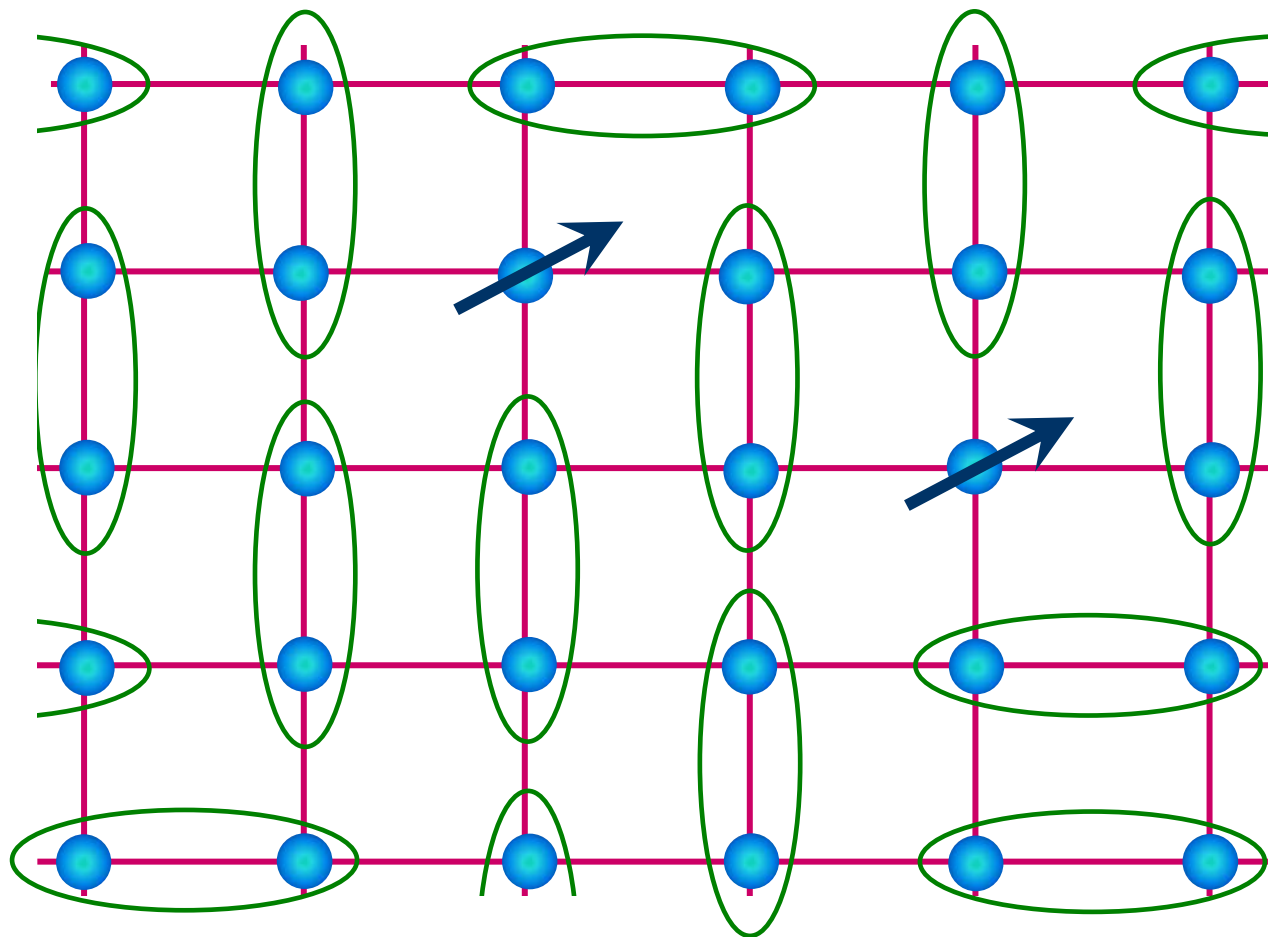
$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



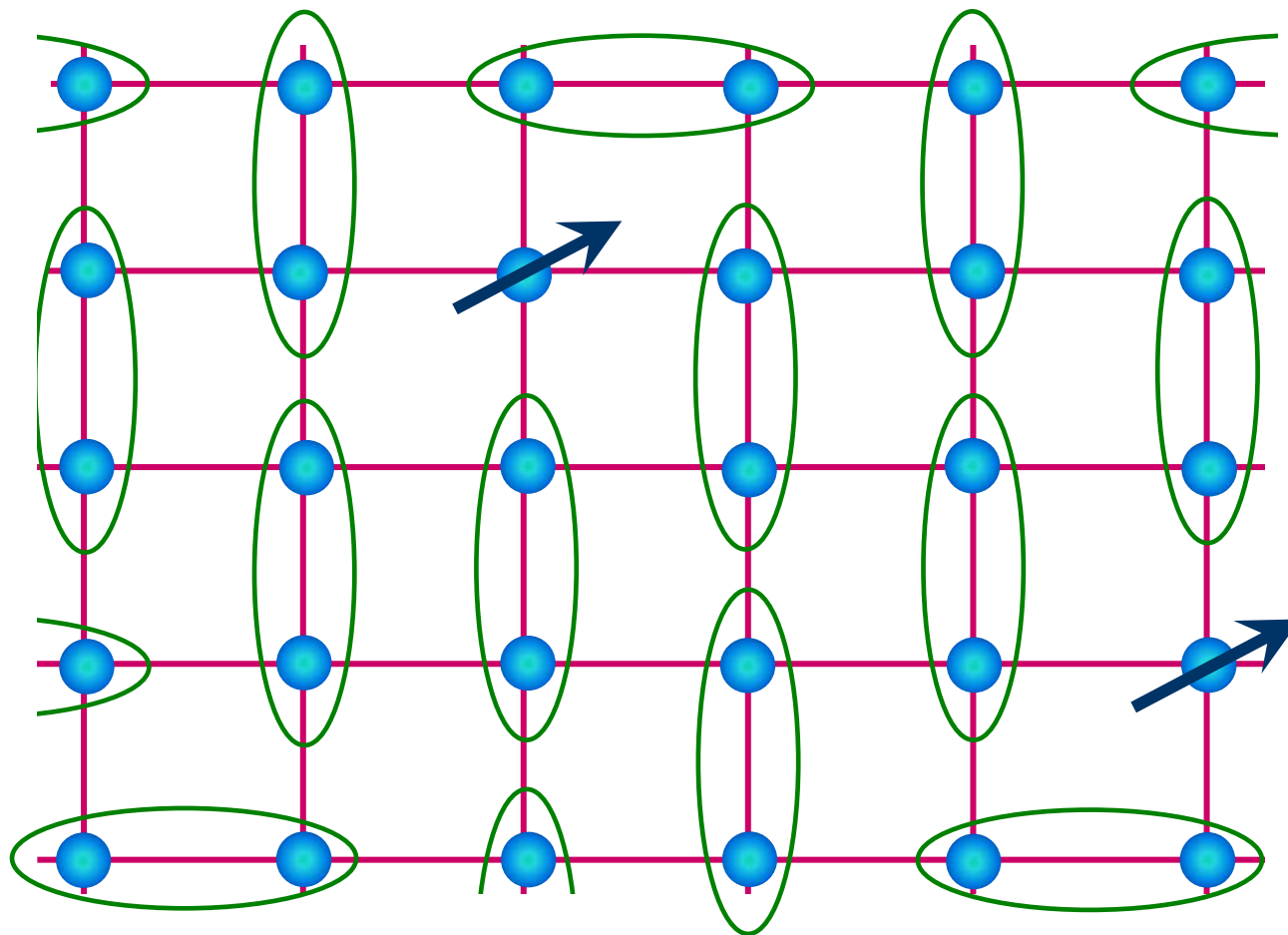
$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



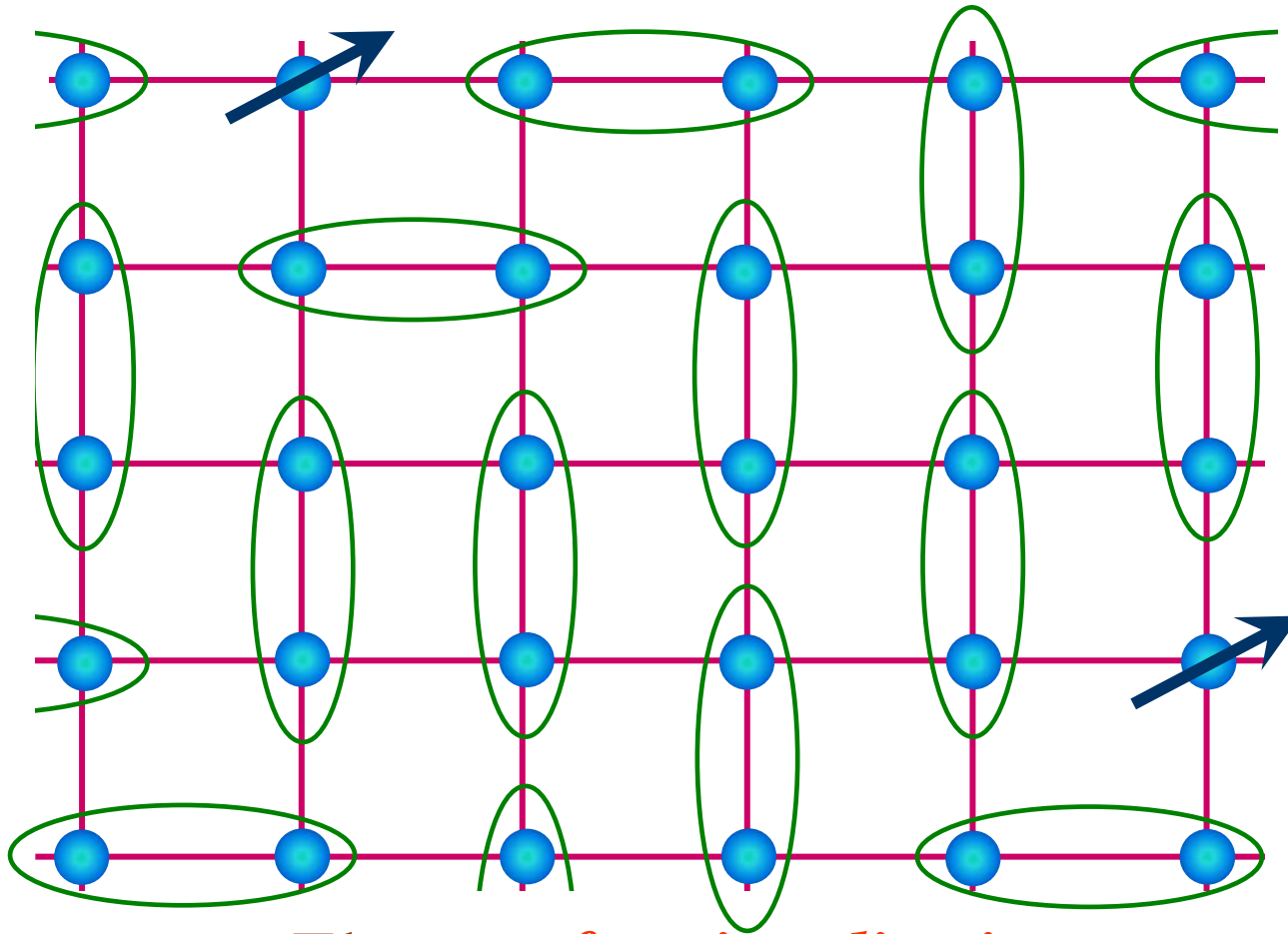
$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid

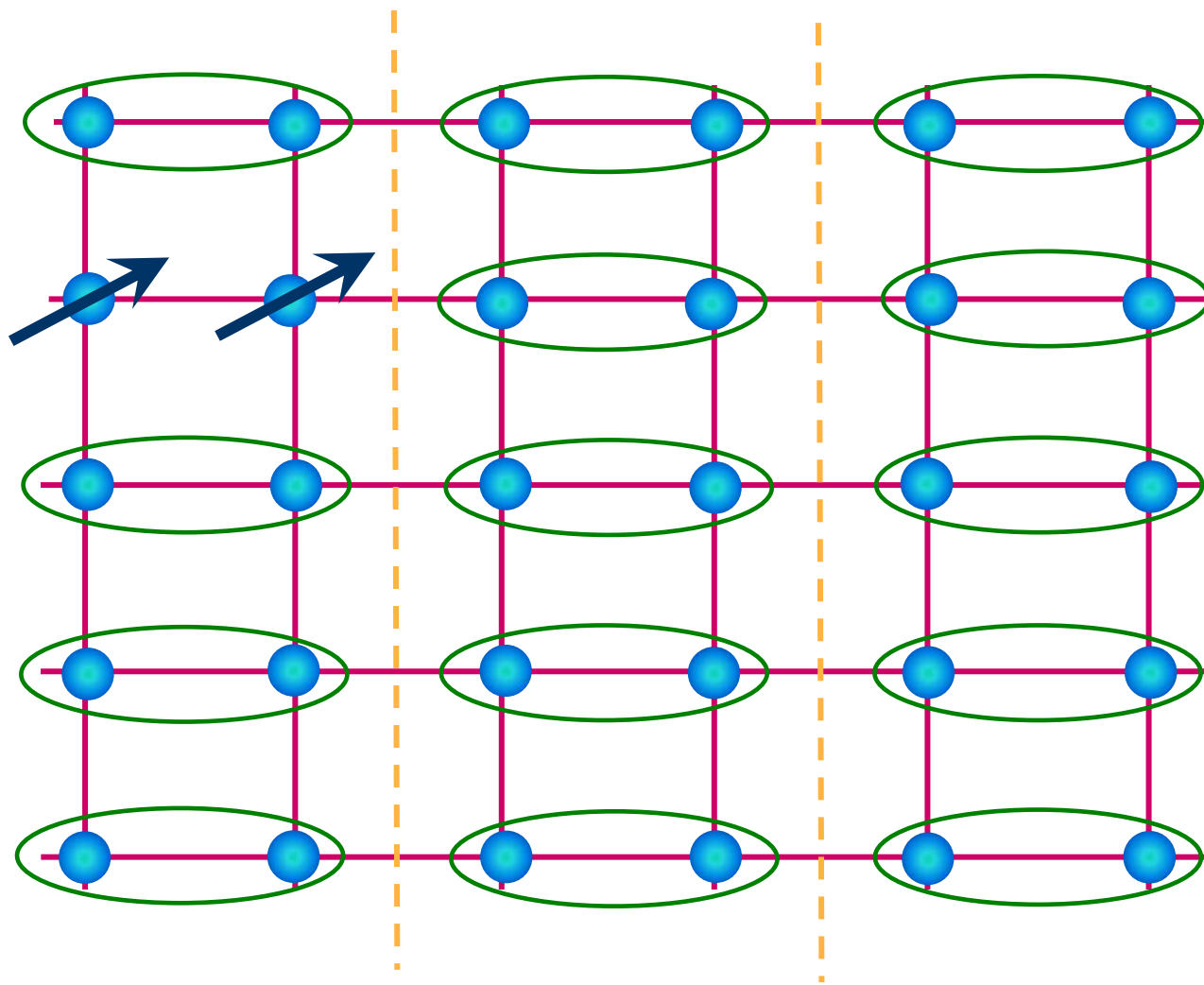


$$\begin{array}{c} \text{Oval with two blue spheres} \\ \hline = \\ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \end{array}$$

Electron *fractionalization*:

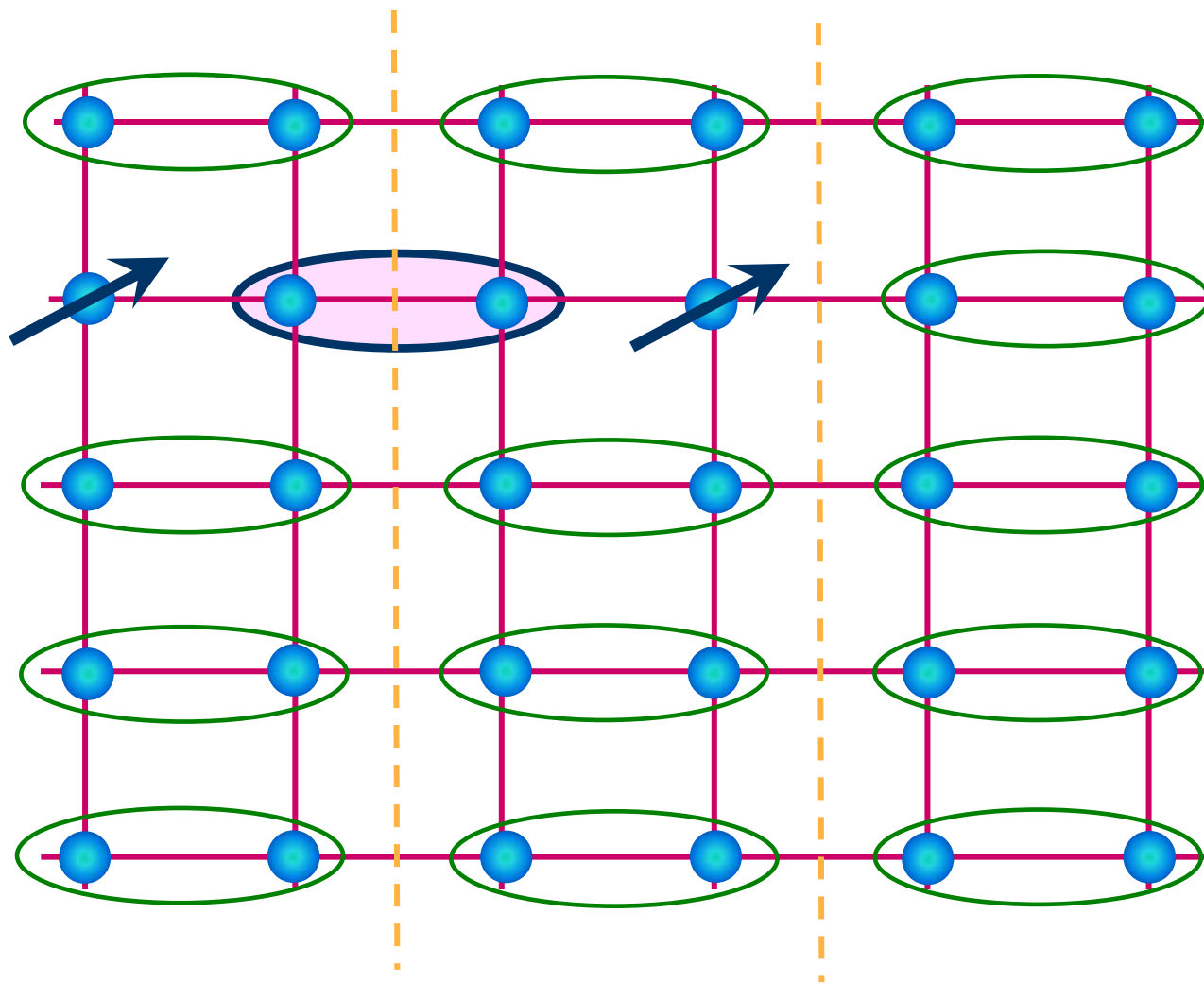
Excitations carry spin $S=1/2$ but no charge

Excitations of the VBS



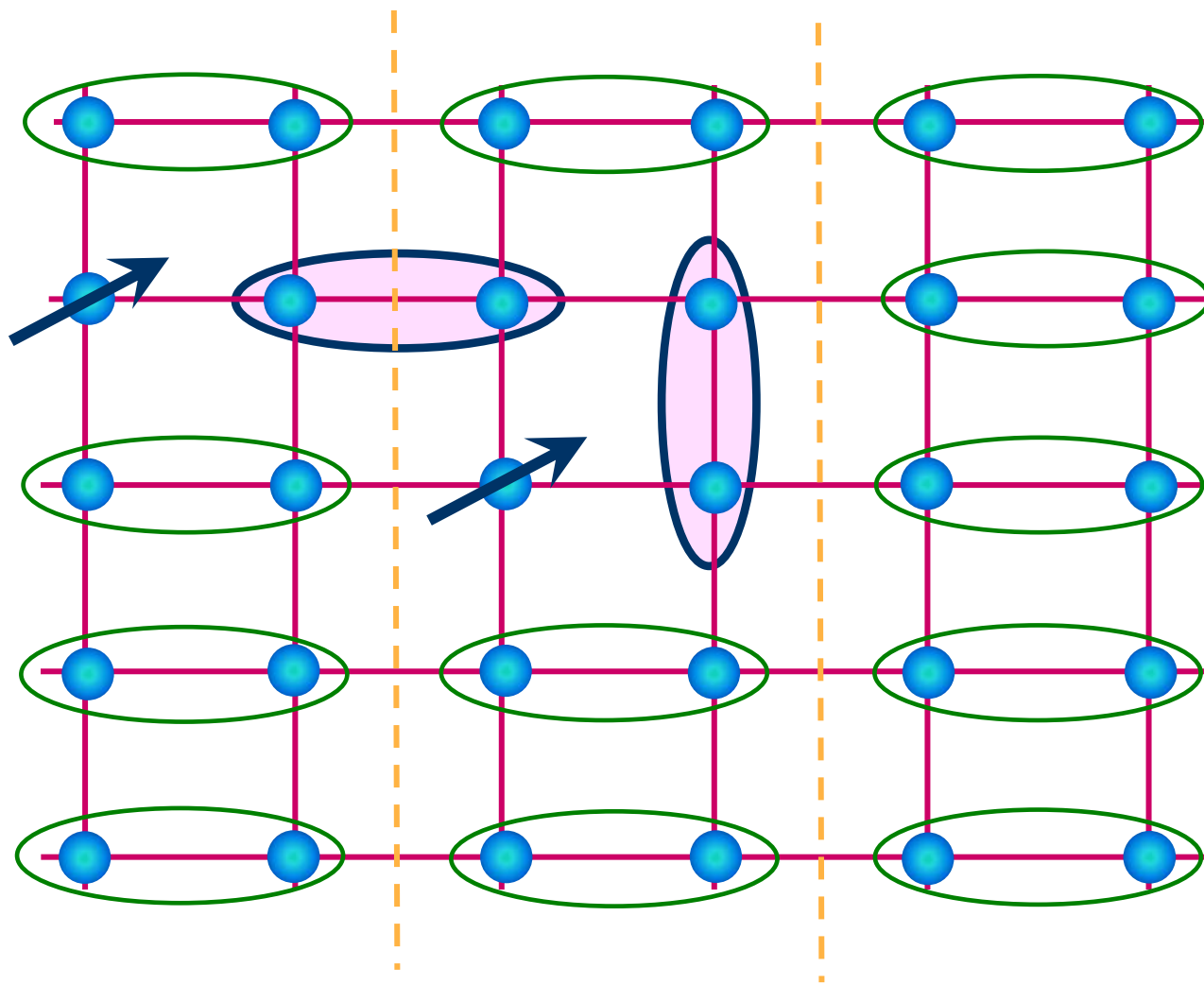
$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the VBS



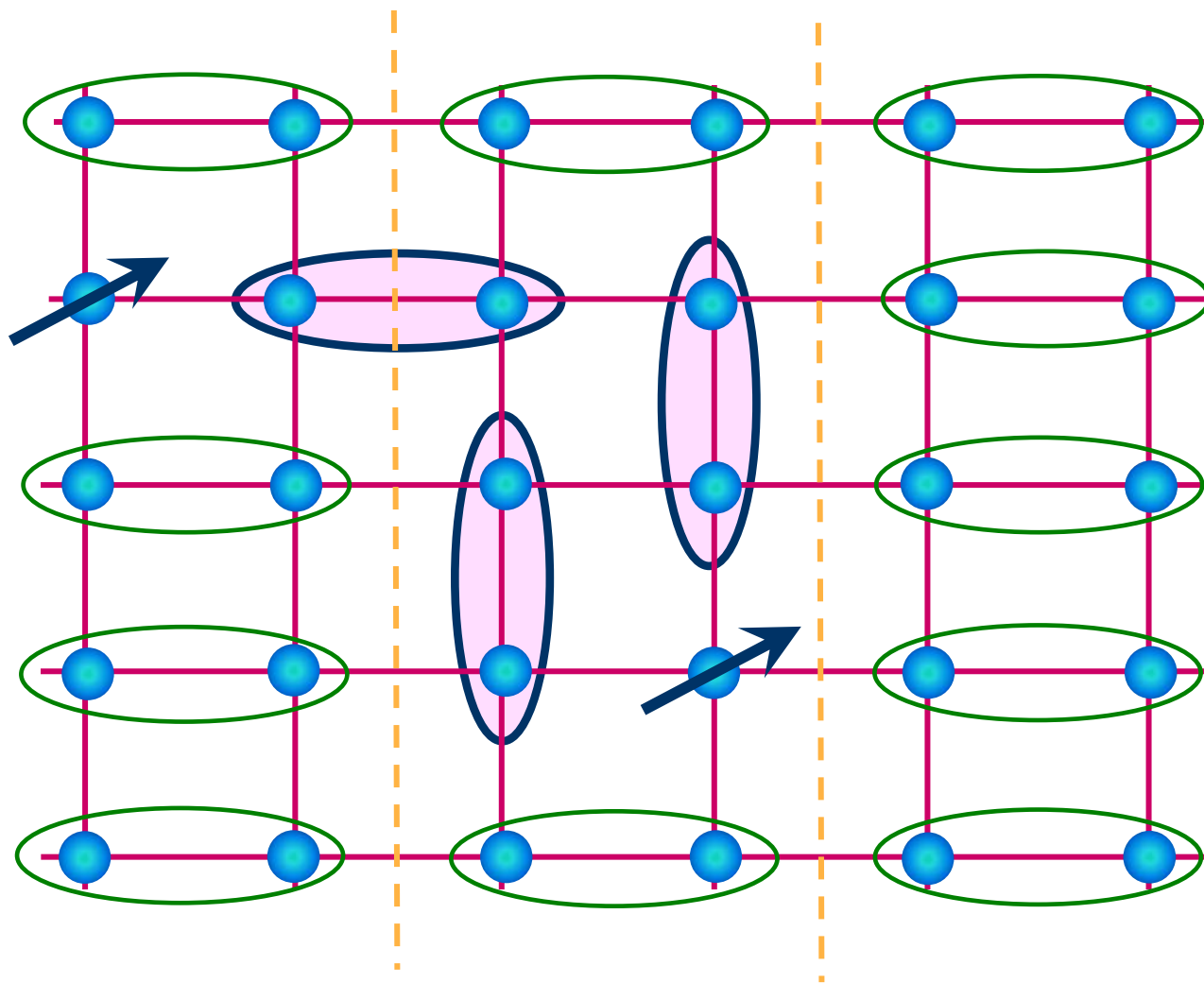
$$\text{[Green oval with two blue spheres]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the VBS



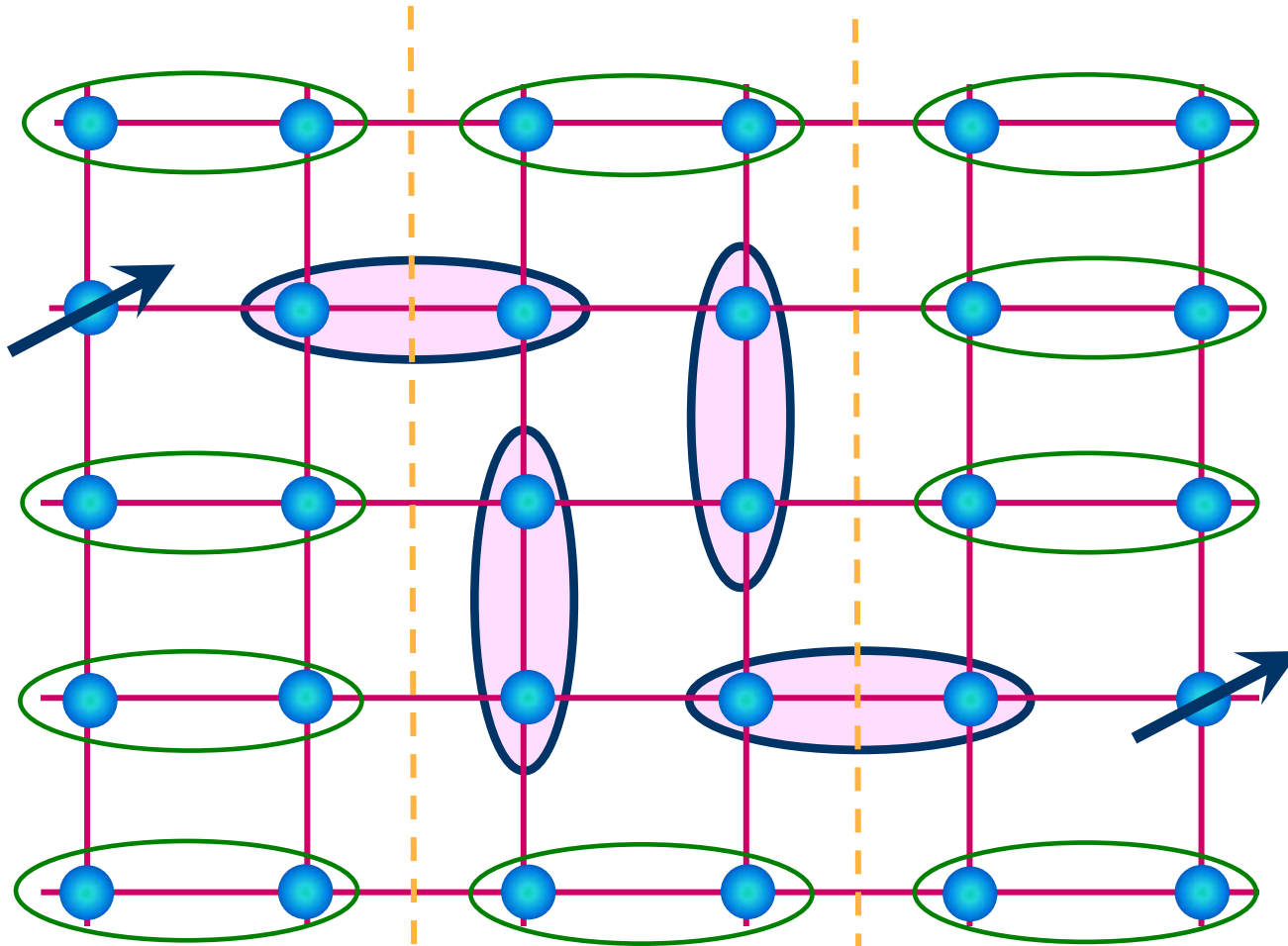
$$\text{green oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the VBS



$$\text{Green oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

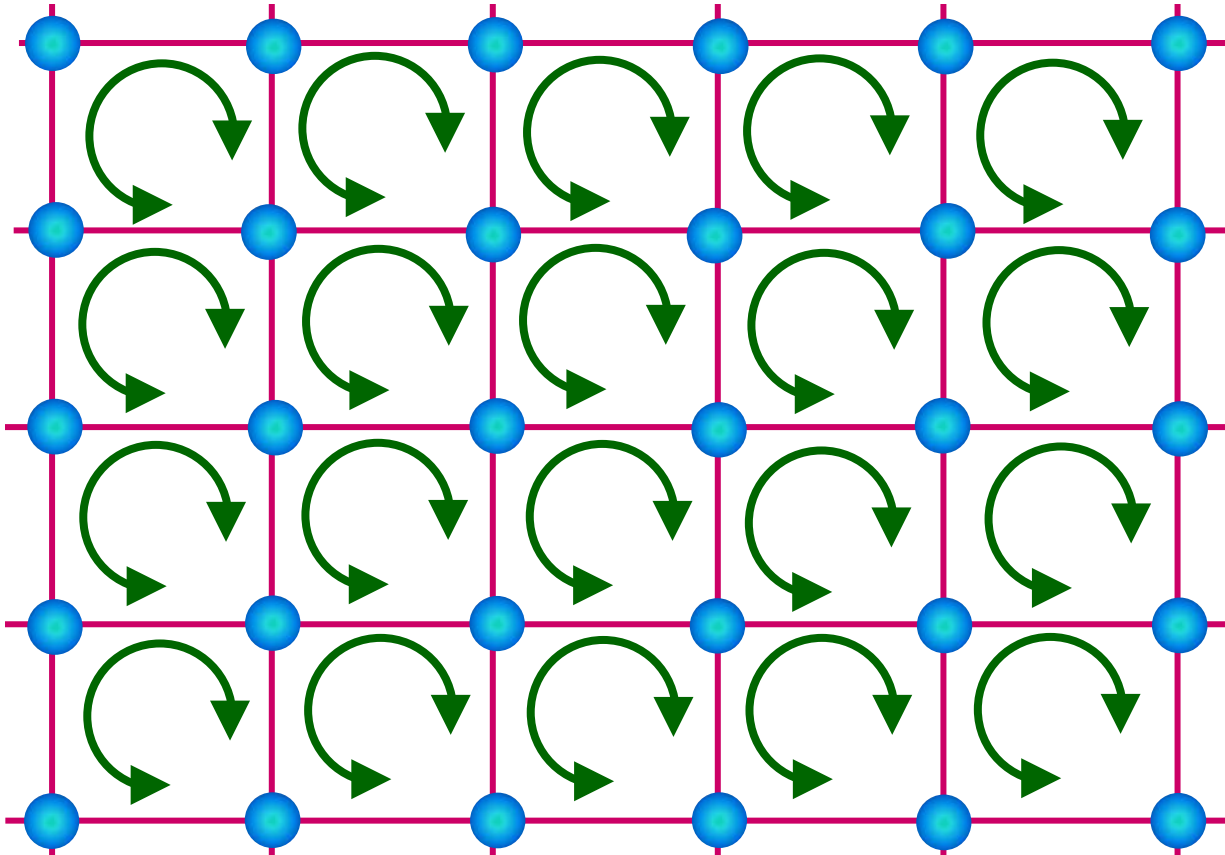
Excitations of the VBS



$$\text{green oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

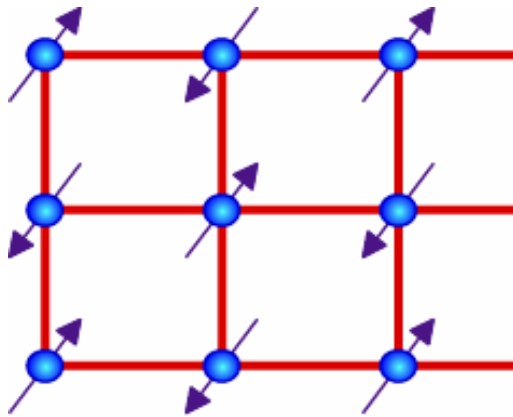
Free spins are unable to move apart:
no fractionalization, but *confinement*

Phase diagram of square lattice antiferromagnet

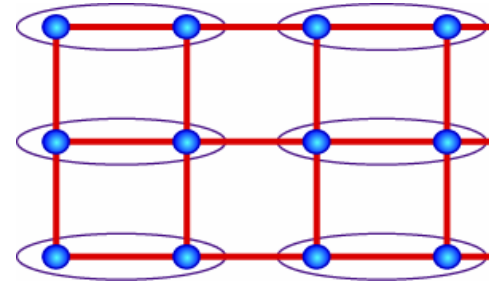


$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} 4\text{-spin exchange}$$

Phase diagram of square lattice antiferromagnet



Neel order

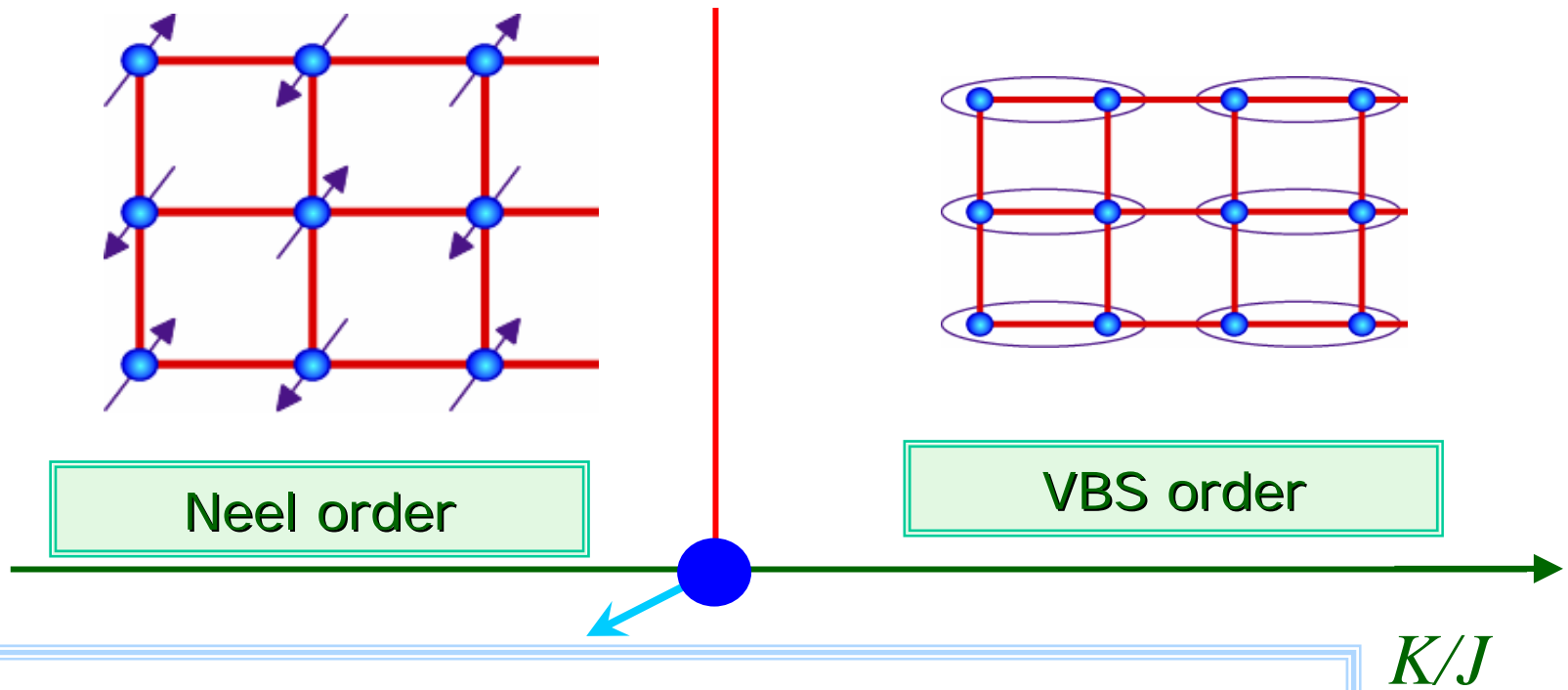


VBS order

K/J

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} 4\text{-spin exchange}$$

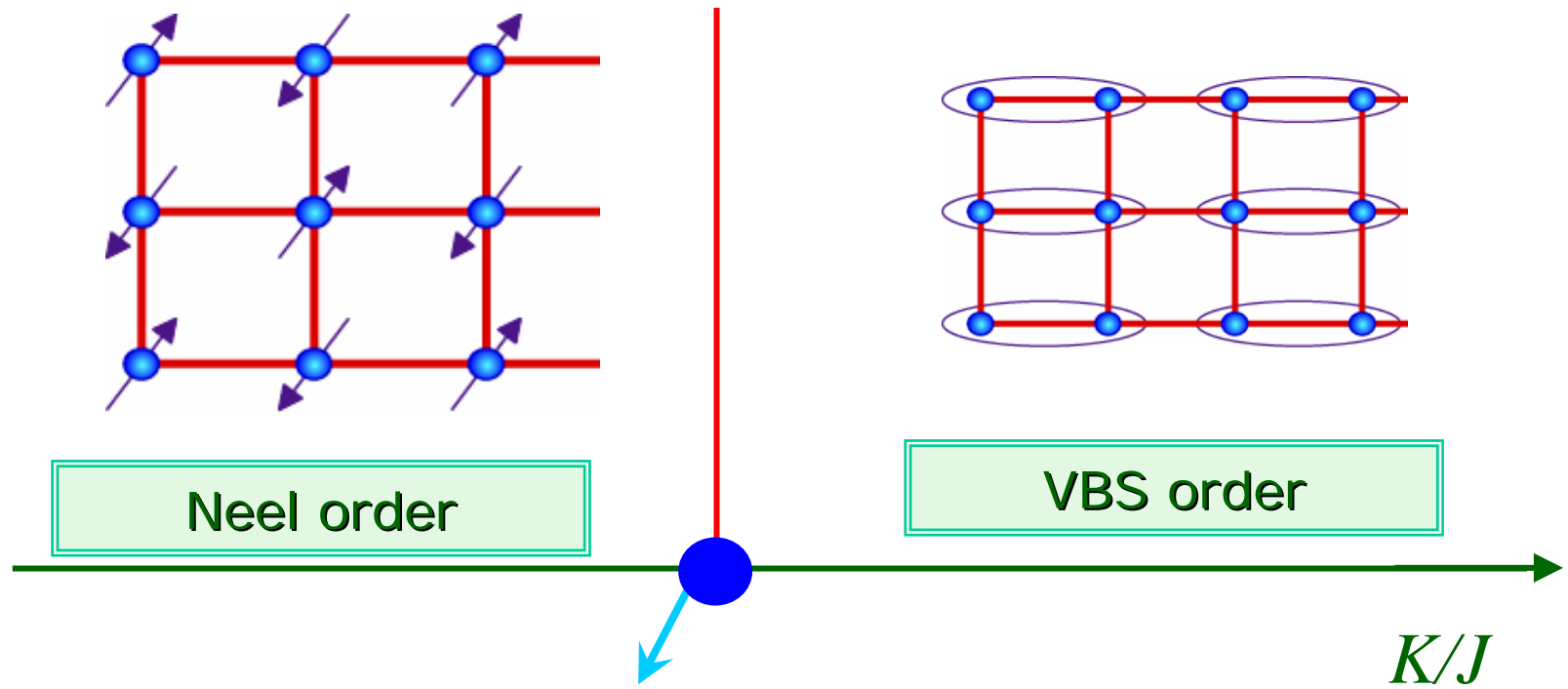
Phase diagram of square lattice antiferromagnet



RVB physics appears at the quantum critical point which has fractionalized excitations: “deconfined criticality”

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} 4\text{-spin exchange}$$

Phase diagram of square lattice antiferromagnet



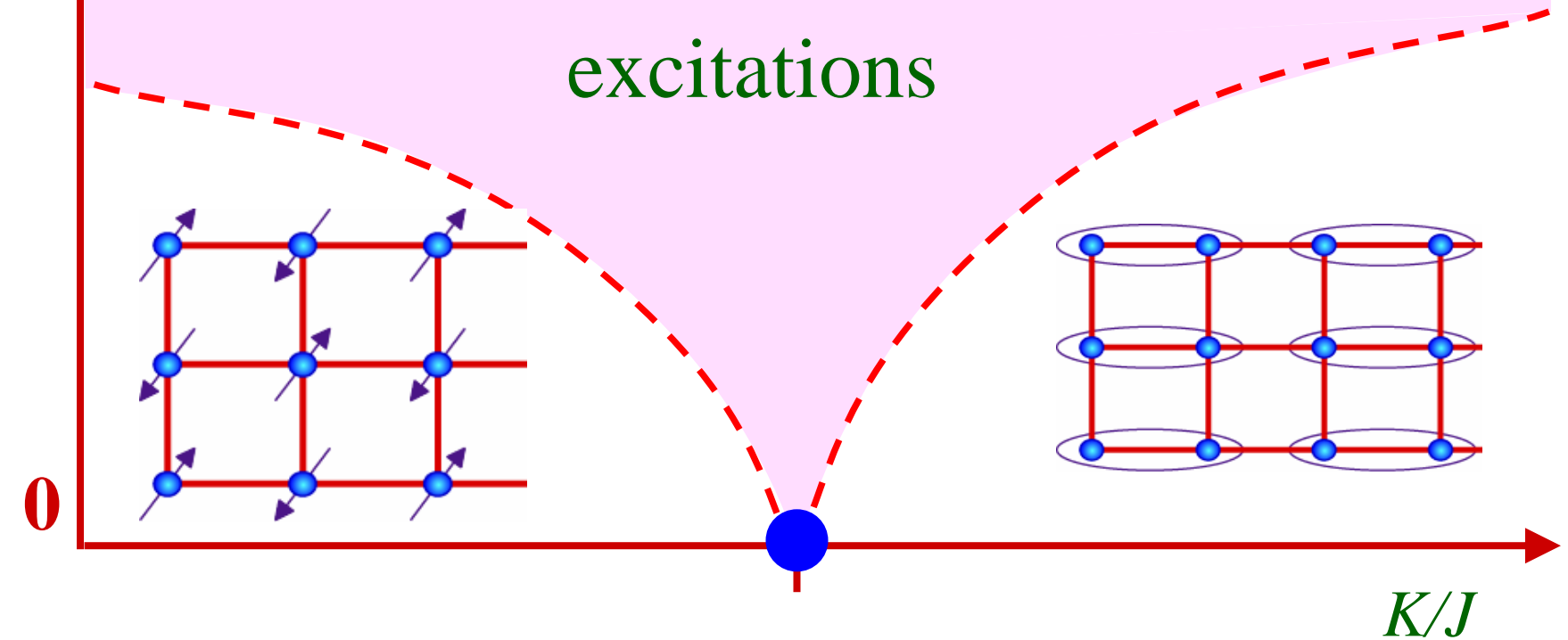
Second-order critical point described by

$$\mathcal{S}_{\text{critical}} = \int d^2x d\tau \left[|(\partial_\mu - iA_\mu)z_\alpha|^2 + r |z_\alpha|^2 + \frac{u}{2} (|z_\alpha|^2)^2 + \frac{1}{4e^2} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2 \right]$$

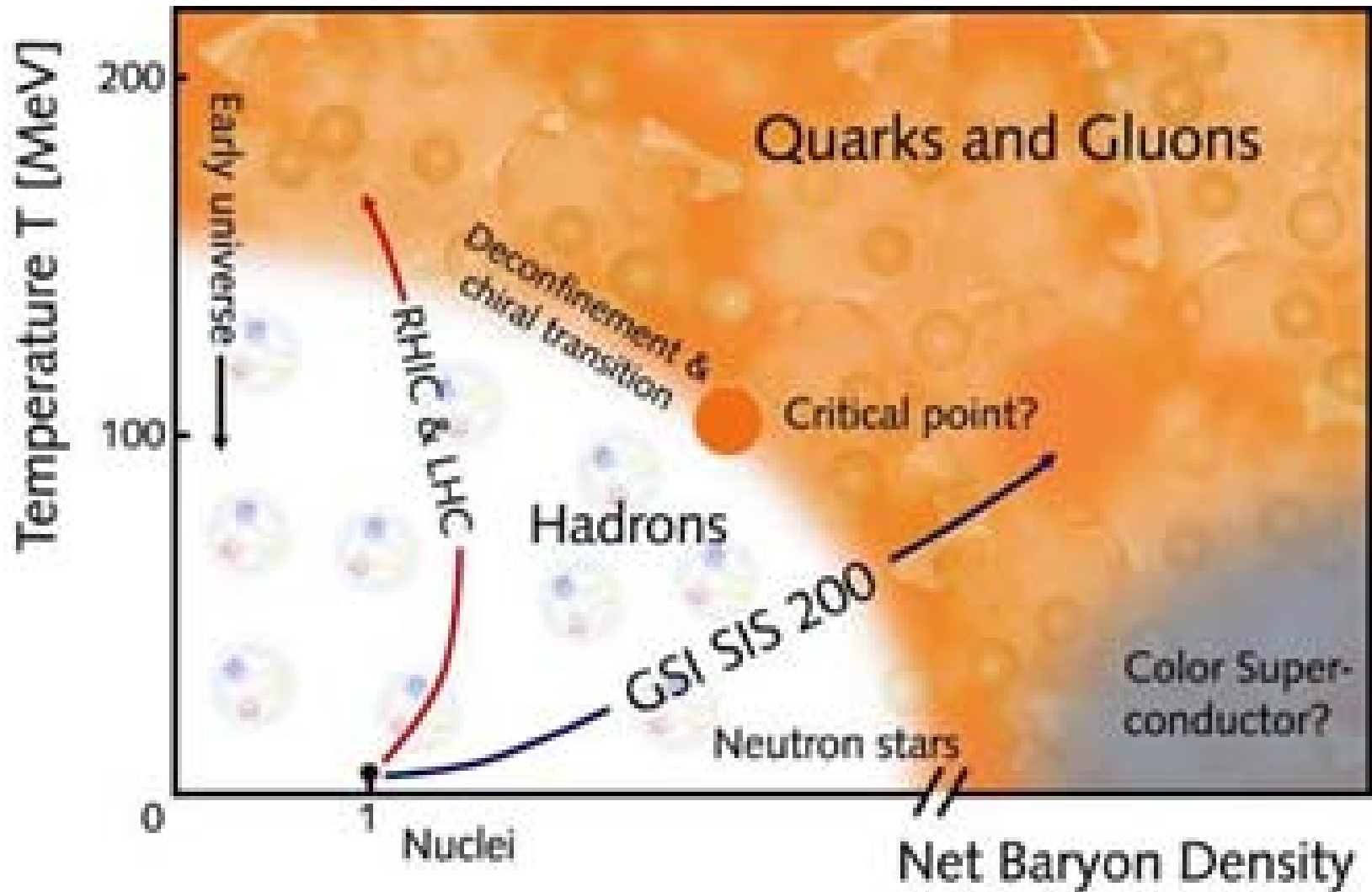
at its critical point $r = r_c$, where z_α are the neutral $S = 1/2$ spinons and A_μ is a *non-compact* U(1) gauge field.

Temperature, T

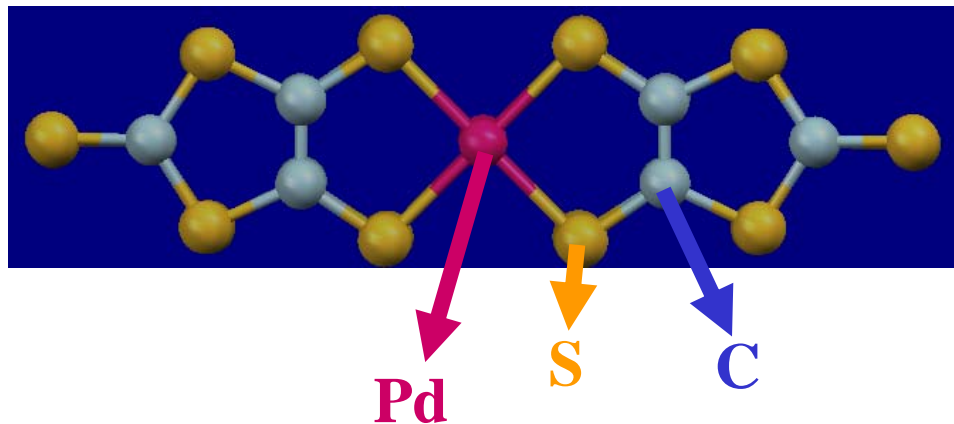
Quantum criticality of
fractionalized
excitations



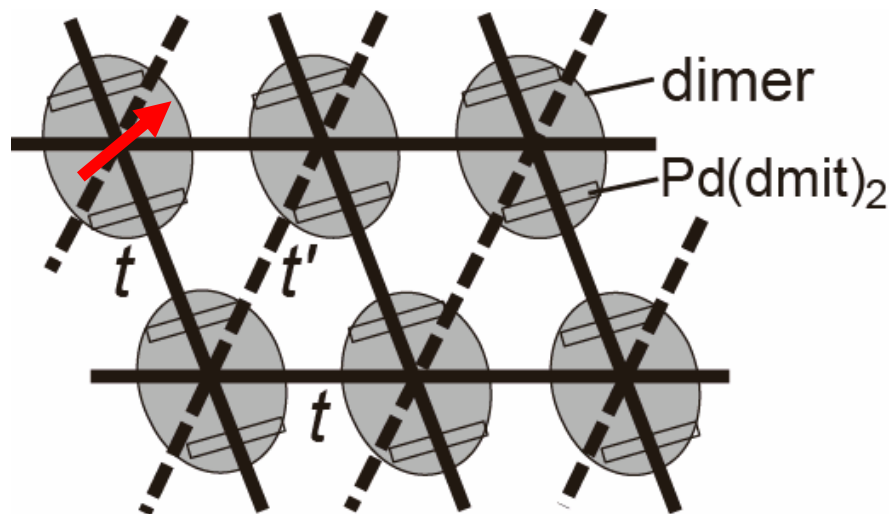
Phases of nuclear matter



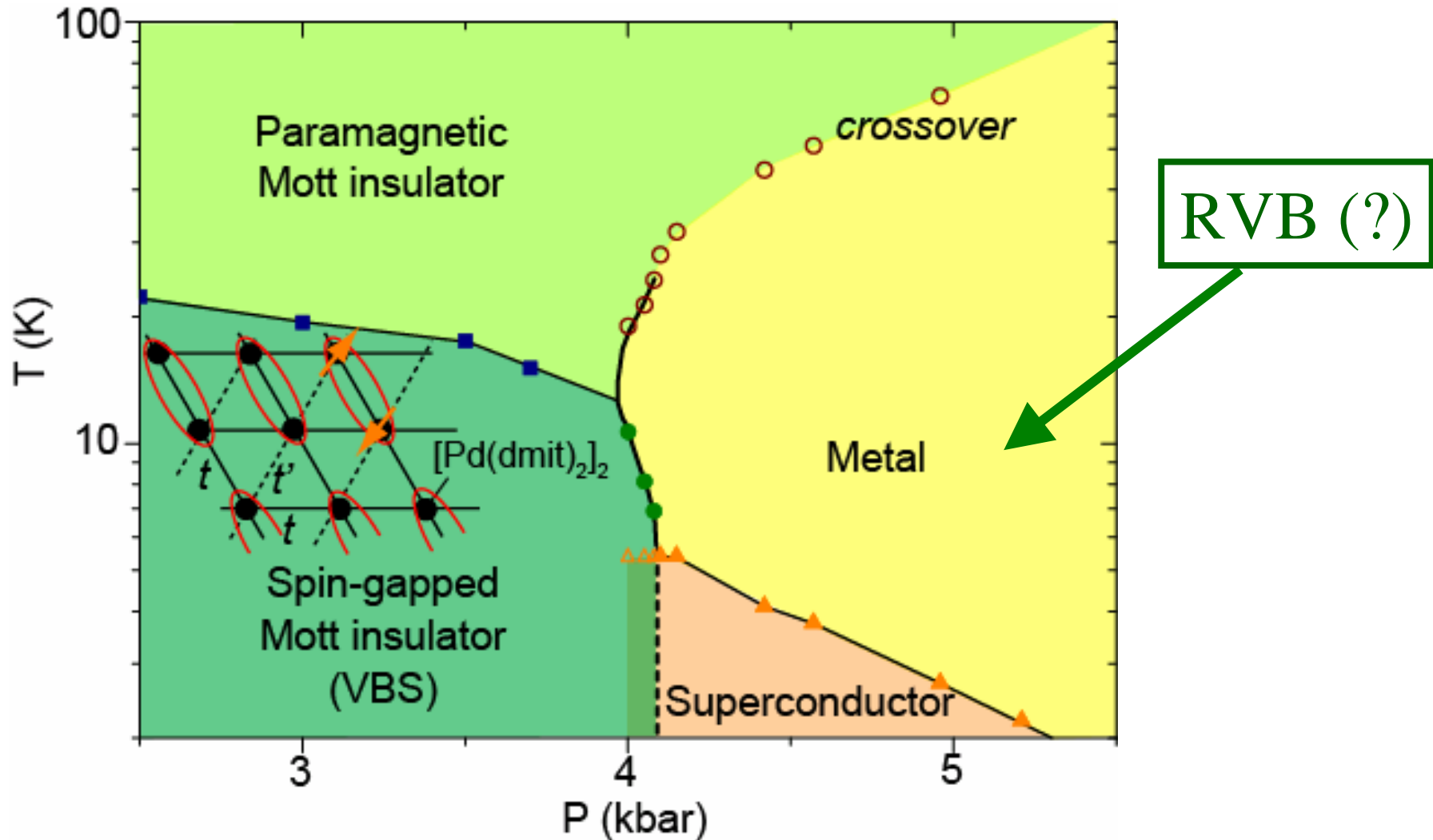
Observation of a valence bond solid (VBS)



One free electron
spin on each
vertex of a
triangular lattice

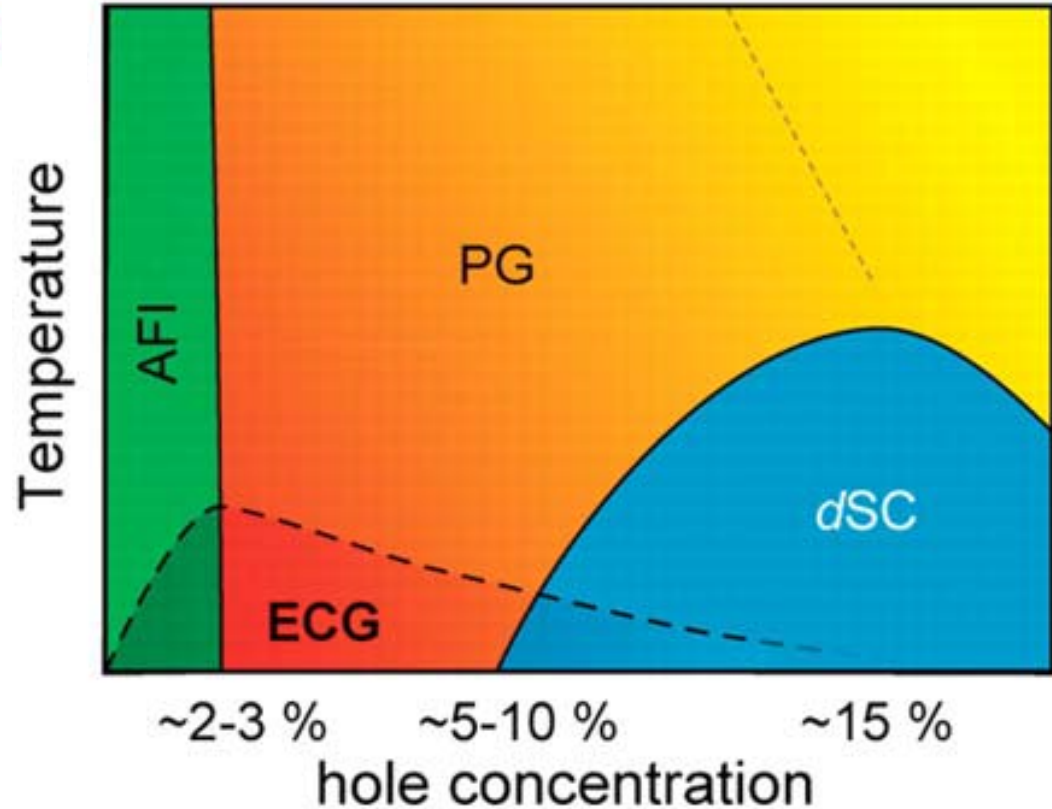
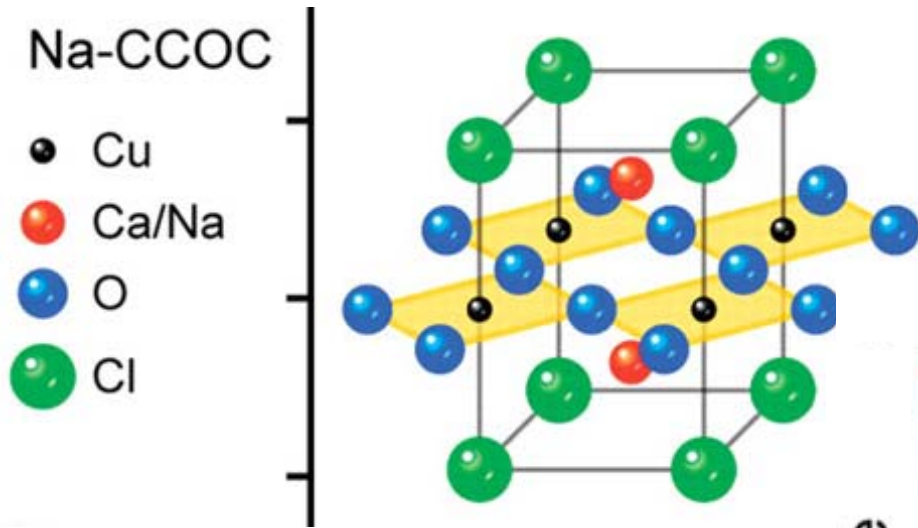


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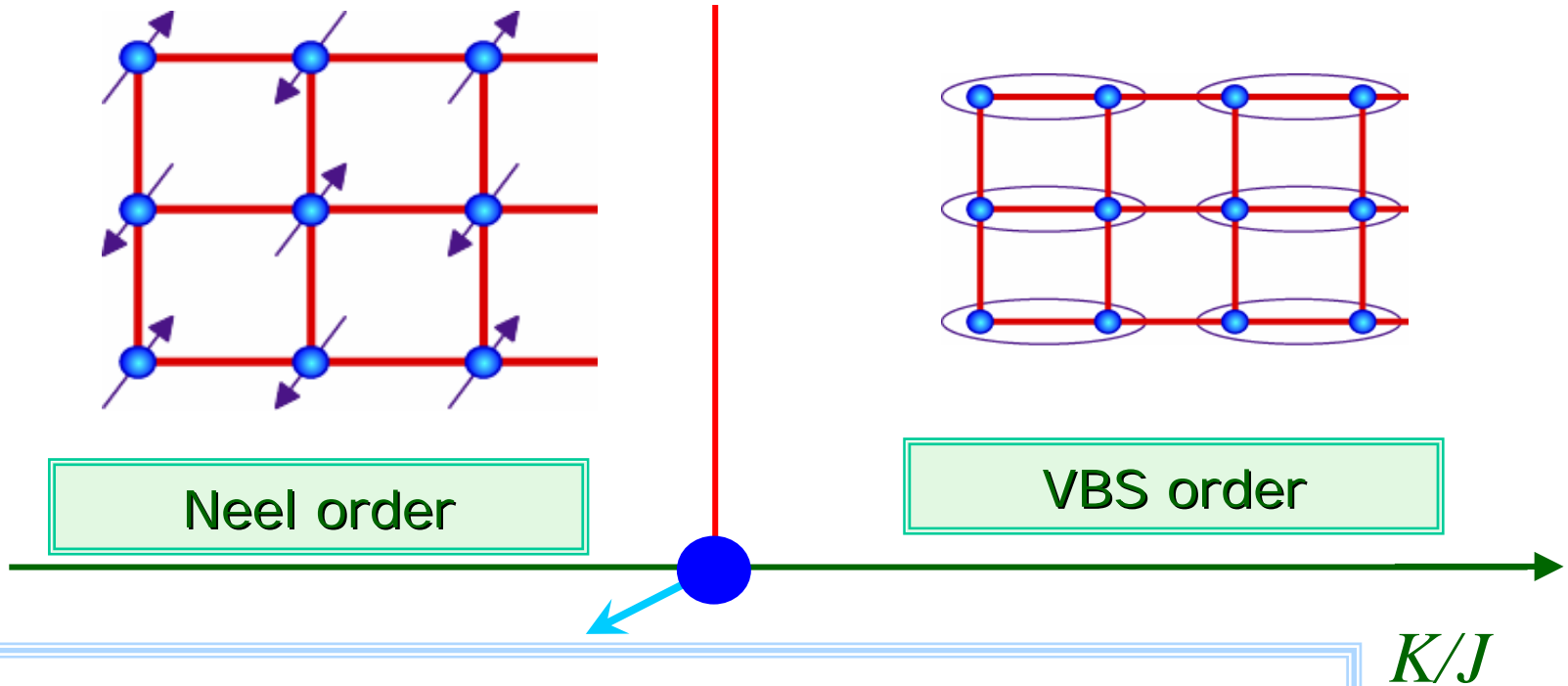


Pressure-temperature phase diagram of $\text{ETMe}_3\text{P}[\text{Pd}(\text{dmit})_2]_2$

Temperature-doping phase diagram of the cuprate superconductors

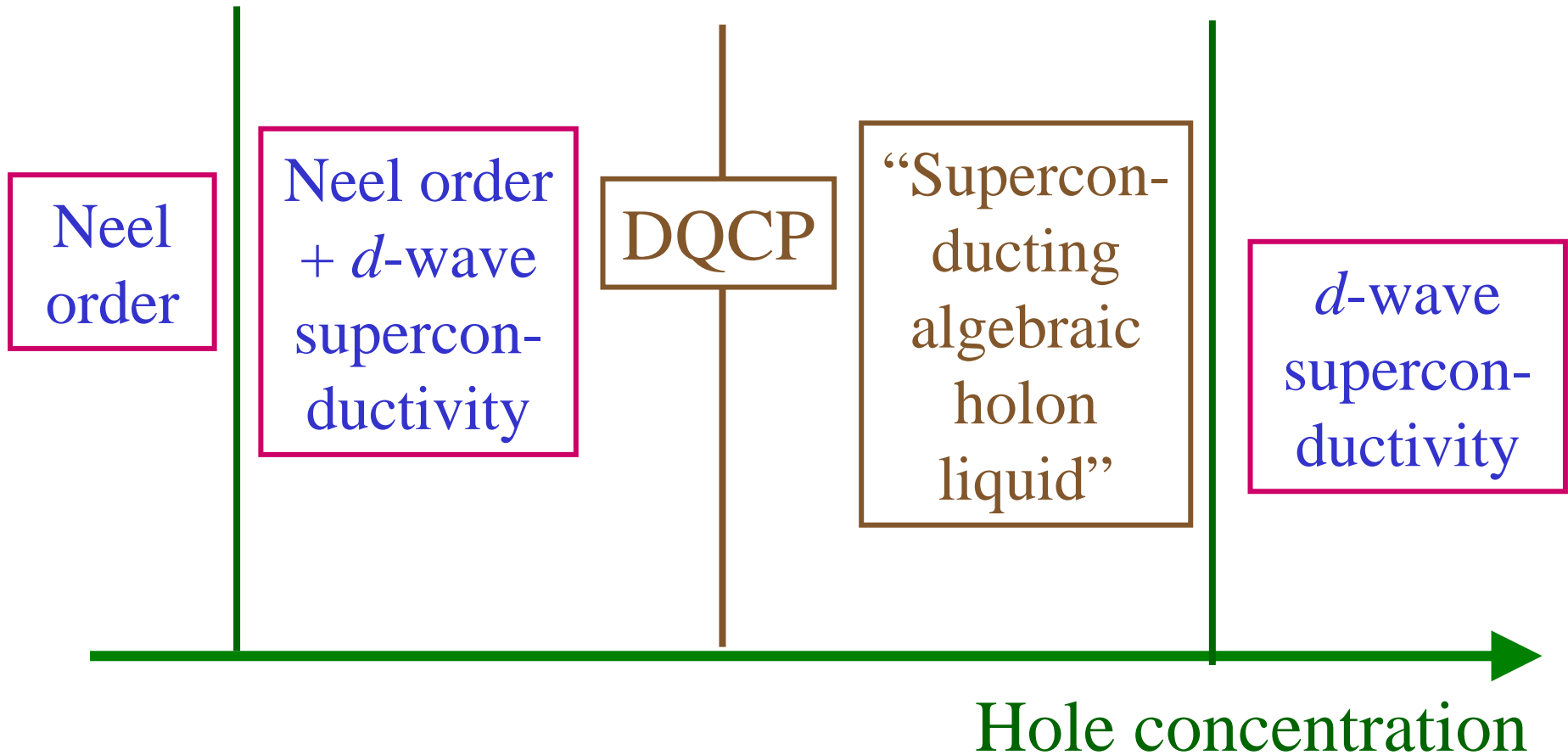


Temperature-doping phase diagram of the cuprate superconductors

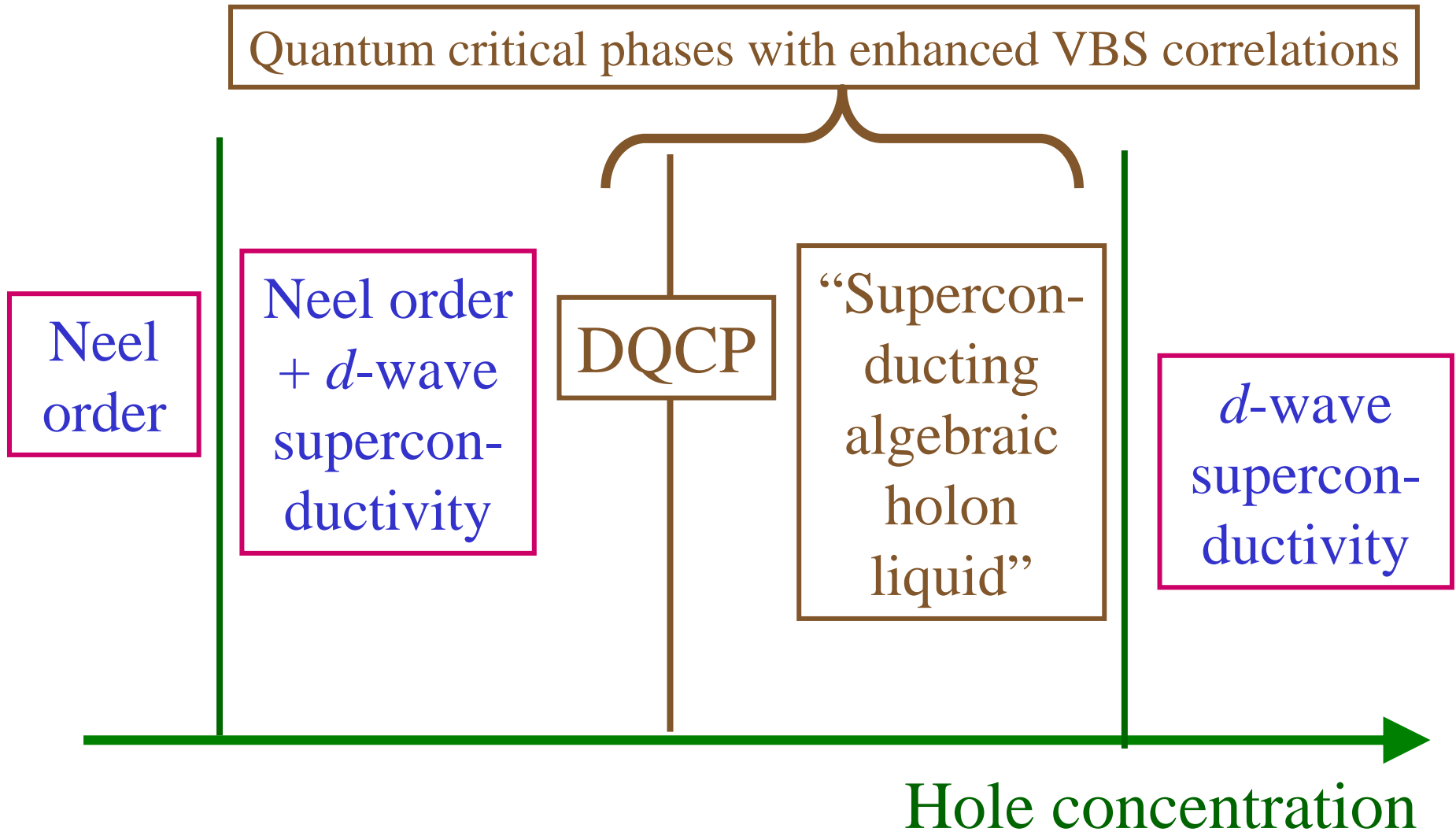


Deconfined quantum critical point
(DQCP)

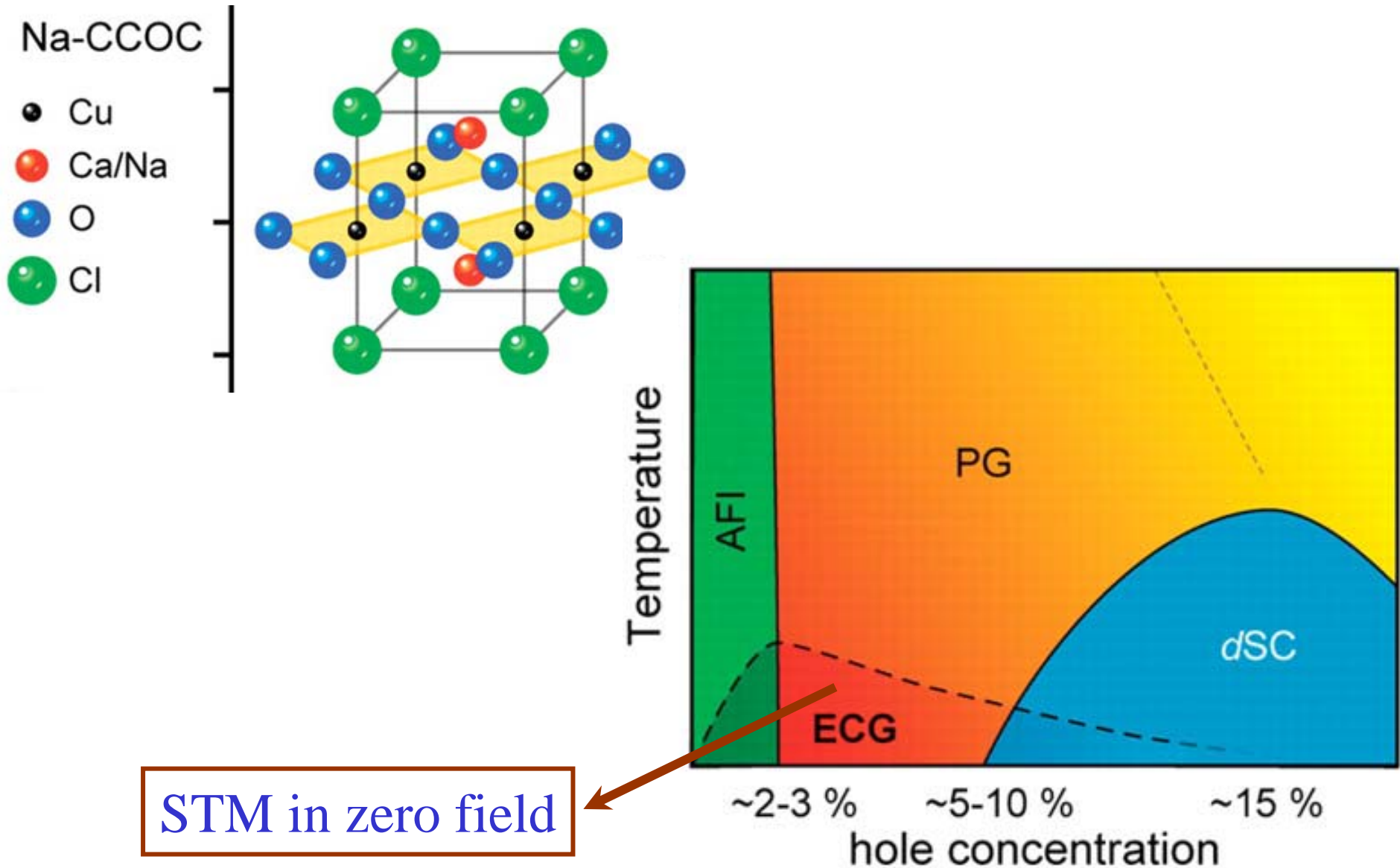
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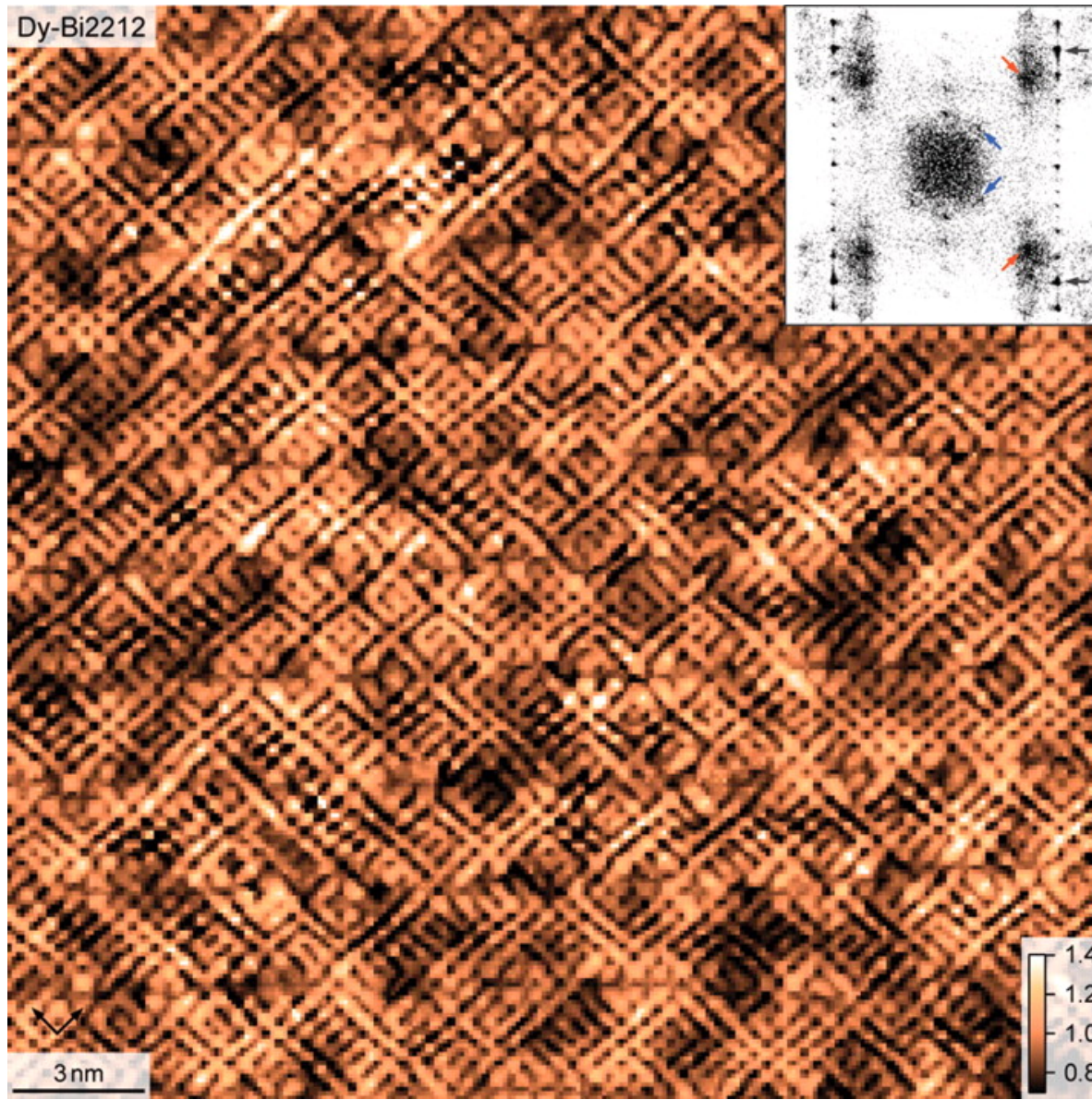


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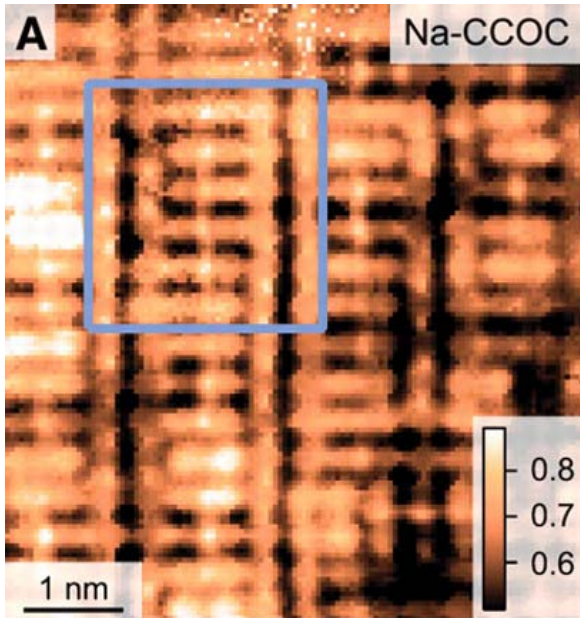


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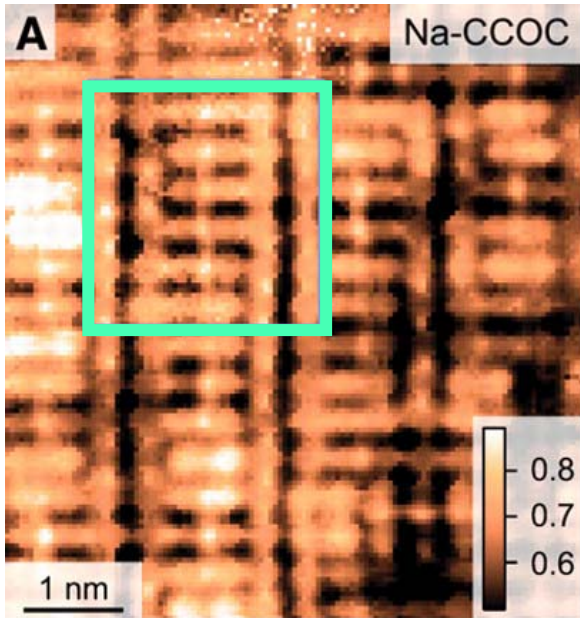




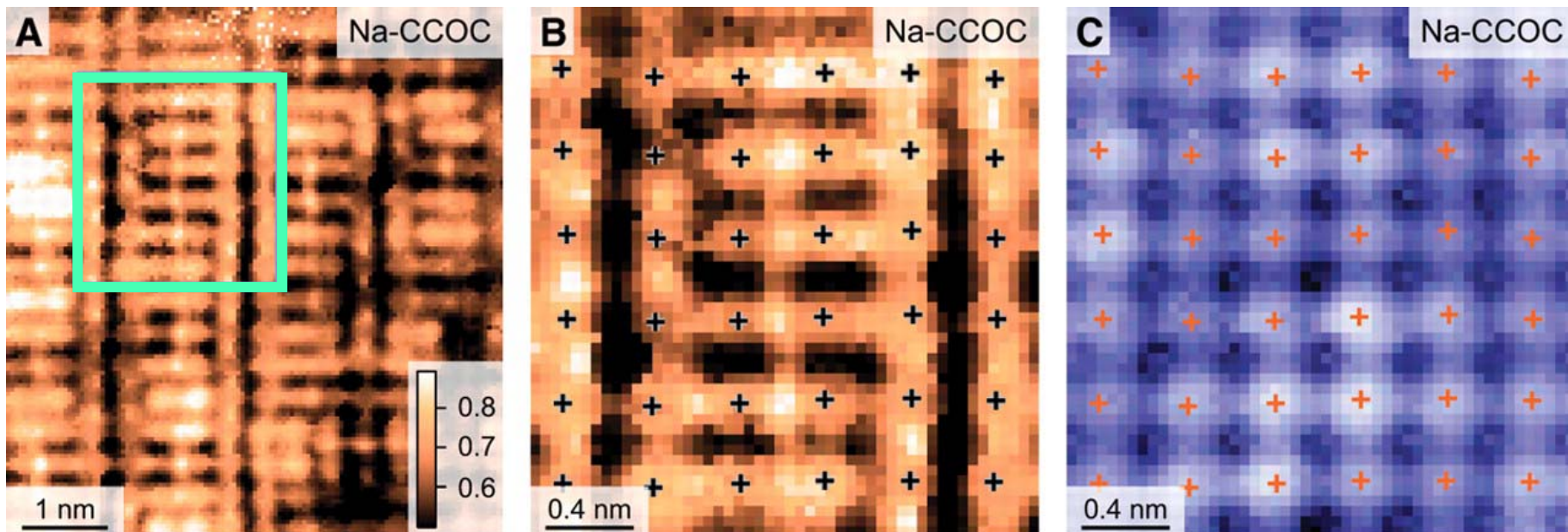
Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)



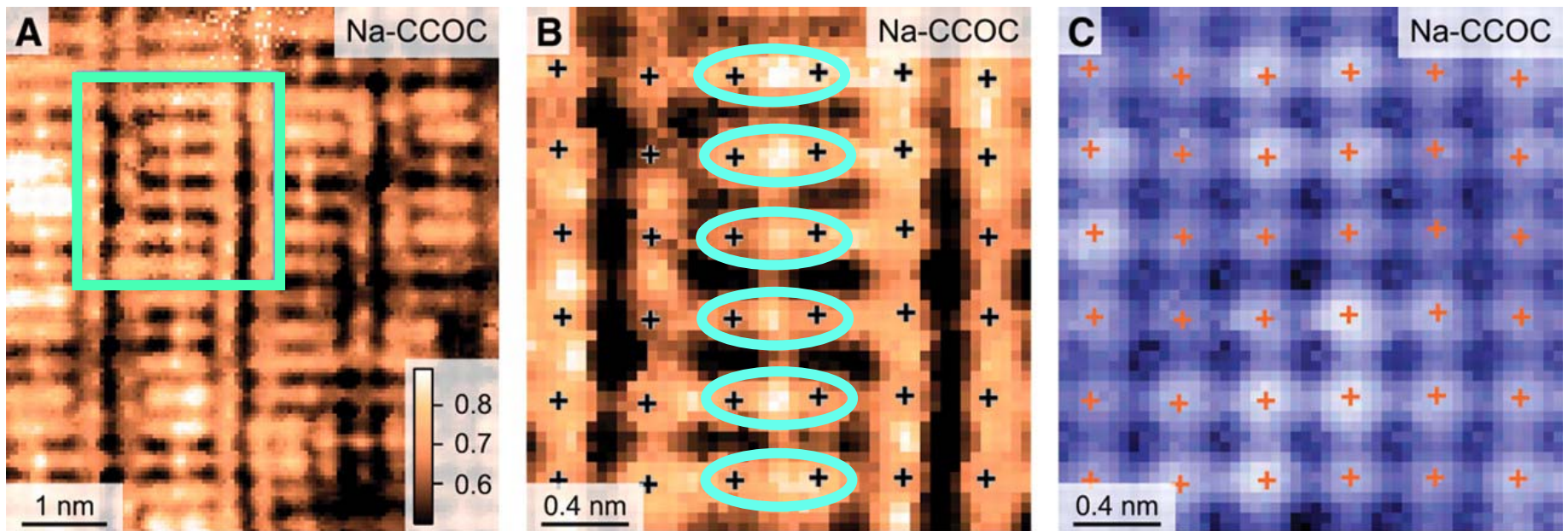
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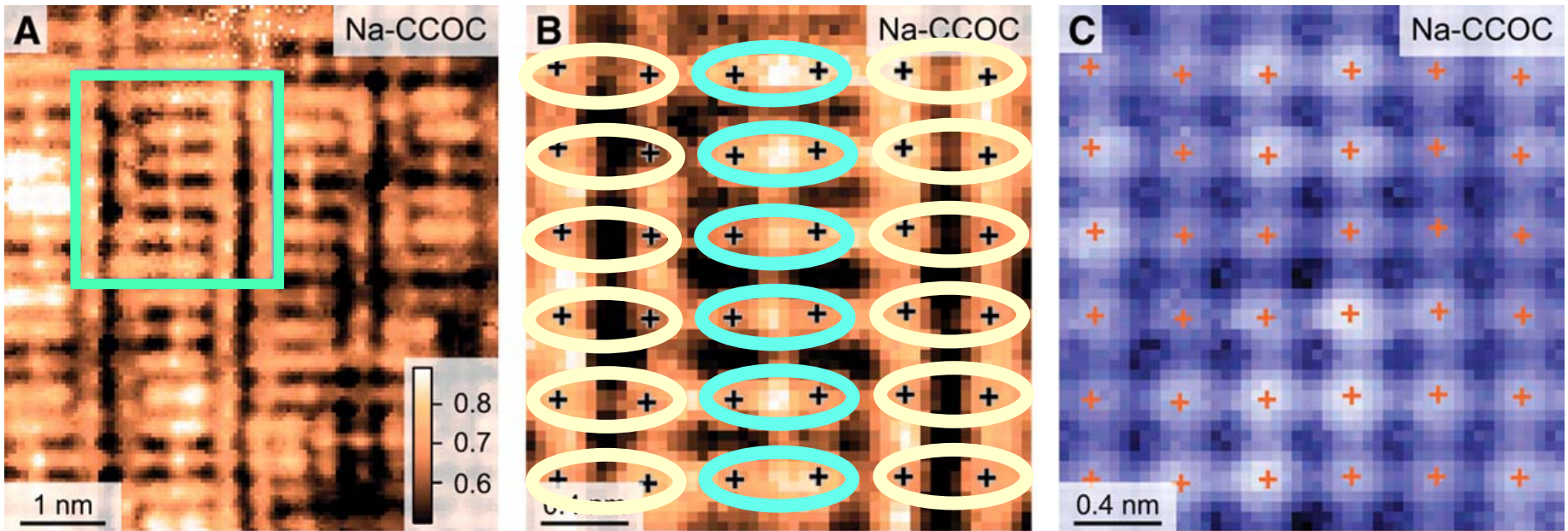
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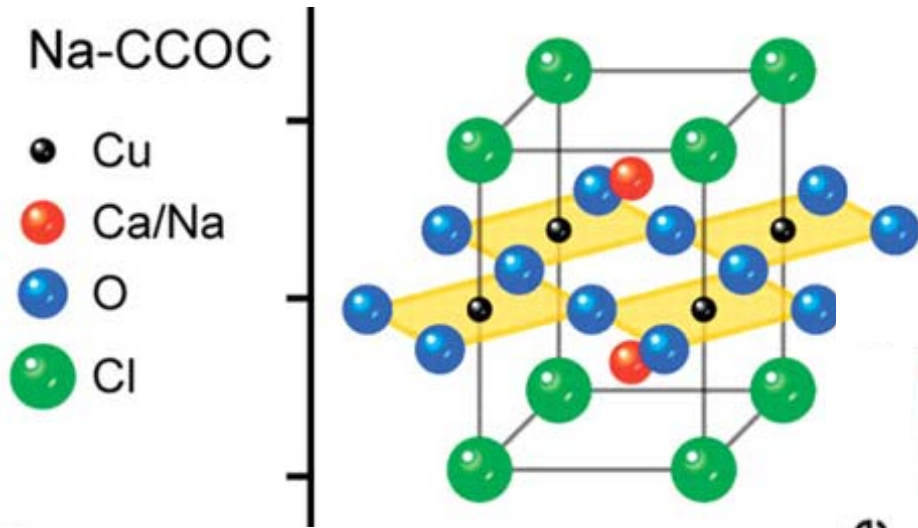
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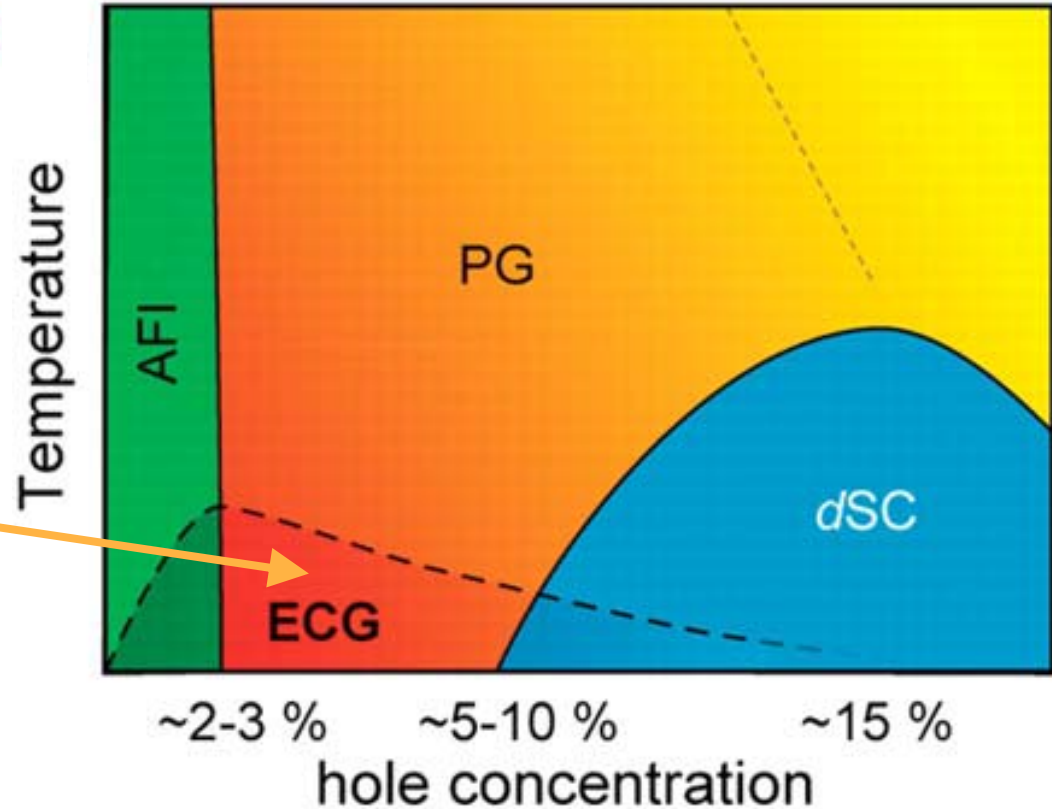
“Glassy” Valence Bond Solid (VBS)

Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)

Temperature-doping phase diagram of the cuprate superconductors



“Glassy” Valence Bond Solid (VBS)



Outline

Quantum phase transitions

1. Spin ordering in “Han purple”
2. Entanglement at the critical point: physical consequences at non-zero temperatures
 - (a) Double-layer antiferromagnet
 - (b) Superfluid-insulator transition
 - (c) Hydrodynamics via mapping to quantum theory of black holes.
3. Entanglement of valence bonds
4. Conclusions

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4. Conclusions

Conclusions

- Studies of new materials and trapped ultracold atoms are yielding new quantum phases, with novel forms of quantum entanglement.
- Some materials are of technological importance: *e.g.* high temperature superconductors.
- Real-world studies on the entanglement of large numbers of qubits: insights may be important for quantum cryptography and quantum computing.
- Tabletop “laboratories for the entire universe”: quantum mechanics of black holes, quark-gluon plasma, neutrons stars, and big-bang physics.