

From the SYK model,
to a theory of the strange metal,
and of quantum gravity in
two spacetime dimensions

ARO MURI review,
University of Maryland
October 13, 2017

Subir Sachdev

Talk online: sachdev.physics.harvard.edu



Gapless spin-fluid ground state in a random quantum Heisenberg magnet

[\[PDF\]](#) from [aps.org](#)

Authors **Subir Sachdev, Jinwu Ye**

Publication date **1993/5/24**

Journal **Physical review letters**

Volume **70**

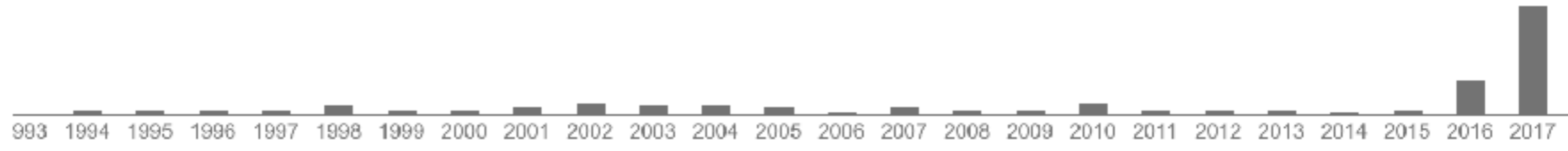
Issue **21**

Pages **3339**

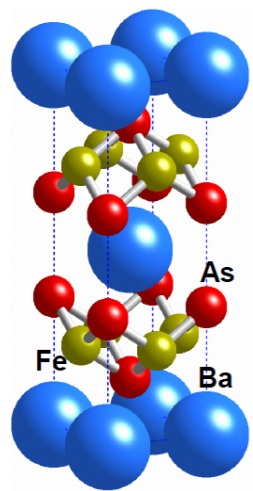
Publisher **American Physical Society**

Description **Abstract** We examine the spin-S quantum Heisenberg magnet with Gaussian-random, infinite-range exchange interactions. The quantum-disordered phase is accessed by generalizing to SU (M) symmetry and studying the large M limit. For large S the ground state is a spin glass, while quantum fluctuations produce a spin-fluid state for small S. The spin-fluid phase is found to be generically gapless—the average, zero temperature, local dynamic spin susceptibility obeys $\chi(\omega) \sim \ln(1/|\omega|) + i(\pi/2) \text{sgn}(\omega)$ at low frequencies.

Total citations [Cited by 320](#)

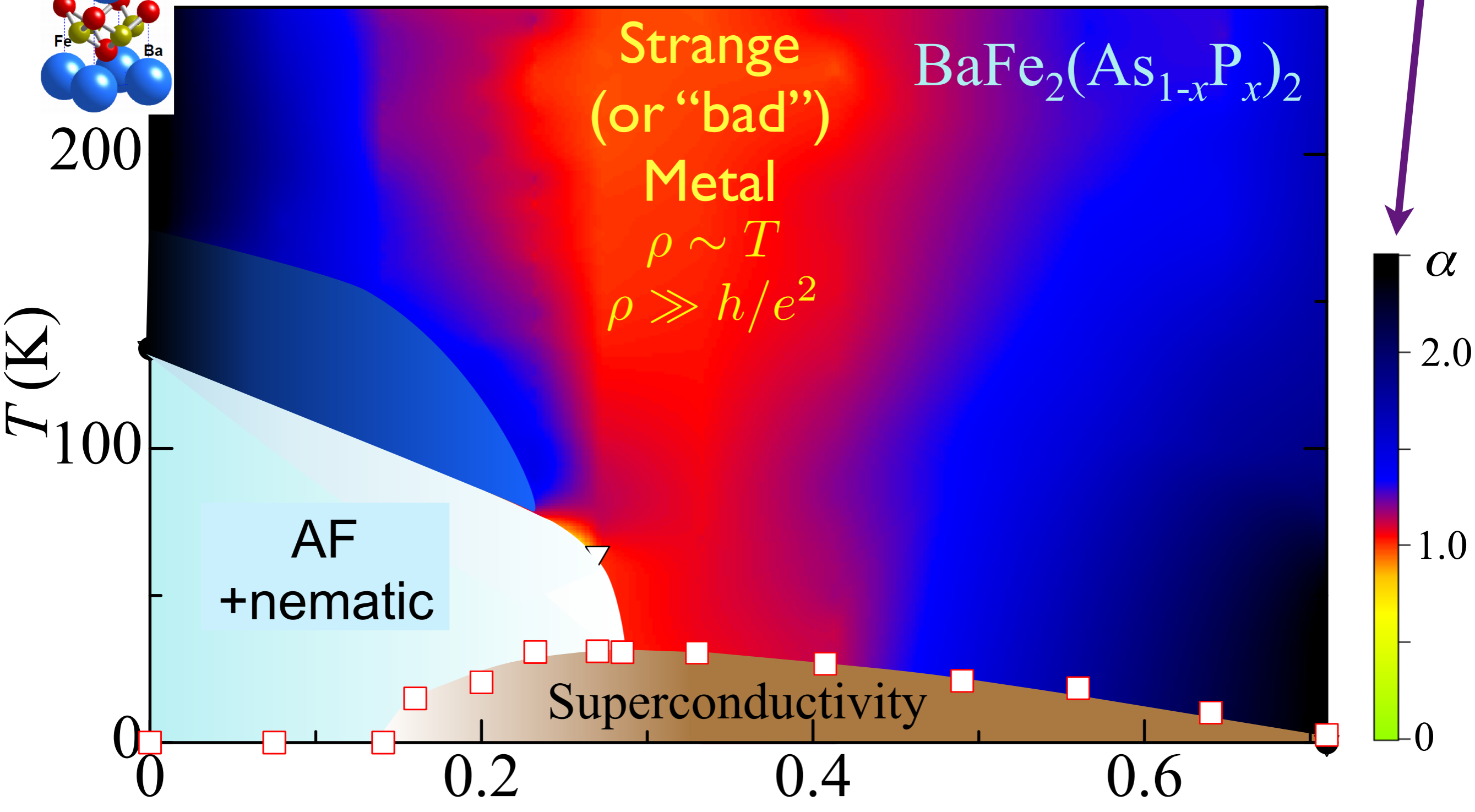


Scholar articles [Gapless spin-fluid ground state in a random quantum Heisenberg magnet](#)
S Sachdev, J Ye - Physical review letters, 1993
[Cited by 320](#) [Related articles](#) [All 11 versions](#)



Quantum matter without quasiparticles

Resistivity
 $\sim \rho_0 + AT^\alpha$

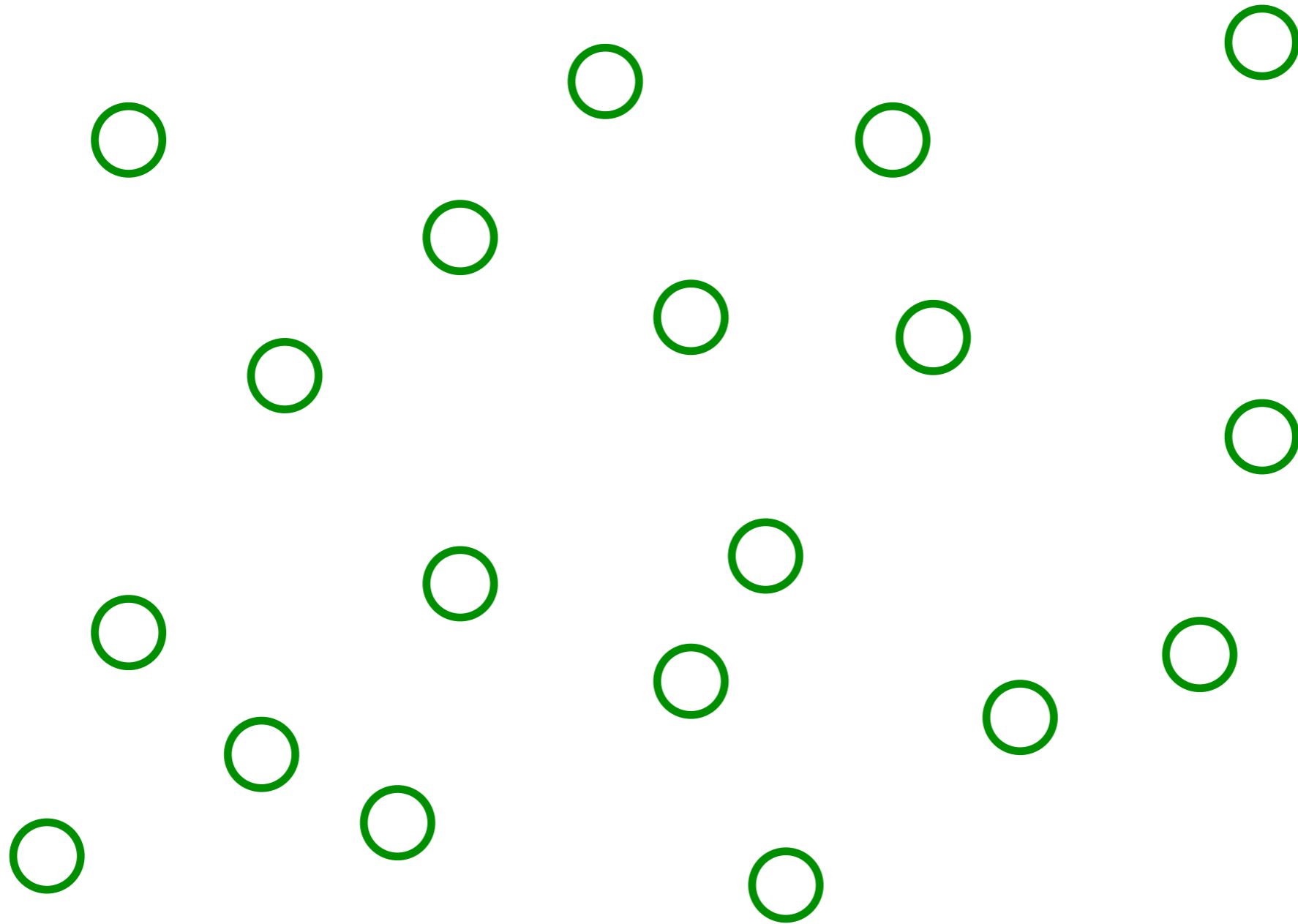


Superconductivity in Bad Metals

V. J. Emery and S. A. Kivelson
 Phys. Rev. Lett. **74**, 3253 – Published 17 April 1995

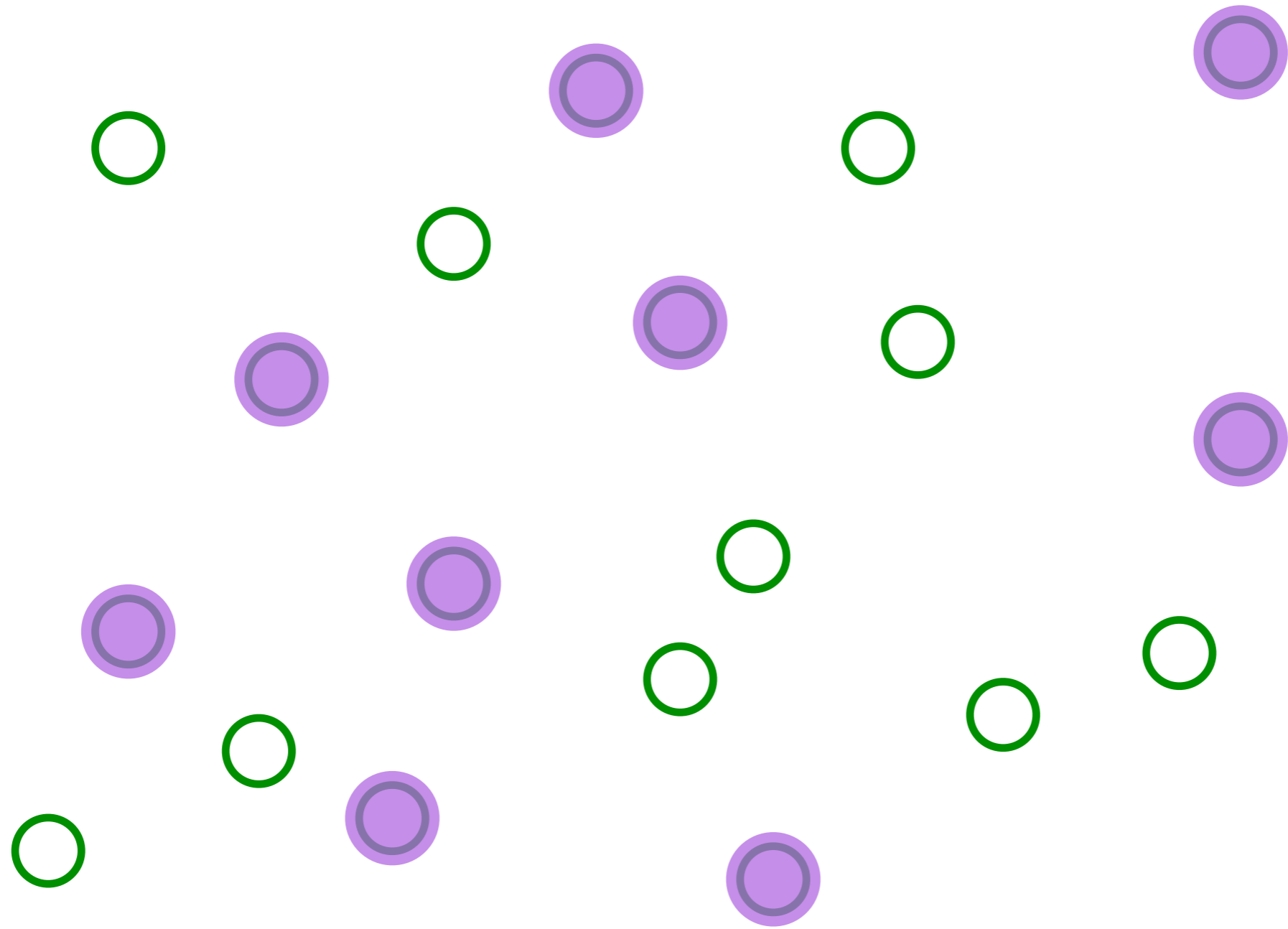
S. Kasahara, T. Shibauchi, K. Hashimoto, K. Ikada, S. Tonegawa, R. Okazaki, H. Shishido, H. Ikeda, H. Takeya, K. Hirata, T. Terashima, and Y. Matsuda, *PRB* **81**, 184519 (2010)

The Sachdev-Ye-Kitaev (SYK) model



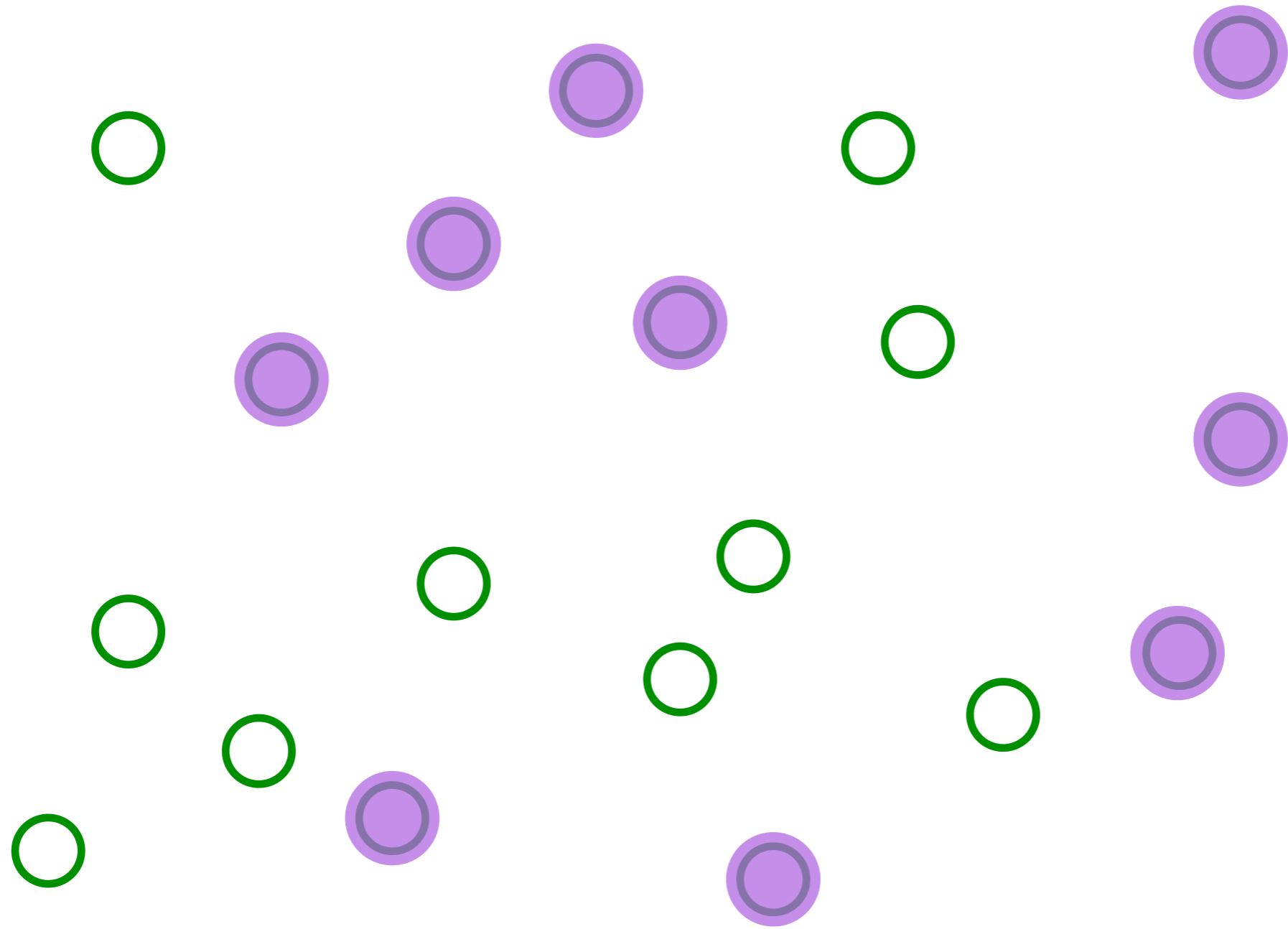
Pick a set of random sites/orbitals

The SYK model



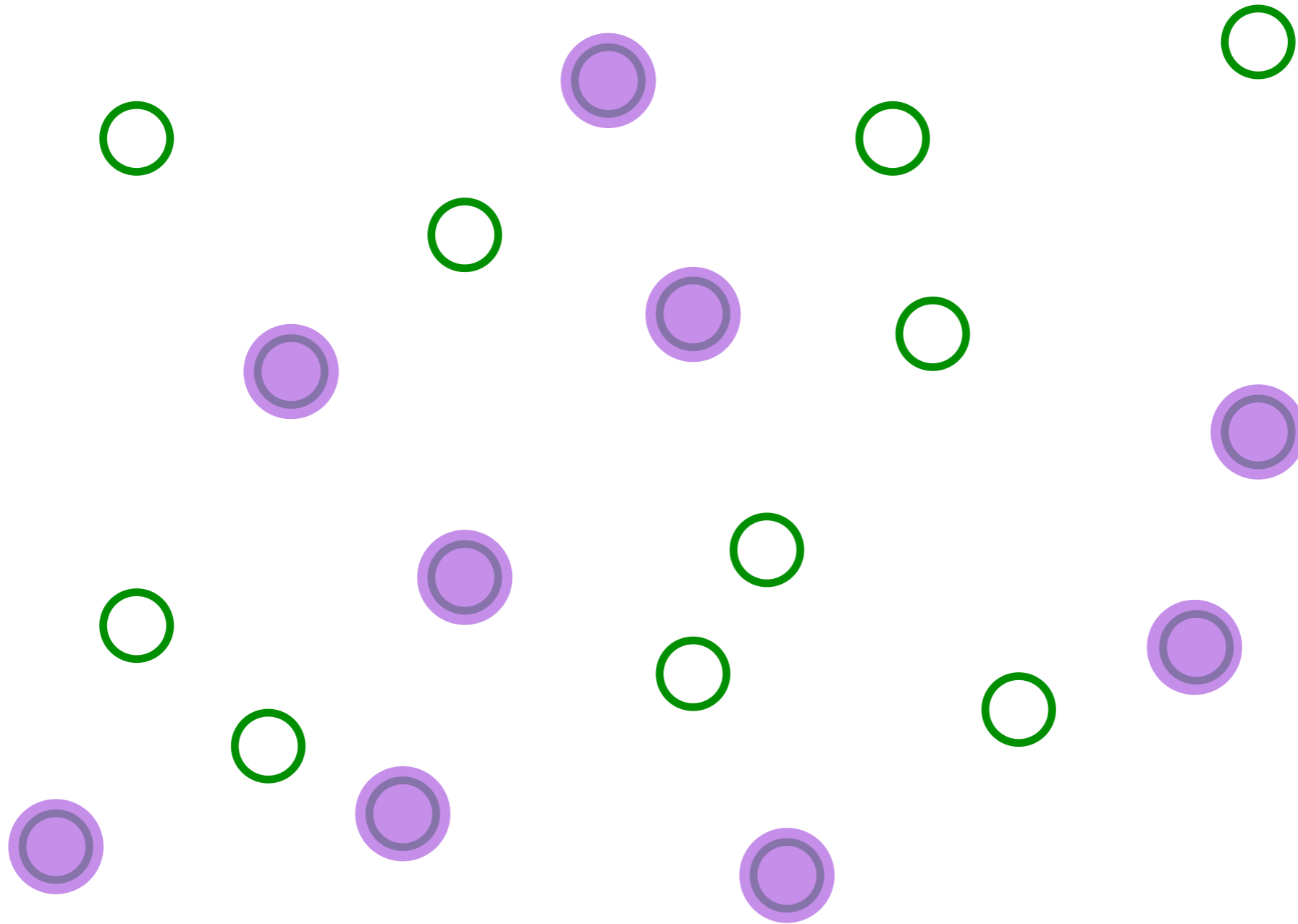
Place electrons randomly on some sites/orbitals

The SYK model



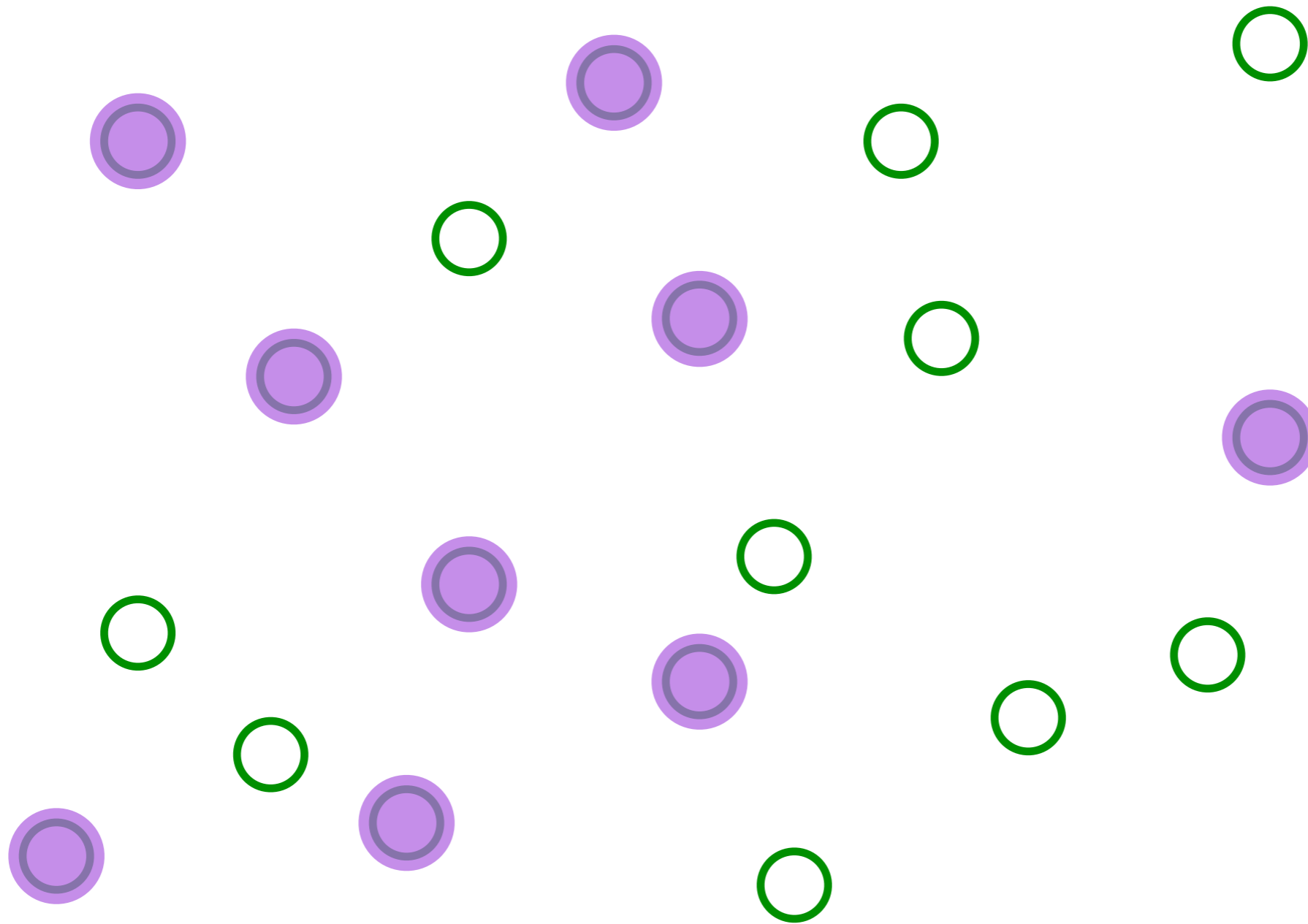
Entangle electrons pairwise randomly

The SYK model



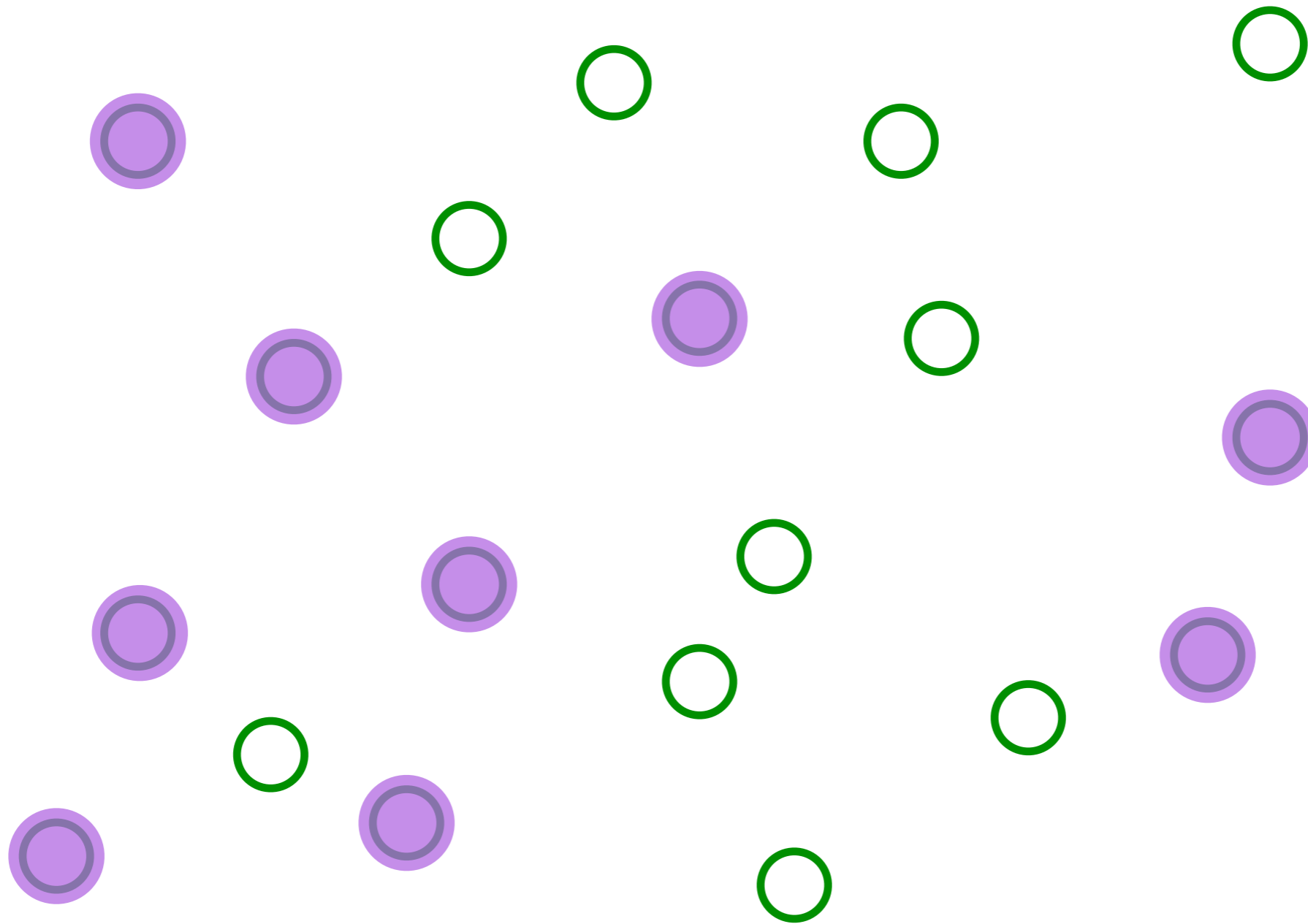
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The SYK model



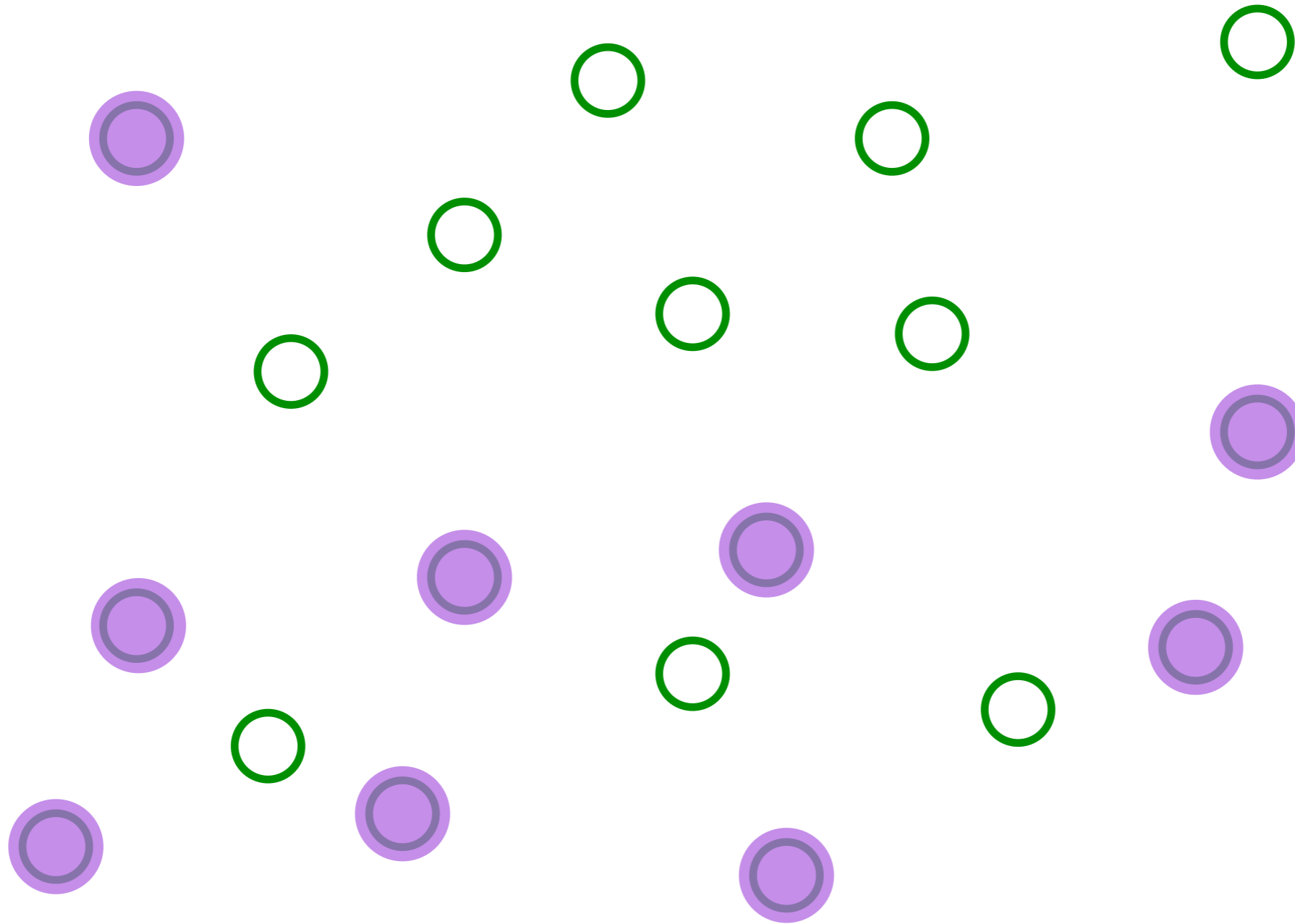
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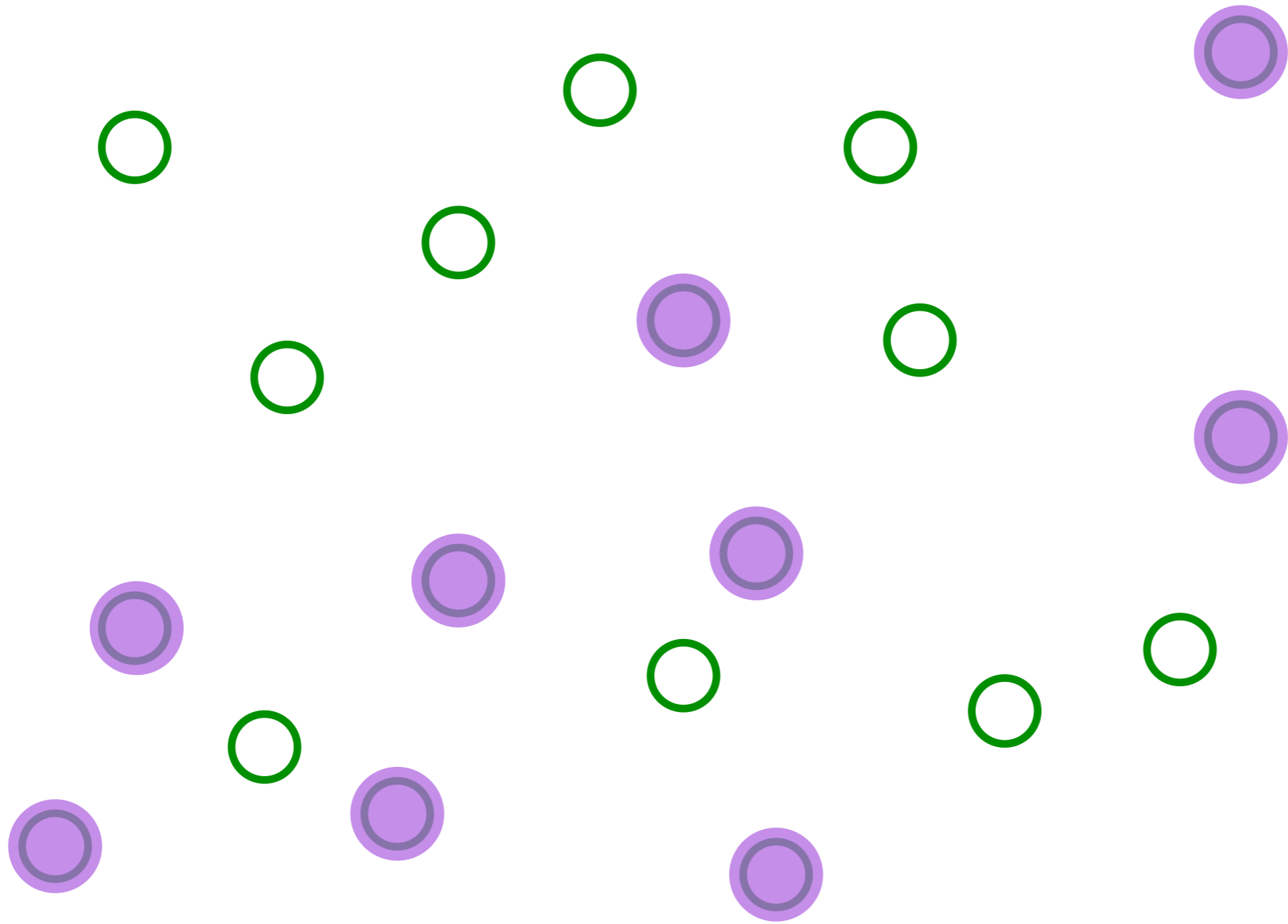
Entangle electrons pairwise randomly

The SYK model



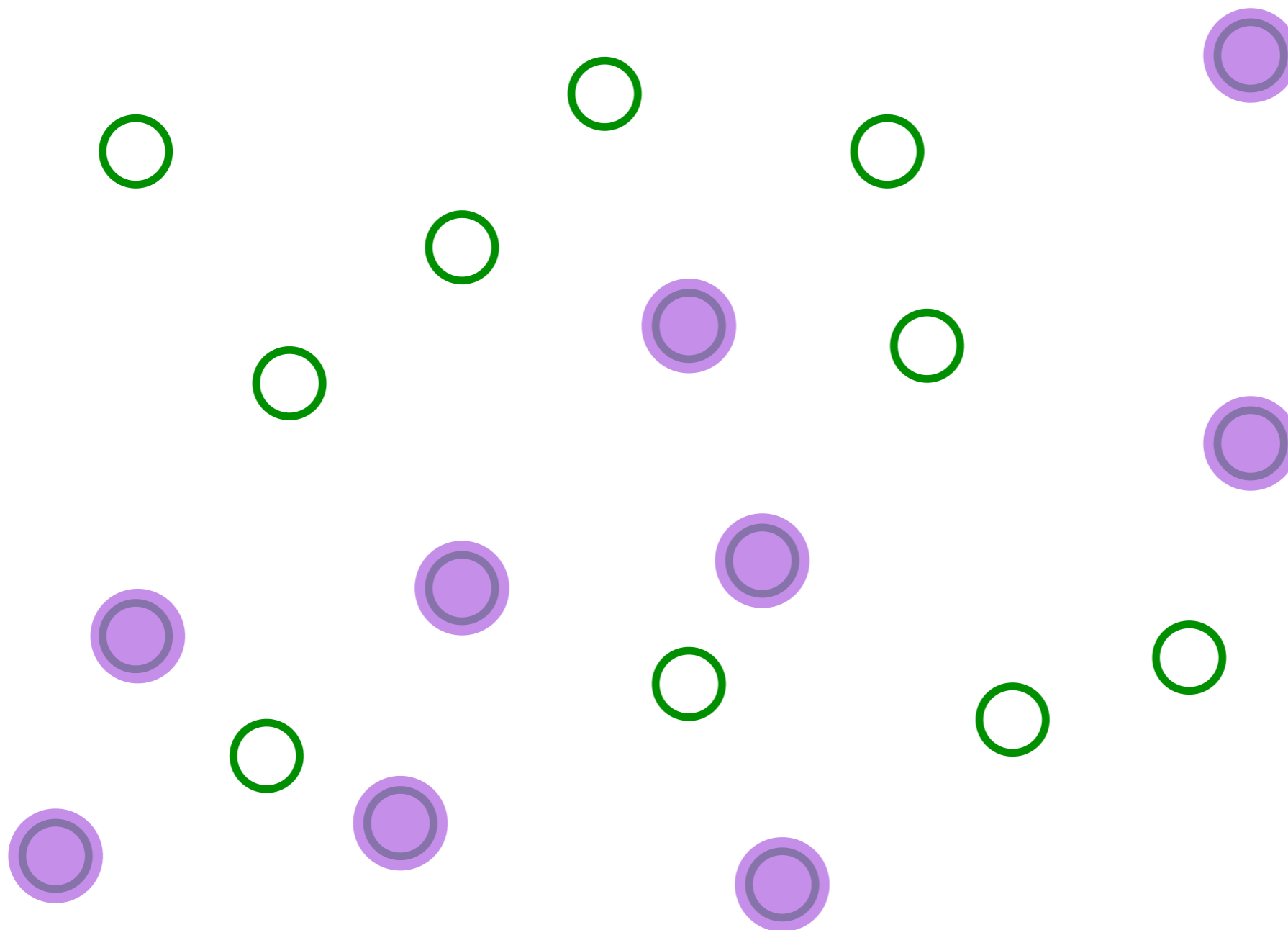
Entangle electrons pairwise randomly

The SYK model



Entangle electrons pairwise randomly

The SYK model



This describes both a strange metal and a black hole!

The SYK model

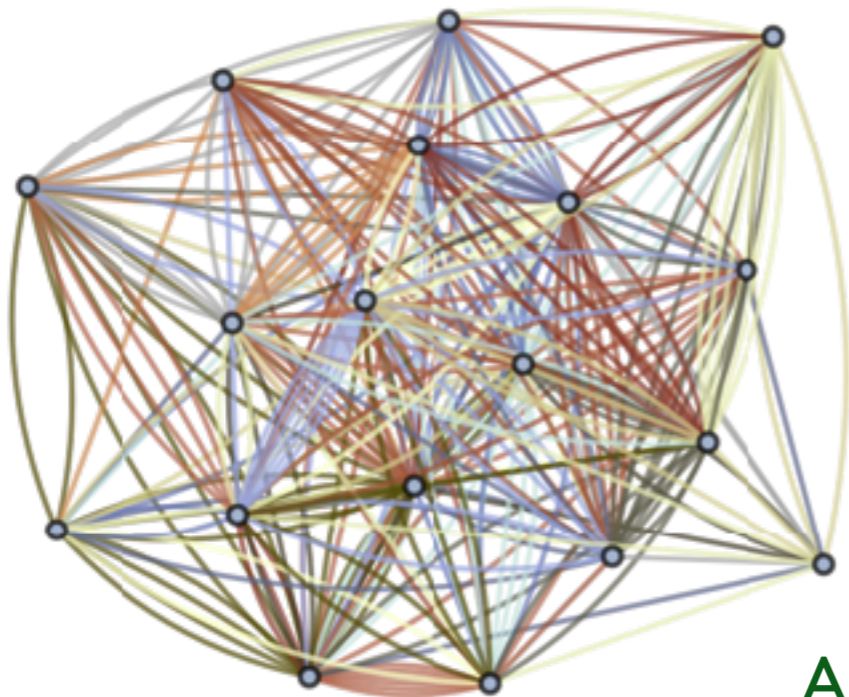
(See also: the “2-Body Random Ensemble” in nuclear physics; did not obtain the large N limit; T.A. Brody, J. Flores, J.B. French, P.A. Mello, A. Pandey, and S.S.M. Wong, Rev. Mod. Phys. **53**, 385 (1981))

$$H = \frac{1}{(2N)^{3/2}} \sum_{i,j,k,\ell=1}^N U_{ij;k\ell} c_i^\dagger c_j^\dagger c_k c_\ell - \mu \sum_i c_i^\dagger c_i$$

$$c_i c_j + c_j c_i = 0 \quad , \quad c_i c_j^\dagger + c_j^\dagger c_i = \delta_{ij}$$

$$Q = \frac{1}{N} \sum_i c_i^\dagger c_i$$

$U_{ij;k\ell}$ are independent random variables with $\overline{U_{ij;k\ell}} = 0$ and $\overline{|U_{ij;k\ell}|^2} = U^2$
 $N \rightarrow \infty$ yields critical strange metal.



S. Sachdev and J. Ye, PRL **70**, 3339 (1993)

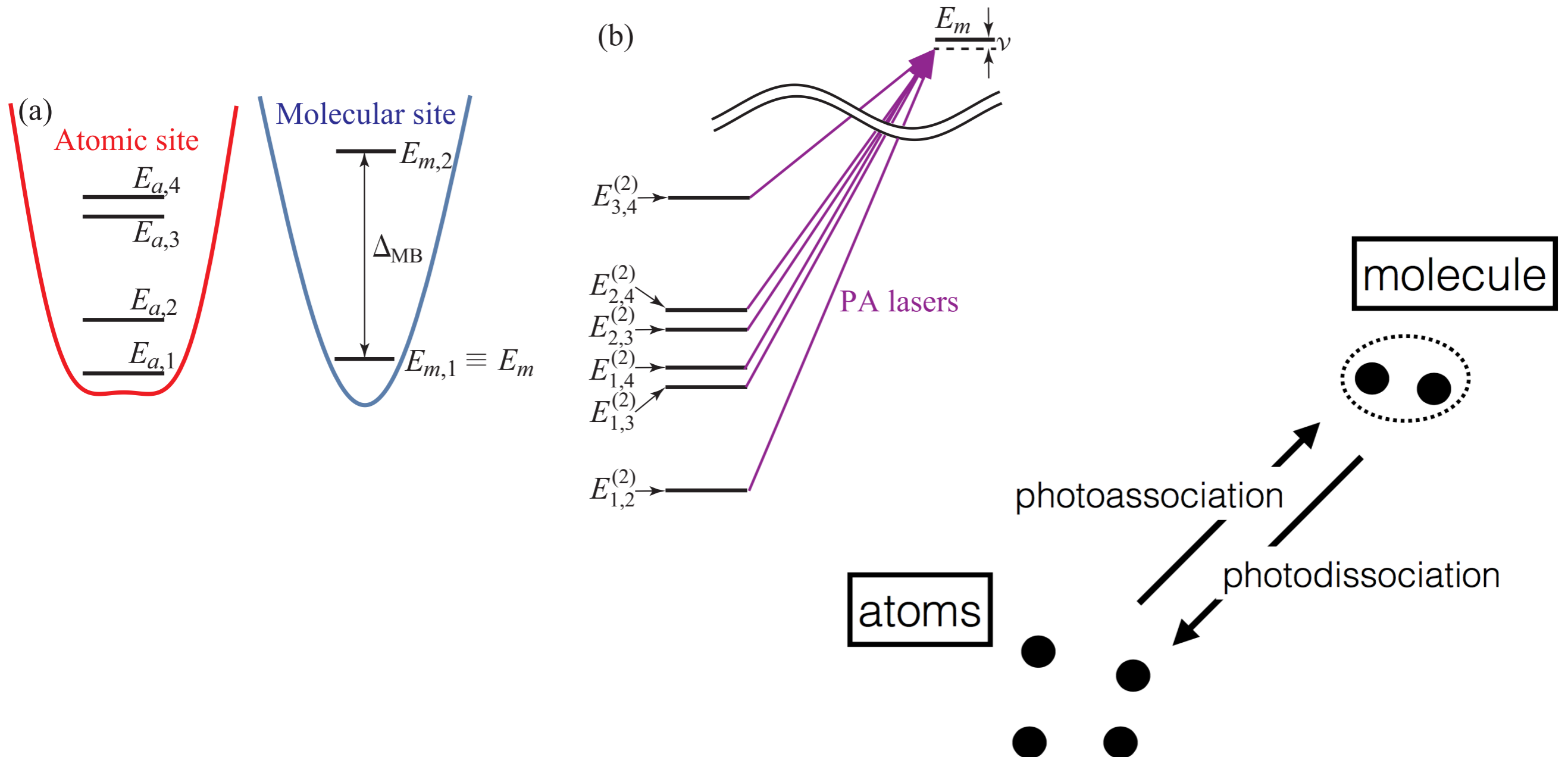
A. Kitaev, unpublished; S. Sachdev, PRX **5**, 041025 (2015)

Creating and probing the Sachdev-Ye-Kitaev model with ultracold gases:

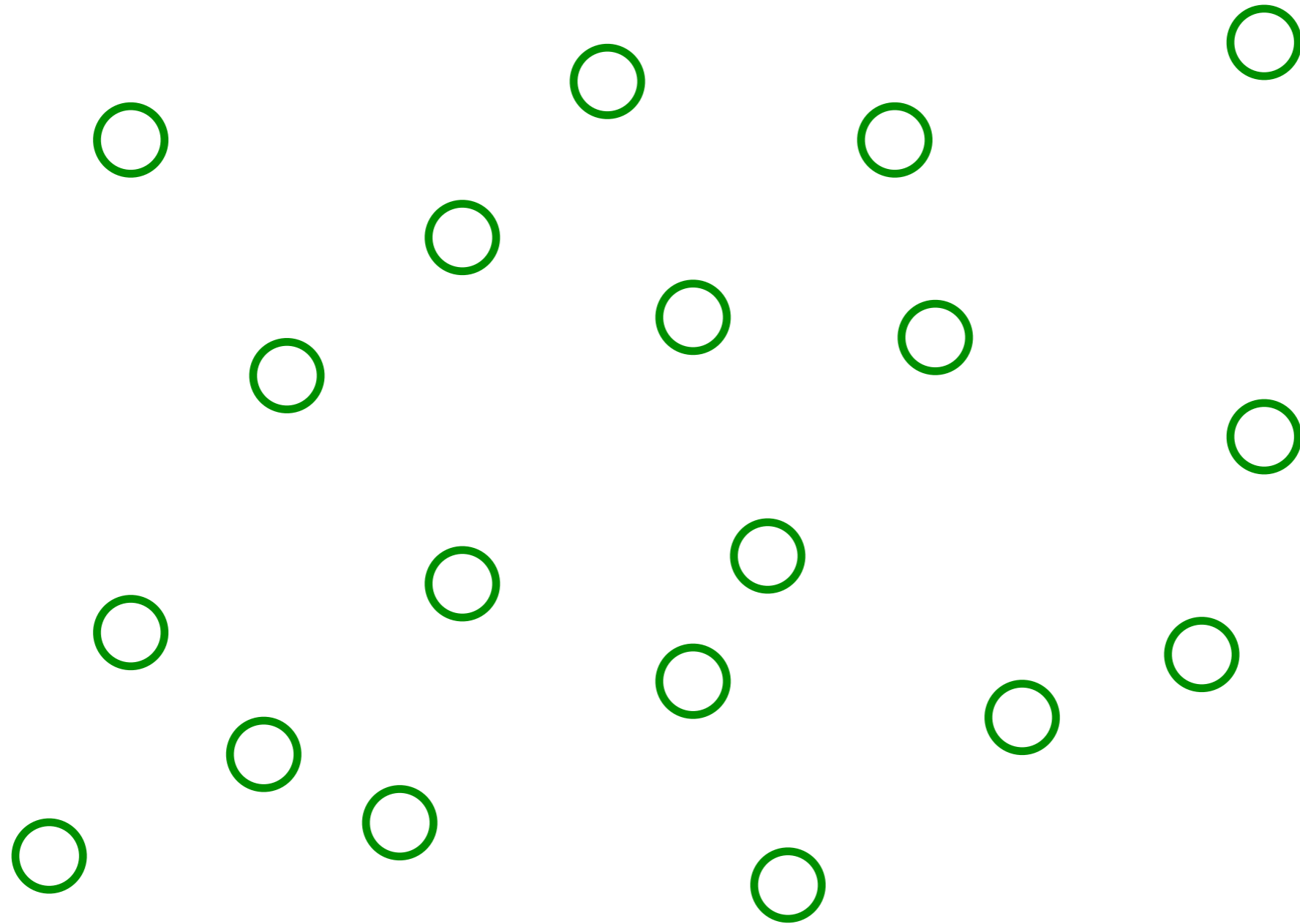
Towards experimental studies of quantum gravity

Prog.Theor. Exp. Phys. 2017, 083101

Ippei Danshita^{1,*}, Masanori Hanada^{1,2,3}, and Masaki Tezuka⁴

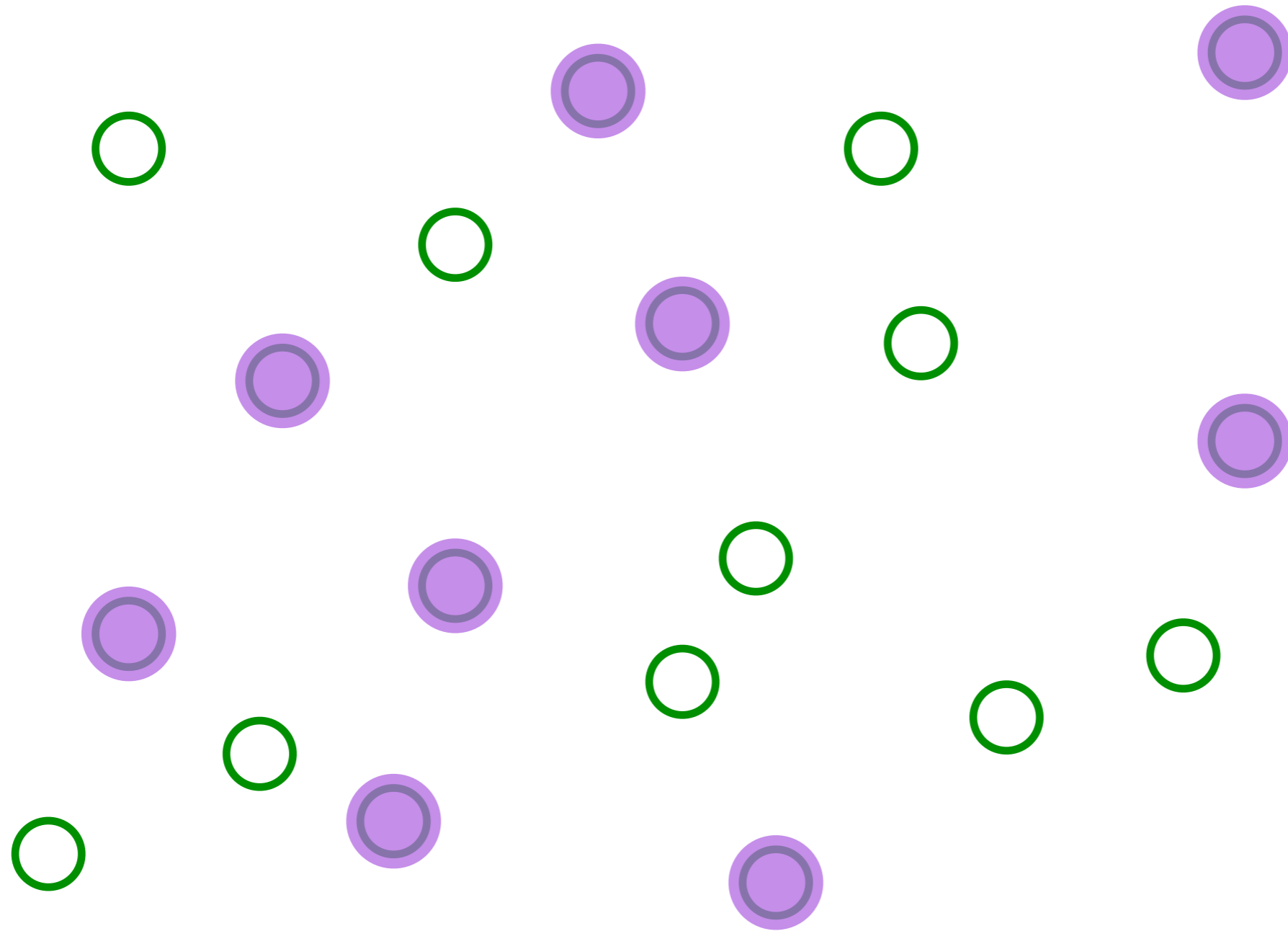


A simple model of a metal with quasiparticles



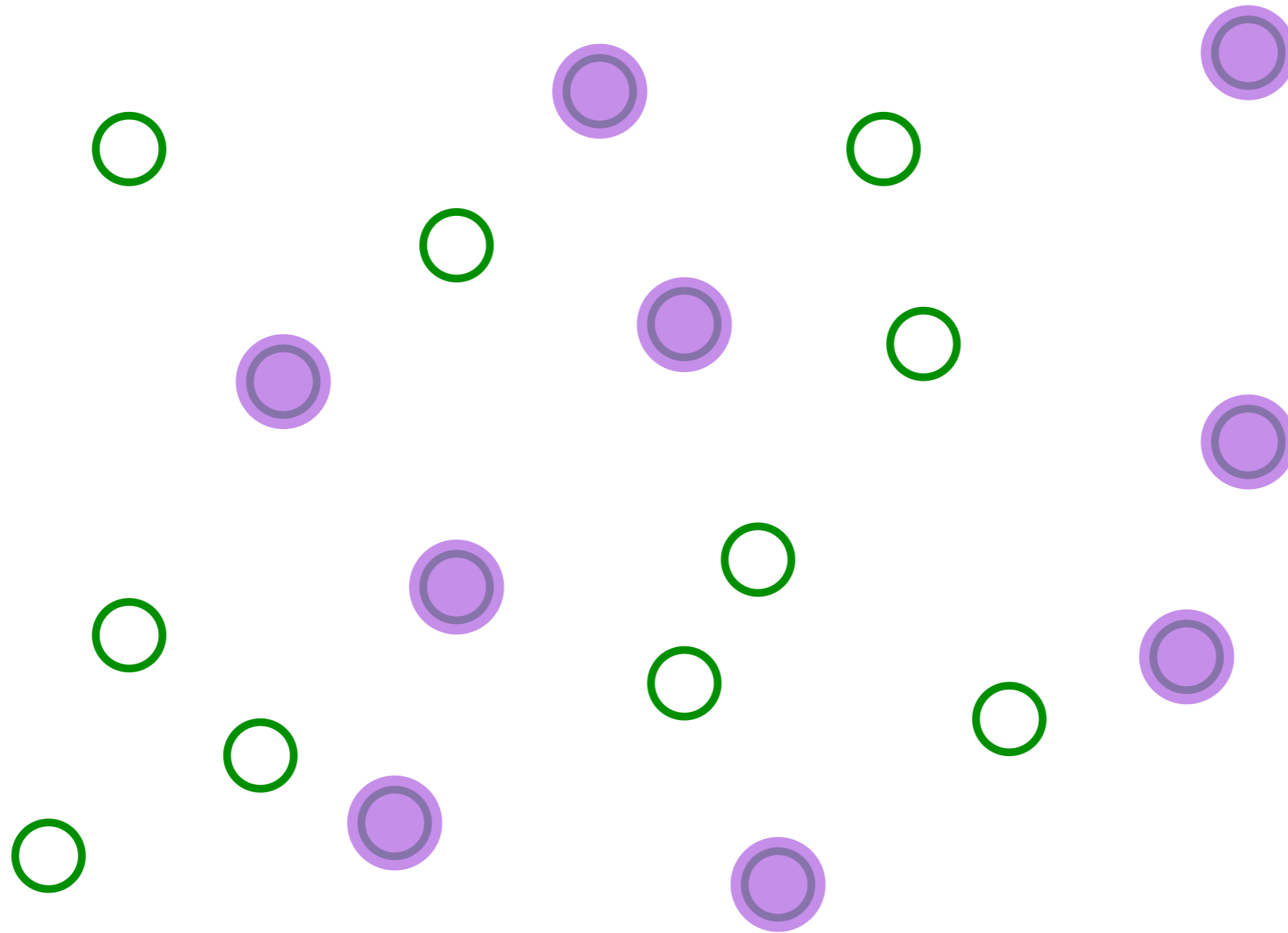
Pick a set of random positions

A simple model of a metal with quasiparticles



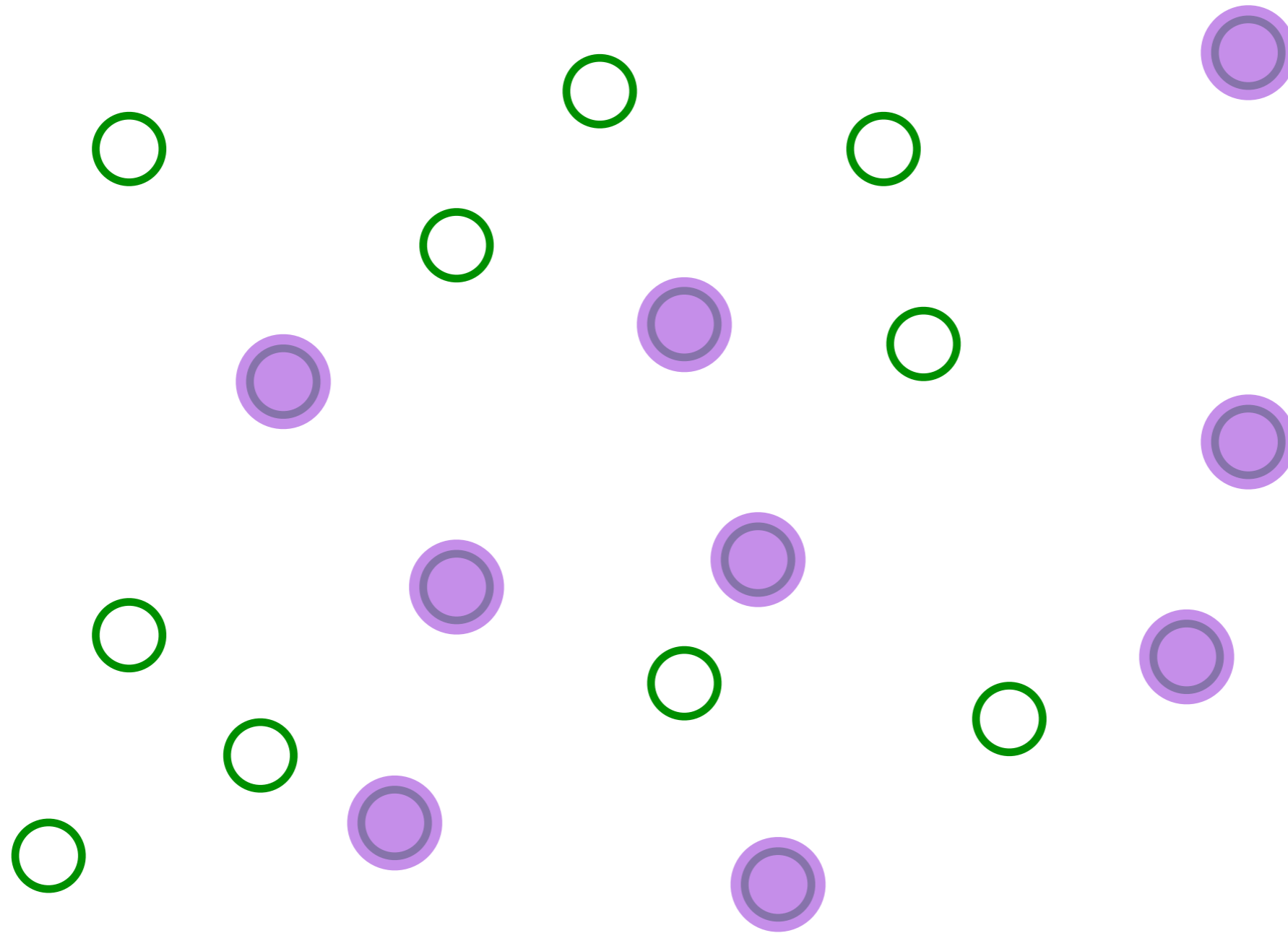
Place electrons randomly on some sites

A simple model of a metal with quasiparticles



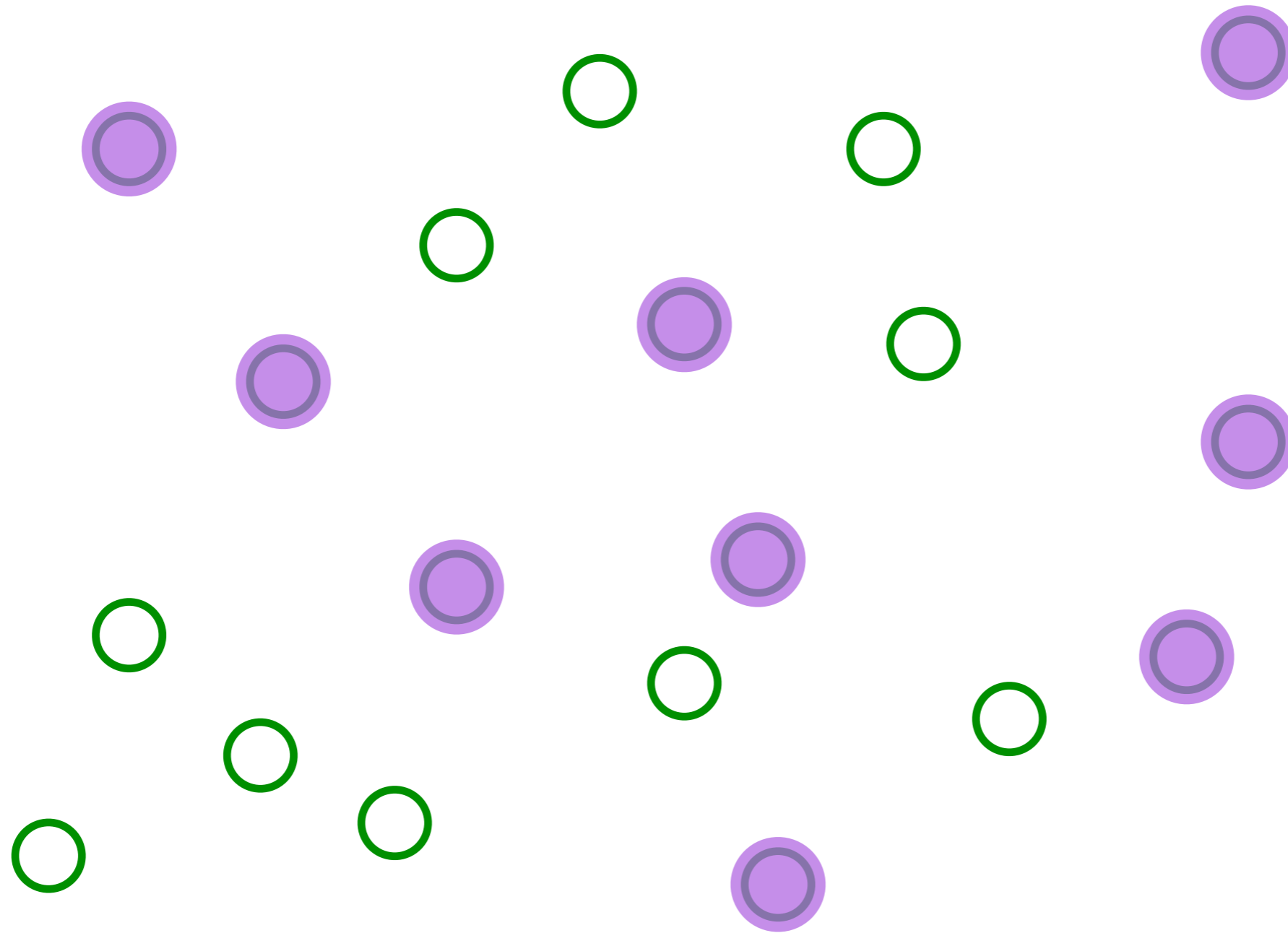
Electrons move one-by-one randomly

A simple model of a metal with quasiparticles



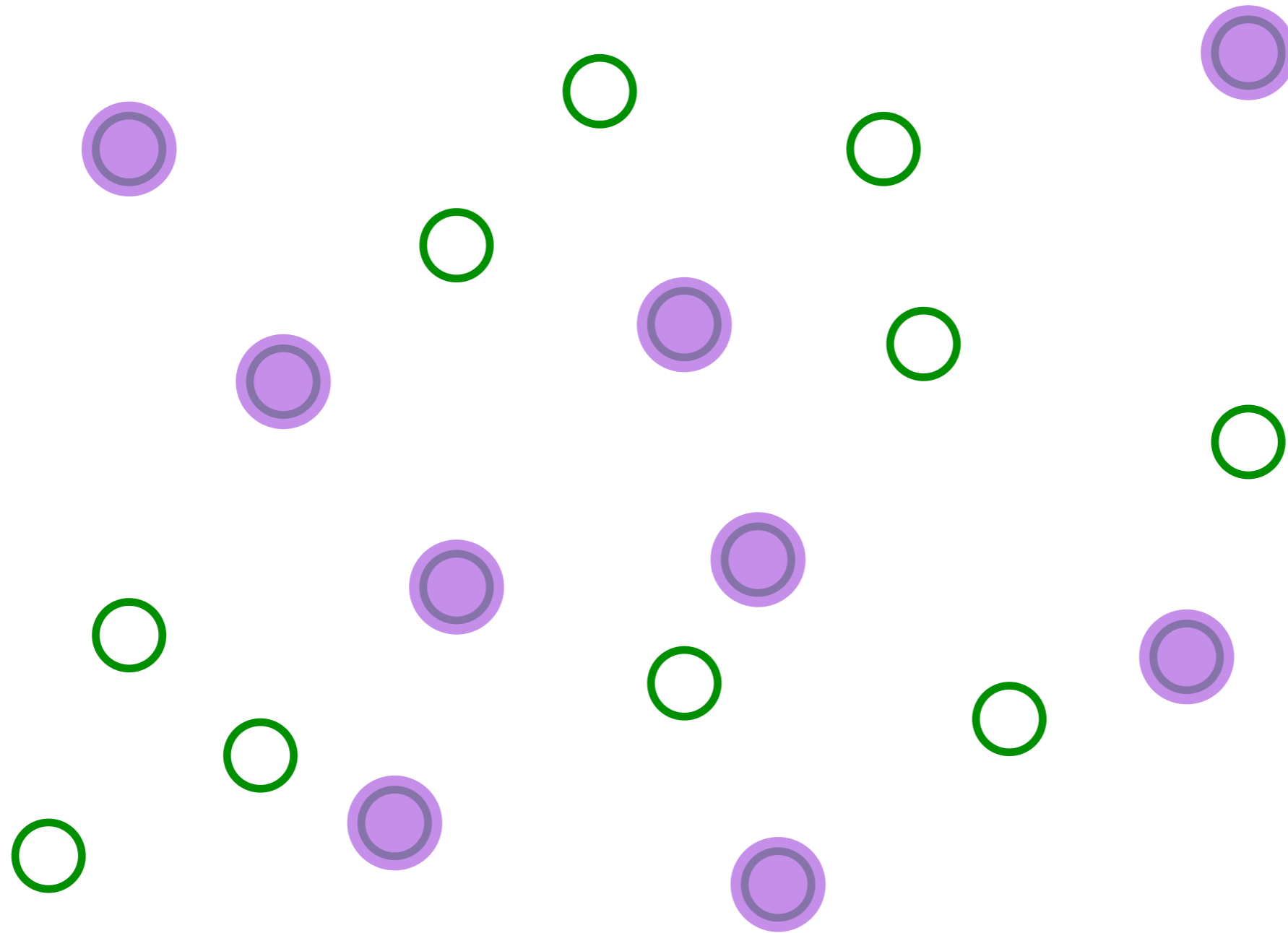
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Electrons move one-by-one randomly

A simple model of a metal with quasiparticles



Electrons move one-by-one randomly

A simple model of a metal with quasiparticles

$$H = \frac{1}{(N)^{1/2}} \sum_{i,j=1}^N t_{ij} c_i^\dagger c_j + \dots$$

$$c_i c_j + c_j c_i = 0 \quad , \quad c_i c_j^\dagger + c_j^\dagger c_i = \delta_{ij}$$

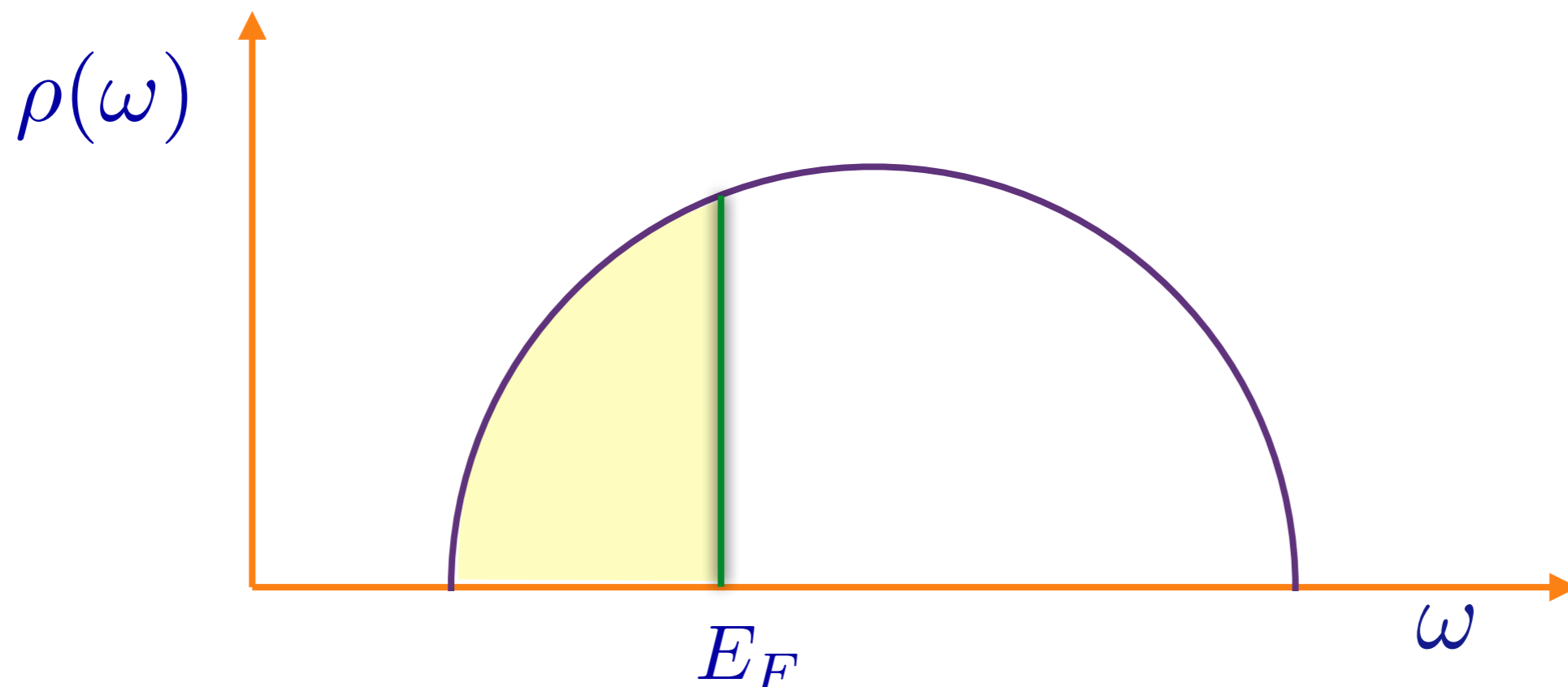
$$\frac{1}{N} \sum_i c_i^\dagger c_i = Q$$

t_{ij} are independent random variables with $\overline{t_{ij}} = 0$ and $\overline{|t_{ij}|^2} = t^2$

**Fermions occupying the eigenstates of a
 $N \times N$ random matrix**

A simple model of a metal with quasiparticles

Let ε_α be the eigenvalues of the matrix t_{ij}/\sqrt{N} . The fermions will occupy the lowest NQ eigenvalues, upto the Fermi energy E_F . The density of states is $\rho(\omega) = (1/N) \sum_\alpha \delta(\omega - \varepsilon_\alpha)$.



A simple model of a metal with quasiparticles



Many-body
level spacing
 $\sim 2^{-N}$

Quasiparticle
excitations with
spacing $\sim 1/N$

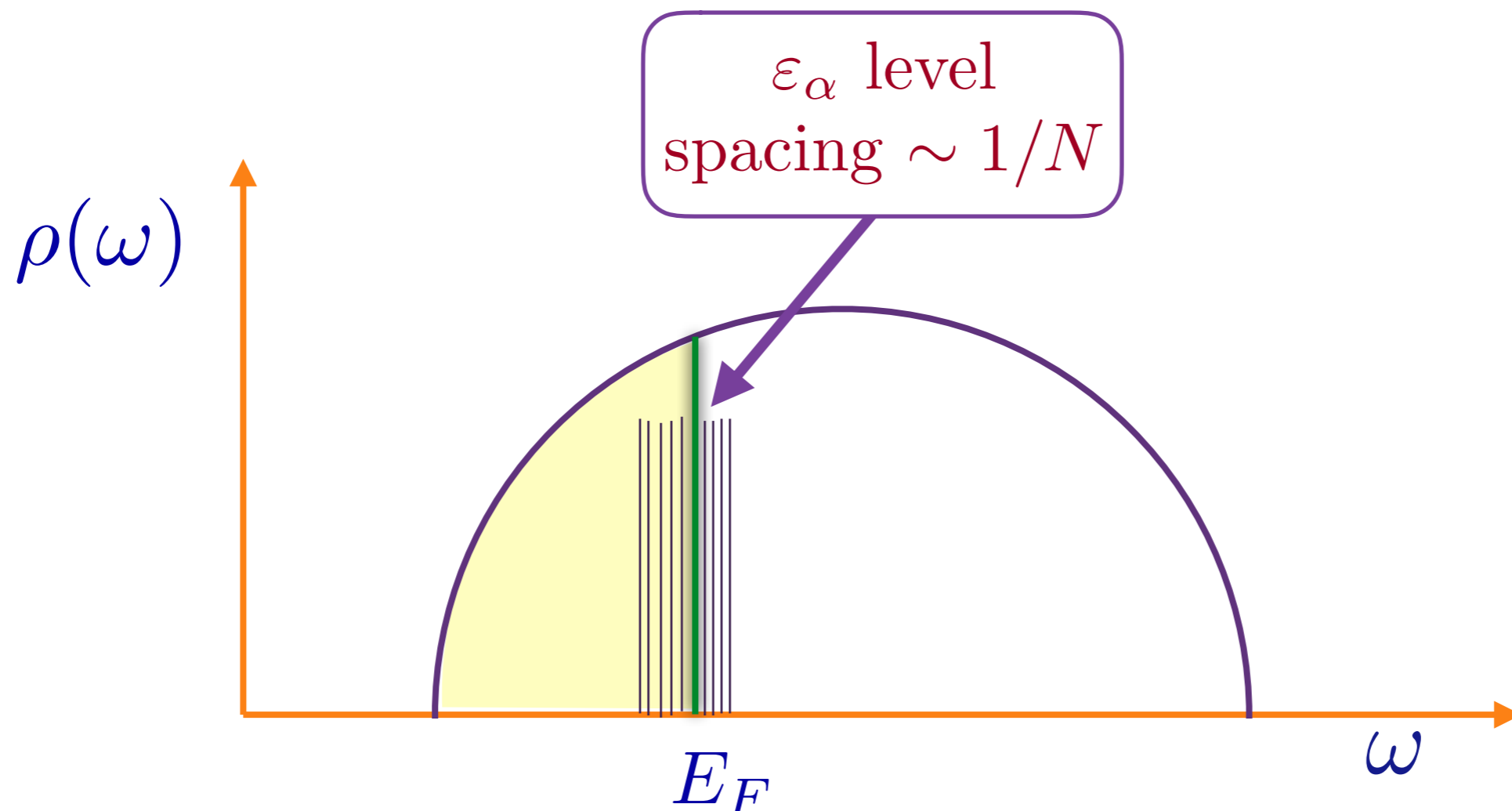
There are 2^N many
body levels with energy

$$E = \sum_{\alpha=1}^N n_{\alpha} \varepsilon_{\alpha},$$

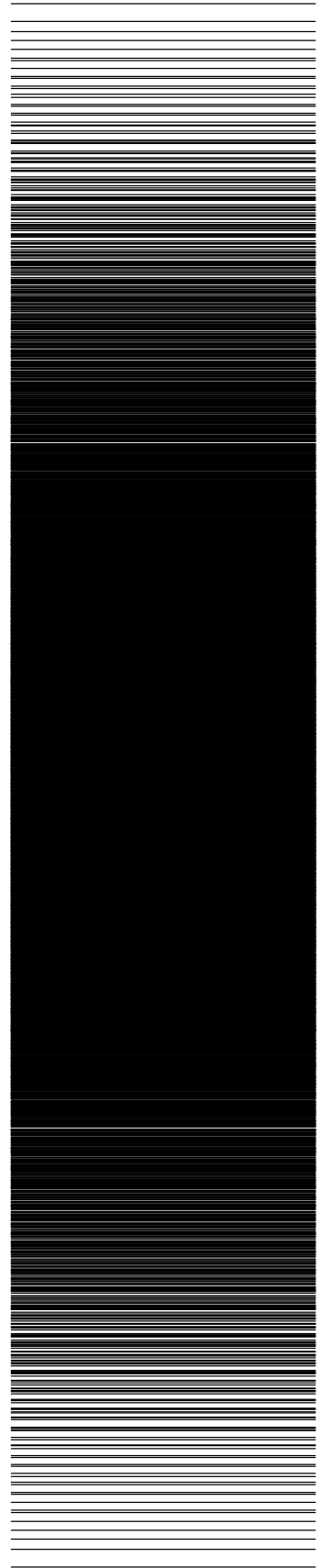
where $n_{\alpha} = 0, 1$. Shown
are all values of E for a
single cluster of size
 $N = 12$. The ε_{α} have a
level spacing $\sim 1/N$.

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The SYK model

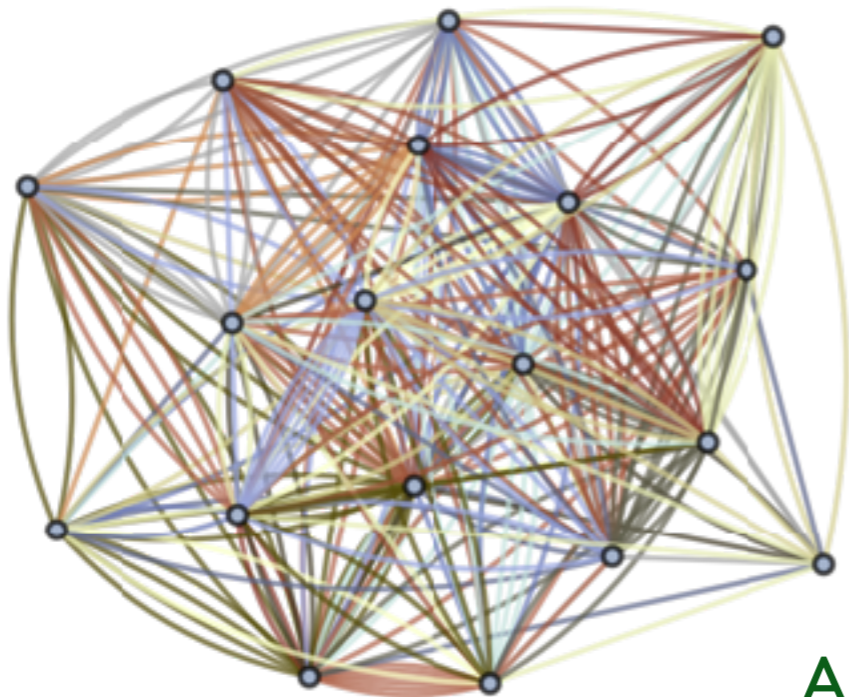
(See also: the “2-Body Random Ensemble” in nuclear physics; did not obtain the large N limit; T.A. Brody, J. Flores, J.B. French, P.A. Mello, A. Pandey, and S.S.M. Wong, Rev. Mod. Phys. **53**, 385 (1981))

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S. Sachdev and J. Ye, PRL **70**, 3339 (1993)

A. Kitaev, unpublished; S. Sachdev, PRX **5**, 041025 (2015)

The SYK model

There are 2^N many body levels with energy E , which do not admit a quasiparticle decomposition. Shown are all values of E for a single cluster of size $N = 12$. The $T \rightarrow 0$ state has an entropy $S_{GPS} = N s_0$ with

$$s_0 = \frac{G}{\pi} + \frac{\ln(2)}{4} = 0.464848\dots$$
$$< \ln 2$$

where G is Catalan's constant, for the half-filled case $Q = 1/2$.

GPS: A. Georges, O. Parcollet, and S. Sachdev, PRB **63**, 134406 (2001)

Many-body level spacing $\sim 2^{-N} = e^{-N \ln 2}$

Non-quasiparticle excitations with spacing $\sim e^{-N s_0}$

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Non-quasiparticle excitations with spacing $\sim e^{-N s_0}$

No quasiparticles !

$$E \neq \sum_{\alpha} n_{\alpha} \varepsilon_{\alpha} + \sum_{\alpha, \beta} F_{\alpha\beta} n_{\alpha} n_{\beta} + \dots$$

PRB **63**, 134406 (2001)

The SYK model

- Low energy, many-body density of states

$$\rho(E) \sim e^{Ns_0} \sinh(\sqrt{2(E - E_0)N\gamma})$$

D. Stanford and E. Witten, 1703.04612

A. M. Garcia-Garcia, J.J.M. Verbaarschot, 1701.06593

D. Bagrets, A. Altland, and A. Kamenev, 1607.00694

A. Georges, O. Parcollet, and S. Sachdev, PRB **63**, 134406 (2001)

- Low temperature entropy $S = Ns_0 + N\gamma T + \dots$

A. Kitaev, unpublished

J. Maldacena and D. Stanford, 1604.07818

- $T = 0$ fermion Green's function $G(\tau) \sim \tau^{-1/2}$ at large τ . (Fermi liquids with quasiparticles have $G(\tau) \sim 1/\tau$)

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- $T > 0$ Green's function has conformal invariance

$$G \sim (T / \sin(\pi k_B T \tau / \hbar))^{1/2}$$

A. Georges and O. Parcollet PRB **59**, 5341 (1999)

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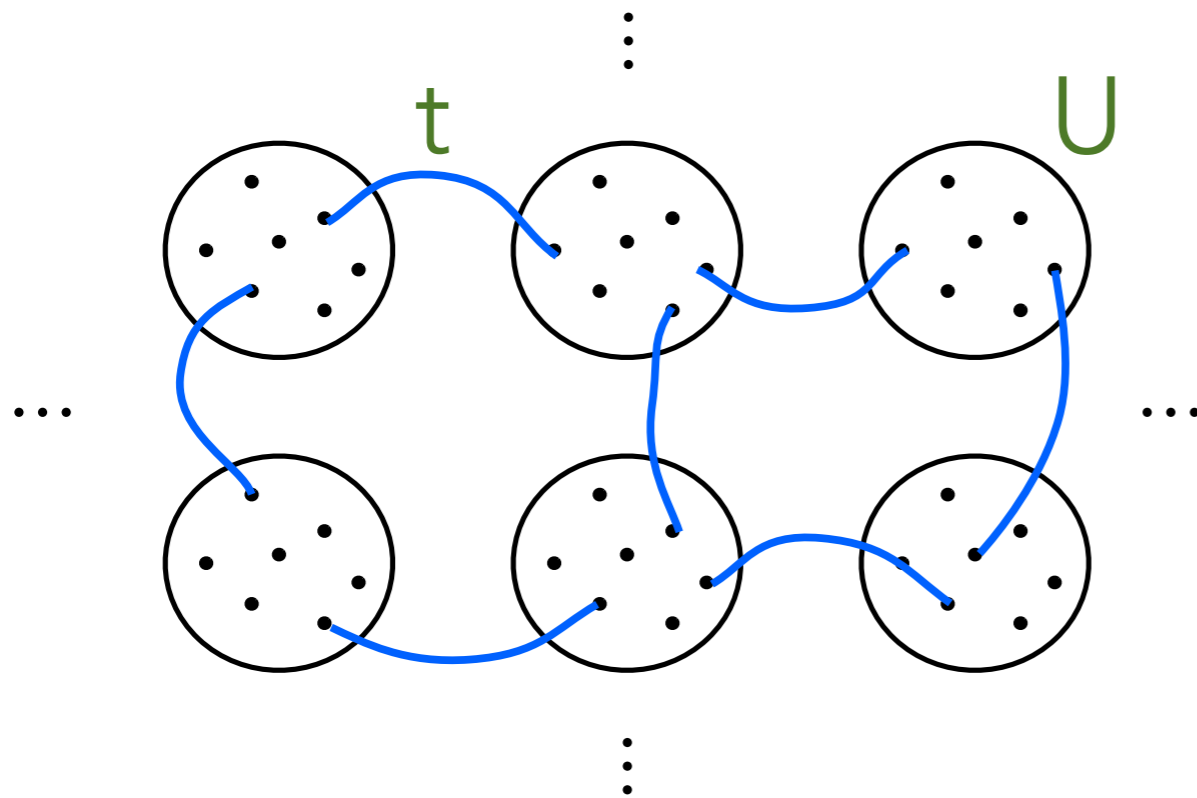
A. Georges and O. Parcollet PRB **59**, 5341 (1999)

- Study of non-equilibrium quench dynamics shows $\tau_{\text{eq}} \sim \hbar / (k_B T)$.

A. Eberlein, V. Kasper, S. Sachdev, and J. Steinberg, arXiv:1706.07803

[arXiv:1705.00117](https://arxiv.org/abs/1705.00117)

Title: A strongly correlated metal built from Sachdev-Ye-Kitaev models

Authors: [Xue-Yang Song](#), [Chao-Ming Jian](#), [Leon Balents](#)

$$H = \sum_x \sum_{i < j, k < l} U_{ijkl,x} c_{ix}^\dagger c_{jx}^\dagger c_{kx} c_{lx} + \sum_{\langle xx' \rangle} \sum_{i,j} t_{ij,xx'} c_{i,x}^\dagger c_{j,x'}$$

$$\overline{|U_{ijkl}|^2} = \frac{2U^2}{N^3}$$

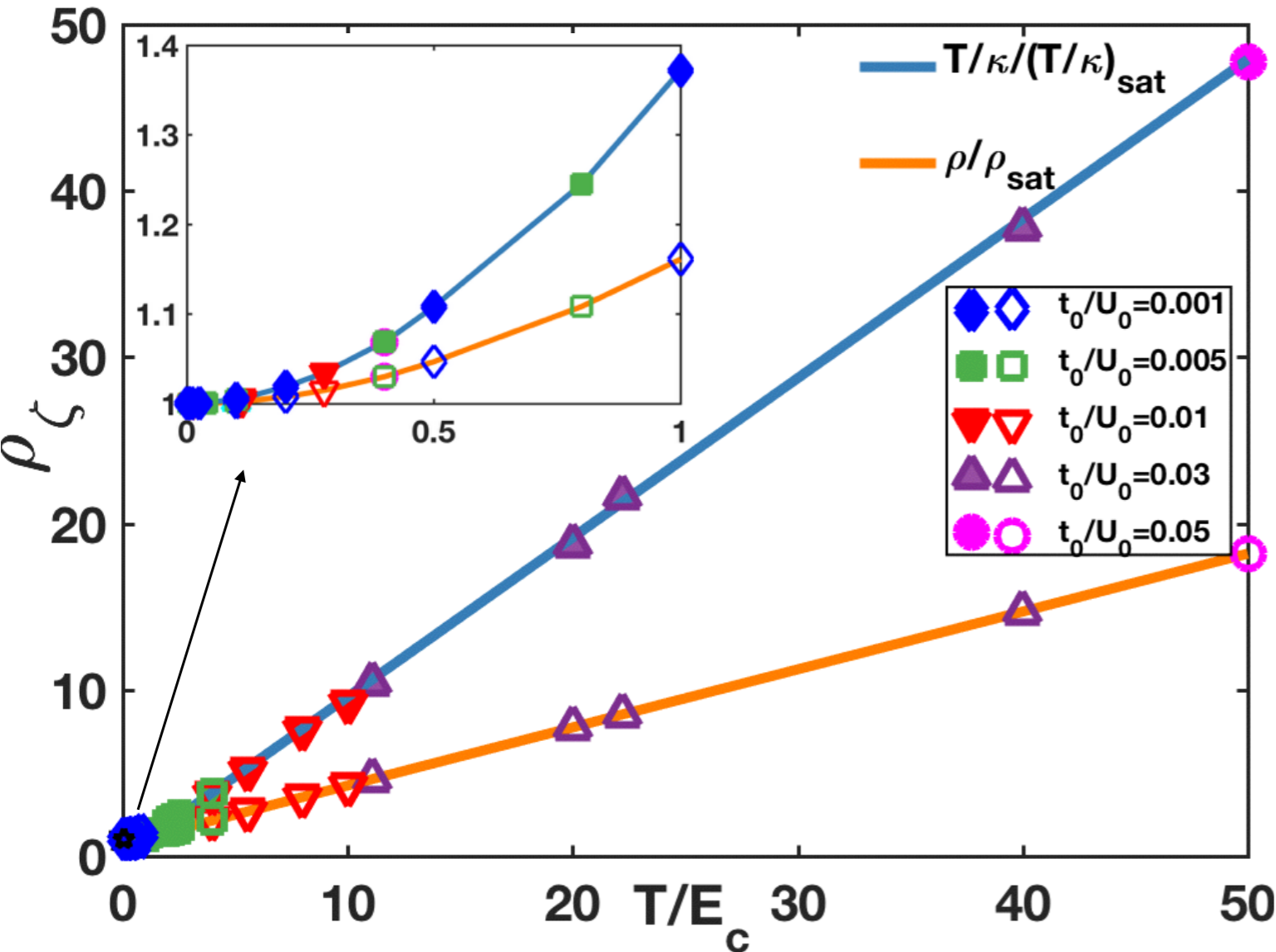
$$\overline{|t_{ij,xx'}|^2} = t_0^2/N$$

[arXiv:1705.00117](https://arxiv.org/abs/1705.00117)

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Low ‘coherence’ scale



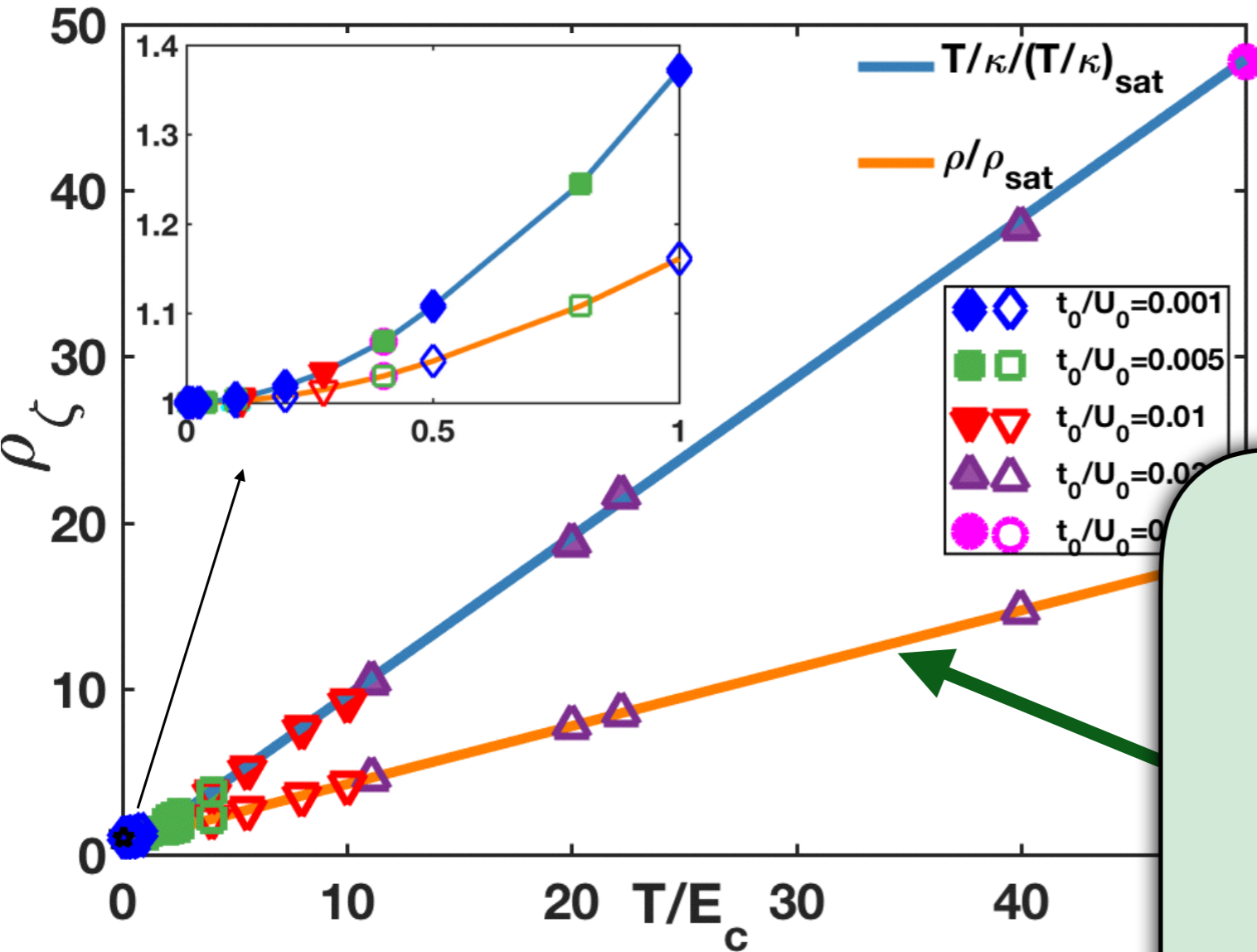
$$E_c \sim \frac{t_0^2}{U}$$

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$$E_c \sim \frac{t_0^2}{U}$$

For $E_c < T < U$, the resistivity, ρ , and entropy density, s , are

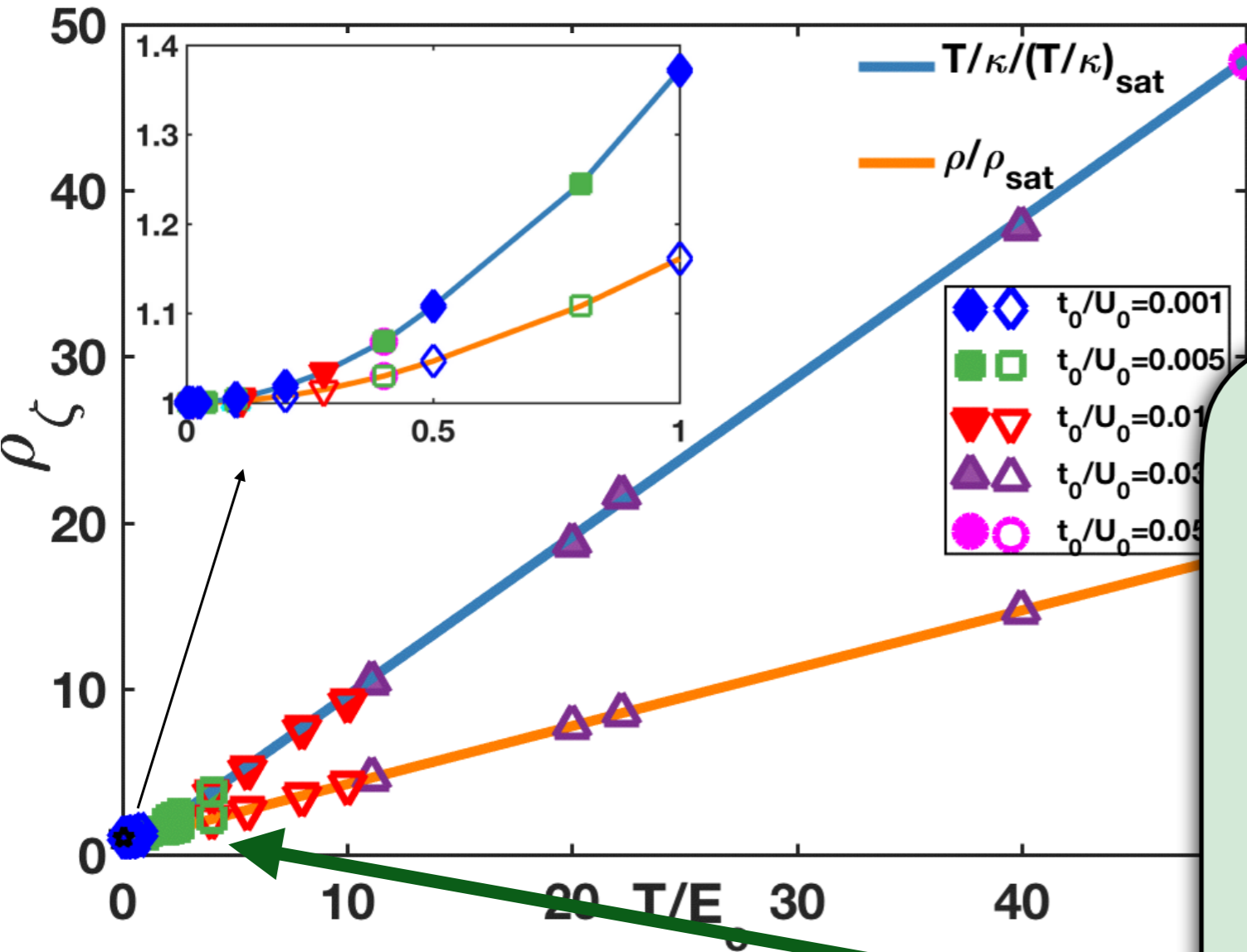
$$\rho \sim \frac{h}{e^2} \left(\frac{T}{E_c} \right), \quad s = s_0$$

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Low ‘coherence’ scale



$$E_c \sim \frac{t_0^2}{U}$$

For $T < E_c$, the resistivity, ρ , and entropy density, s , are

$$\rho = \frac{h}{e^2} \left[c_1 + c_2 \left(\frac{T}{E_c} \right)^2 \right]$$

$$s \sim s_0 \left(\frac{T}{E_c} \right)$$

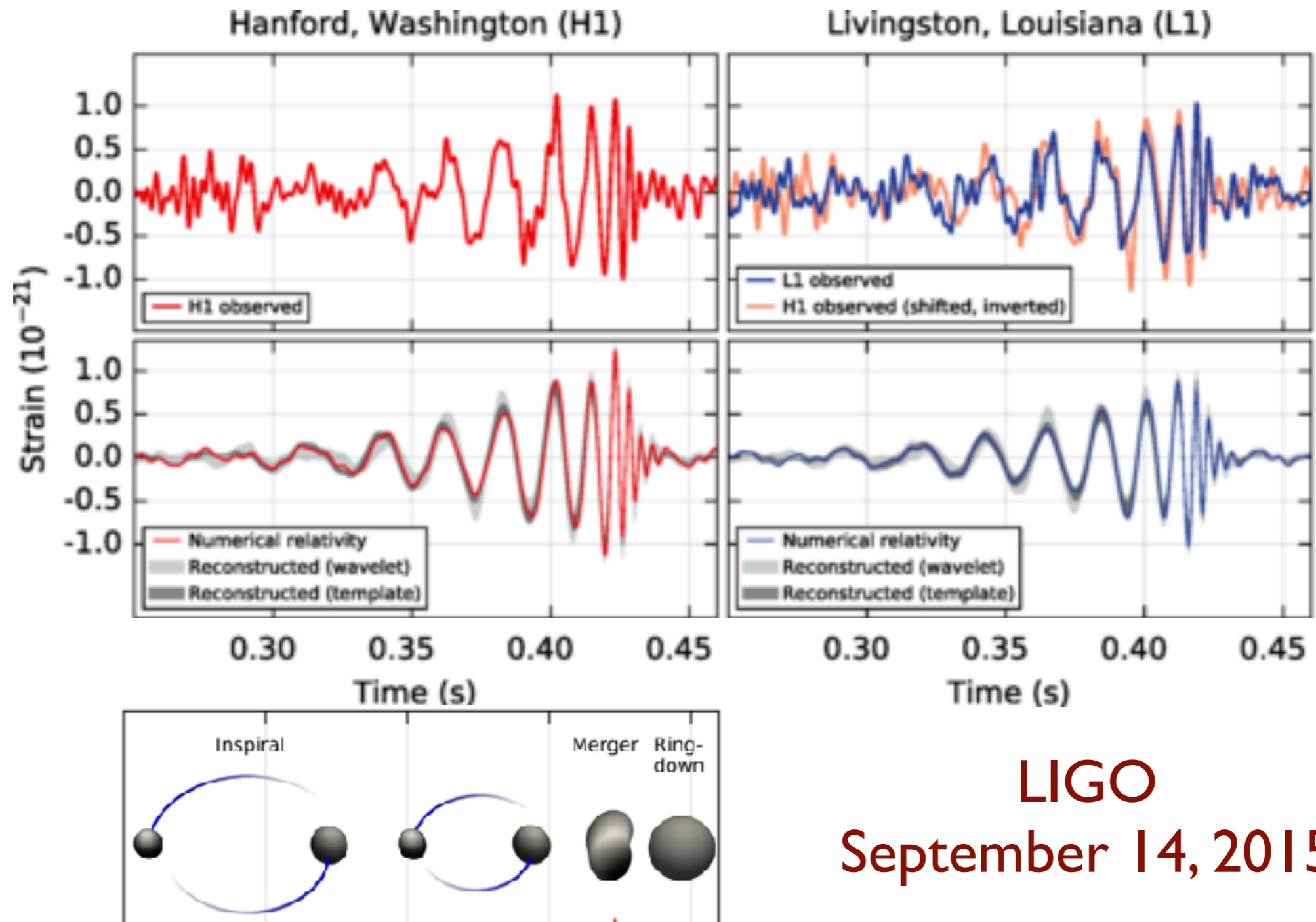
Quantum matter without quasiparticles:

- No quasiparticle decomposition of low-lying states:
$$E \neq \sum_{\alpha} n_{\alpha} \varepsilon_{\alpha} + \sum_{\alpha, \beta} F_{\alpha\beta} n_{\alpha} n_{\beta} + \dots$$
- Thermalization and many-body chaos in the shortest possible time of order $\hbar/(k_B T)$.
- These are also characteristics of black holes in quantum gravity.

Black holes

- Black holes have an entropy and a temperature, T_H .
- The entropy is proportional to their surface area.
- They relax to thermal equilibrium in a time $\sim \hbar / (k_B T_H)$.

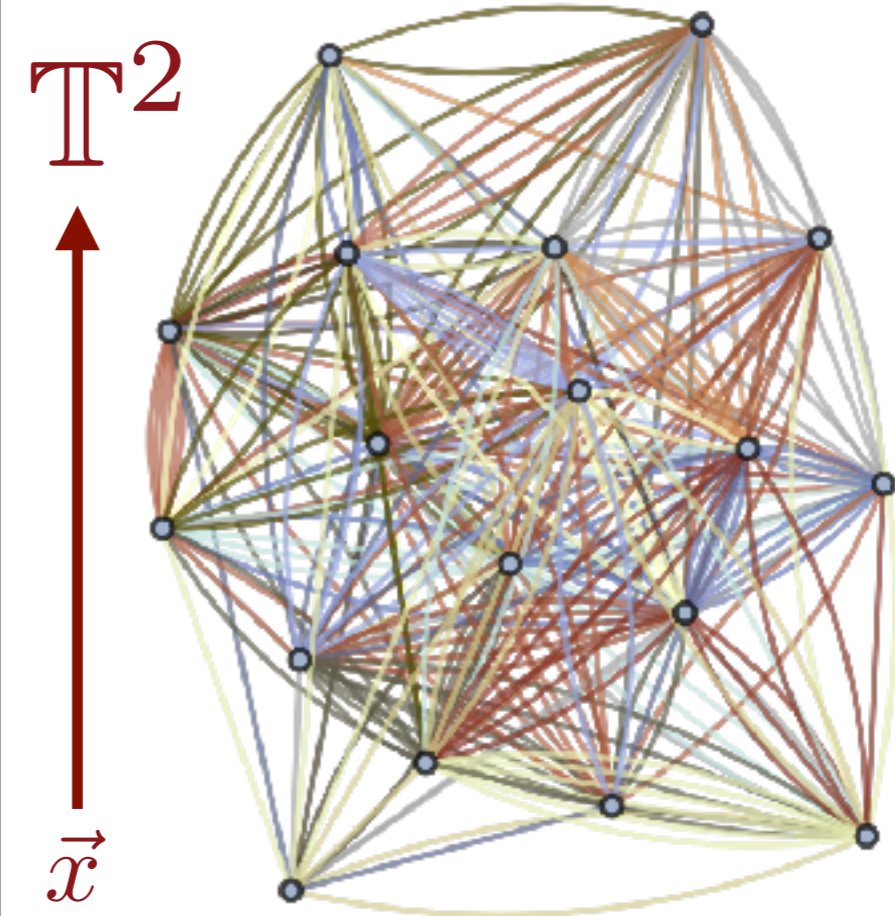
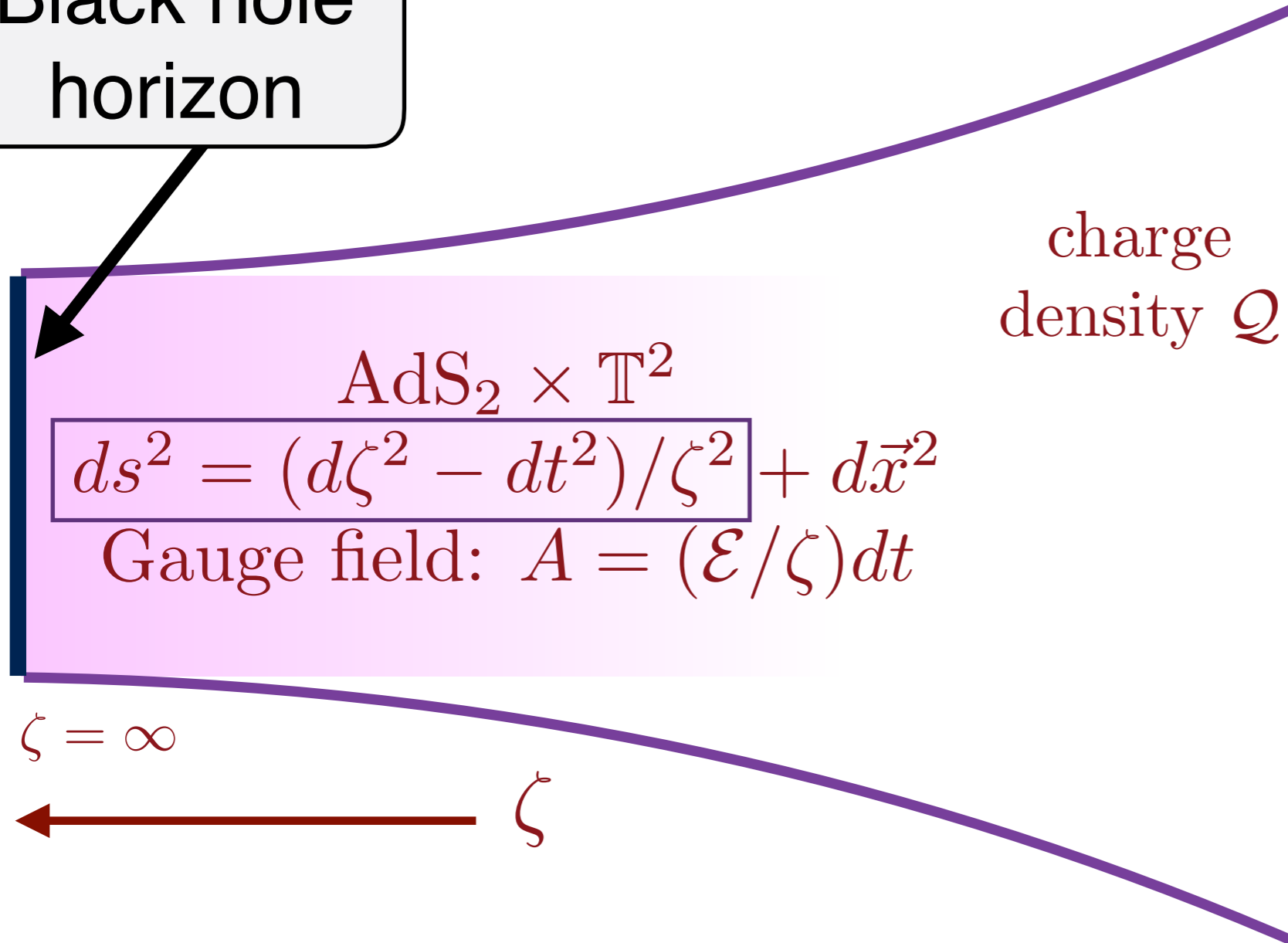




- The Hawking temperature, T_H influences the radiation from the black hole at the very last stages of the ring-down (not observed so far). The ring-down (approach to thermal equilibrium) happens very rapidly in a time $\sim \frac{\hbar}{k_B T_H} = \frac{8\pi GM}{c^3} \sim 8$ milliseconds.

SYK and black holes

Black hole horizon



Quantum gravity on the $1+1$ dimensional spacetime AdS_2 (when embedded in AdS_4) is holographically matched to the $0+1$ dimensional SYK model