

A model of d-wave superconductivity, antiferromagnetism, and charge order on the square lattice

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Subir Sachdev

F30.00008

Maine Christos, Zhu-Xi Luo, Henry Shackleton, Ya-Hui Zhang,
Mathias Scheurer, and S. S., arXiv:2302.07885

Alexander Nikolaenko, Jonas v. Milczewski, Darshan G. Joshi,
and S.S., arXiv:2211.10452

Talk online: sachdev.physics.harvard.edu

PHYSICS

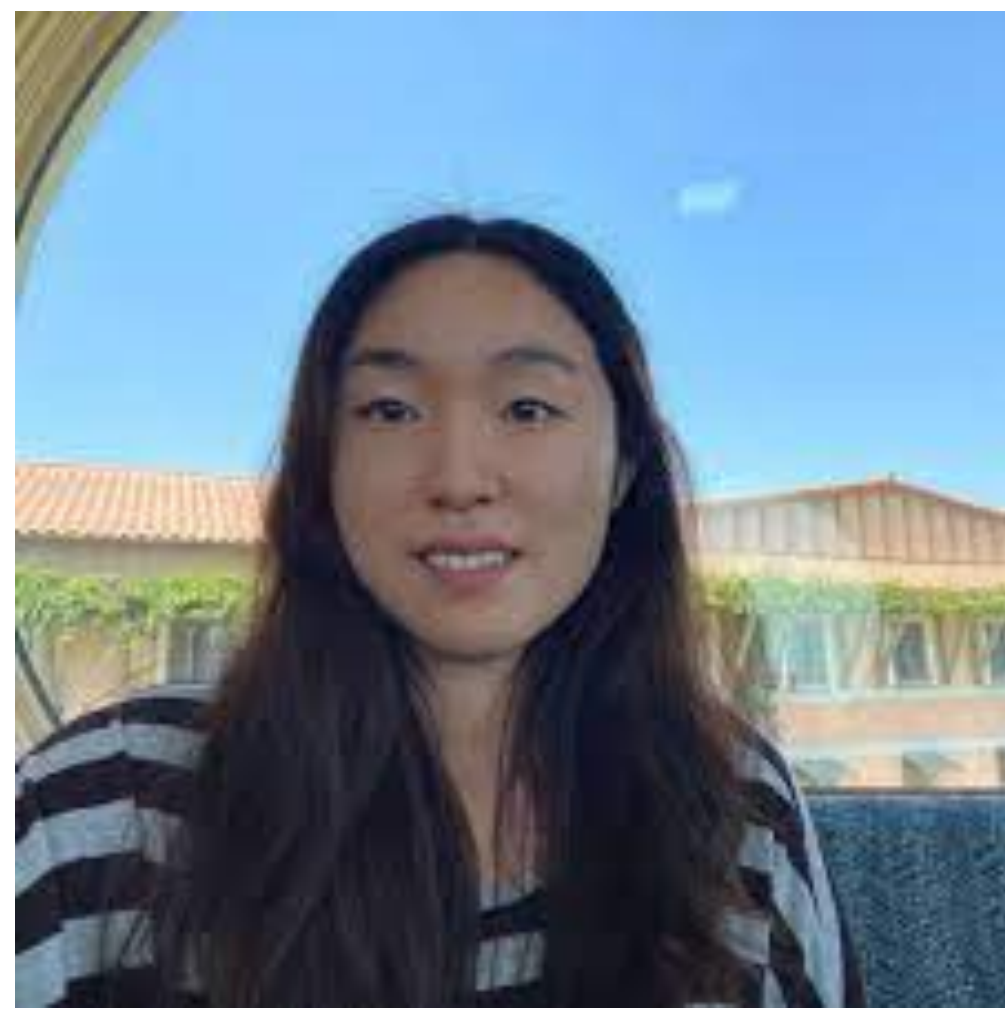


HARVARD





Maine Christos



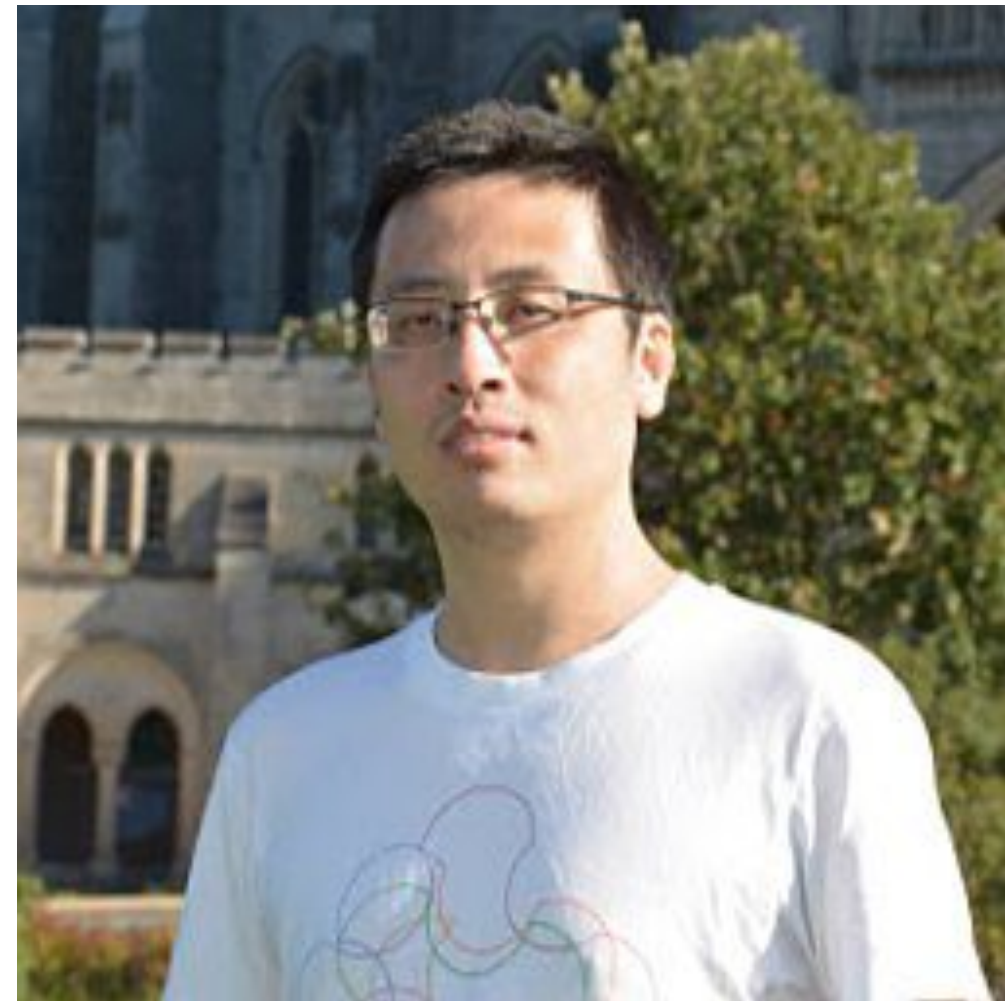
Zhu-Xi Luo



Henry Shackleton



**Mathias
Scheurer**



**Ya-Hui
Zhang**



**Alexander
Nikolaenko**

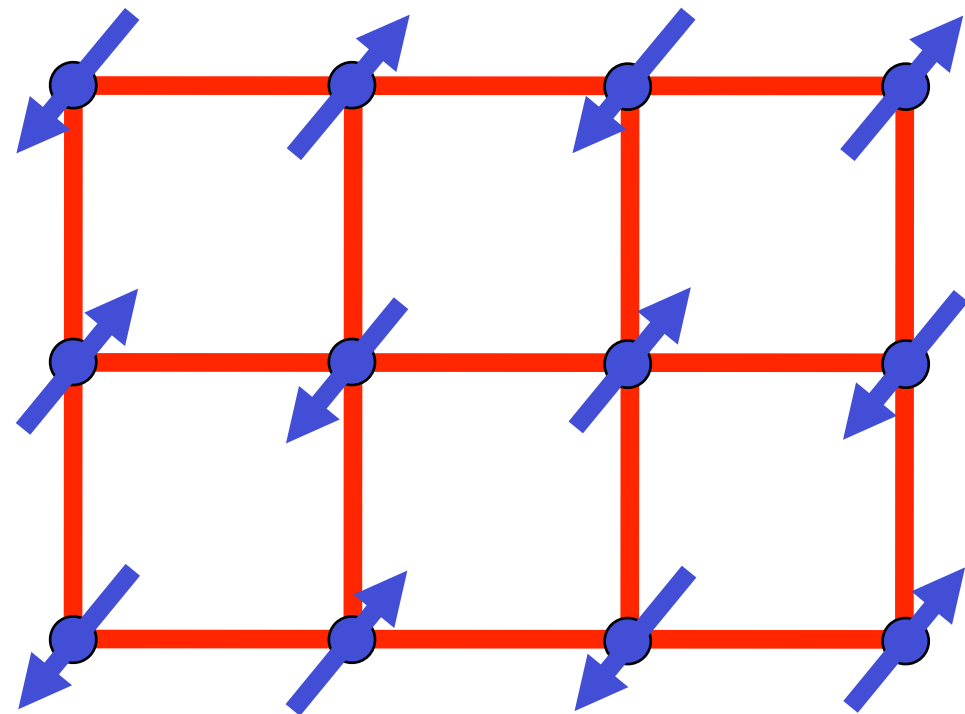


Darshan Joshi



Jonas von Milczewski

Insulating $S=1/2$ antiferromagnet

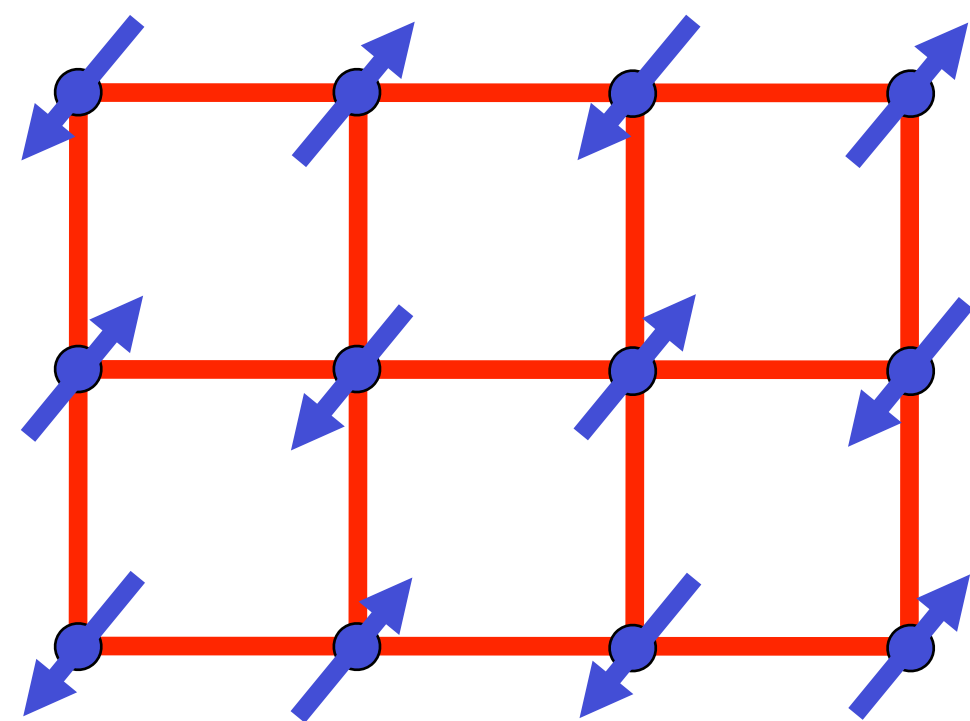


$$H = \sum_{i < j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

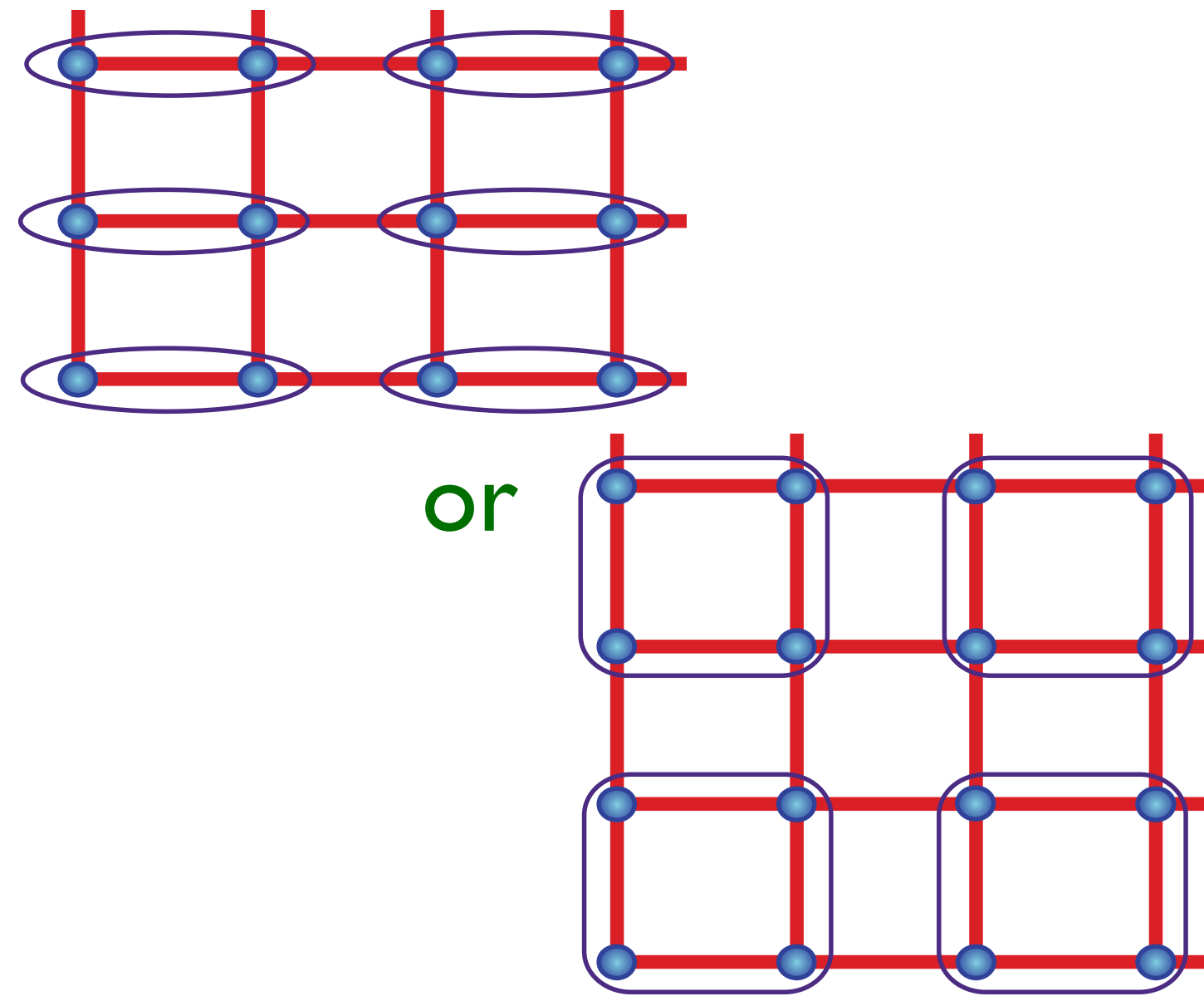
Schwinger bosons

$$\mathbf{S}_i = \frac{1}{2} b_{i\alpha}^\dagger \boldsymbol{\sigma}_{\alpha\beta} b_{i\beta}, \quad \sum_{\alpha=\uparrow,\downarrow} b_{i\alpha}^\dagger b_{i\alpha} = 1$$

Insulating $S=1/2$ antiferromagnet



Higgs phase, $\langle z_\alpha \rangle \neq 0$:
Néel order



Confining phase, $\langle z_\alpha \rangle = 0$:
VBS order

s

$$H = \sum_{i < j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

Schwinger bosons

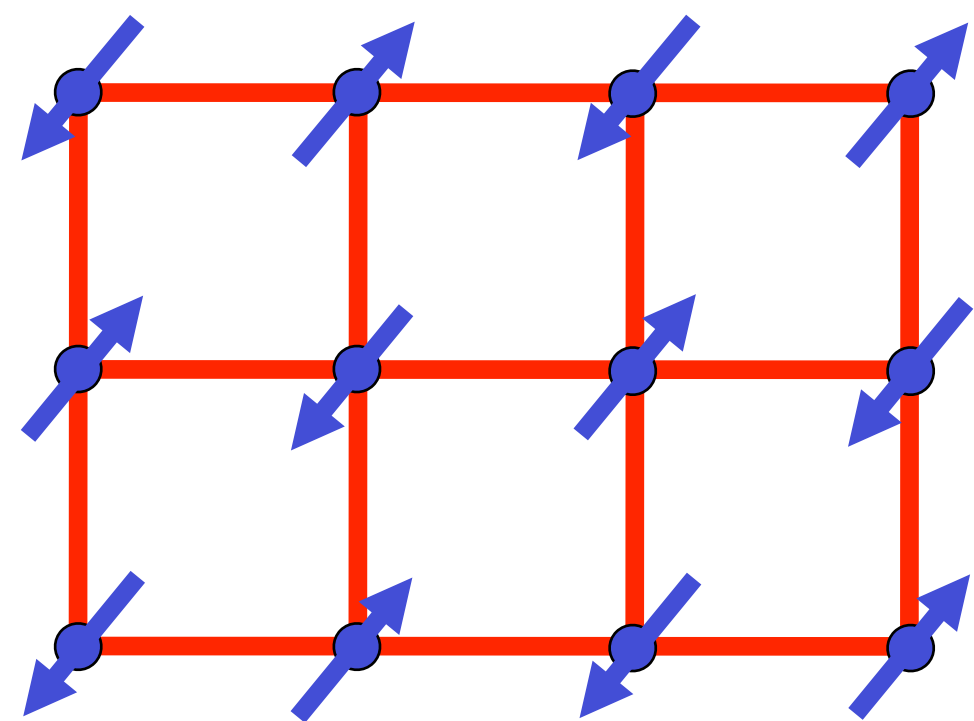
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Low energy $\mathbb{C}\mathbb{P}^1$ U(1) gauge theory

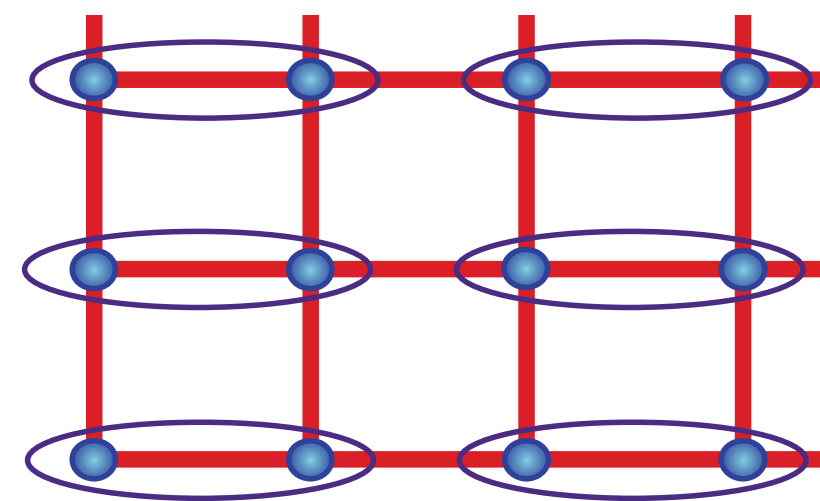
$$z_\alpha \sim b_{A\alpha} + \varepsilon_{\alpha\beta} b_{B\beta}$$

$$\mathcal{L} = |(\partial_\mu - ia_\mu) z_\alpha|^2 + s |z_\alpha|^2 + u |z_\alpha|^4 + \mathcal{L}_{\text{monopole}}$$

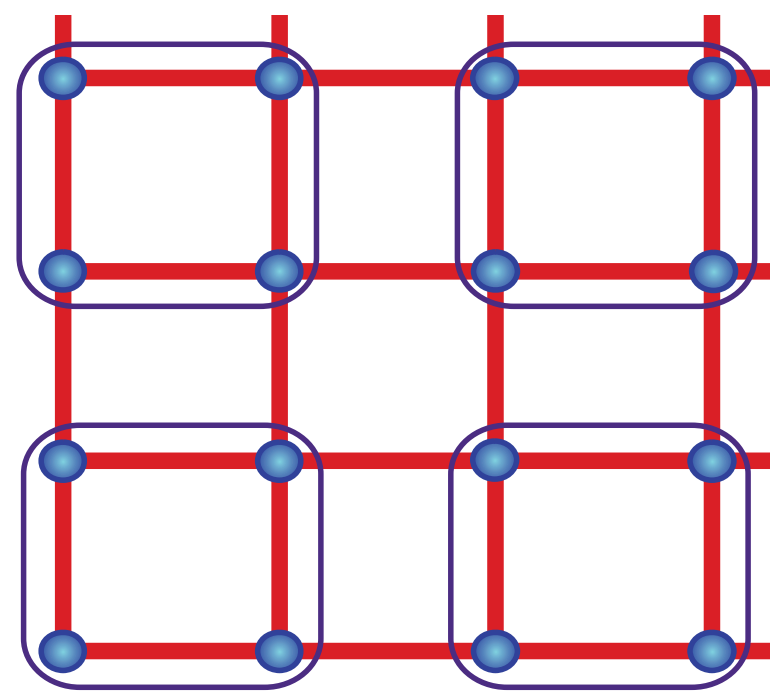
Insulating $S=1/2$ antiferromagnet



Néel order



or



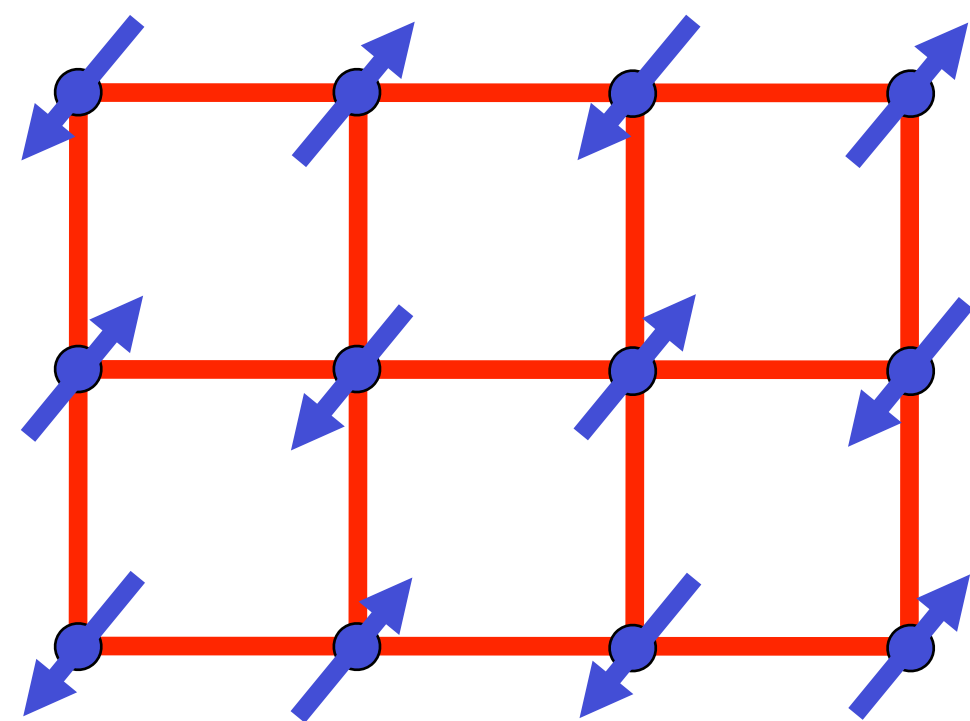
VBS order

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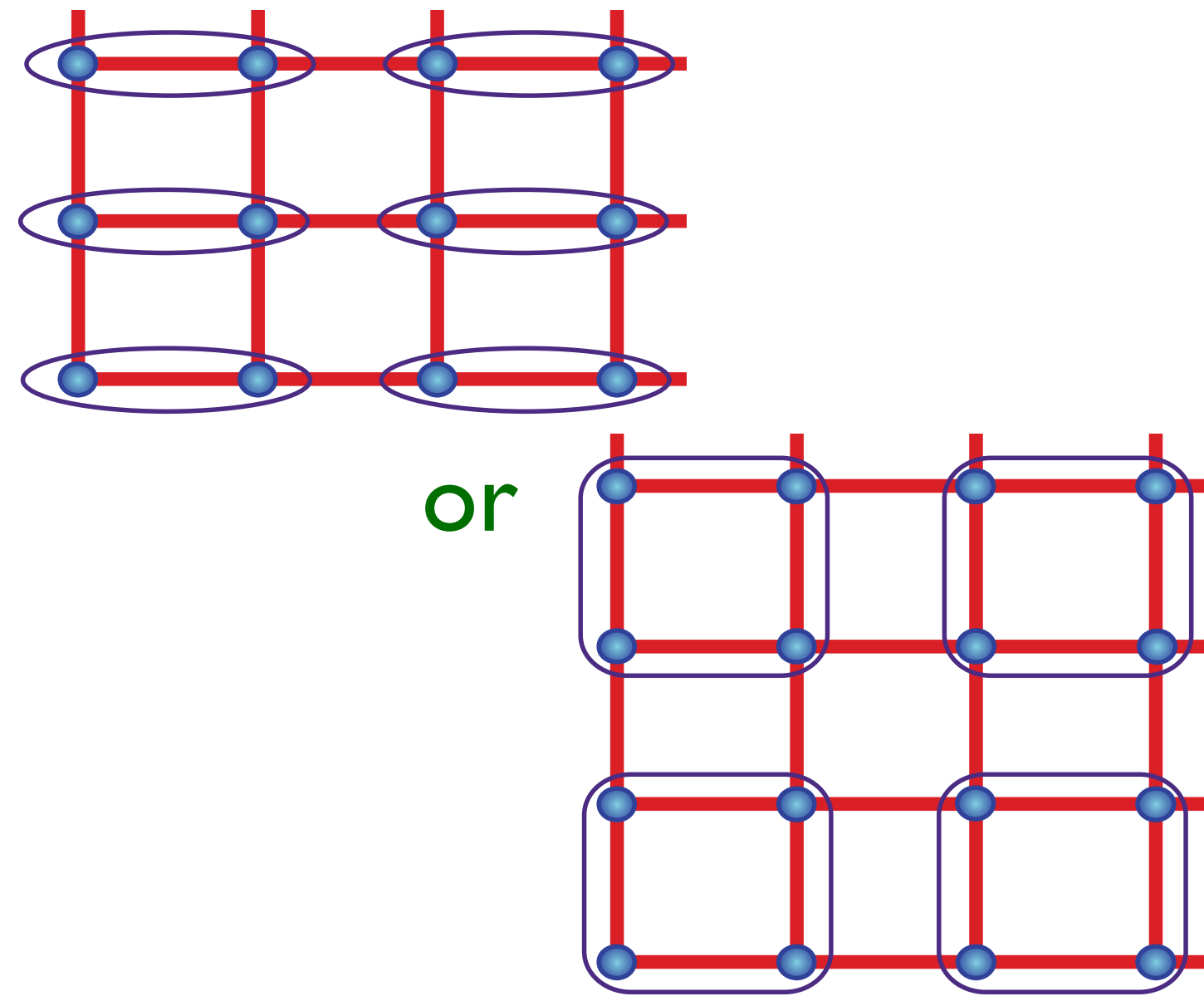
Schwinger fermions

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Insulating $S=1/2$ antiferromagnet



Confining phase:
Néel order



Confining phase:
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π -flux mean-field leads to a
low energy theory of
 $N_f = 2$ Dirac fermions Ψ_s
coupled to
an emergent SU(2) gauge field.

Dual to $\mathbb{C}\mathbb{P}^1$ U(1) gauge theory!

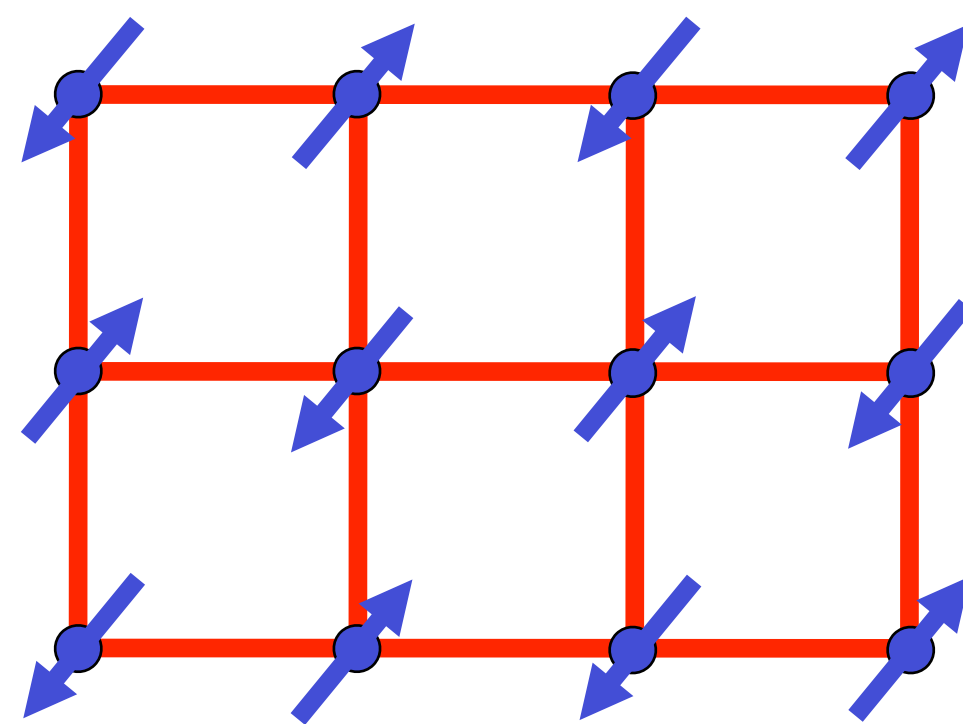
$$\mathcal{L} = i\bar{\Psi}_s D_\mu \Psi_s + \dots$$

Include charge fluctuations at half-filling: repulsive Hubbard model with pair-hopping

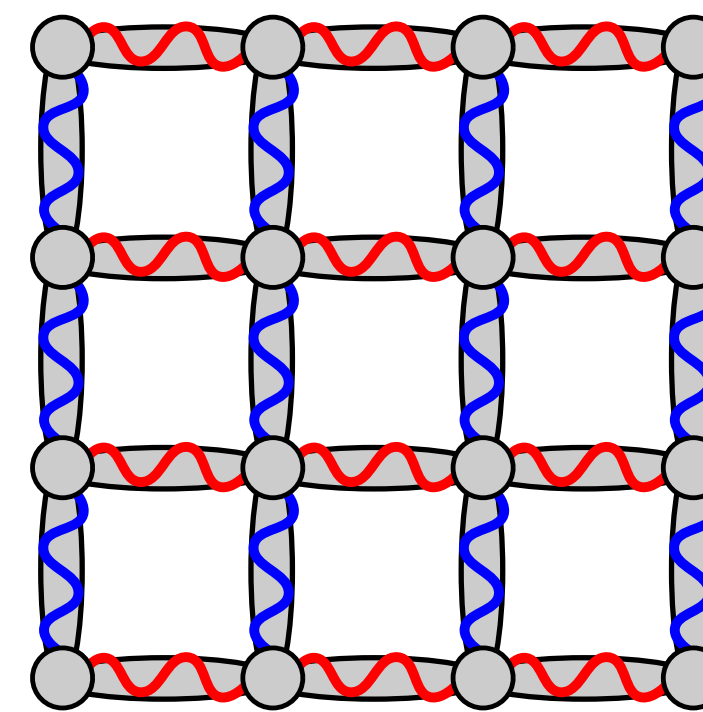
$$H = -\frac{t}{2} \sum_i K_i + U \sum_i (n_{i\uparrow} - \frac{1}{2})(n_{i\downarrow} - \frac{1}{2}) - W \sum_i K_i^2,$$

$$K_i = \sum_{\hat{e}=\pm\hat{x},\pm\hat{y}} \left(c_{i\alpha}^\dagger c_{i+\hat{e},\alpha} + c_{i+\hat{e},\alpha}^\dagger c_{i\alpha} \right)$$

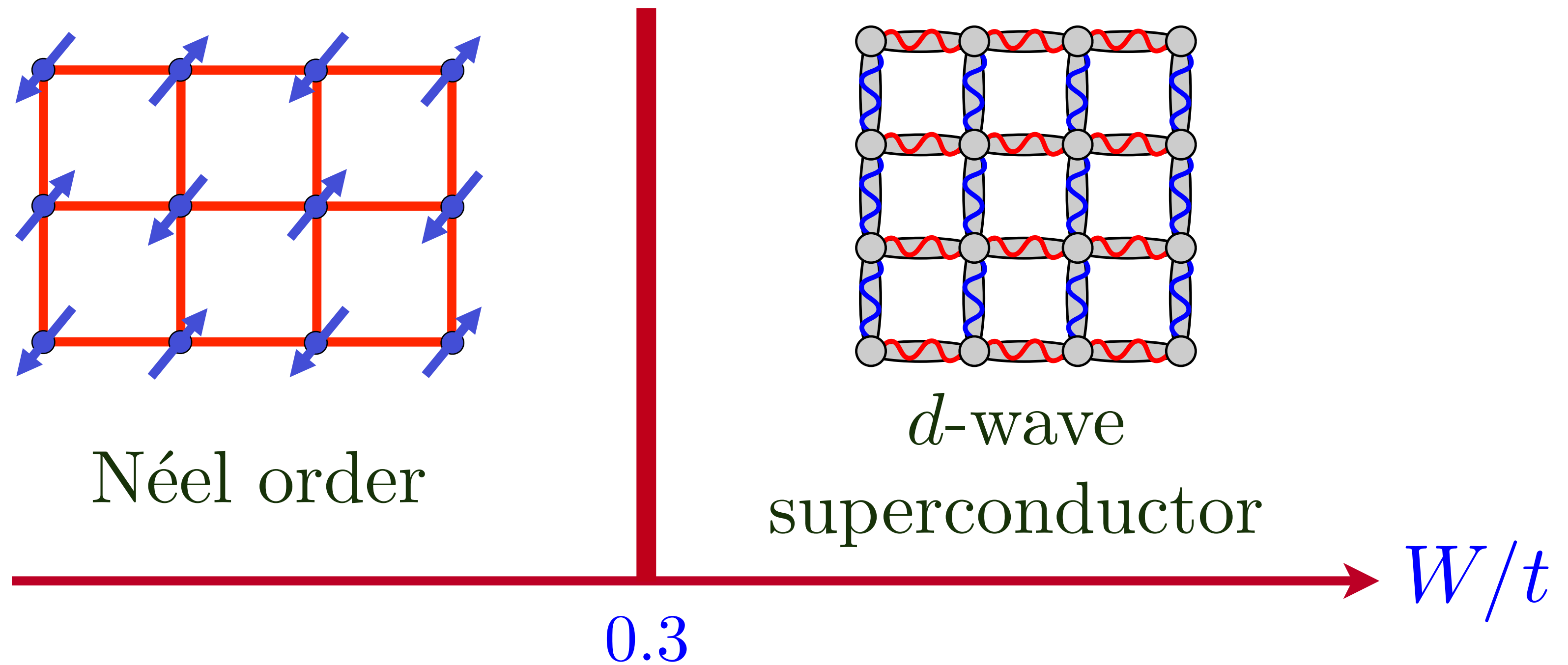
$$U/t = 4$$



Néel order



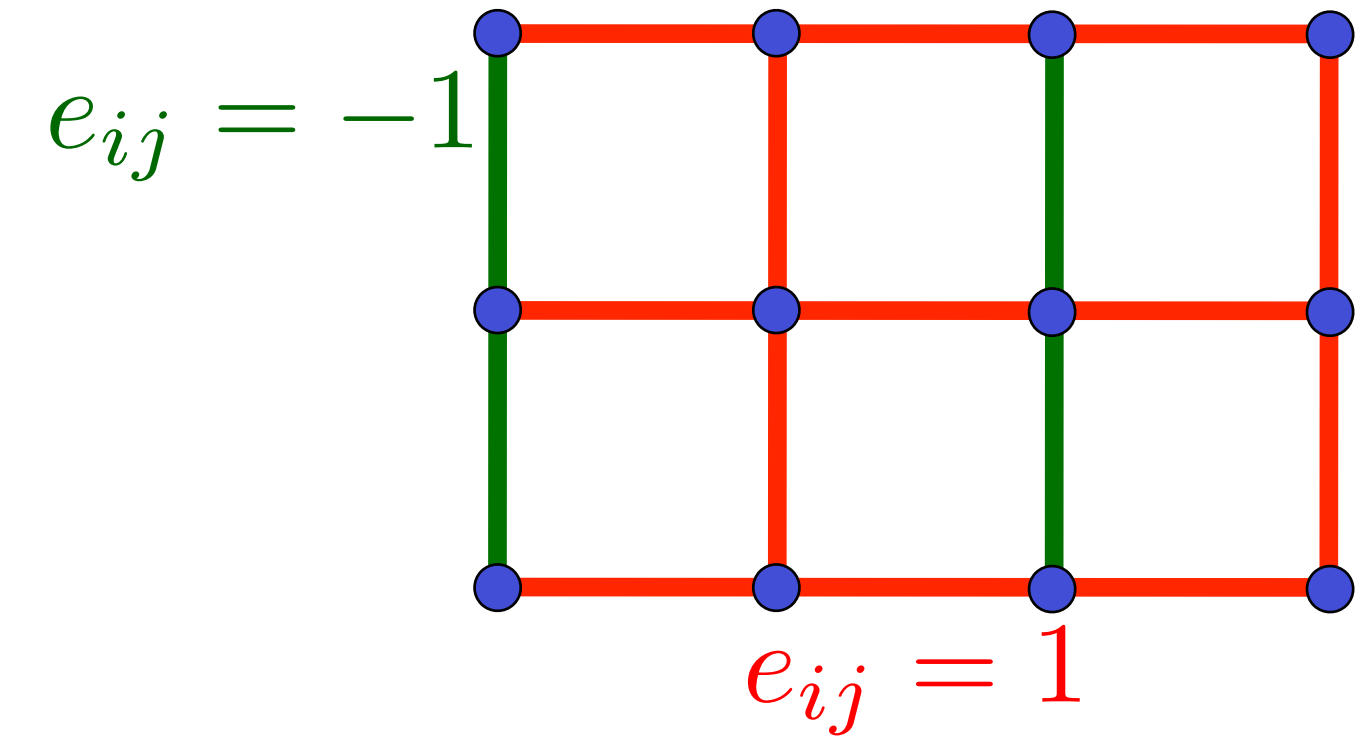
d-wave
superconductor



Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

- Begin with the π -flux spin liquid in the fermionic spinon description.

$$H_f = iJ \sum_{\langle ij \rangle} e_{ij} \left(f_{i\alpha}^\dagger f_{j\alpha} - f_{j\alpha}^\dagger f_{i\alpha} \right)$$

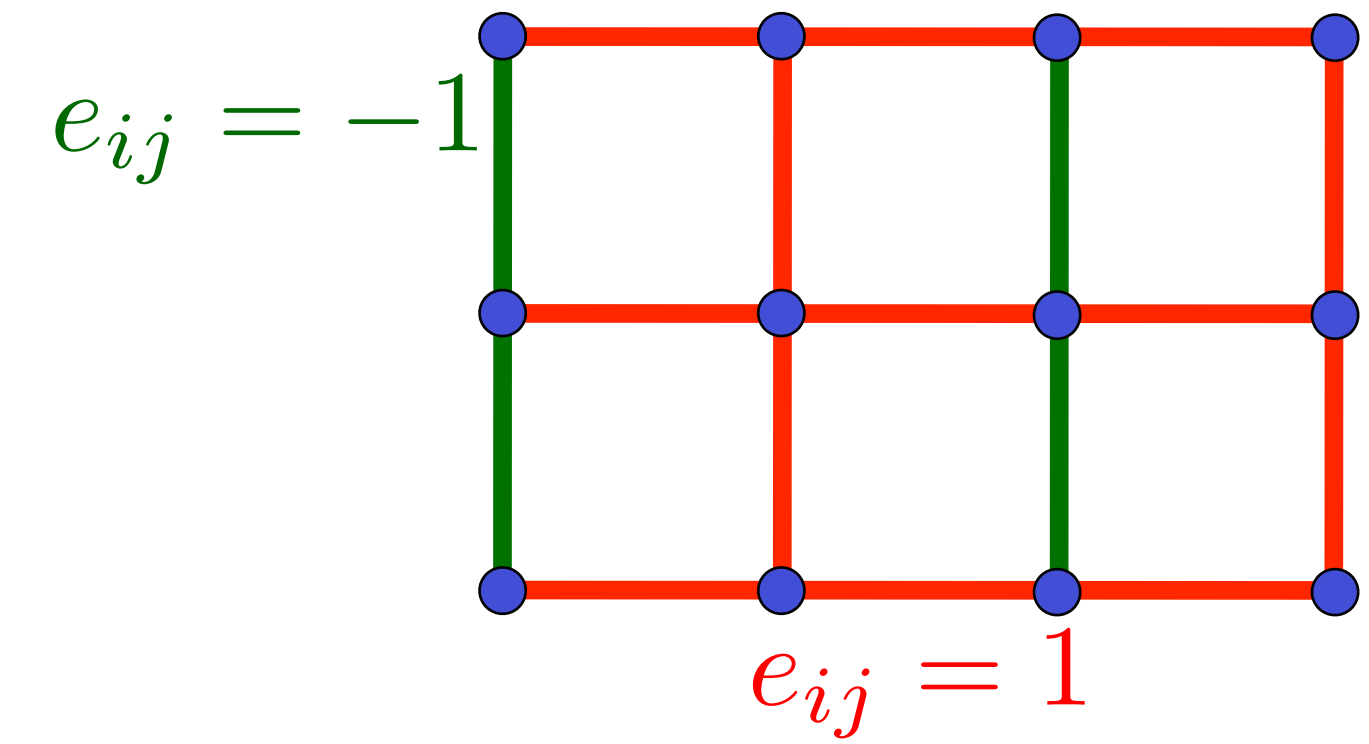


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H_f is invariant under distinct SU(2) rotations in spin and Nambu space.

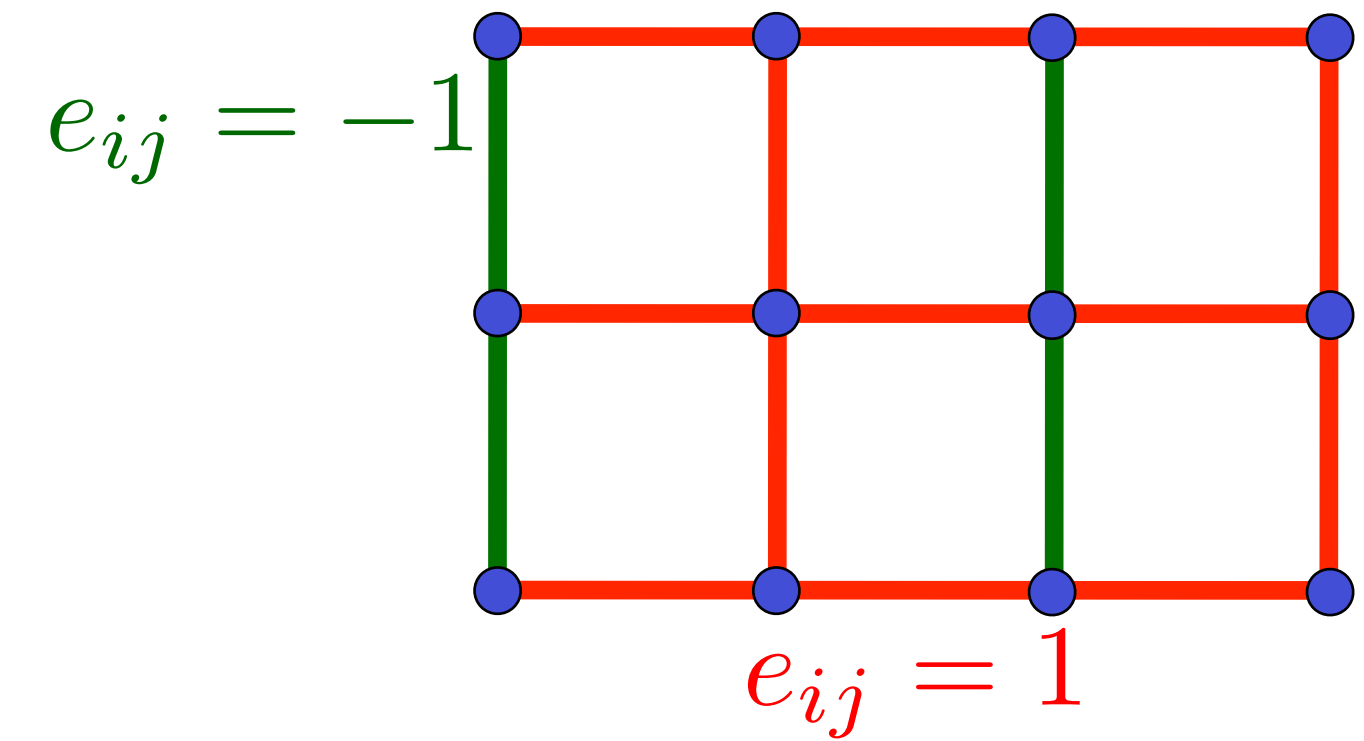


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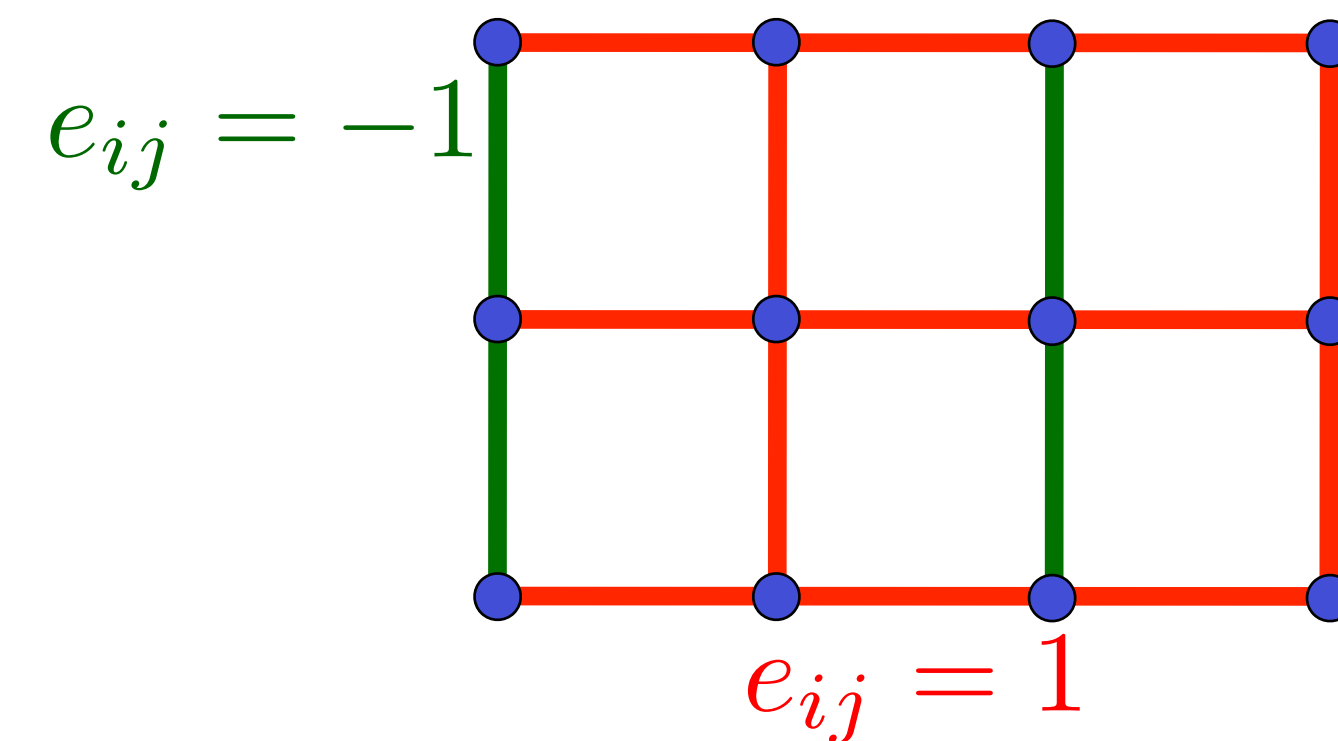


- We can fully confine the SU(2) gauge field by condensing a boson, B_i , which is a fundamental of gauge SU(2). To obtain superconductivity with charge $2e$ pairs in the confining phase, B_i should also carry electromagnetic charge e .

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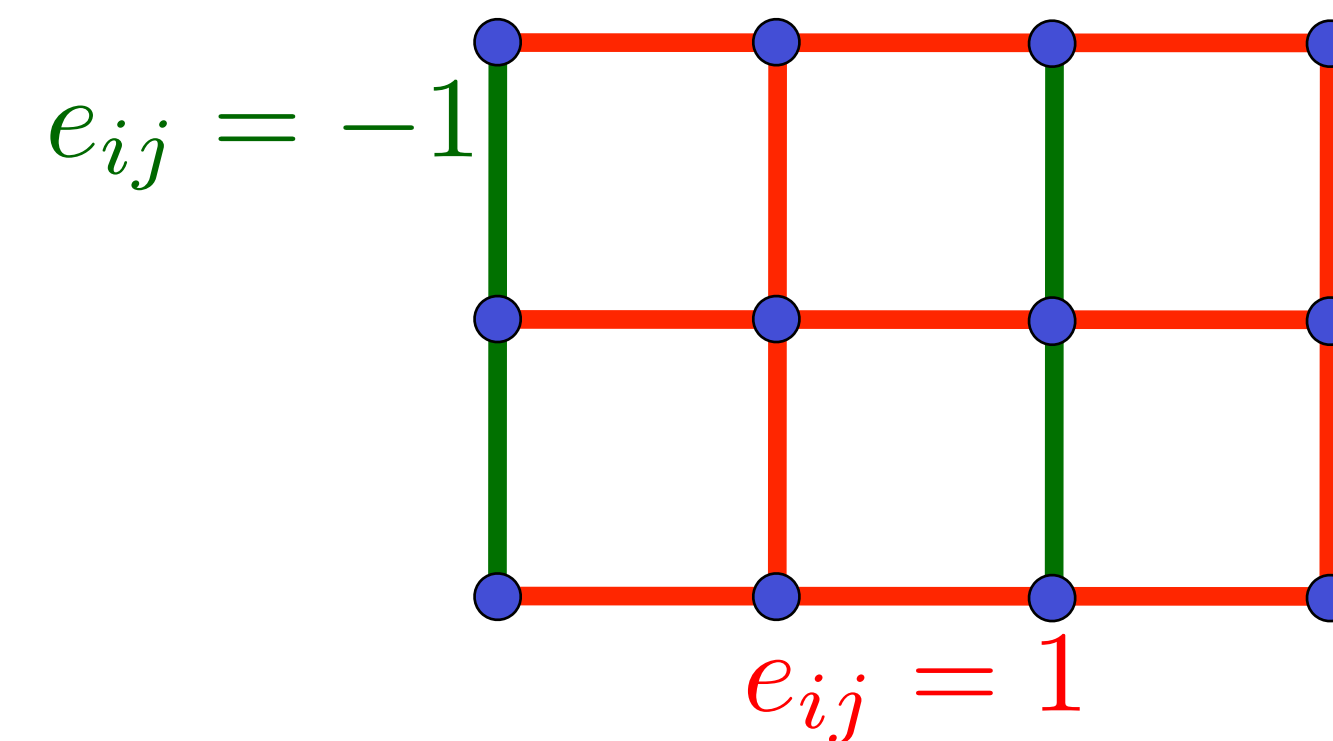
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- This uniquely identifies B_i as the ‘chargon’ of X.-G. Wen and P.A. Lee, PRL **76**, 503 (1996), related to the electrons $c_{i\sigma}$ by

$$B_i = \begin{pmatrix} B_{1i} \\ B_{2i} \end{pmatrix} ; \quad \begin{pmatrix} c_{i\uparrow} \\ c_{i\downarrow}^\dagger \end{pmatrix} \sim \begin{pmatrix} B_{1i}^* & B_{2i}^* \\ -B_{2i} & B_{1i} \end{pmatrix} \begin{pmatrix} f_{i\uparrow} \\ f_{i\downarrow}^\dagger \end{pmatrix}$$

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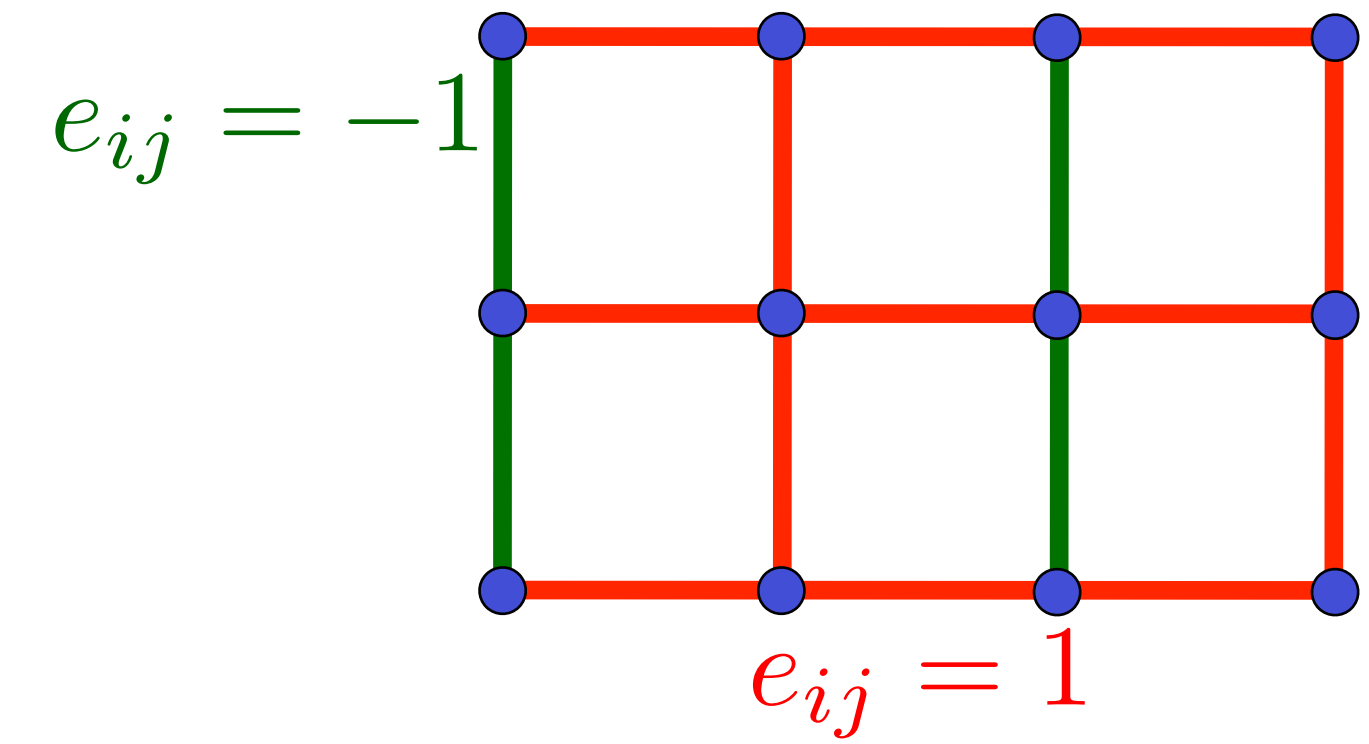
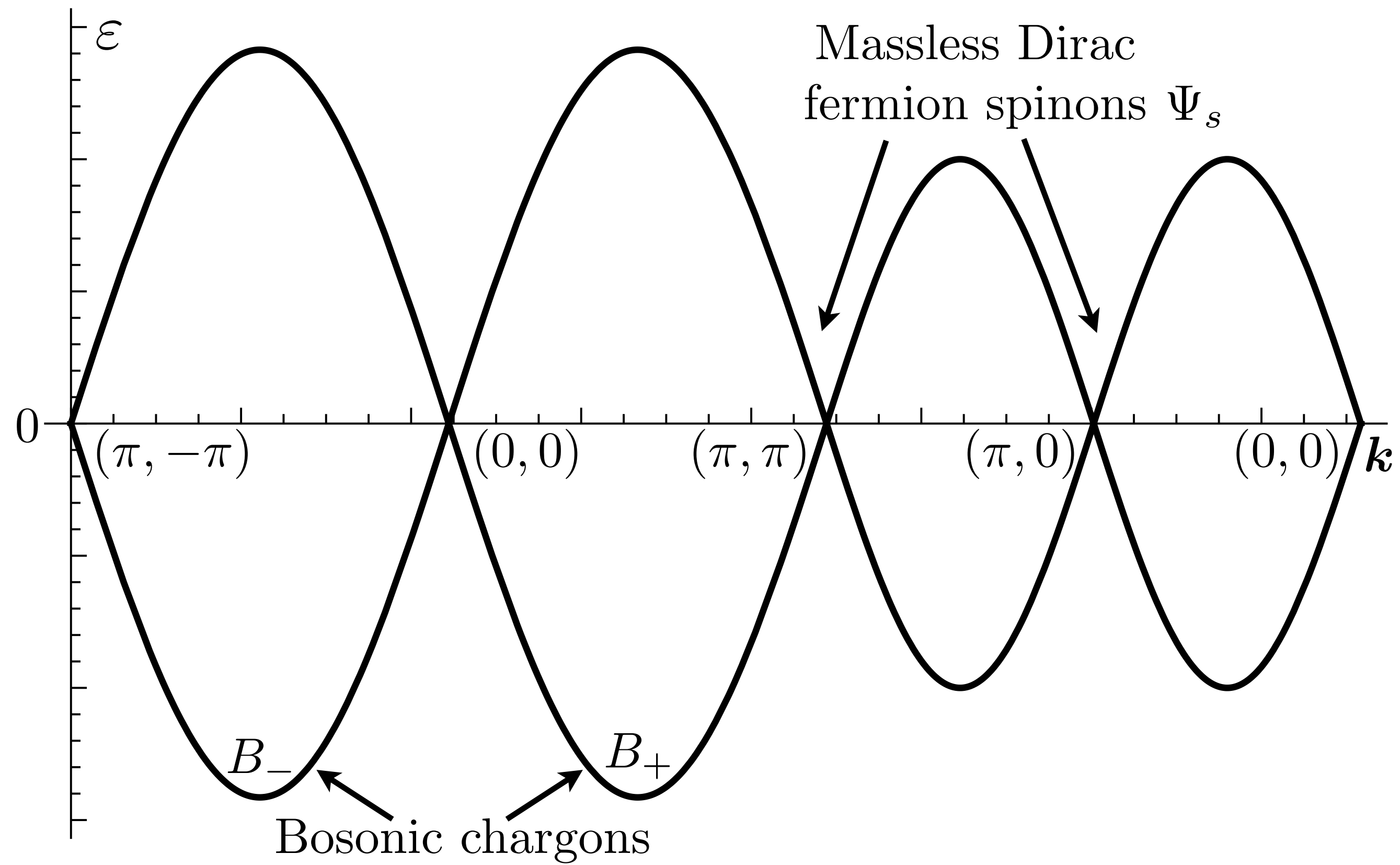
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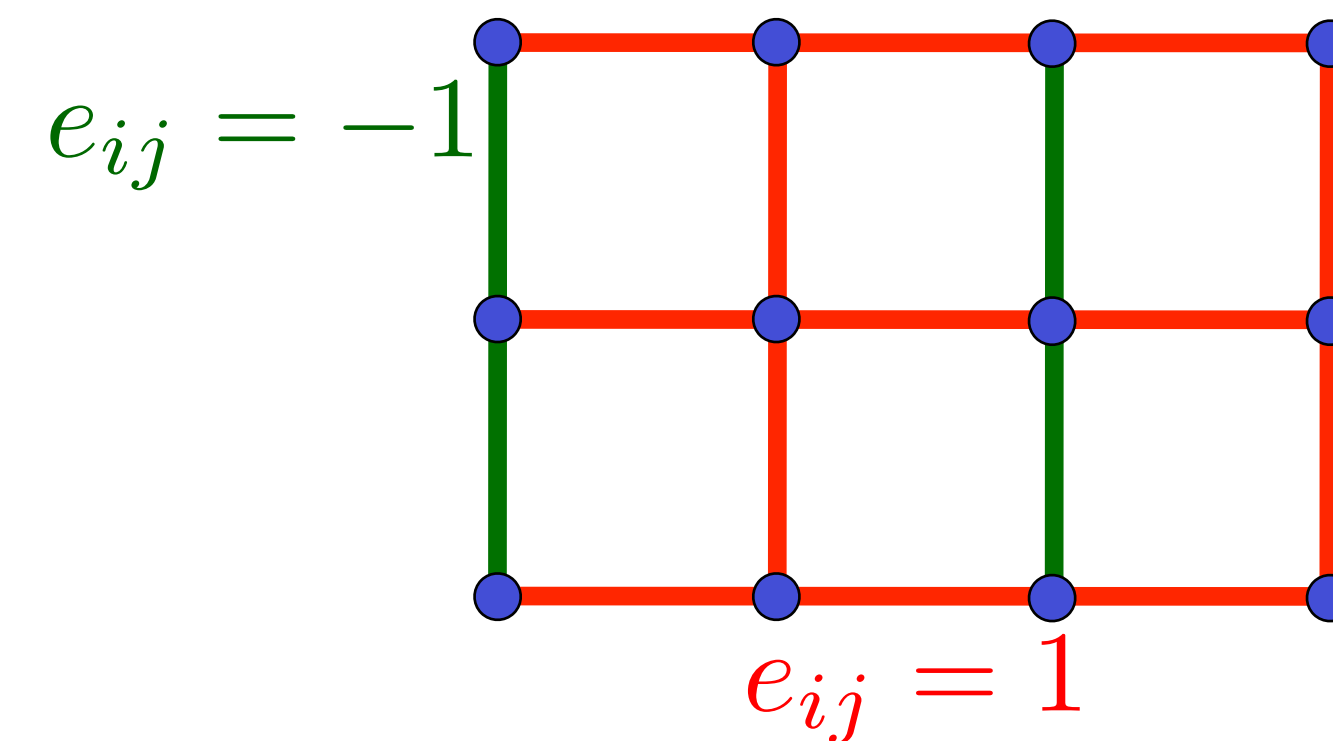
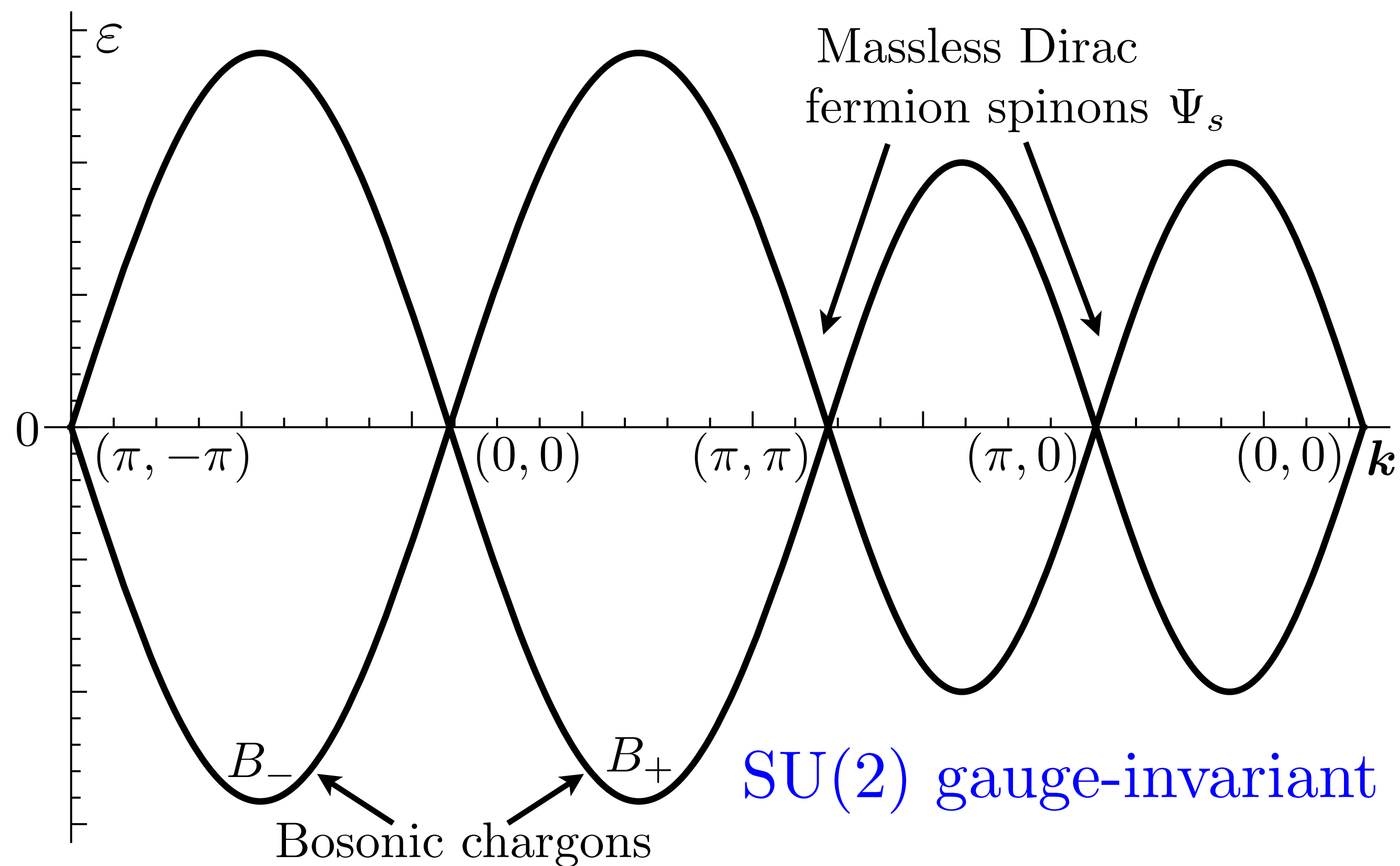
- Knowing the projective symmetry transformations of Ψ_i , we can deduce those of the B_i , and obtain the effective Hamiltonian for B_i

$$H_B = r \sum_i B_i^\dagger B_i + iw \sum_{\langle ij \rangle} e_{ij} \left(B_i^\dagger B_j - B_j^\dagger B_i \right) + \dots$$

Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



SU(2) gauge-invariant order parameters of Higgs phases:

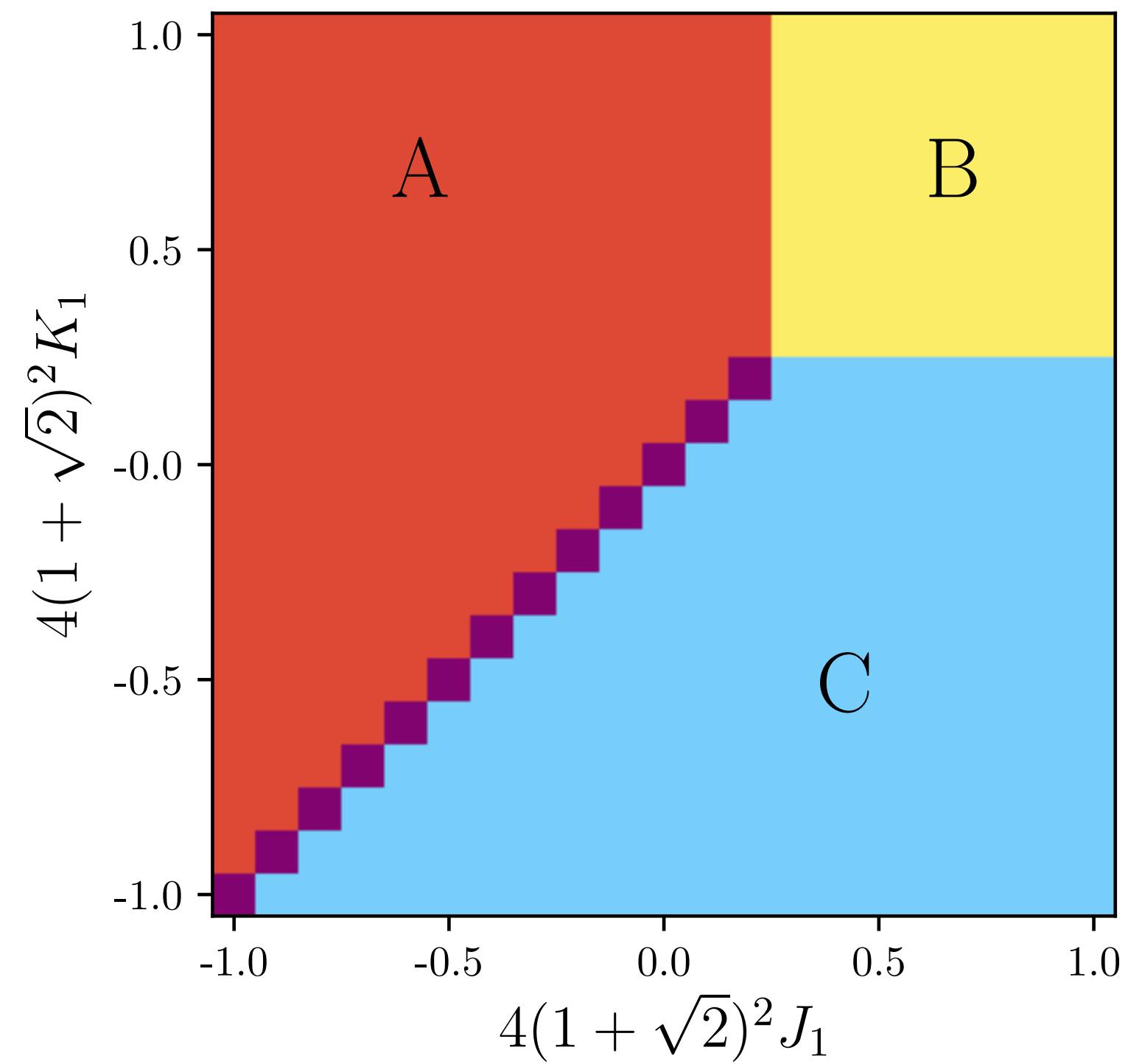
$$x\text{-CDW} : \rho_{(\pi,0)} = B_{a+}^* B_{a+} - B_{a-}^* B_{a-}$$

$$y\text{-CDW} : \rho_{(0,\pi)} = B_{a+}^* B_{a-} + B_{a-}^* B_{a+}$$

$$d\text{-density wave} : D = i (B_{a+}^* B_{a-} - B_{a-}^* B_{a+})$$

$$d\text{-wave superconductor} : \Delta = \varepsilon_{ab} B_{a+} B_{b-}$$

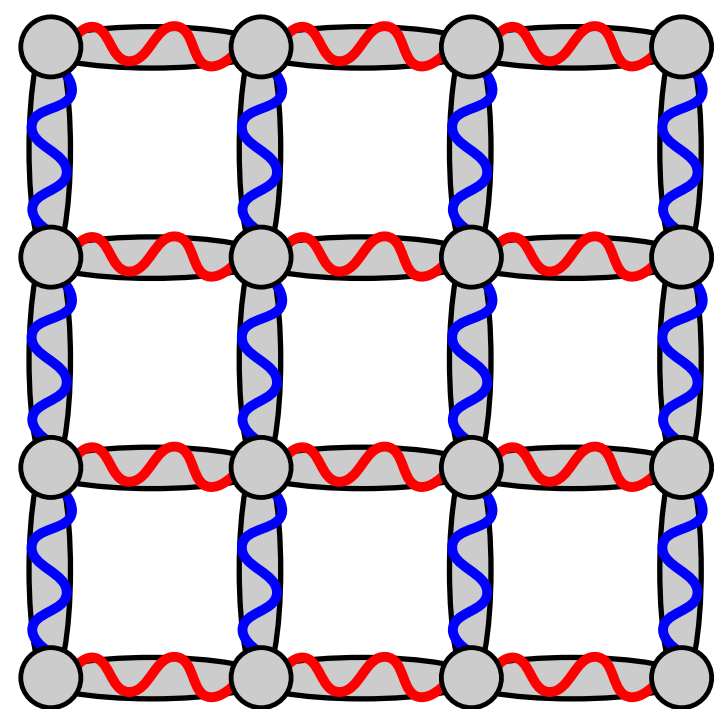
Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



$$\langle B \rangle \neq 0$$

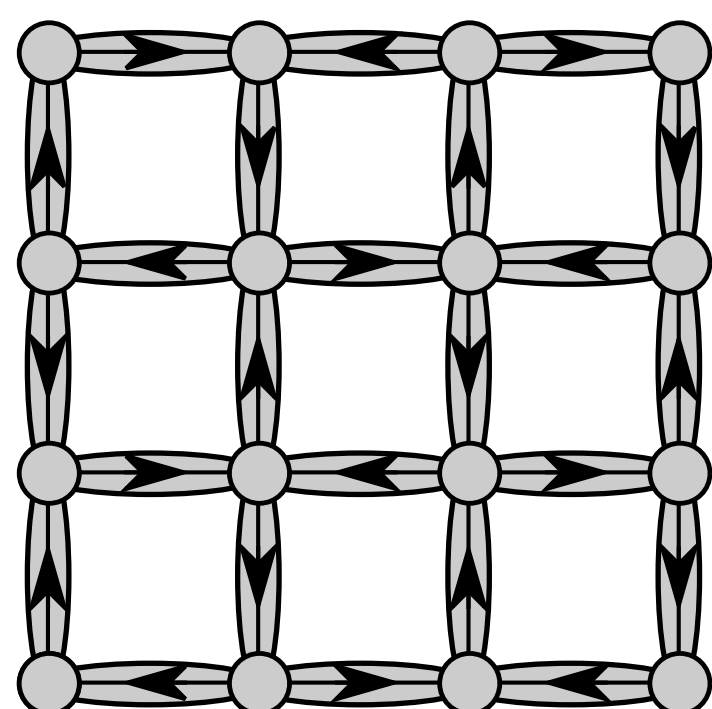
Phase B

d-wave SC



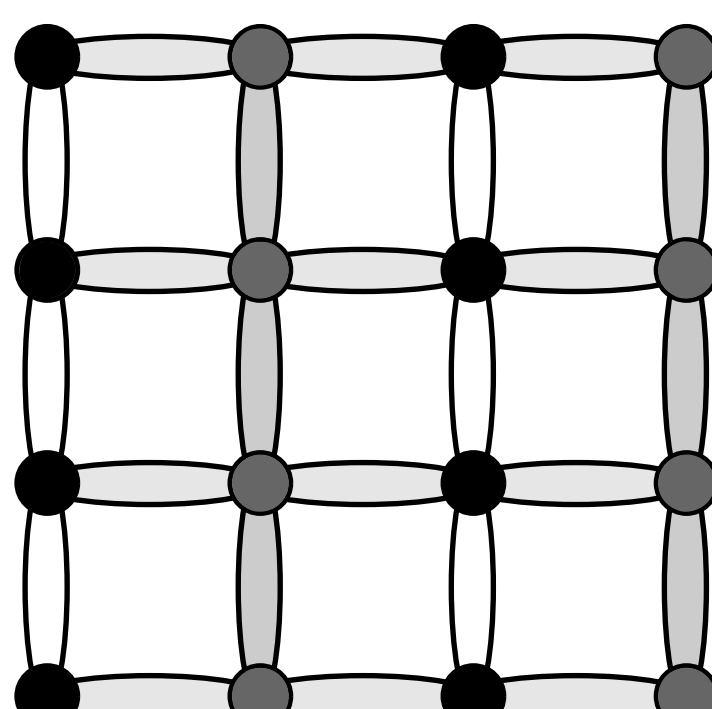
Phase C

d-density



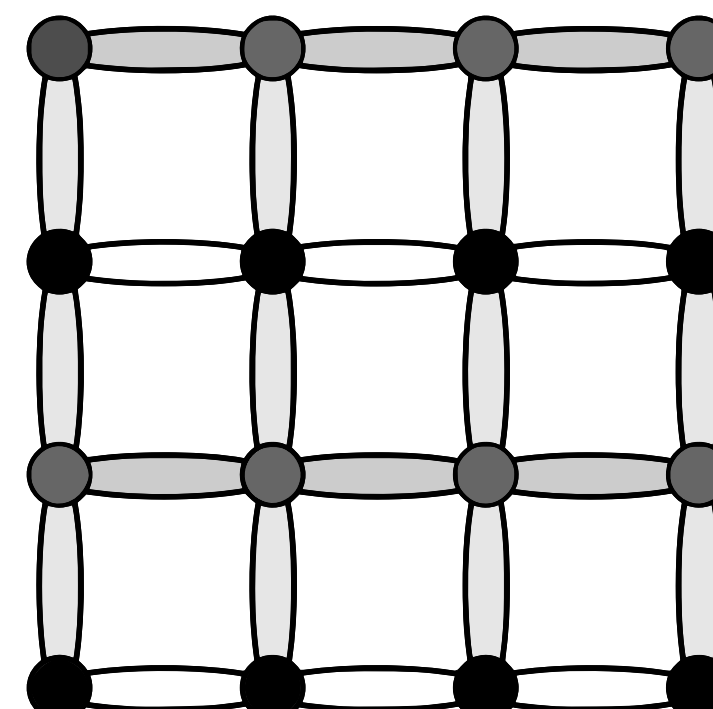
Phase A

$(\pi, 0)$ stripe

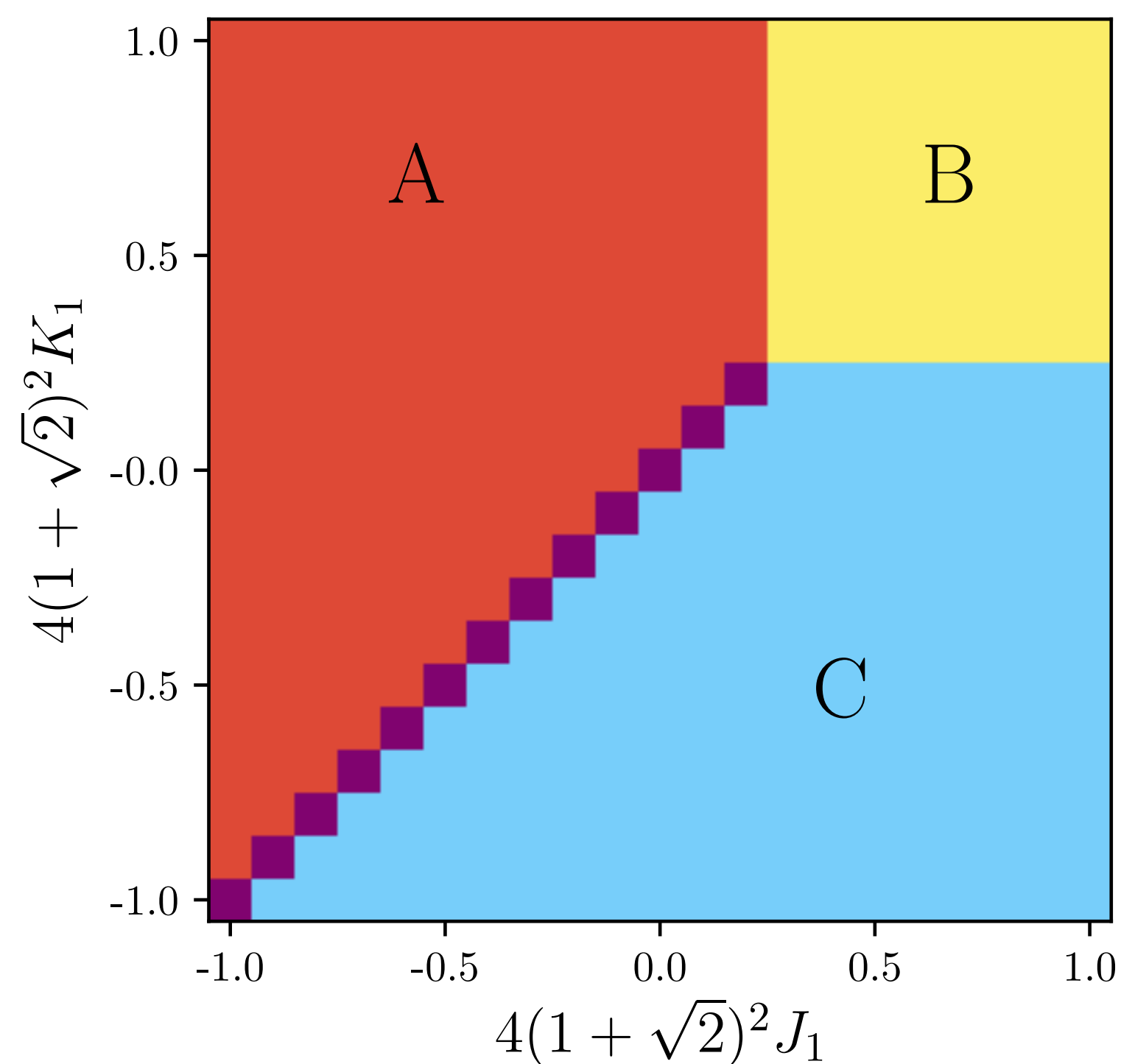


Phase A

$(0, \pi)$ stripe

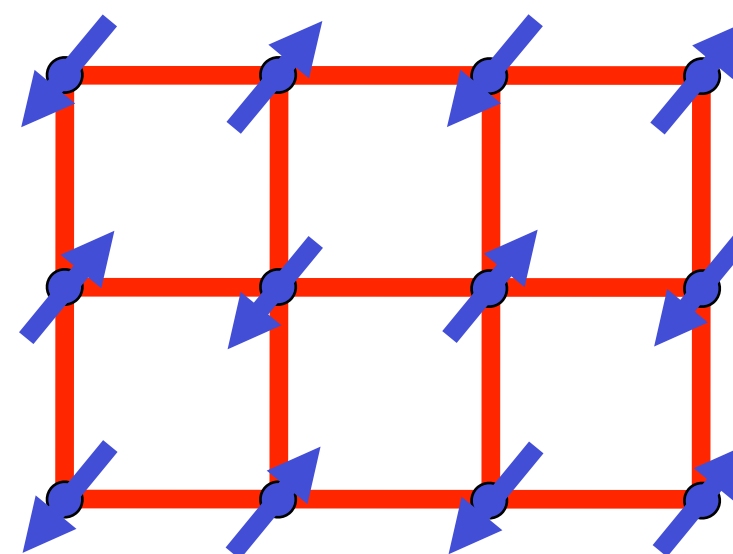


Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

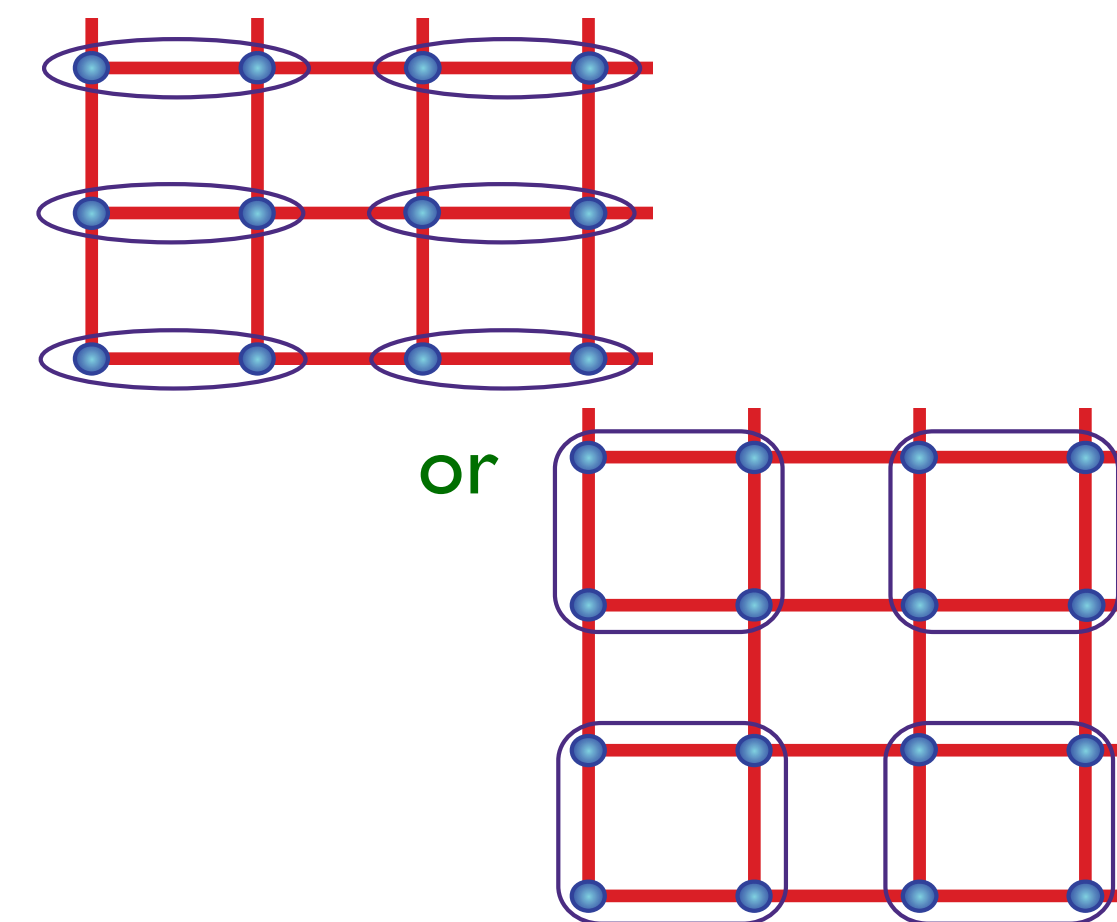


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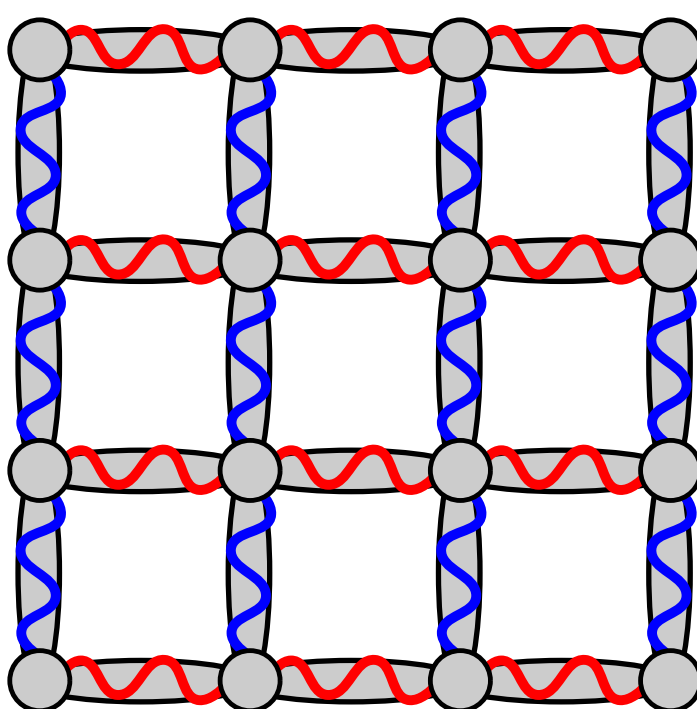


Confining phase:
Néel order

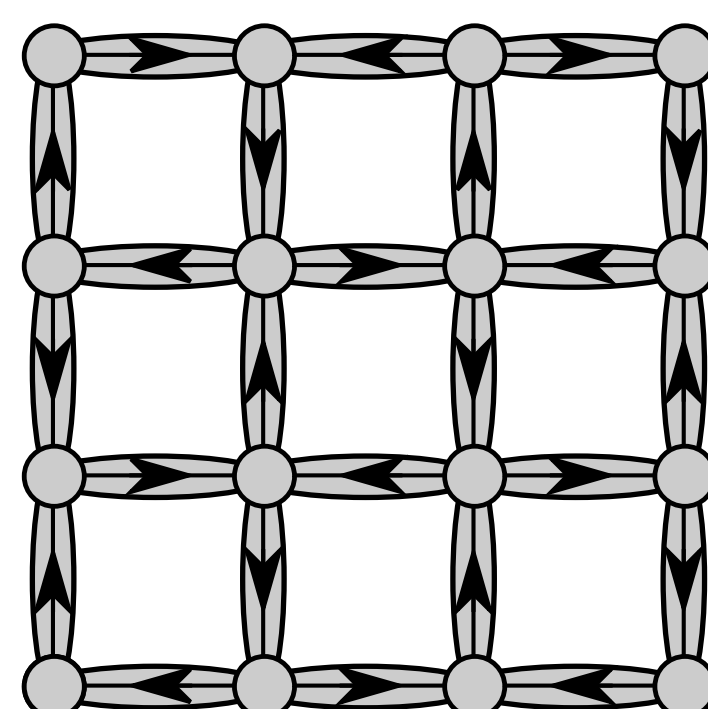


Confining phase:
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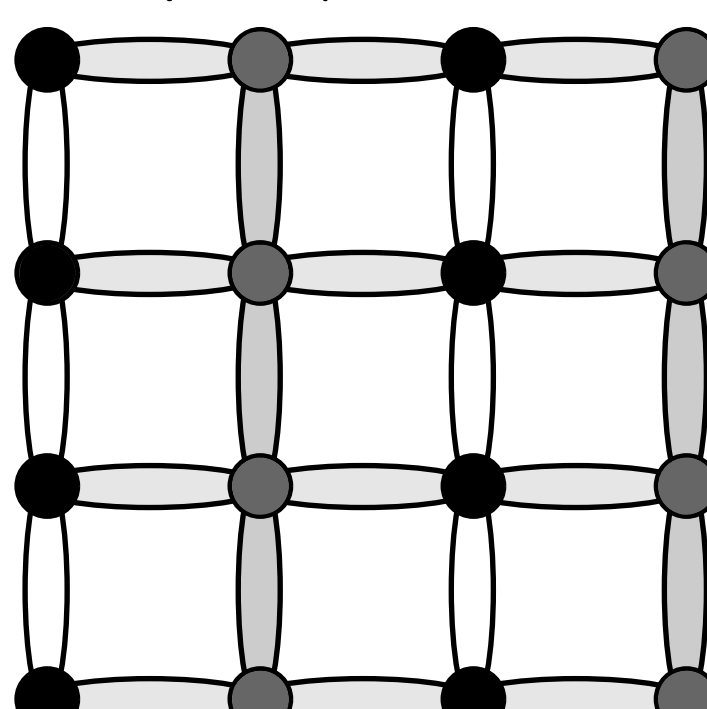
Phase B
d-wave SC



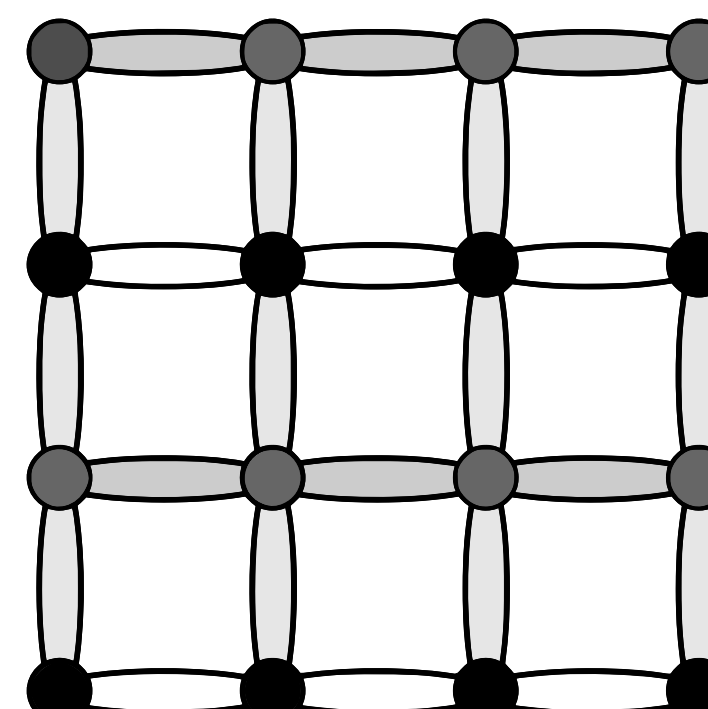
Phase C
d-density



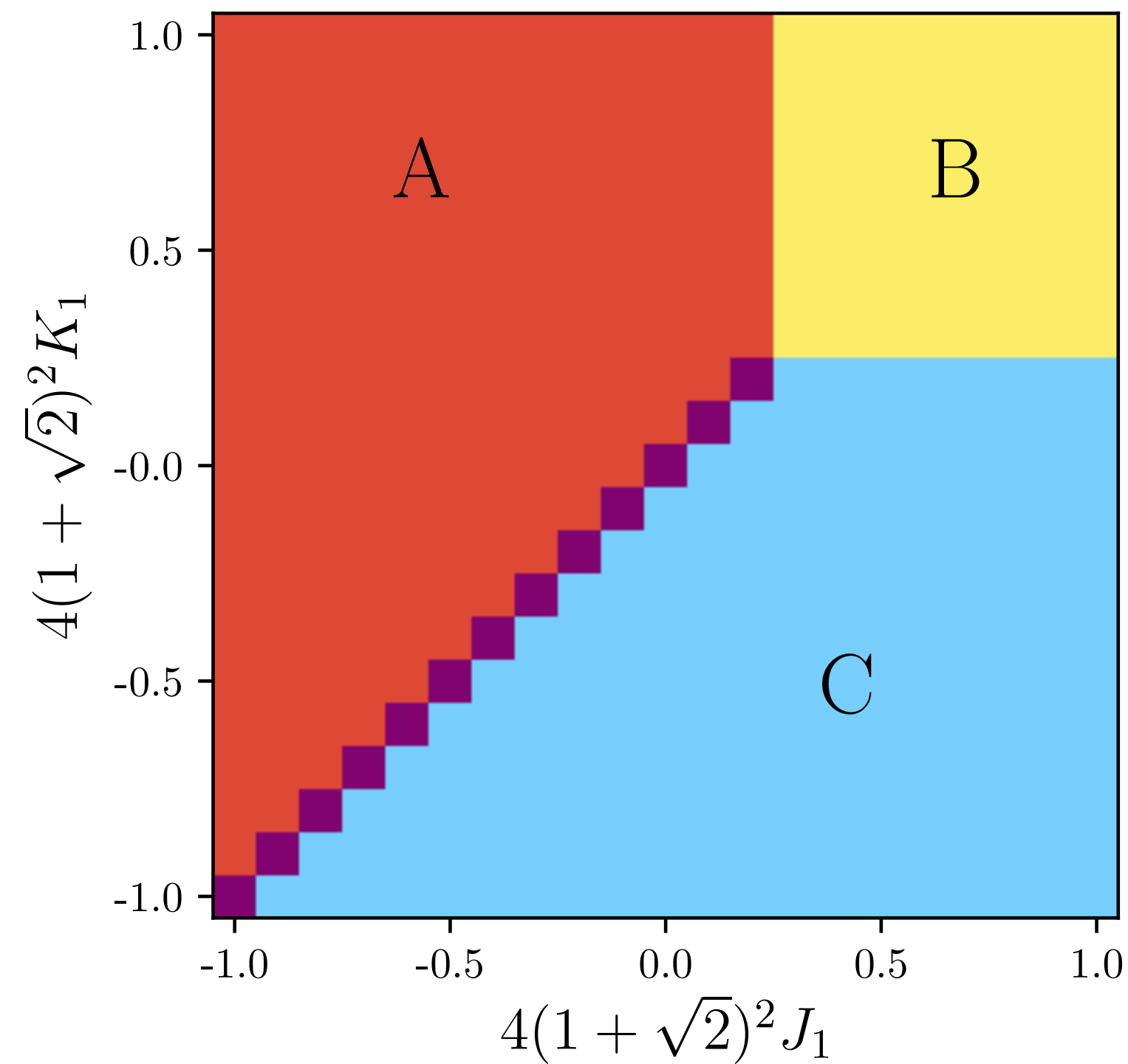
Phase A
 $(\pi, 0)$ stripe



Phase A
 $(0, \pi)$ stripe



Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



$$\langle B \rangle = 0$$

Confining phase.
 $SO(5)_f$ broken.
 Néel or
 valence bond solid
 order.

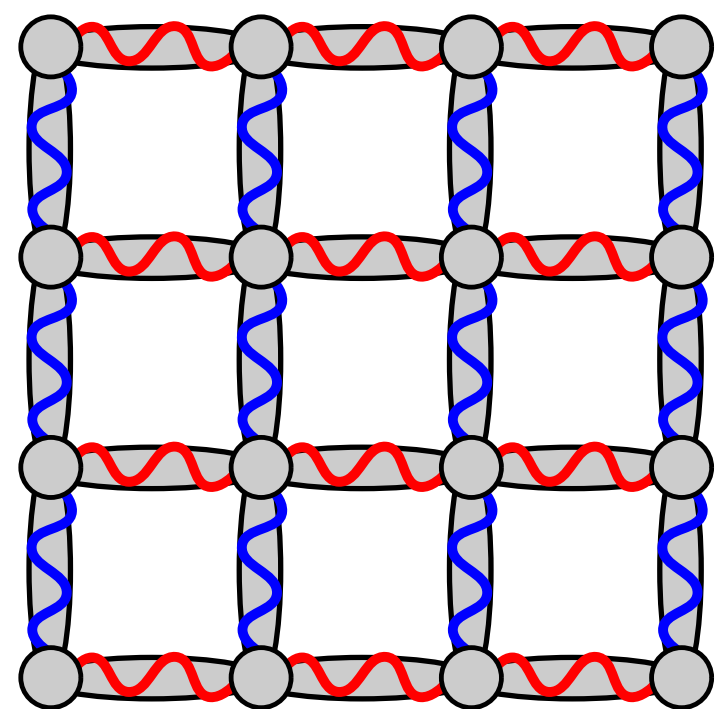
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r ← r_c

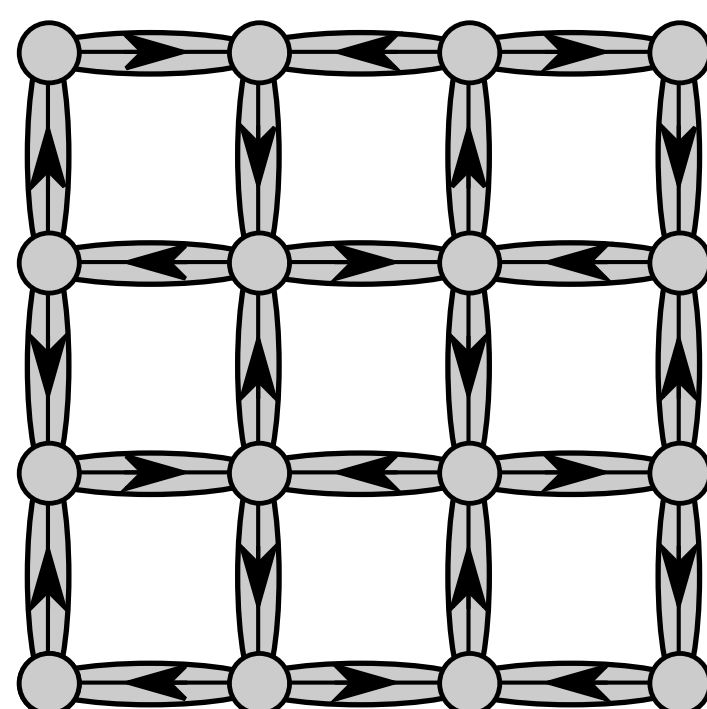
Phase B

d-wave SC



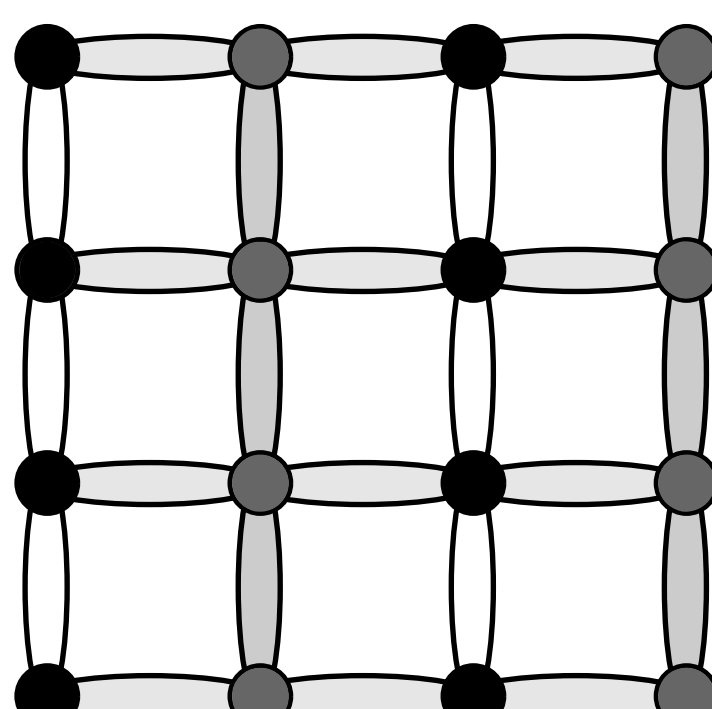
Phase C

d-density



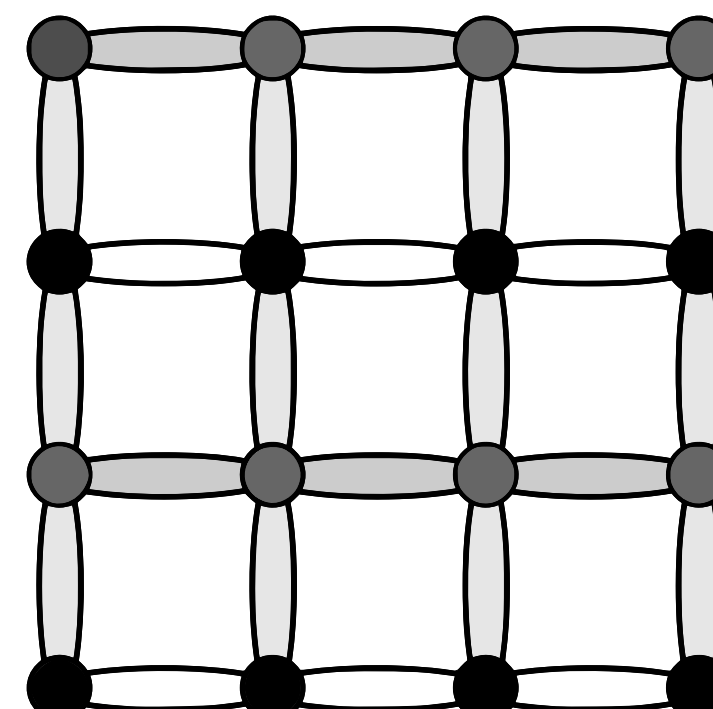
Phase A

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Phase A

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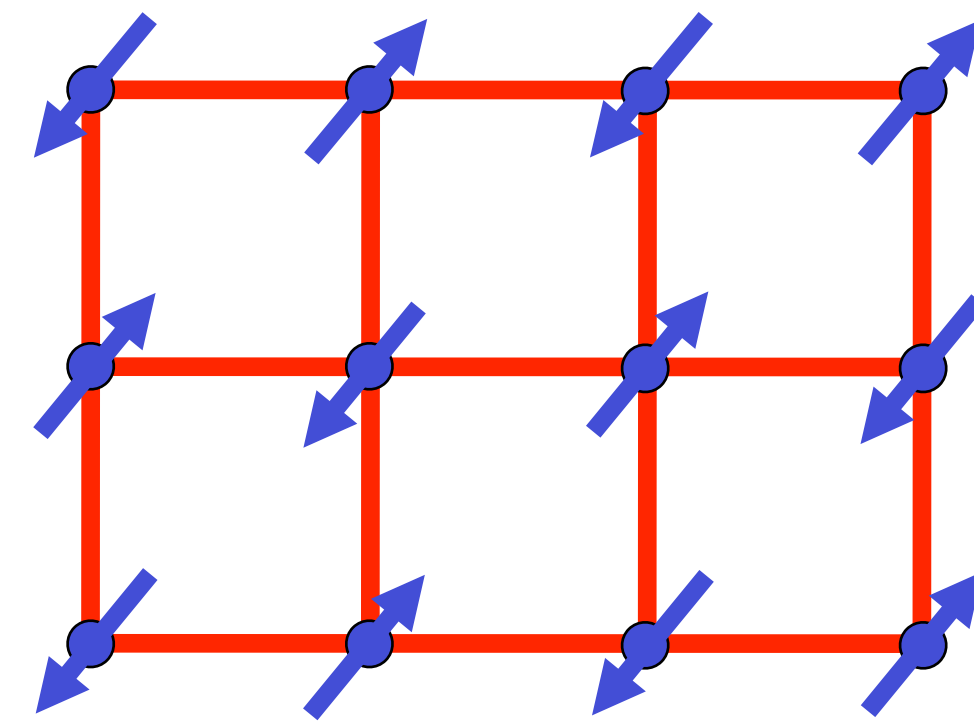
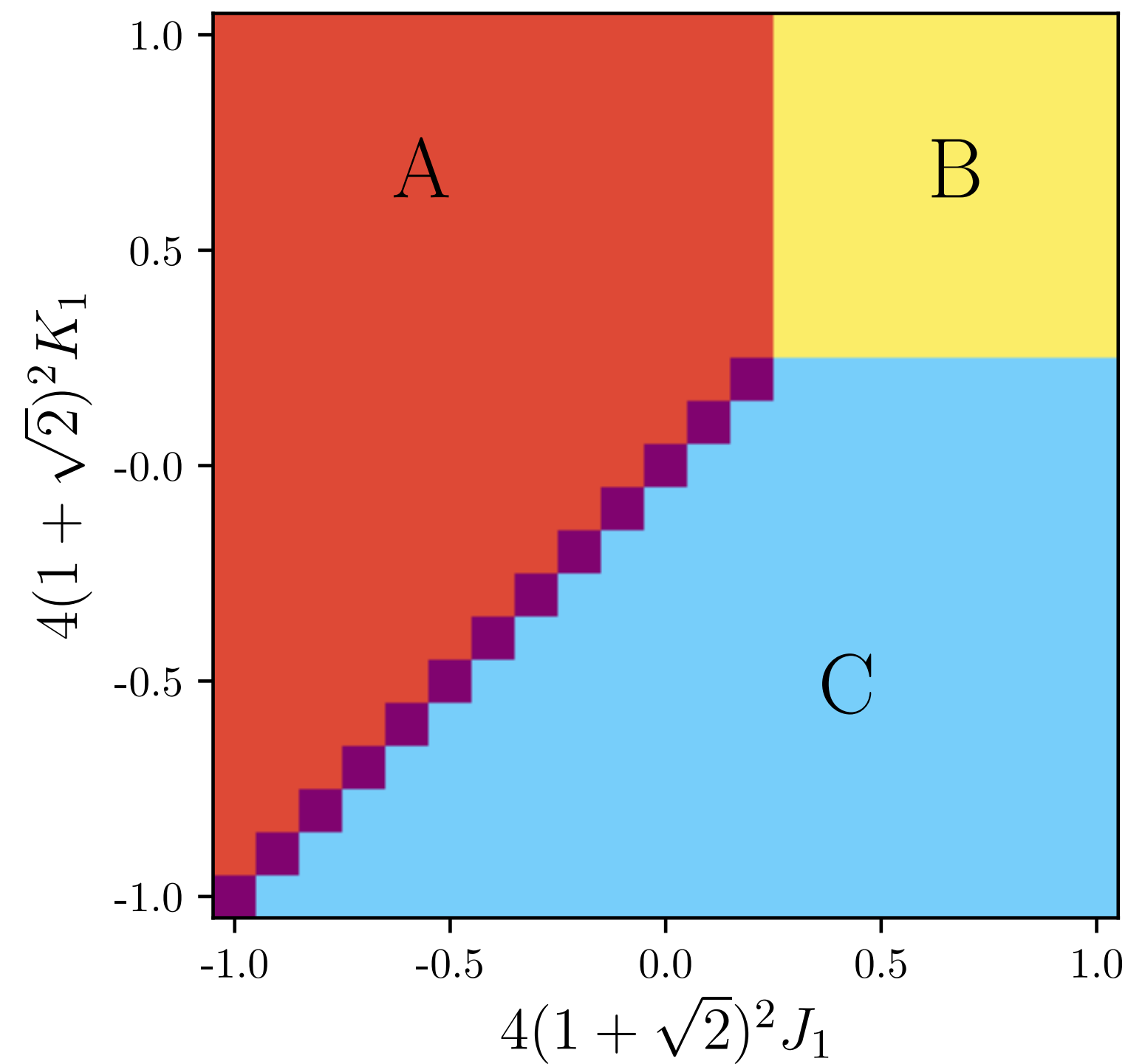


$$\mathcal{L} = i\bar{\Psi}_s D_\mu \Psi_s$$

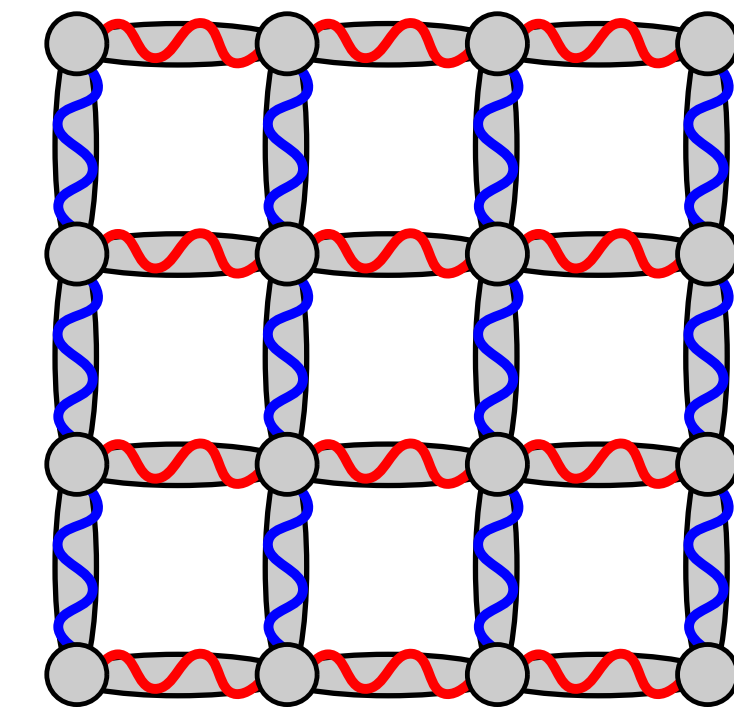
$$+ |D_\mu B_s|^2 + r|B_s|^2 \dots$$

Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

F. F. Assaad, M. Imada, and D. J. Scalapino, PRL 77, 4592 (1996)



Néel order



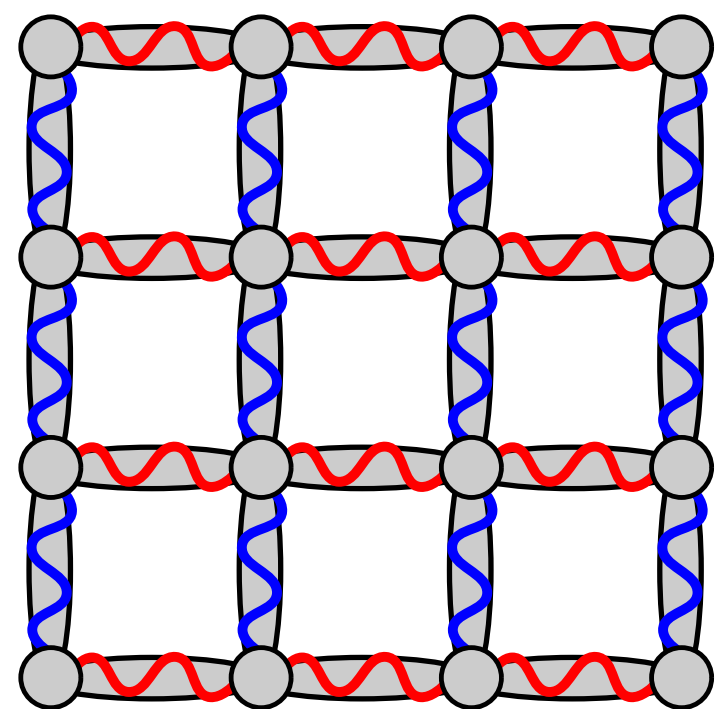
d -wave superconductor

0.3

W/t

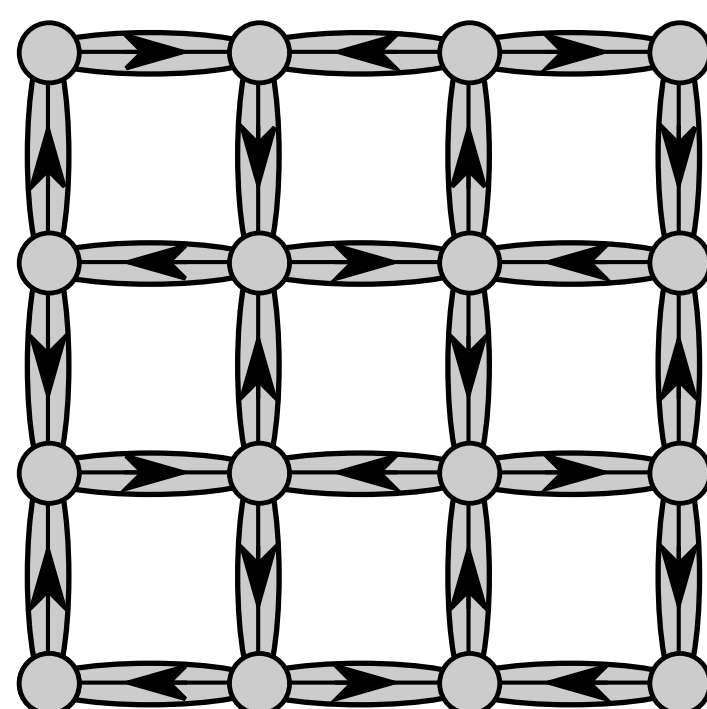
Phase B

d -wave SC



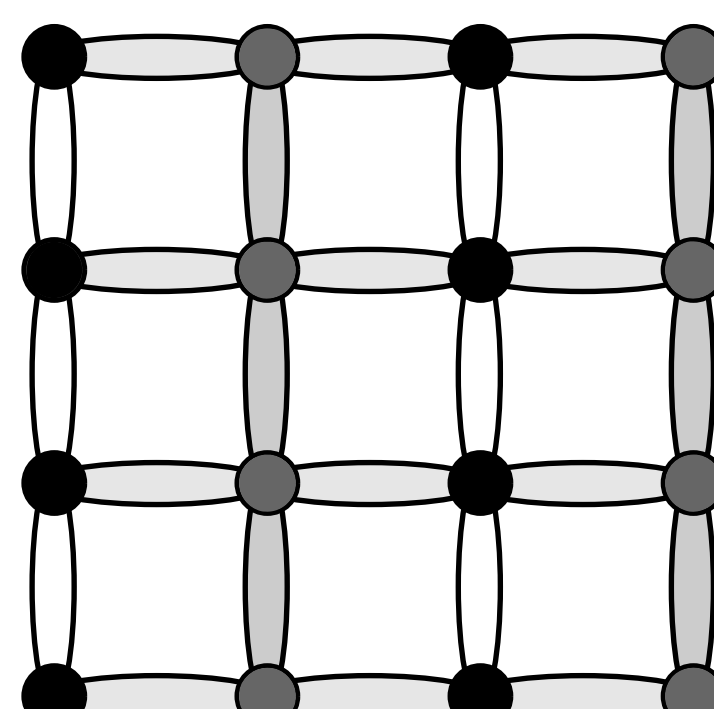
Phase C

d -density



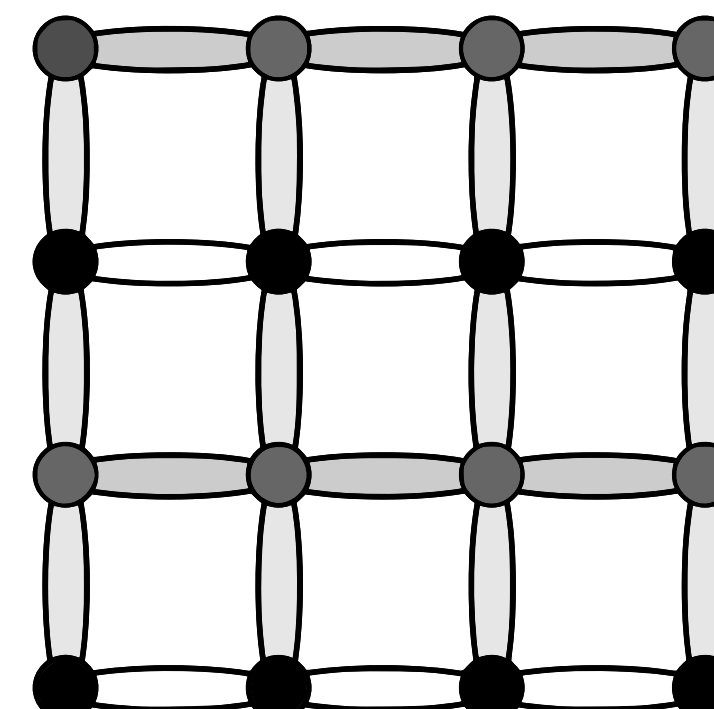
Phase A

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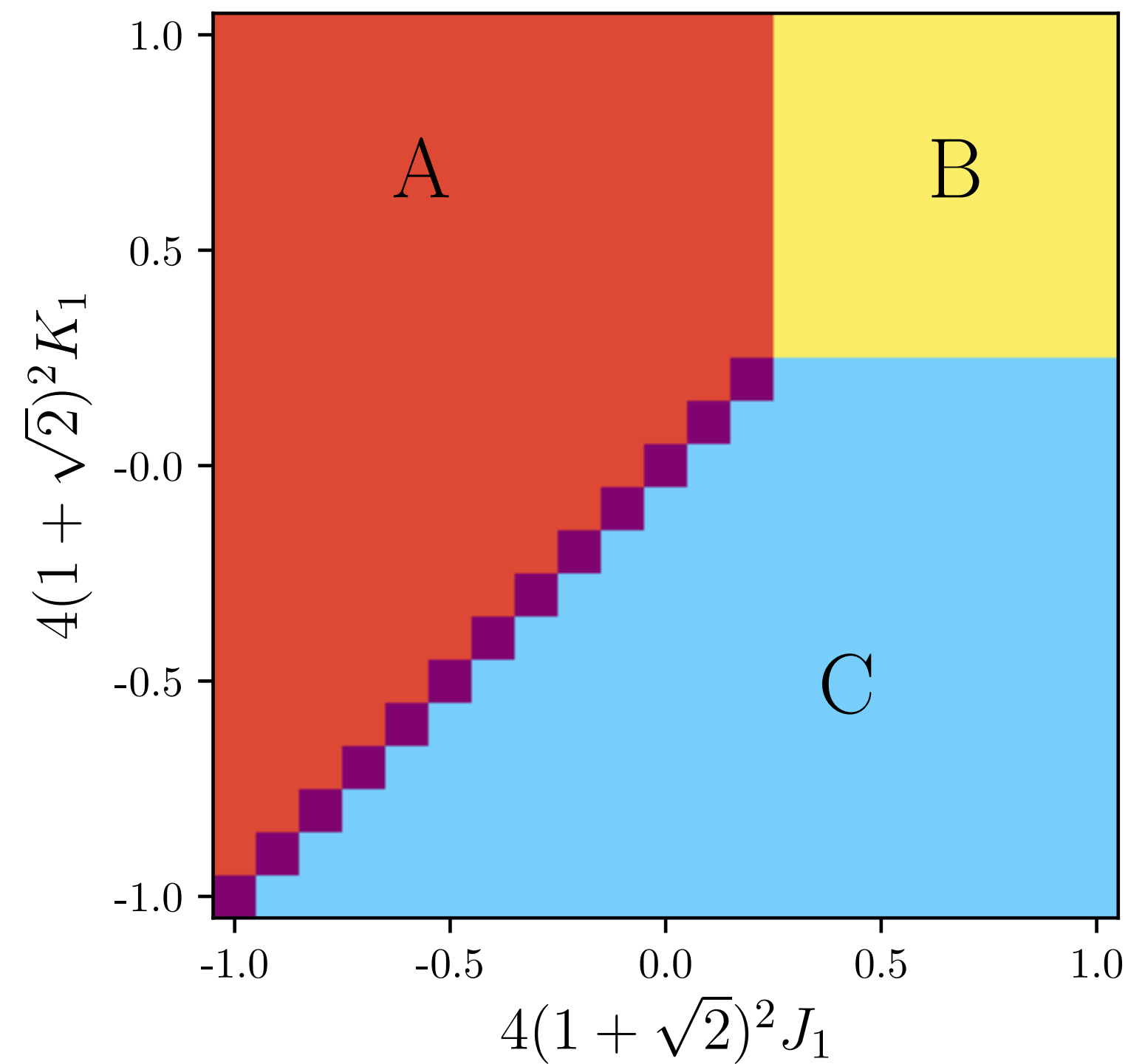
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Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



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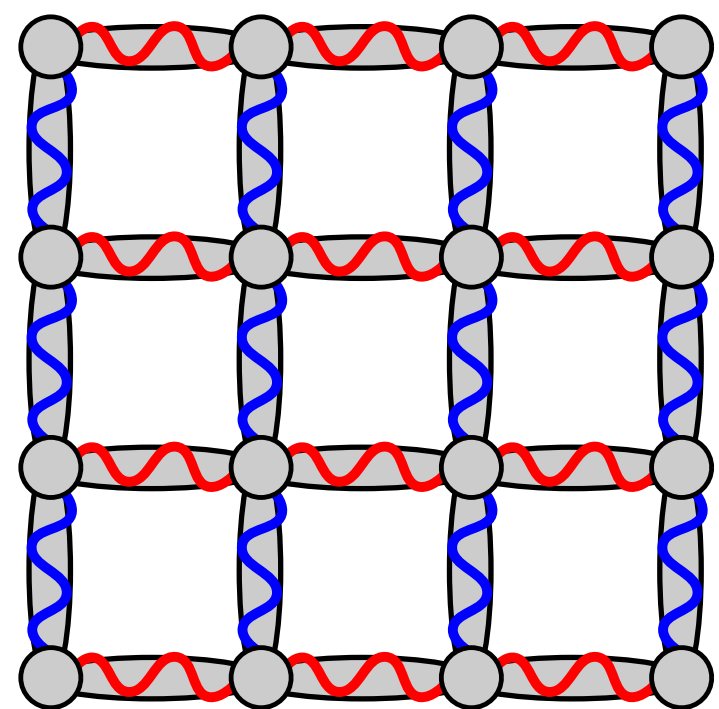
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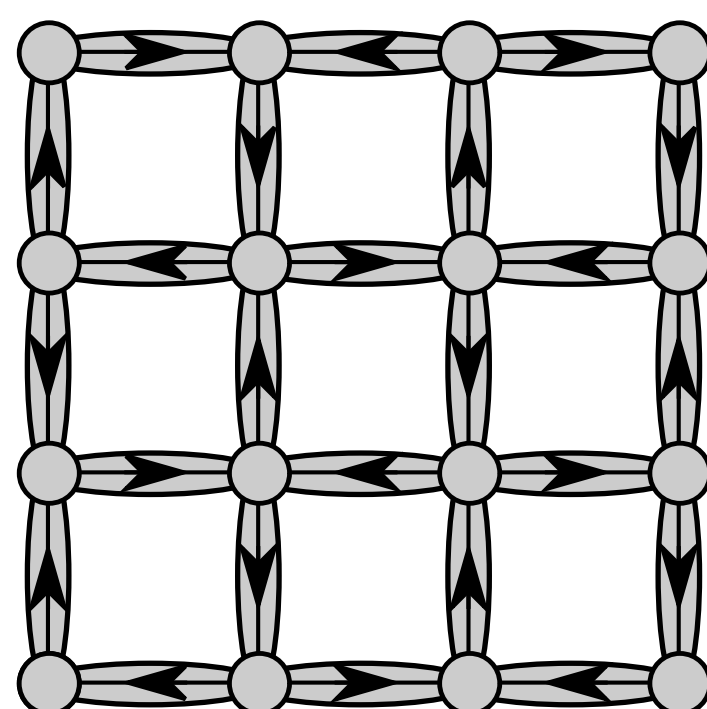
Phase B

d-wave SC



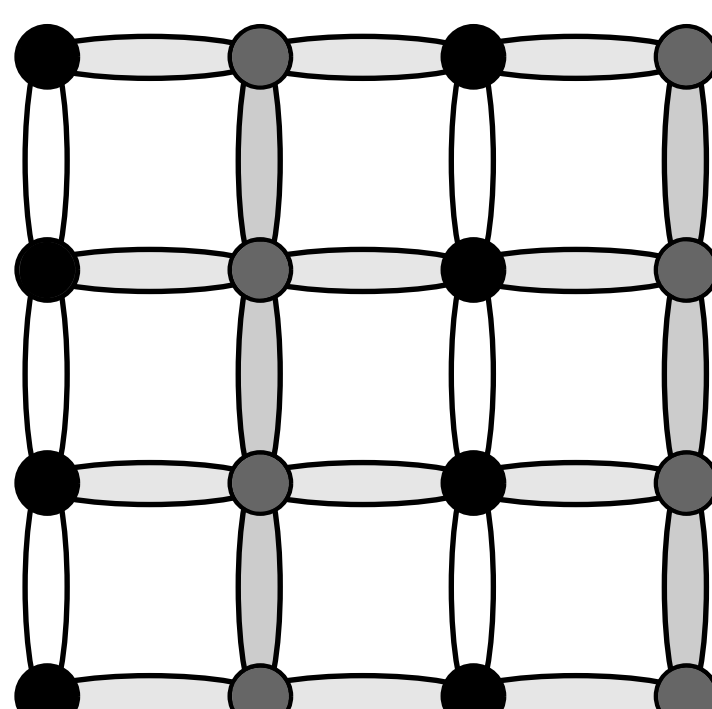
Phase C

d-density



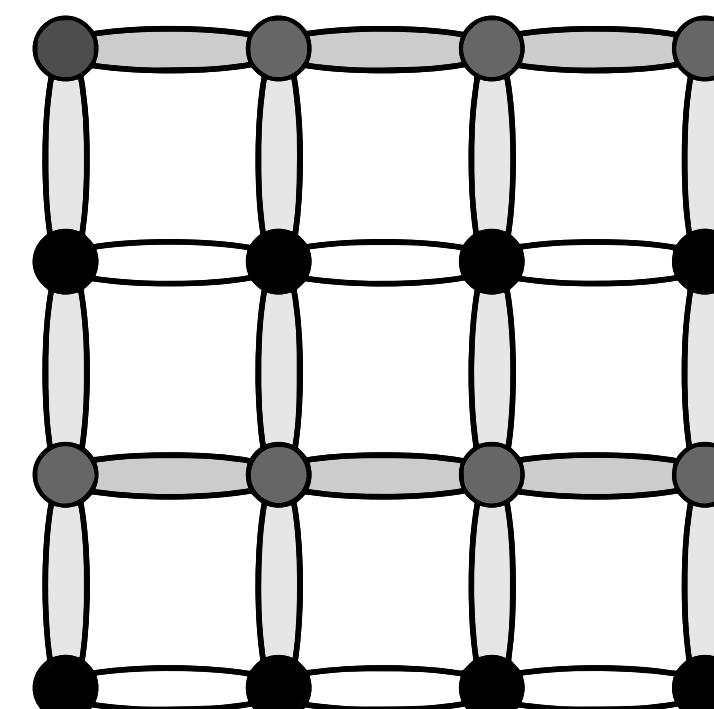
Phase A

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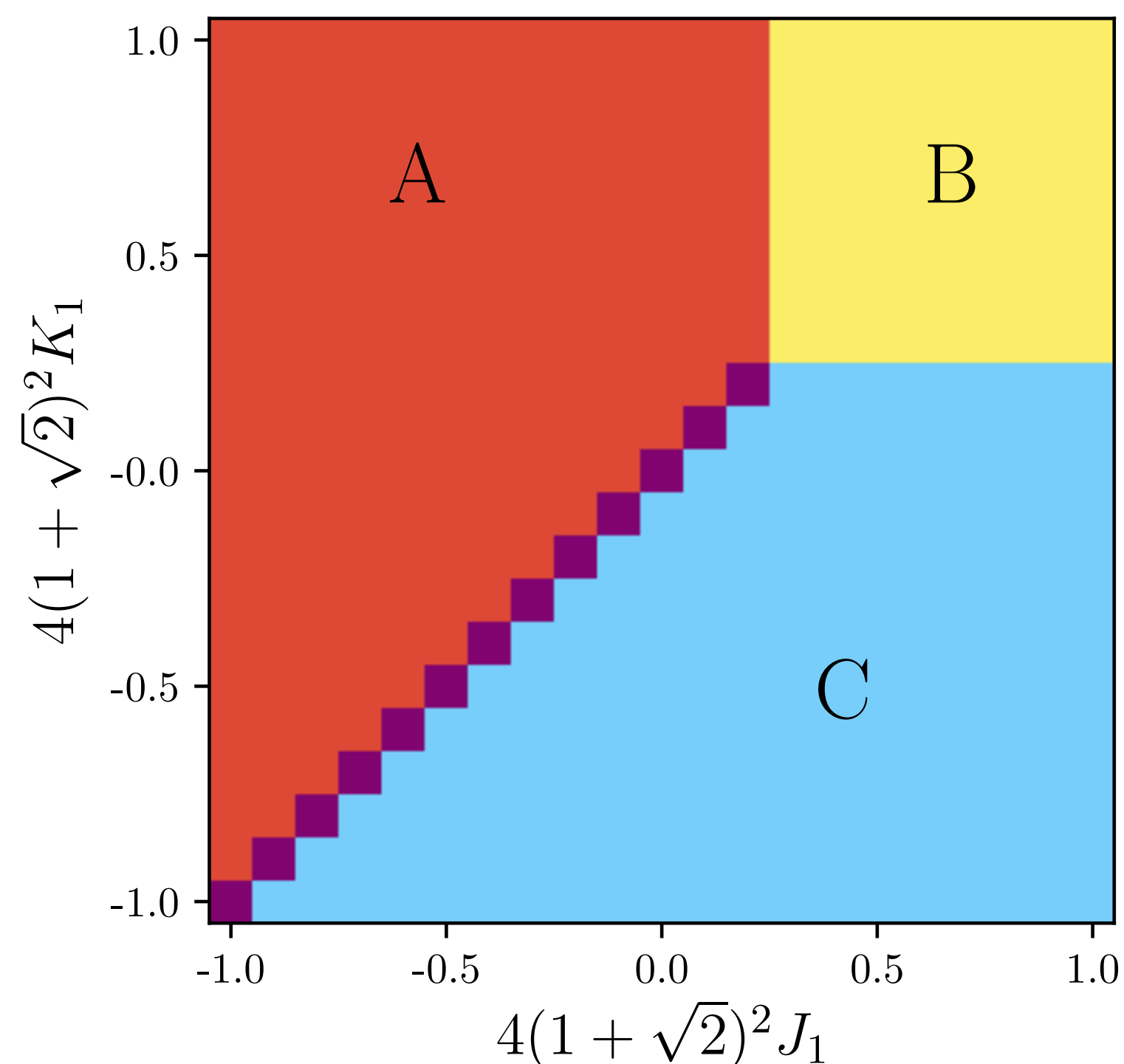
Phase A

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Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

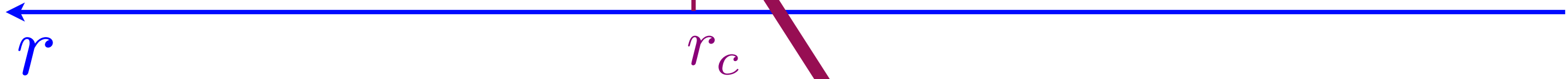


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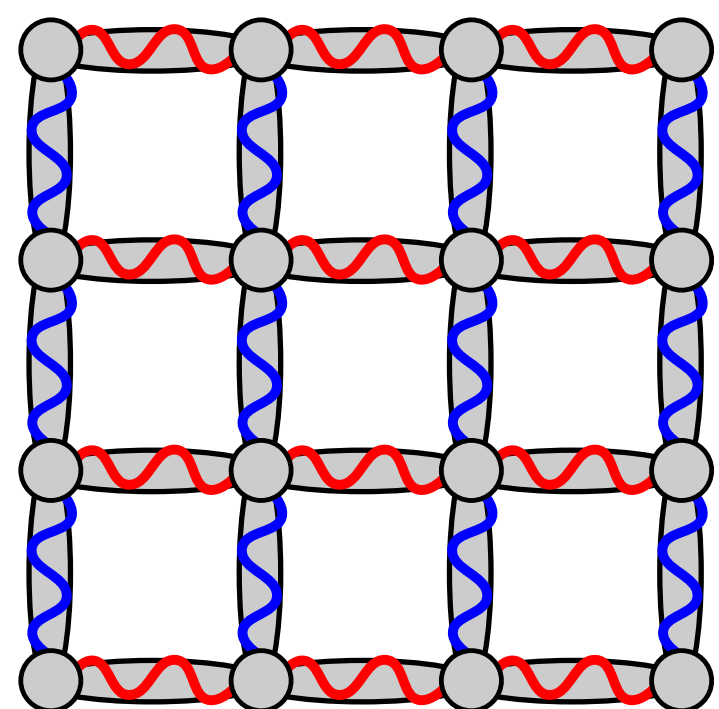


$$\mathcal{L} = i\bar{\Psi}_s D_\mu \Psi_s + |D_\mu B_s|^2 + r|B_s|^2 \dots$$

Possible CFT.
 DQCP with
 $SO(5)_f \times SO(5)_b$
 symmetry.

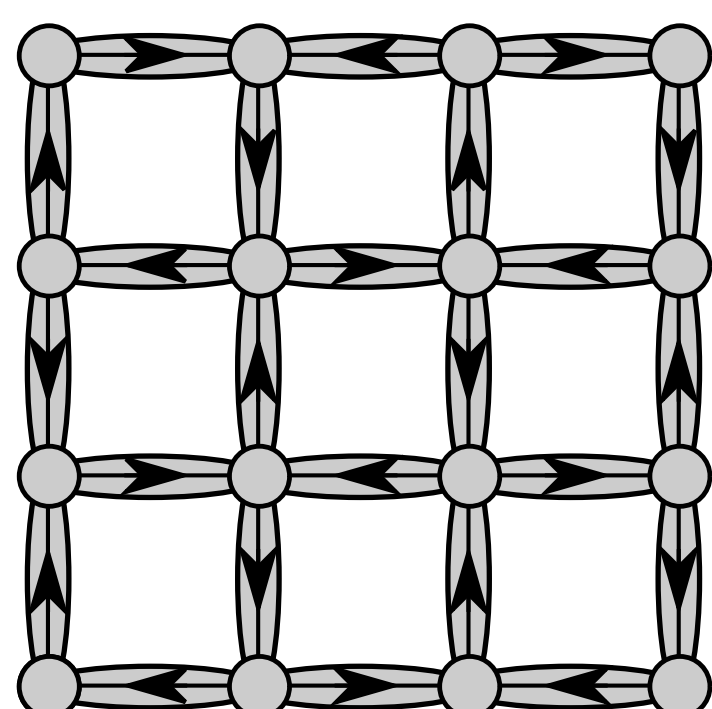
Phase B

d-wave SC



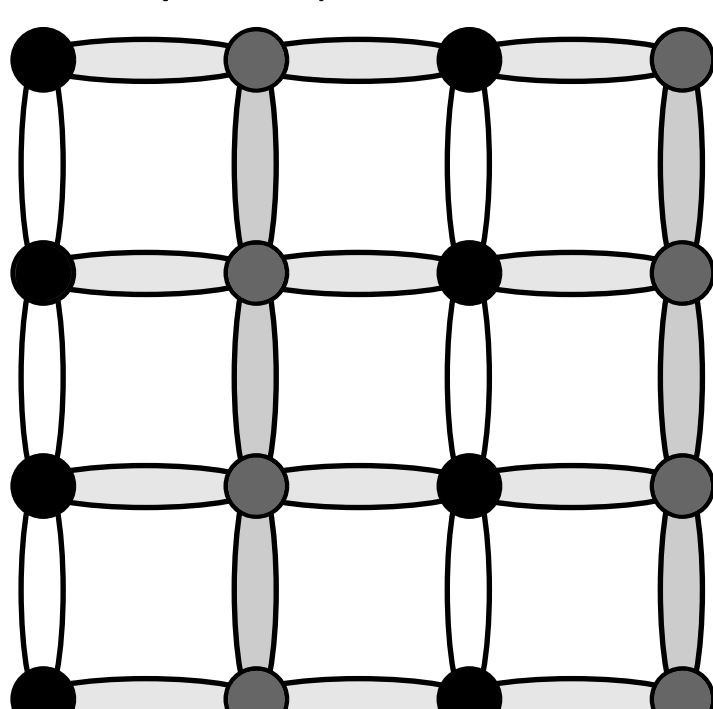
Phase C

d-density



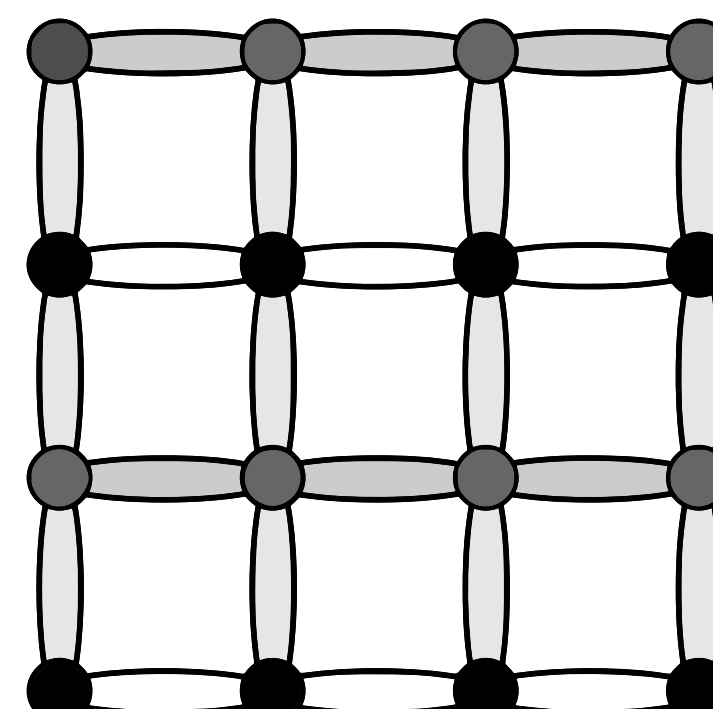
Phase A

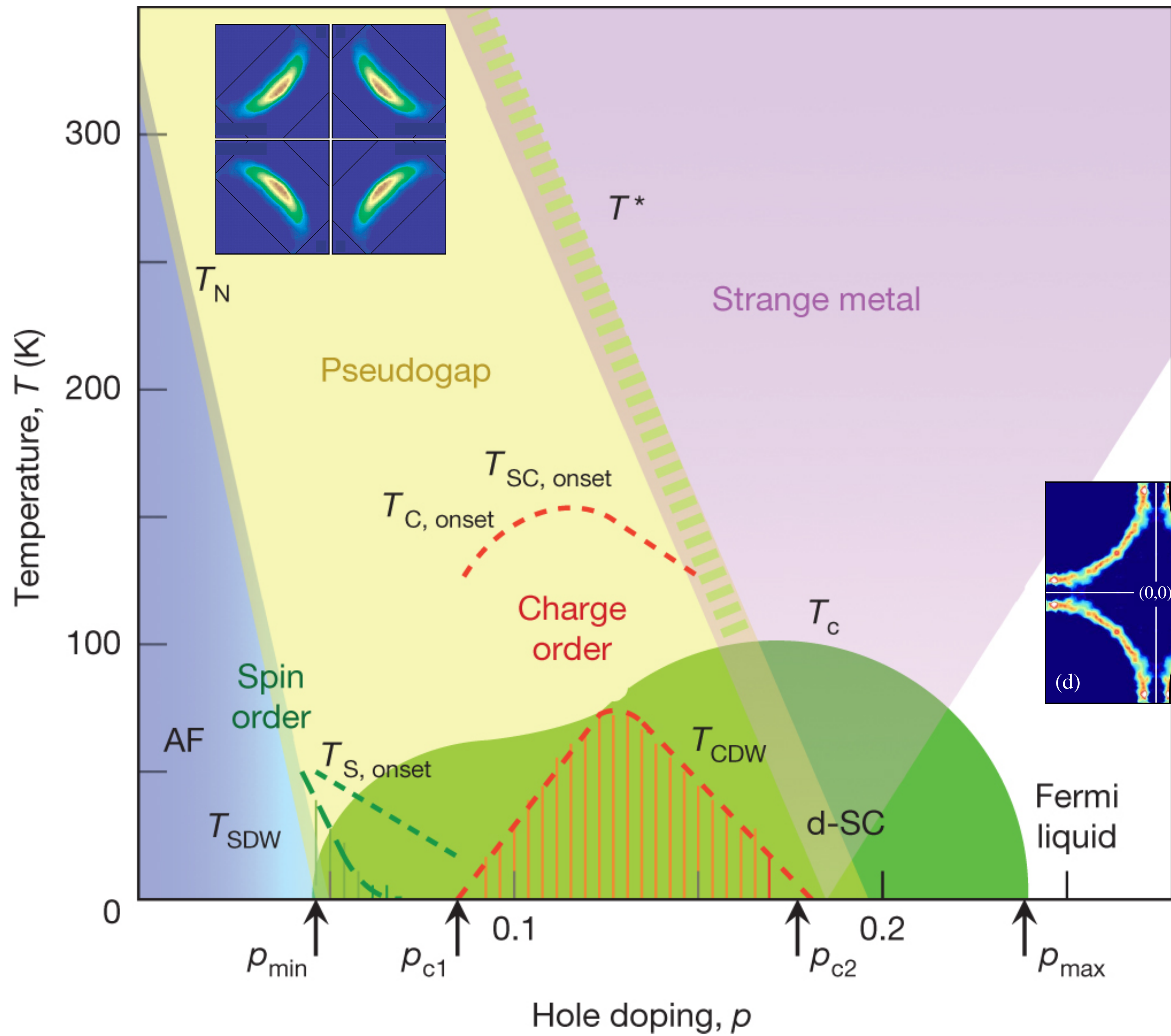
$(\pi, 0)$ stripe

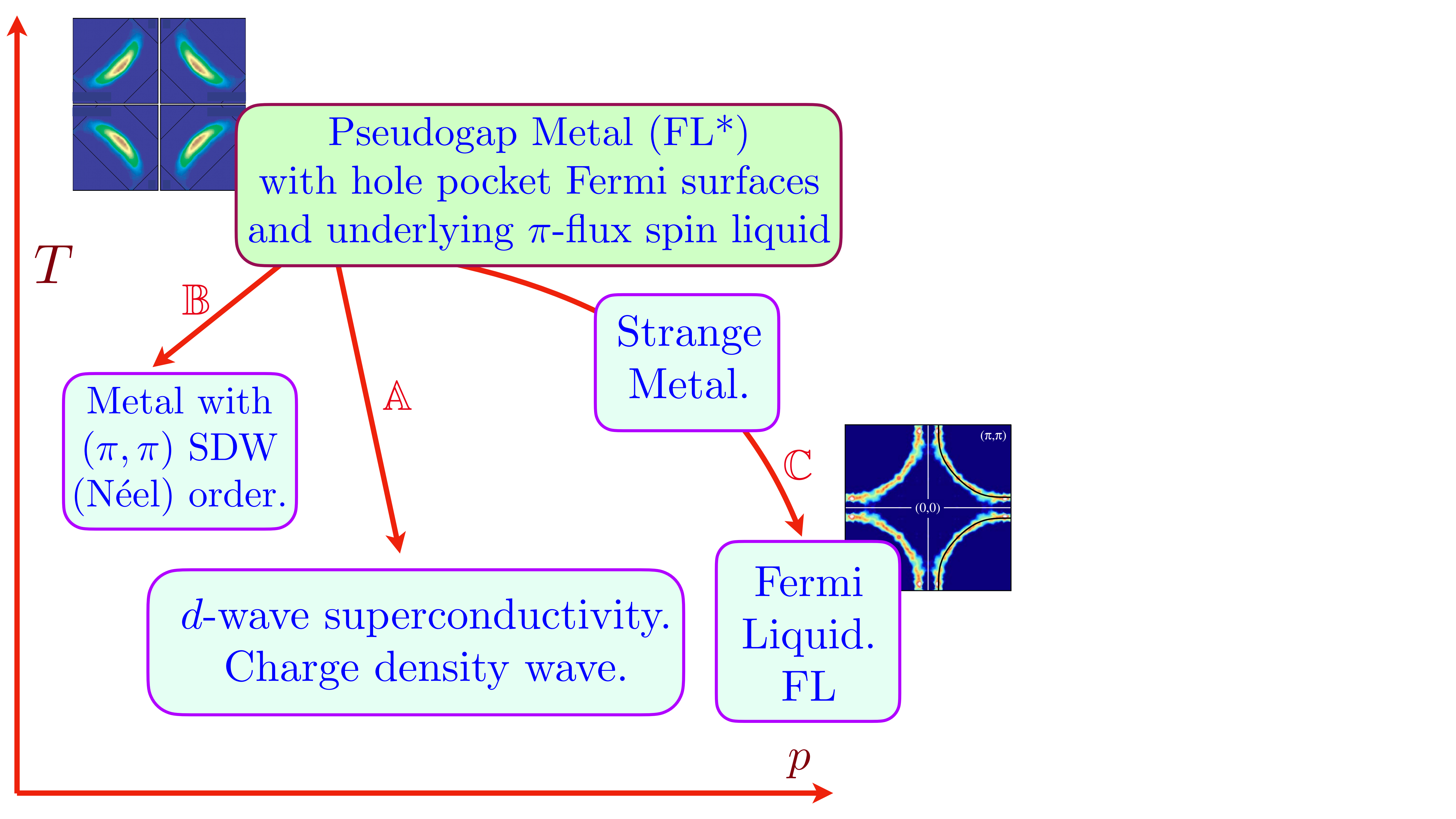


Phase A

$(0, \pi)$ stripe







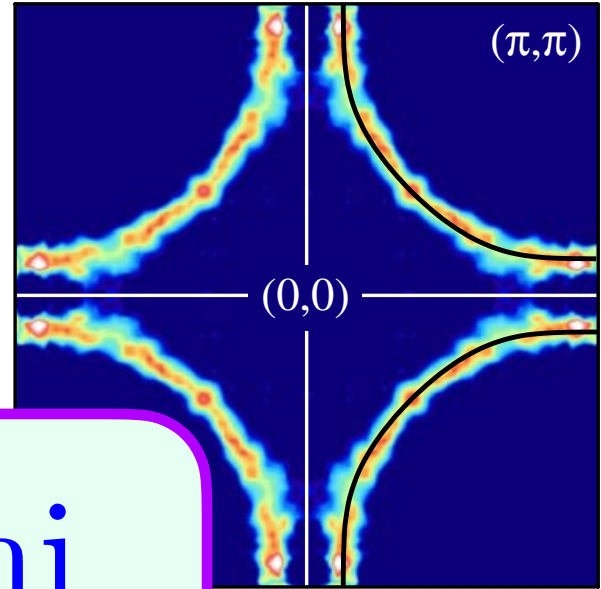
Pseudogap Metal (FL*)
with hole pocket Fermi surfaces
and underlying π -flux spin liquid

Strange
Metal.

Metal with
 (π, π) SDW
(Néel) order.

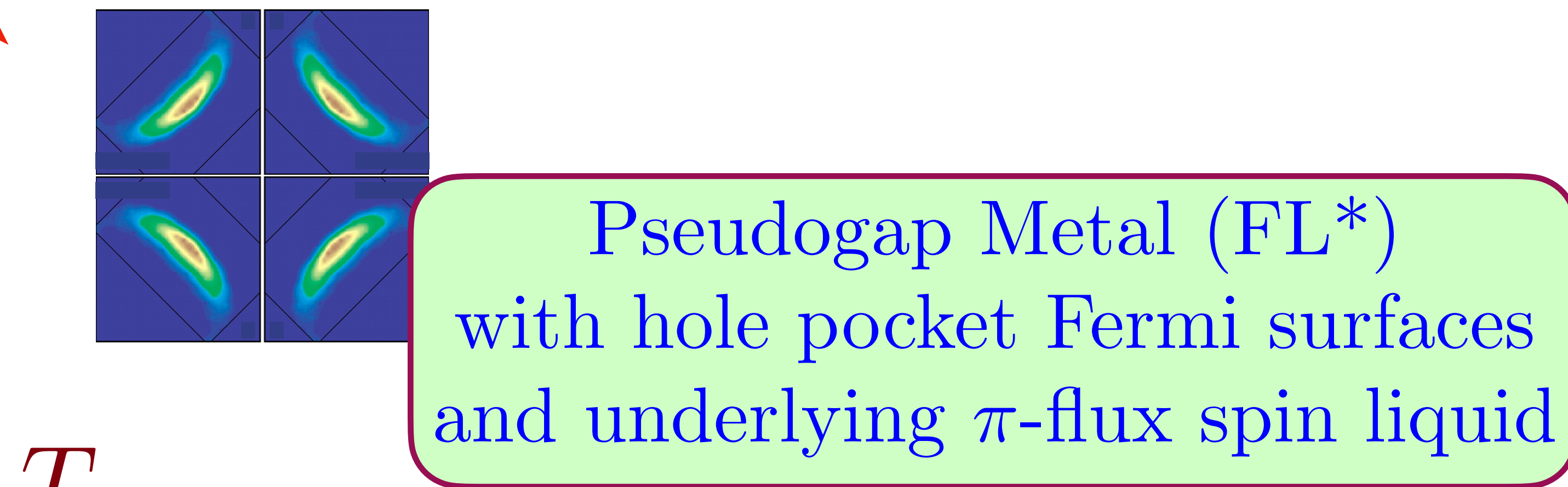
d -wave superconductivity.
Charge density wave.

Fermi
Liquid.
FL



Arrow B

Condensation of z_α in dual $\mathbb{C}P^1$
U(1) gauge theory.



T

B

Metal with
 (π, π) SDW
(Néel) order.

A

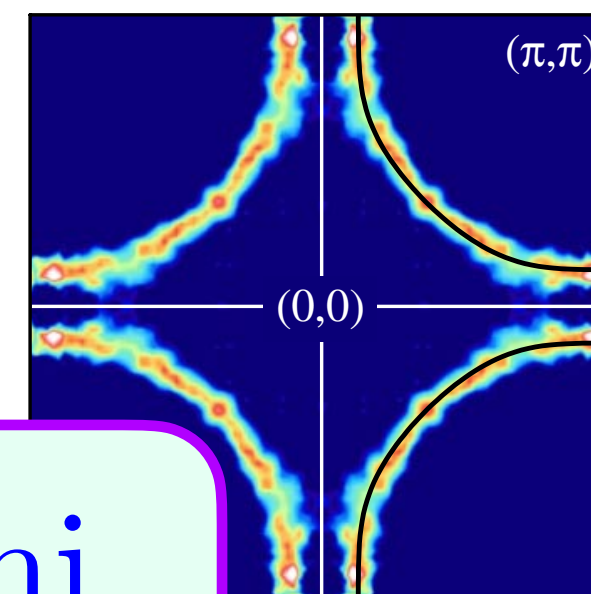
d -wave superconductivity.
Charge density wave.

Strange
Metal.

C

Fermi
Liquid.
FL

p



Arrow A

Condensation of B in SU(2) gauge theory.

Longer-range couplings in H_B can lead to charge order with other periods

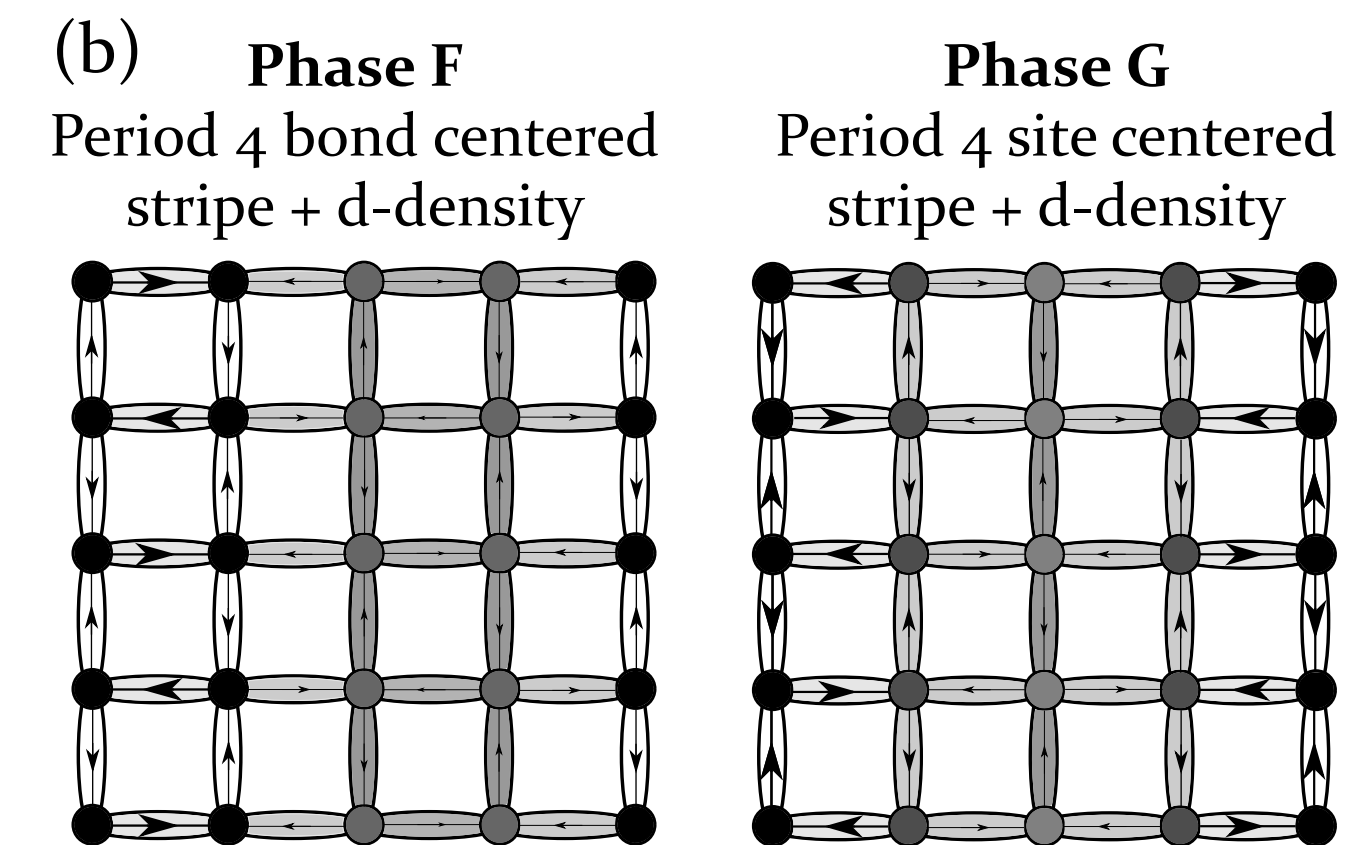
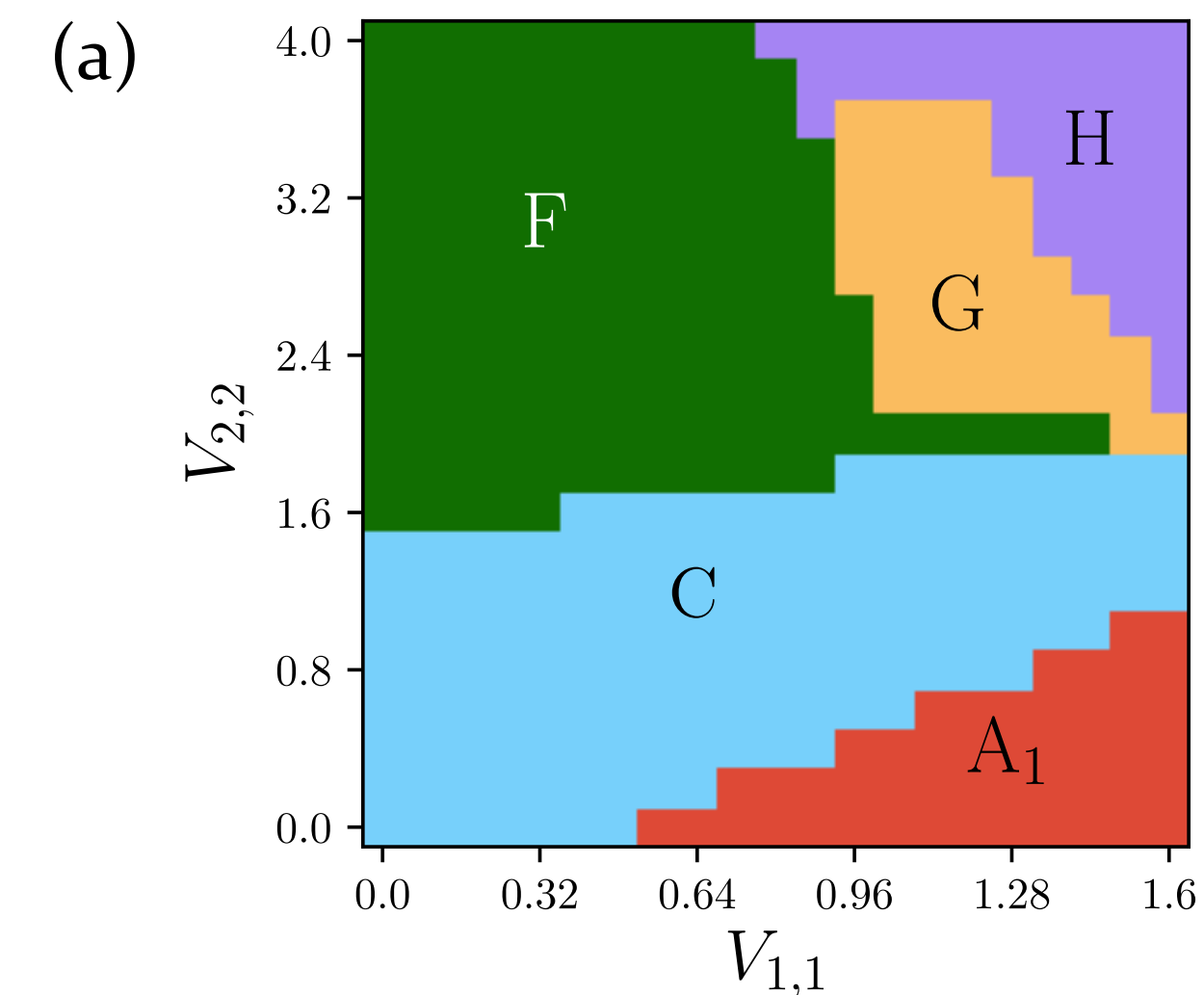
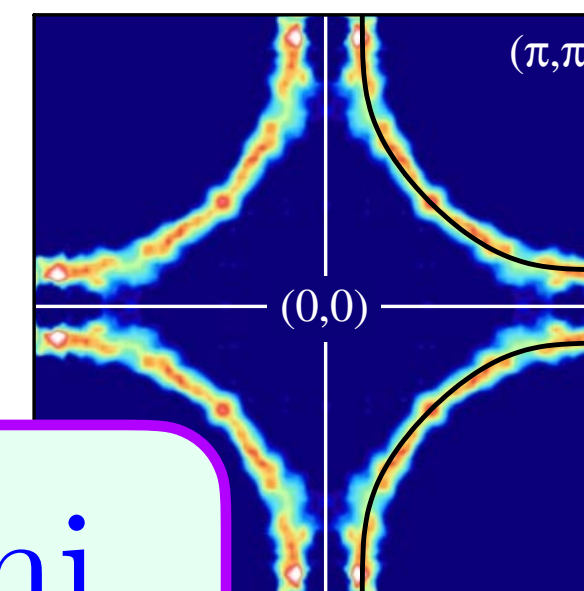
Pseudogap Metal (FL*)
with hole pocket Fermi surfaces
and underlying π -flux spin liquid

Strange
Metal.

Metal with
 (π, π) SDW
(Néel) order.

d -wave superconductivity.
Charge density wave.

Fermi
Liquid.
FL



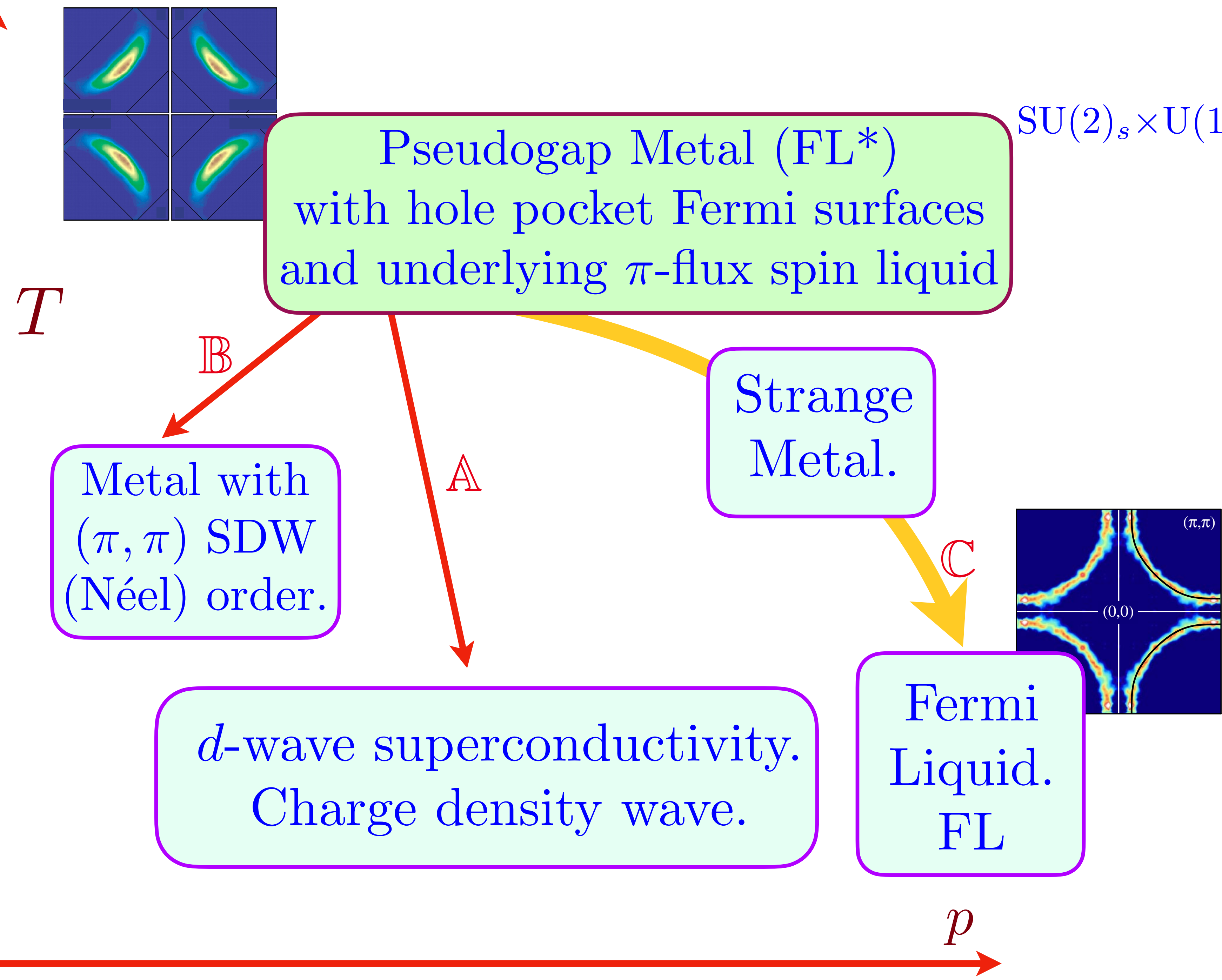
T

B

A

C

p



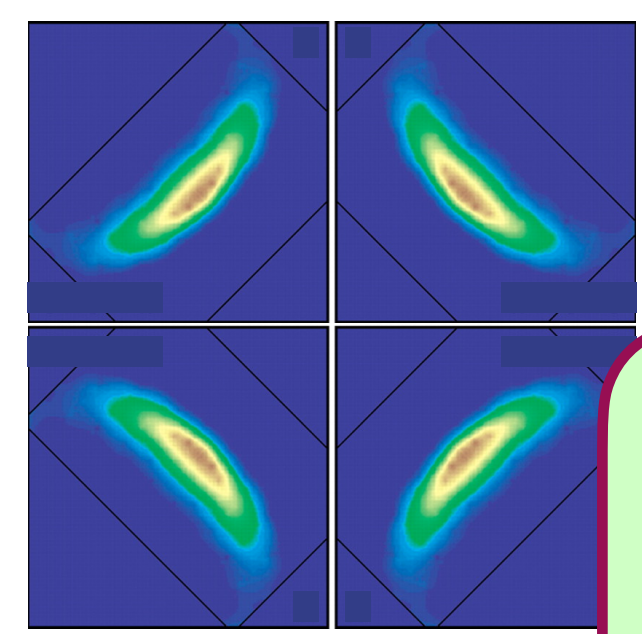
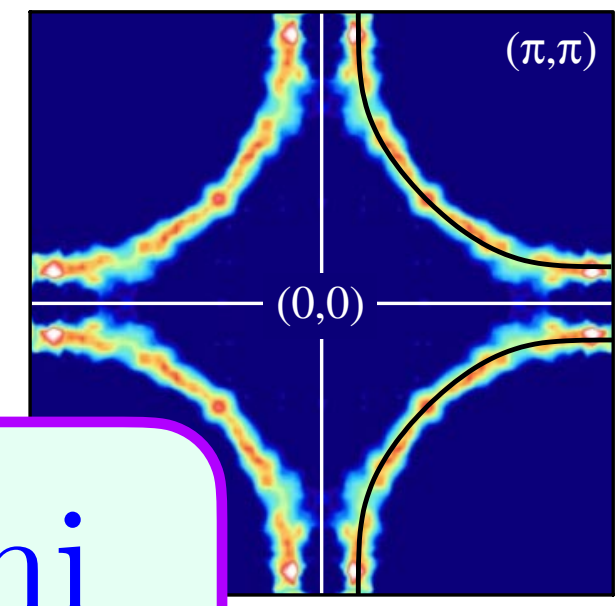
Pseudogap Metal (FL*)
with hole pocket Fermi surfaces
and underlying π -flux spin liquid

Strange
Metal.

Metal with
 (π, π) SDW
(Néel) order.

d -wave superconductivity.
Charge density wave.

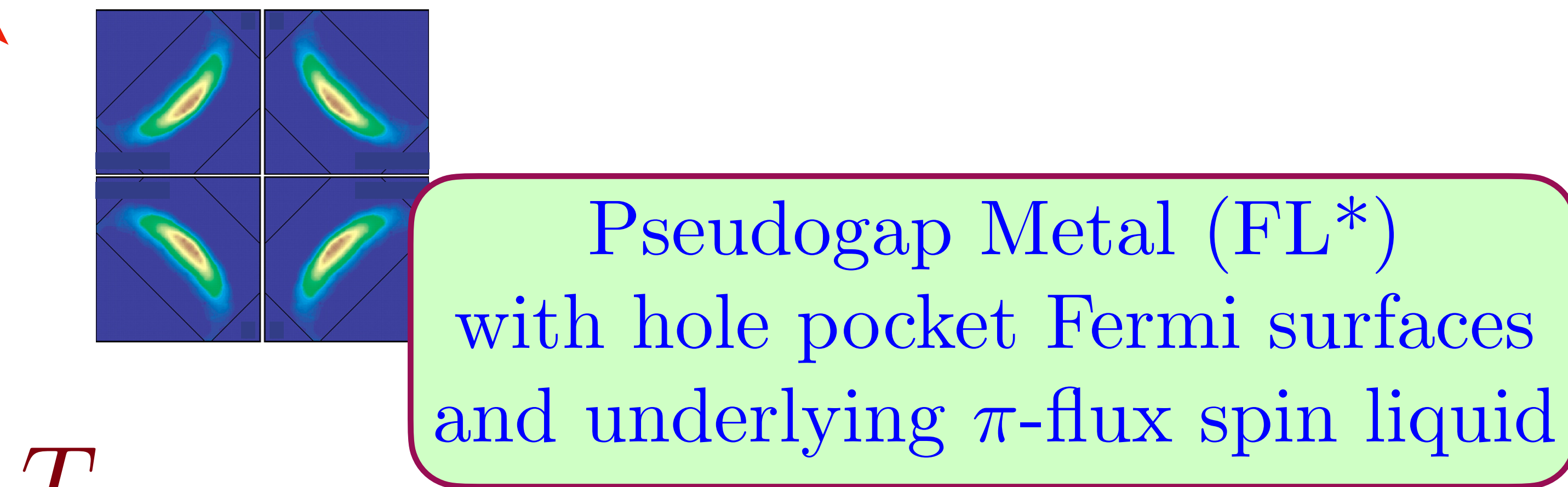
Fermi
Liquid.
FL



Arrow A

Condensation of B in SU(2) gauge theory.

Longer-range couplings in H_B can lead to charge order with other periods



T

Pseudogap Metal (FL*)
with hole pocket Fermi surfaces
and underlying π -flux spin liquid

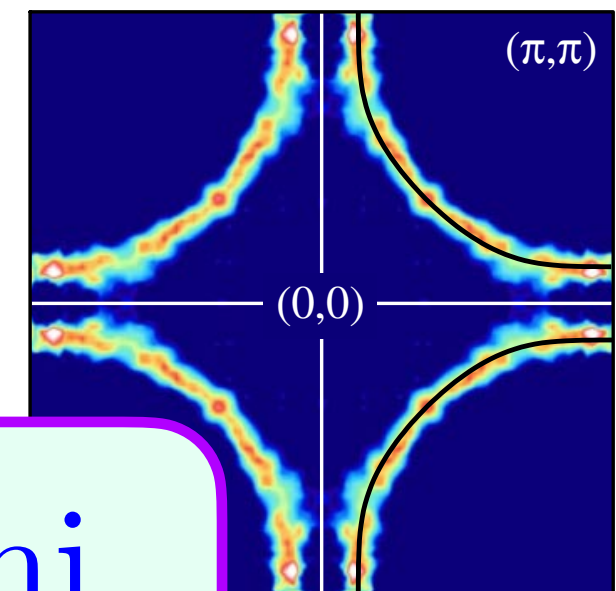
B

Strange
Metal.

A

Metal with
 (π, π) SDW
(Néel) order.

C



d -wave superconductivity.
Charge density wave.

Fermi
Liquid.
FL

p

